



Erratum to: “The Schrödinger–Maxwell system with Dirac mass” [Ann. Inst. H. Poincaré Anal. Non Linéaire 24 (5) (2007) 773–793] [☆]

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Abstract

We correct the proof of [G.M. Coclite, H. Holden, The Schrödinger–Maxwell system with Dirac mass, *Ann. Inst. H. Poincaré Anal. Non Linéaire* 24 (5) (2007) 773–793, Lemma 4.1].

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1. Introduction

The proof of [1, Lemma 4.1] is incorrect. We here present a new proof. We employ the notation and assumptions of [1].

Lemma 4.1. *Assume α and β are positive constants. There exists a subsequence $\{v_{n_k}\}_{k \in \mathbb{N}}$ and a map $v \in H^2(\Omega) \cap H_0^1(\Omega)$ such that*

$$v_{n_k} \rightharpoonup v \quad \text{weakly in } H^2(\Omega) \cap H_0^1(\Omega). \quad (1)$$

In particular,

$$v_{n_k} \rightarrow v \quad \text{uniformly in } \Omega. \quad (2)$$

Proof. We split the proof in two steps.

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Step 1. We claim that the following inequality

$$J_n(v_n) = \min_{v \in H_0^1(\Omega)} J_n(v) \leq 0 \quad (3)$$

holds for infinitely many $n \in \mathbb{N}$.

Let φ_0 be the normalized positive first eigenfunction of $-\Delta$ on Ω , namely φ_0 is the unique smooth map satisfying the following conditions

$$\begin{cases} -\Delta\varphi_0 = \omega_0\varphi_0, & \text{in } \Omega, \\ \varphi_0 > 0, & \text{in } \Omega, \\ \varphi_0 = 0, & \text{on } \partial\Omega, \\ \|\varphi_0\|_{L^2(\Omega)} = 1. \end{cases} \quad (4)$$

Since

$$\int_{\Omega} |\nabla\varphi_0(x)|^2 dx = \omega_0 \int_{\Omega} \varphi_0^2(x) dx = \omega_0,$$

we find, by evaluating J_n in¹ $\lambda \operatorname{sign}(v_{n-1}(x_0))\varphi_0$, $\lambda > 0$, that (writing $v_{n-1} = \operatorname{sign}(v_{n-1}(x_0))$)

$$\begin{aligned} J_n(\lambda v_{n-1}\varphi_0) &= \frac{\lambda^2}{2} \int_{\Omega} |\nabla\varphi_0|^2 dx + \lambda^4 \frac{\alpha}{4} \int_{\Omega \times \Omega} G(x, y)\varphi_0^2(y)\varphi_0^2(x) dx dy \\ &\quad - |v_{n-1}(x_0)|\lambda^3 \frac{\alpha}{\beta} \int_{\Omega \times \Omega} G(x, y)G(y, x_0)\varphi_0(y)\varphi_0^2(x) dx dy \\ &\quad + v_{n-1}^2(x_0)\lambda^2 \frac{\alpha}{\beta^2} \int_{\Omega \times \Omega} G(x, y)G(y, x_0)G(x, x_0)\varphi_0(y)\varphi_0(x) dx dy \\ &\quad + v_{n-1}^2(x_0)\lambda^2 \frac{\alpha}{2\beta^2} \int_{\Omega \times \Omega} G(x, y)G^2(y, x_0)\varphi_0^2(x) dx dy \\ &\quad - |v_{n-1}^3(x_0)|\lambda \frac{\alpha}{\beta^3} \int_{\Omega \times \Omega} G(x, y)G^2(y, x_0)G(x, x_0)\varphi_0(x) dx dy \\ &\quad - \frac{\omega}{2}\lambda^2 \int_{\Omega} \varphi_0^2(x) dx - |v_{n-1}(x_0)|\lambda \frac{\omega}{\beta} \int_{\Omega} G(x, x_0)\varphi_0(x) dx \\ &= \lambda^2 \frac{\omega_0 - \omega}{2} + \lambda^4 \kappa_1 - |v_{n-1}(x_0)|\lambda^3 \kappa_2 + v_{n-1}^2(x_0)\lambda^2 \kappa_3 - |v_{n-1}^3(x_0)|\lambda \kappa_4 - |v_{n-1}(x_0)|\lambda \kappa_5. \end{aligned}$$

Due to the positivity and the boundedness of φ_0 , we have $\kappa_1, \dots, \kappa_5 > 0$, hence

$$J_n(\lambda v_{n-1}\varphi_0) \leq \lambda^2 \frac{\omega_0 - \omega}{2} + \lambda^4 \kappa_1 + v_{n-1}^2(x_0)\lambda^2 \kappa_3 - |v_{n-1}(x_0)|\lambda \kappa_5. \quad (5)$$

We have the following two cases.

(i) If

$$\liminf_n |v_{n-1}(x_0)| = 0,$$

then there exists an n_0 such that, by passing to a subsequence and using, e.g., [1, (2.6)], we find

$$J_n(v_n) = \min_{v \in H_0^1(\Omega)} J_n(v) \leq \min_{\lambda > 0} J_n(\lambda v_{n-1}\varphi_0) \leq \min_{\lambda > 0} \left(\lambda^2 \frac{\omega_0 - \omega}{2} + \lambda^4 \kappa_1 \right) < 0, \quad n > n_0. \quad (6)$$

¹ The point x_0 is where the point interaction is located, cf. [1].

(ii) If

$$0 < \liminf_n |v_{n-1}(x_0)| \leq \infty,$$

then there exists n_0 and $c_1 \in (0, |v_{n-1}(x_0)|)$ for $n > n_0$ such that

$$\begin{aligned} J_n(v_n) &= \min_{v \in H_0^1(\Omega)} J_n(v) \leq \min_{\lambda > 0} J_n(\lambda v_{n-1} \varphi_0) \\ &\leq \min_{\lambda > 0} \left(\lambda^2 \frac{\omega_0 - \omega}{2} + \lambda^4 \kappa_1 + |v_{n-1}(x_0)|^2 \lambda^2 \kappa_3 - c_1 \lambda \kappa_5 \right) < 0, \quad n > n_0. \end{aligned} \quad (7)$$

Clearly, (6) and (7) prove (3). So *Step 1* is concluded.

Step 2. We prove (1) and (2). Clearly it suffices to prove that

$$\text{the sequence } \{v_n\}_{n \in \mathbb{N}} \text{ is bounded in } H^2(\Omega) \cap H_0^1(\Omega). \quad (8)$$

Due to [1, Lemmas 2.1 and 3.5] we only have to show that

$$\text{the sequence } \{v_n\}_{n \in \mathbb{N}} \text{ is bounded in } H_0^1(\Omega), \quad (9)$$

$$\text{the sequence } \{v_n(x_0)\}_{n \in \mathbb{N}} \text{ is bounded in } \mathbb{R}. \quad (10)$$

If, by contradiction, (9) does not hold, we have that

$$\limsup_n \|v_n\|_{H_0^1(\Omega)} = \infty. \quad (11)$$

Therefore, [1, Lemma 3.3] and a diagonal argument guarantee

$$\limsup_n J_n(v_n) = \infty. \quad (12)$$

Since, by construction, v_n is a minimizer for J_n , Eq. (3) says

$$J_n(v_n) = \min_{f \in H_0^1(\Omega)} J_n(f) \leq 0, \quad n \in \mathbb{N}, \quad (13)$$

which contradicts (12). This proves (9).

We conclude by proving (10). Assume by contradiction that $\{v_n(x_0)\}_{n \in \mathbb{N}}$ is not bounded, namely (passing to a subsequence)

$$\lim_n |v_n(x_0)| = \infty. \quad (14)$$

Multiplying [1, (2.21)] by u_n , which is defined by [1, (2.20)], and integrating over Ω we get

$$\int_{\Omega} |\nabla v_n(x)|^2 dx + \frac{v_{n-1}(x_0)}{\beta} \int_{\Omega} G(x, x_0) \Delta v_n(x) dx + \alpha \int_{\Omega \times \Omega} G(x, y) u_n^2(y) u_n^2(x) dx dy = \omega \int_{\Omega} u_n^2(x) dx.$$

Integration by parts gives

$$-\frac{v_{n-1}(x_0)}{\beta} \int_{\Omega} G(x, x_0) \Delta v_n(x) dx = -\frac{v_{n-1}(x_0)}{\beta} \int_{\Omega} v_n(x) \Delta G(x, x_0) dx = \frac{v_{n-1}(x_0) v_n(x_0)}{\beta},$$

thus

$$\int_{\Omega} |\nabla v_n(x)|^2 dx + \alpha \int_{\Omega \times \Omega} G(x, y) u_n^2(y) u_n^2(x) dx dy = \omega \int_{\Omega} u_n^2(x) dx + \frac{v_n(x_0) v_{n-1}(x_0)}{\beta}. \quad (15)$$

Introducing the notation

$$a_n = \alpha \int_{\Omega \times \Omega} G(x, y) \left(\frac{v_n(y)}{v_{n-1}(x_0)} - \frac{G(y, x_0)}{\beta} \right)^2 \left(\frac{v_n(x)}{v_{n-1}(x_0)} - \frac{G(x, x_0)}{\beta} \right)^2 dx dy,$$

$$b_n = \omega \int_{\Omega} \left(\frac{v_n(x)}{v_{n-1}(x_0)} - \frac{G(x, x_0)}{\beta} \right)^2 dx,$$

Eq. (15) reads (cf. [1, (2.20)])

$$\int_{\Omega} |\nabla v_n|^2 dx + v_{n-1}^4(x_0) a_n = v_{n-1}^2(x_0) b_n + \frac{v_n(x_0) v_{n-1}(x_0)}{\beta}. \quad (16)$$

Due to (9) and (14)

$$a_n \rightarrow \alpha \int_{\Omega \times \Omega} G(x, y) \frac{G^2(y, x_0)}{\beta^2} \frac{G^2(x, x_0)}{\beta^2} dx dy, \quad b_n \rightarrow \omega \int_{\Omega} \frac{G^2(x, x_0)}{\beta^2} dx. \quad (17)$$

Passing to a subsequence we can assume that the sequence $\{v_n(x_0)/v_{n-1}^3(x_0)\}$ has a limit as $n \rightarrow \infty$. We have the following three cases.

(i) If

$$\lim_n \frac{v_n(x_0)}{v_{n-1}^3(x_0)} = \infty, \quad (18)$$

we divide (16) by $v_{n-1}^4(x_0)$

$$\int_{\Omega} \left| \frac{\nabla v_n}{v_{n-1}^2(x_0)} \right|^2 dx + a_n = \frac{b_n}{v_{n-1}^2(x_0)} + \frac{v_n(x_0)}{v_{n-1}^3(x_0) \beta}. \quad (19)$$

Using (9), (14), and (17) in (19), we have

$$\alpha \int_{\Omega \times \Omega} G(x, y) \frac{G^2(y, x_0)}{\beta^2} \frac{G^2(x, x_0)}{\beta^2} dx dy = \infty,$$

which is a contradiction.

(ii) If

$$\lim_n \frac{v_n(x_0)}{v_{n-1}^3(x_0)} = 0, \quad (20)$$

we divide (16) by $v_{n-1}(x_0) v_n(x_0)$

$$\int_{\Omega} \frac{|\nabla v_n|^2}{v_{n-1}(x_0) v_n(x_0)} dx + \frac{a_n v_{n-1}^3(x_0) - b_n v_{n-1}(x_0)}{v_n(x_0)} = \frac{1}{\beta}. \quad (21)$$

Since by (9), (14), and (17),

$$\lim_n \frac{a_n v_{n-1}^3(x_0) - b_n v_{n-1}(x_0)}{v_n(x_0)} = \lim_n \frac{v_{n-1}^3(x_0)}{v_n(x_0)} \left(a_n - \frac{b_n}{v_{n-1}^2(x_0)} \right) = \infty,$$

which implies, using (21), that $\infty = \frac{1}{\beta}$, which is a contradiction.

(iii) Finally, if

$$\lim_n \frac{v_n(x_0)}{v_{n-1}^3(x_0)} = \ell \in (0, \infty),$$

we observe that (see (14))

$$\lim_n \frac{v_{n+1}(x_0)}{v_{n-1}^3(x_0)} = \lim_n \frac{v_{n+1}(x_0)}{v_n^3(x_0)} \frac{v_n(x_0)}{v_{n-1}^3(x_0)} v_n^2(x_0) = \ell^2 \lim_n v_n^2(x_0) = \infty,$$

therefore the subsequence $\{v_{2n}(x_0)\}_{n \in \mathbb{N}}$ satisfies (18) and we get a contradiction.

Therefore (10) is proved. \square

References

- [1] G.M. Coclite, H. Holden, The Schrödinger–Maxwell system with Dirac mass, Ann. Inst. H. Poincaré Anal. Non Linéaire 24 (5) (2007) 773–793.