

Corrections to "Finiteness of the Number of Discrete Eigenvalues of the Schrödinger Operator for a Three Particle System"

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By

Jun UCHIYAMA*

Remark 4, p. 59, was an error due to the negligence of the fact that R depends on Z . The correct assertion is the following:

There exists some constant $Z_0 \left(\frac{Z_3}{2} \geq Z_0 > 0 \right)$ depending only on Z_3 such that for any positive constants $Z_1, Z_2 (Z_0 \geq Z_1 \geq Z_2 > 0)$ the operator of the form

$$(1) \quad H = -\Delta_1 - \Delta_2 - \frac{Z_1}{r_1} - \frac{Z_2}{r_2} + \frac{Z_3}{|r_1 - r_2|}$$

has no discrete eigenvalues.

In fact let $\frac{Z_3}{2} \geq Z_1 \geq Z_2 > 0$. Then taking into consideration Remark 1 and the fact that μ given to (1) by (2.7) and (2.8) equals $-\frac{Z_1^2}{4}$, (3.1) and (3.2) are satisfied by $R = \frac{c_1}{Z_1}$, where c_1 is sufficiently large constant depending only on Z_3 . On the other hand, we can show the following fact which is more precise than Lemma 5.

There exist an "extension operator" \mathcal{O}_1 which maps $\mathcal{E}_{L^2}^1(\mathcal{Q}_1)$ to $\mathcal{D}_{L^2}^1(R^6)$ and some constant c_2 depending only on Z_3 such that for any $\varphi \in \mathcal{E}_{L^2}^1(\mathcal{Q}_1)$

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* Mathematical Institute, Kyoto University of Industrial Arts and Textile Fibres.

$$(2) \quad (\Phi_1\varphi)(x) = \varphi(x) \quad \text{for } x \in \Omega_1$$

and

$$(3) \quad \begin{cases} \|\nabla(\Phi_1\varphi)\|_{R^6}^2 \leq c_2(\|\nabla\varphi\|_{\Omega_1}^2 + Z_1^2\|\varphi\|_{\Omega_1}^2), \\ \|\Phi_1\varphi\|_{R^6}^2 \leq c_2\Delta\|\varphi\|_{\Omega_1}^2. \end{cases}$$

Indeed, for $\varphi \in \mathcal{E}_{L^2}^1(\Omega_1)$ we define $f \in \mathcal{E}_{L^2}^1(\Omega_0)$ by

$$(4) \quad f(x) = \varphi(Rx) \quad \text{for } x \in \Omega_0,$$

where $\Omega_0 = \{x \in R^6; r_1 < 1, r_2 < 1\}$. Then we have

$$(5) \quad \begin{cases} \|f\|_{\mathcal{E}_0}^2 = \frac{1}{R^6}\|\varphi\|_{\Omega_1}^2, \\ \|\nabla f\|_{\Omega_0}^2 = \frac{R^2}{R^6}\|\nabla\varphi\|_{\Omega_1}^2. \end{cases}$$

Let Φ_2 and $c_3 = c_3(\Omega_0, \Phi_2)$ be the Φ and \tilde{c} satisfying the relations given in Lemma 5 with Ω_1 replaced by Ω_0 . Now we define $\Phi_1\varphi \in \mathcal{D}_{L^2}^1(R^6)$ by

$$(6) \quad (\Phi_1\varphi)(x) = (\Phi_2f)\left(\frac{x}{R}\right) \quad \text{for } x \in R^6.$$

Then by Lemma 5 Φ_1 satisfies (2), and

$$(7) \quad \frac{1}{R^6}\|\Phi_1\varphi\|_{R^6}^2 = \|\Phi_2f\|_{R^6}^2 \leq c_3\|f\|_{\Omega_0}^2 = \frac{c_3}{R^6}\|\varphi\|_{\Omega_1}^2$$

and

$$(8) \quad \begin{aligned} \frac{R^2}{R^6}\|\nabla(\Phi_1\varphi)\|_{R^6}^2 &= \|\nabla(\Phi_2f)\|_{R^6}^2 \leq c_3(\|\nabla f\|_{\Omega_0}^2 + \|f\|_{\Omega_0}^2) \\ &= \frac{c_3R^2}{R^6}\left(\|\nabla\varphi\|_{\Omega_1}^2 + \frac{1}{R^2}\|\varphi\|_{\Omega_1}^2\right). \end{aligned}$$

By (7) and (8) we have (3).

Then using the well-known inequality

$$(9) \quad \int_{R^3} \frac{|\varphi|^2}{r^2} dx \leq 4 \int_{R^3} |\nabla\varphi|^2 dx \quad \text{for } \varphi \in \mathcal{D}_{L^2}^1(R^3),$$

and Schwartz's inequality, we have by (2) and (3) for any $\psi \in \mathcal{D}_{L^2}^2(R^6)$

$$(10) \quad \int_{\Omega_1} \frac{|\psi|^2}{r_1} dx \leq \int_{R^6} \frac{|\Phi_1\psi|^2}{r_1} dx \leq \left(\int_{R^6} \frac{|\Phi_1\psi|^2}{r_1^2} dx\right)^{1/2} \cdot \|\Phi_1\psi\|_{R^6}$$

$$\begin{aligned}
 &= \left(\int_{R^3} d\mathbf{r}_2 \int_{R^3} \frac{|\Phi_1 \psi|^2}{r_1^2} d\mathbf{r}_1 \right)^{1/2} \cdot \|\Phi_1 \psi\|_{R^6} \leq 2 \left(\int_{R^3} d\mathbf{r}_2 \int_{R^3} |\nabla_1(\Phi_1 \psi)|^2 d\mathbf{r}_1 \right)^{1/2} \cdot \|\Phi_1 \psi\|_{R^6} \\
 &\leq \frac{\eta}{Z_1} \|\nabla_1(\Phi_1 \psi)\|_{R^6}^2 + \frac{Z_1}{\eta} \|\Phi_1 \psi\|_{R^6}^2 \\
 &\leq \frac{c_2 \eta}{Z_1} \|\nabla \psi\|_{\Omega_1}^2 + c_2 \left(\eta + \frac{1}{\eta} \right) Z_1 \|\psi\|_{\Omega_1}^2
 \end{aligned}$$

and similarly

$$(11) \quad \int_{\Omega_1} \frac{|\psi|^2}{r_2} dx \leq \frac{c_2 \eta}{Z_1} \|\nabla \psi\|_{\Omega_1}^2 + c_2 \left(\eta + \frac{1}{\eta} \right) Z_1 \|\psi\|_{\Omega_1}^2,$$

where η is an arbitrary positive constant. Let $\eta = \frac{1}{2c_2}$. Then if we take into consideration $\frac{Z_3}{|\mathbf{r}_1 - \mathbf{r}_2|} > \frac{Z_3}{2R} = \frac{Z_1 Z_3}{2c_1}$ in Ω_1 , we have by (10) and (11) for any $\psi \in \mathcal{D}_{L^2}^2(R^6)$

$$\begin{aligned}
 (12) \quad L_1[\psi] &\geq \left(1 - c_2 \eta - \frac{Z_2}{Z_1} c_2 \eta \right) \|\nabla \psi\|_{\Omega_1}^2 + \\
 &\quad + \left\{ \frac{Z_1 Z_3}{2c_1} - c_2 \left(\eta + \frac{1}{\eta} \right) Z_1 (Z_1 + Z_2) \right\} \|\psi\|_{\Omega_1}^2 \\
 &\geq N_1 \left\{ \frac{Z_3}{2c_1} - 2c_2 \left(\eta + \frac{1}{\eta} \right) Z_1 \right\} \|\psi\|_{\Omega_1}^2.
 \end{aligned}$$

Then there exists some constant Z_0 such that for any Z_1 and Z_2 ($Z_0 \leq Z_1 \leq Z_2 > 0$) we have

$$(13) \quad L_1[\psi] \geq \mu \|\psi\|_{\Omega_1}^2 = -\frac{Z_1^2}{4} \|\psi\|_{\Omega_1}^2$$

for any $\psi \in \mathcal{D}_{L^2}^2(R^6)$. By Lemma 1, Lemma 3 and (13), we have the assertion.

Remark. There exists some constant Z'_0 depending on Z_1 ($Z'_0 > Z_1 \geq Z_2 > 0$) such that for any $Z_3 > Z'_0$ the operator of the form (1) has no discrete eigenvalues.

In fact R satisfying (3.1) and (3.2) is independent of Z_3 ($Z_3 \geq 2Z_1$). Then by the same calculation as (3.16) and (3.17) we have

$$(14) \quad L_1[\psi] \geq (1 - 2\tilde{c}\eta Z_1) \|\nabla \psi\|_{\Omega_1}^2 + \left(\frac{Z_3}{2R} - 2Z_1(\eta + c(\eta))\tilde{c} \right) \|\psi\|_{\Omega_1}^2,$$

Take $\eta = \frac{1}{2\tilde{c}Z_1}$ and Z_3 sufficiently large. We have for any $\psi \in \mathcal{D}_{L^2}^2(\mathbb{R}^8)$

$$(15) \quad L_1[\psi] \geq -\frac{Z_1^2}{4} \|\psi\|_{\mathfrak{h}_1}^2 = \mu \|\psi\|_{\mathfrak{h}_1}^2.$$

Then by Lemma 1, Lemma 3 and (15), we have the assertion.