

# $L_2$ and $L_\infty$ Estimates of the Solutions for the Compressible Navier-Stokes Equations in a 3D Exterior Domain

*Dedicated to Professors Takaaki Nishida and Masayasu Mimura  
on their sixtieth birthdays*

By

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## Abstract

We consider the boundary value problem of the equation of motion of viscous compressible fluid in a 3D exterior domain. We shall give the  $L_\infty$  estimates of the solutions and  $L_2$ -estimates of the derivative with respect to space variable of the solutions.

## §1. Introduction

In this paper, we consider the equation of the motion of compressible viscous fluid in a 3D exterior domain. The equation is given by the following system for the density  $\rho(t, x)$  and the velocity  $v(t, x) = (v_1(t, x), v_2(t, x), v_3(t, x))$ ,

(1.1)

$$\begin{aligned} \rho_t + \nabla \cdot (\rho v) &= 0 && \text{in } (0, \infty) \times \Omega, \\ \rho(v_t + (v \cdot \nabla)v) + \nabla P(\rho) &= \mu \Delta v + (\mu + \nu) \nabla(\nabla \cdot v) && \text{in } (0, \infty) \times \Omega, \\ v|_{\partial\Omega} &= 0 && \text{on } (0, \infty) \times \partial\Omega, \\ \rho(0, x) &= \rho_0(x), \quad v(0, x) = v_0(x) && \text{in } \Omega, \end{aligned}$$

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where  $\Omega$  is an exterior domain in  $\mathbf{R}^3$  with the compact smooth boundary  $\partial\Omega$ ,  $P = P(\rho)$  the pressure,  $\mu \geq 0, (2/3)\mu + \nu > 0$  the viscosity coefficients. The unique existence of smooth solutions globally in time near constant state  $(\bar{\rho}_0, 0)$ , where  $\bar{\rho}_0$  is a positive constant, was proved by the employing the same argument as in Matsumura and Nishida [11], [12] for the Cauchy problem in  $\mathbf{R}^3$ ; Matsumura and Nishida [13], [14], [15] for the exterior domain in  $\mathbf{R}^3$ . Concerning the decay property of solutions  $(\rho(t, x), v(t, x))$ , if the initial data  $(\rho_0(x) - \bar{\rho}_0, v_0(x))$  belongs to  $H^4$  and  $L_1$ , then as  $t \rightarrow \infty$

$$\begin{aligned} \|(\rho(t, \cdot) - \bar{\rho}_0, v(t, \cdot))\|_{L_\infty} &= O(t^{-3/2}), \\ \|(\rho(t, \cdot) - \bar{\rho}_0, v(t, \cdot))\|_{L_2} &= O(t^{-3/4}), \\ \|(\rho(t, \cdot) - \bar{\rho}_0, v(t, \cdot))\|_{L_2} &= O(t^{1/2}). \end{aligned}$$

This fact was investigated by Hoff and Zumbrun [3], [4], Liu and Wang [9], Matsumura and Nishida [11], [12], Ponce [16] and Weike [17] for the Cauchy problem case; Kobayashi [7], Kobayashi and Shibata [8] for the exterior domain case. On the other hand, if the initial data belongs to  $H^3$  only, namely we do not assume that the initial data belongs to  $L_1$ , then, Deckelnick [1], [2] showed that as  $t \rightarrow \infty$

$$\begin{aligned} \|(\rho_t(t, \cdot), v_t(t, \cdot))\|_{L_2(\Omega)} &= O(t^{-1/2}), \\ \|\partial_x(\rho(t, \cdot), v(t, \cdot))\|_{L_2(\Omega)} &= O(t^{-1/4}), \\ \|v(t, \cdot)\|_{C^0(\bar{\Omega})} &= O(t^{-1/4}), \\ \|\rho(t, \cdot) - \bar{\rho}_0\|_{C^0(\bar{\Omega})} &= O(t^{-1/8}), \end{aligned}$$

in the exterior domain case; Matsumura [10] showed that as  $t \rightarrow \infty$

$$\begin{aligned} \|(\rho_t(t, \cdot), v_t(t, \cdot))\|_{L_2(\mathbf{R}^3)} &= O(t^{-1/2}), \\ \|\partial_x(\rho(t, \cdot), v(t, \cdot))\|_{L_2(\mathbf{R}^3)} &= O(t^{-1/2}), \\ \|\partial_x^2(\rho(t, \cdot), v(t, \cdot))\|_{L_2(\mathbf{R}^3)} &= O(t^{-1}), \\ \|(\rho(t, \cdot) - \bar{\rho}_0, v(t, \cdot))\|_{L_\infty(\mathbf{R}^3)} &= O(t^{-3/4}), \end{aligned}$$

in the Cauchy problem case. In this paper, we shall investigate the exterior problem of the system (1.1) and give the better decay rate than the rate obtained by Deckelnick [1], [2] in the case that the initial data belongs to  $H^3$  or  $H^4$  only. In particular, the  $L_2$ -decay rate of the first derivative with respect to the spacial variable  $x$  for the solutions corresponds to the rate obtained by Matsumura [10].

## §2. Notation and Main Results

Let  $L_p$  denotes the usual  $L_p$  space on  $\Omega$  with norm  $\|\cdot\|_{L_p}$ . Put

$$W_p^m = \{u \in L_p \mid \|u\|_{W_p^m} < \infty\}, \quad \|u\|_{W_p^m} = \sum_{|\alpha| \leq m} \|\partial_x^\alpha u\|_{L_p},$$

$$H^m = W_2^m, \quad W_p^0 = L_p, \quad H^0 = L_2.$$

Set

$$W_p^{k,m} = \{(\rho, v) = (\rho, v_1, v_2, v_3) \mid \rho \in W_p^k, v_j \in W_p^m, j = 1, 2, 3\},$$

$$\|(\rho, v)\|_{W_p^{k,m}} = \|\rho\|_{W_p^k} + \|v\|_{W_p^m},$$

and

$$H^{k,m} = W_2^{k,m}, \quad \|u\|_{H^{k,m}} = \|u\|_{W_2^{k,m}}.$$

First, in order to state the existence of solutions according to Matsumura and Nishida [13], [14], [15], we introduce the assumptions and notations. Let  $\bar{\rho}_0$  be a positive constant. We assume that

A1.  $P$  is a smooth function in a neighborhood of  $\bar{\rho}_0$  and  $\partial P / \partial \rho > 0$ .

A2. The initial data  $(\rho_0, v_0)$  satisfies the compatibility condition and regularity, namely  $(\rho_0 - \bar{\rho}_0, v_0) \in H^3$ ,  $v_0|_{\partial\Omega} = 0$  and  $(\rho_1, v_1) = (\rho_t, v_t)|_{t=0}$  satisfies

$$\rho_1 = -\nabla \cdot (\rho_0 v_0),$$

$$v_1 = -(v_0 \cdot \nabla) v_0 + \frac{\mu}{\rho_0} \Delta v_0 + \frac{\nu}{\rho_0} \nabla (\nabla \cdot v_0) - \frac{\nabla P(\rho_0)}{\rho_0},$$

and

$$\rho_1 \in H^2, \quad v_1 \in H^1, \quad v_1|_{\partial\Omega} = 0.$$

Put

$$X(0, \infty) = \left\{ U = (\rho, v) \mid \rho - \bar{\rho}_0 \in \bigcap_{j=0}^1 C^j([0, \infty); H^{3-j}), \right.$$

$$\left. \begin{aligned} \partial_x \rho \in L_2((0, \infty); H^2), \quad \rho_t, v_t \in L_2((0, \infty); H^2), \\ v \in \bigcap_{j=0}^1 C^j([0, \infty); H^{3-2j}), \quad \partial_x v \in L_2((0, \infty); H^3) \end{aligned} \right\},$$

and

$$N(0, \infty)^2 = \sup_{0 \leq t < \infty} (\|U(t) - \bar{U}_0\|_{H^3}^2 + \|U_t(t)\|_{H^{2,1}}^2)$$

$$+ \int_0^\infty (\|\partial_x U(s)\|_{H^{2,3}}^2 + \|U_s(s)\|_{H^2}^2) ds,$$

where  $\bar{U}_0 = (\bar{\rho}_0, 0)$ . Then, we have

**Proposition 2.1** (Matsumura and Nishida [13], [14], [15]). *Assume that the assumptions A.1 and A.2 hold. Then, there exists an  $\epsilon_0$  such that if  $\|(\rho_0 - \bar{\rho}_0, v_0)\|_{H^3} \leq \epsilon_0$ , then (1.1) admits a unique solution  $(\rho, v) \in X(0, \infty)$ .*

*Moreover, there exists a constant  $C$  such that*

$$(2.1) \quad N(0, \infty) \leq C \|(\rho_0 - \bar{\rho}_0, v_0)\|_{H^3}.$$

*Remark.* If the initial data  $(\rho_0 - \bar{\rho}_0, v_0) \in H^4$  and satisfies the second order compatibility condition and regularity, namely  $(\rho_2, v_2) = (\rho_{tt}, v_{tt})|_{t=0}$  is determined successively by the initial data  $(\rho_0, v_0)$  through the system (1.1), then we have

$$(2.2) \quad \begin{aligned} \tilde{N}(0, \infty)^2 &= \sup_{0 \leq t < \infty} (\|U(t) - \bar{U}_0\|_{H^4}^2 + \|U_t(t)\|_{H^{3,2}}^2) \\ &\quad + \int_0^\infty (\|\partial_x U(s)\|_{H^{3,4}}^2 + \|U_s(s)\|_{H^3}^2) ds, \\ &\leq C \|(\rho_0 - \bar{\rho}_0, v_0)\|_{H^4}^2. \end{aligned}$$

Now, we shall state our main results.

**Theorem 2.1.** *Assume that the assumptions A.1 and A.2 hold. Then, there exists an  $\epsilon_1$  such that if  $\|(\rho_0 - \bar{\rho}_0, v_0)\|_{3,2} \leq \epsilon_1$ , the solution  $(\rho, v)$  of the system (1.1) has the following asymptotic behavior as  $t \rightarrow \infty$ :*

$$\begin{aligned} \|(\rho_t(t, \cdot), v_t(t, \cdot))\|_{L_2} &= O(t^{-1/2}), \\ \|\partial_x v(t, \cdot)\|_{H^1} &= O(t^{-1/2}), \\ \|\partial_x \rho(t, \cdot)\|_{L_2} &= O(t^{-1/2}), \\ \|\partial_x^2 \rho(t, \cdot)\|_{L_2} &= O(t^{-3/4} \log t), \\ \|(\rho(t, \cdot) - \bar{\rho}_0, v(t, \cdot))\|_{L_\infty} &= O(t^{-3/4} \log t). \end{aligned}$$

**Corollary 2.1.** *The assumptions in Theorem 2.1 hold. Moreover, if the initial data  $(\rho_0 - \bar{\rho}_0, v_0) \in H^4$  satisfies the second order compatibility condition and regularity in Remark, then we have*

$$\|(\rho(t, \cdot) - \bar{\rho}_0, v(t, \cdot))\|_{L_\infty} = O(t^{-3/4}) \quad \text{as } t \rightarrow \infty.$$

### §3. Linearized Problem

In this section, we shall consider the following linearized problem associated to the problem (1.1) (see Section 4)

$$(3.1) \quad \begin{aligned} \rho_t + \gamma \nabla \cdot v &= 0 & \text{in } [0, \infty) \times \Omega, \\ v_t - \alpha \Delta v - \beta \nabla (\nabla \cdot v) + \gamma \nabla \rho &= 0 & \text{in } [0, \infty) \times \Omega, \\ v|_{\partial \Omega} &= 0 & \text{on } [0, \infty) \times \partial \Omega, \\ \rho(0, x) = \rho_0(x), \quad v(0, x) &= v_0(x) & \text{in } \Omega, \end{aligned}$$

where  $\alpha > 0, \beta \geq 0$  and  $\gamma > 0$ . Let  $A$  be the  $4 \times 4$  matrix of the differential operator of the form:

$$A = \begin{pmatrix} 0 & \gamma \nabla \\ \gamma \nabla & -\alpha \Delta - \beta \nabla \nabla \end{pmatrix}$$

with the domain:

$$D_p(A) = \{U = (\rho, v) \in W_p^{1,2} \mid v|_{\Omega} = 0\}$$

for  $1 < p < \infty$ . Then, (3.1) is written in the form:

$$U_t + AU = 0 \quad \text{for } t > 0, \quad U|_{t=0} = U_0,$$

where  $U_0 = (\rho_0, v_0)$  and  $U = (\rho, v)$ . Then,

**Proposition 3.1** (*Kobayashi [5], [6], [7], Kobayashi and Shibata [8]*).

*The operator  $-A$  generates an analytic semigroup  $\{e^{-tA}\}_{t \geq 0}$  on  $W_p^{1,0}$ ,  $1 < p < \infty$  and the following properties hold.*

(I) *Let  $1 < p < \infty$ . For  $0 < t \leq 2$ , we have*

$$(3.2) \quad \|e^{-tA}U\|_{W_p^{1,0}} \leq C\|U\|_{W_p^{1,0}} \quad \text{for } U \in W_p^{1,0},$$

$$(3.3) \quad \|e^{-tA}U\|_{W_p^1} \leq Ct^{-1/2}\|U\|_{W_p^{1,0}} \quad \text{for } U \in W_p^{1,0},$$

$$(3.4) \quad \|(\mathbf{I} - \mathbf{P})e^{-tA}U\|_{W_p^2} \leq Ct^{-1/2}\|U\|_{W_p^{2,1}} \quad \text{for } U \in W_p^{2,1}.$$

*Here and hereafter, we shall use the notations:*

$$\mathbf{I}U = U, \quad \mathbf{P}U = v \quad \text{and} \quad (\mathbf{I} - \mathbf{P})U = \rho \quad \text{for } U = (\rho, v).$$

(II) *Let  $1 \leq q \leq 2 \leq p < \infty$ . For  $U \in W_p^{1,0} \cap L_q$  and  $t \geq 1$*

(3.5)

$$\|e^{-tA}U\|_{L_p} \leq Ct^{-\sigma} \left( \|U\|_{L_q} + \|U\|_{W_p^{1,0}} \right), \quad 2 \leq p < \infty, \quad \sigma = \frac{3}{2} \left( \frac{1}{q} - \frac{1}{p} \right),$$

(3.6)

$$\|\partial_t e^{-tA}U\|_{L_p} + \|\partial_x e^{-tA}U\|_{L_p} \leq Ct^{-\sigma - \frac{1}{2}} \left( \|U\|_{L_q} + \|U\|_{W_p^{1,0}} \right), \quad 2 \leq p \leq 3,$$

(3.7)

$$\|e^{-tA}U\|_{W_\infty^{0,1}} \leq Ct^{-\frac{3}{2q}} \left( \|U\|_{L_q} + \|U\|_{W_p^{1,0}} \right), \quad 3 < p < \infty,$$

*where  $\partial_t = d/dt$  and  $\partial_x^m u = (\partial_x^\alpha u \mid |\alpha| = m)$ . Moreover, if  $q > 1$  and  $U \in W_p^{2,1} \cap W_q^{1,0}$ , then for  $t \leq 1$*

(3.8)

$$\|\partial_x^2(\mathbf{I} - \mathbf{P})e^{-tA}U\|_{L_p} + \|\partial_t \partial_x e^{-tA}U\|_{L_p} \leq Ct^{-\frac{3}{2q}} \left( \|U\|_{W_q^{1,0}} + \|U\|_{W_p^{2,1}} \right), \\ 2 \leq p < \infty.$$

#### §4. Proof of Theorem 2.1

First of all, we shall introduce the linearized equations. By the change of unknown functions:  $(\rho, v) \rightarrow (\rho + \bar{\rho}_0, v)$ , (1.1) is reduced to the following equation:

$$(4.1) \quad \begin{aligned} \rho_t + \bar{\rho}_0 \nabla \cdot v &= f_1 \\ v_t - \hat{\mu} \Delta v - (\hat{\mu} + \hat{\nu}) \nabla (\nabla \cdot v) + p_1 \nabla \rho &= f_2, \\ \rho(0, x) &= \rho_0(x) - \bar{\rho}_0, \quad v(0, x) = v_0(x), \end{aligned}$$

where  $\hat{\mu} = \mu/\bar{\rho}_0, \hat{\nu} = \nu/\bar{\rho}_0, p_1 = P_\rho(\bar{\rho}_0)/\bar{\rho}_0$ ,

$$(4.2) \quad \begin{aligned} f_1 &= -\rho \nabla \cdot v - \nabla \rho \cdot v, \\ f_2 &= -(v \cdot \nabla)v + \left( \frac{\mu}{\rho + \bar{\rho}_0} - \hat{\mu} \right) \Delta v + \left( \frac{\mu + \nu}{\rho + \bar{\rho}_0} - \hat{\mu} - \hat{\nu} \right) \nabla (\nabla \cdot v), \\ &\quad + \left( p_1 - \frac{P_\rho(\rho)}{\rho + \bar{\rho}_0} \right) \nabla \rho. \end{aligned}$$

If we put  $\rho' = (p_1/\bar{\rho}_0)^{1/2} \rho$  and  $v' = v$ , then (4.1) is reduced to the symmetric form

$$\begin{aligned} \rho'_t + \gamma \nabla \cdot v' &= f'_1 \\ v'_t - \alpha \Delta v' - \beta \nabla (\nabla \cdot v') + p_1 \nabla \rho' &= f_2, \\ \rho'(0, x) &= \rho'_0, \quad v'(0, x) = v_0(x), \end{aligned}$$

where  $\alpha = \hat{\mu}, \beta = \hat{\mu} + \hat{\nu}$  and  $\gamma = \sqrt{P_\rho(\bar{\rho}_0)}$ . For the notational simplicity, we write:  $\rho = \rho', v = v', f_1 = f'_1$ , again. If we put  $U = (\rho, v), U_0 = (\rho_0, v_0), F(U) = (f_1, f_2)$  and

$$A = \begin{pmatrix} 0 & \gamma \nabla \\ \gamma \nabla & -\alpha \Delta - \beta \nabla \nabla \end{pmatrix}$$

then (1.1) is reduced to the following equations

$$(4.3) \quad \begin{aligned} U_t + AU &= F(U), \\ U(0) &= U_0. \end{aligned}$$

Here  $F(U) = (f_1, f_2)$  is written as follows:

$$\begin{aligned} f_1 &= -\frac{\gamma}{\bar{\rho}_0} (\rho \nabla \cdot v + \nabla \rho \cdot v), \\ f_2 &= -(v \cdot \nabla)v + a_1(\rho) \rho \Delta v + a_2(\rho) \rho \nabla (\nabla \cdot v) + a_3(\rho) \rho \nabla \rho, \end{aligned}$$

where  $a_j(\rho)$  ( $j = 1, 2, 3$ ) represent identities (4.2). To prove our main results, we shall estimate the following integral equation

$$U(t) = e^{-tA}U_0 - S(t), \quad S(t) = \int_0^t e^{-(t-s)A}F(U)(s)ds.$$

Let  $N(0, \infty)$  be the quantity defined in Section 2. By choosing  $\|(\rho_0 - \bar{\rho}_0, v_0)\|_{H^3}$  small enough, we can make  $N(0, \infty)$  as small we want, and therefore we will state the smallness assumption in term of  $N(0, \infty)$  instead of  $\|(\rho_0 - \bar{\rho}_0, v_0)\|_{H^3}$  in the course of our proof of Theorem 2.1 below.

*Step 1.* Put

$$M_1(t) = \sup_{0 \leq s \leq t} (1+s)^{\frac{1}{2}} \|\partial_s U(s)\|_{L_2},$$

$$M_2(t) = \sup_{0 \leq s \leq t} (1+s)^{\frac{1}{2}} \|\partial_x U(s)\|_{H^1},$$

$$M_3(t) = \sup_{0 \leq s \leq t} (1+s)^{\frac{1}{2}} \|U(s)\|_{L_\infty}.$$

Then, there exists an  $\epsilon > 0$  such that if  $N(0, \infty) \leq \epsilon$ , then

$$M_1(t) + M_2(t) + M_3(t) \leq C\|U_0\|_{H^2}.$$

First, we shall show that

$$(4.4) \quad M_1(t) \leq C((M_2(t) + M_3(t))N(0, \infty) + \|U_0\|_{H^{1,0}}),$$

$$(4.5) \quad \sup_{0 \leq s \leq t} (1+s)^{1/2} \|\partial_x U(s)\|_{L_2} \leq C((M_2(t) + M_3(t))N(0, \infty) + \|U_0\|_{H^{1,0}}),$$

$$(4.6) \quad M_3(t) \leq C((M_2(t) + M_3(t))N(0, \infty) + \|U_0\|_{H^2}).$$

When  $0 \leq t \leq 2$ , by Proposition 2.1 and the inequality

$$(4.7) \quad \|u\|_{L_p} \leq C\|u\|_{H^2}, \quad 6 < p \leq \infty,$$

we have

$$(4.8) \quad M_1(t) + M_2(t) + M_3(t) \leq CN(0, \infty),$$

and therefore we consider the case when  $t \geq 1$ , below. By (3.6) with  $(p, q) = (2, 2)$

$$(4.9) \quad \|\partial_t e^{-tA}U_0\|_{L_2} + \|\partial_x e^{-tA}U_0\|_{L_2} \leq Ct^{-1/2}\|U_0\|_{H^{1,0}},$$

and by (3.7) with  $(p, q) = (4, 2)$  and the inequalities:

$$(4.10) \quad \|u\|_{L_p} \leq C\|u\|_{H^1} \quad (2 \leq p < 6) \quad \text{and} \quad \|u\|_{L_6} \leq C\|\partial_x u\|_{L_2},$$

we have

$$(4.11) \quad \|e^{-tA}U_0\|_{L_\infty} \leq Ct^{-3/4}\|U_0\|_{H^2}.$$

The main task is the estimation of  $S(t)$ , which is divided into the two parts as follows:

$$S(t) = \left\{ \int_{t-1}^t + \int_0^{t-1} \right\} e^{-(t-s)A} F(U)(s) ds = I(t) + II(t).$$

Before going further on the proof of (4.4), (4.5) and (4.6) we prepare the estimates of nonlinear term  $F(U)$ : By (4.7), (4.10) and the Hölder's inequality, we have

$$(4.12) \quad \begin{aligned} \|F(U)(s)\|_{L_1} &\leq C\|U\|_2\|\partial_x U(s)\|_{H^1}, \\ \|F(U)(s)\|_{H^{2,1}} &\leq C(\|U(s)\|_{L_\infty} + \|\partial_x U(s)\|_{H^1})\|\partial_x U(s)\|_{H^2}, \\ \|F(U)(s)\|_{W_4^{1,0}} &\leq C(\|U(s)\|_{L_\infty} + \|\partial_x U(s)\|_{H^1})\|\partial_x U(s)\|_{H^2}, \\ \|\partial_s F(U)(s)\|_{H^{1,0}} &\leq C(\|U(s)\|_{L_\infty} + \|\partial_x U(s)\|_{H^1})\|\partial_s U(s)\|_{H^2}. \end{aligned}$$

Now, we return to estimate  $S(t)$ . By (3.6) with  $(p, q) = (2, 1)$  and (4.12), we have

$$(4.13) \quad \begin{aligned} \|\partial_t II(t)\|_{L_2} + \|\partial_x II(t)\|_{L_2} \\ \leq C \int_0^{t-1} (t-s)^{-5/4} (\|F(U)(s)\|_{L_1} + \|F(U)(s)\|_{H^{1,0}}) ds \\ \leq C \int_0^{t-1} (t-s)^{-5/4} (1+s)^{-1/2} ds (M_2(t) + M_3(t)) N(0, \infty) \\ \leq C(1+t)^{-1/2} (M_2(t) + M_3(t)) N(0, \infty). \end{aligned}$$

On the other hand, by (3.3) and (4.12),

$$(4.14) \quad \begin{aligned} \|\partial_x I(t)\|_{L_2} &\leq C \int_{t-1}^t (t-s)^{-1/2} \|F(U)(s)\|_{H^{1,0}} ds \\ &\leq C \int_{t-1}^t (t-s)^{-1/2} (1+s)^{-1/2} ds (M_2(t) + M_3(t)) N(0, \infty) \\ &\leq C(1+t)^{-1/2} (M_2(t) + M_3(t)) N(0, \infty). \end{aligned}$$

Combining (4.8), (4.9), (4.13) and (4.14), we have (4.5). By integration by parts,

$$\partial_t I(t) = \int_{t-1}^t e^{-(t-s)A} \partial_s F(U)(s) ds$$



and therefore by (3.2) and (4.12),

$$\begin{aligned}
 (4.15) \quad \|\partial_t I(t)\|_{L_2} &\leq C \int_{t-1}^t \|\partial_s F(U)(s)\|_{H^{1,0}} ds \\
 &\leq C(1+t)^{-1/2} \left( \int_{t-1}^t \|\partial_s U(s)\|_{H^2}^2 ds \right)^{1/2} (M_2(t) + M_3(t)) \\
 &\leq C(1+t)^{-1/2} (M_2(t) + M_3(t)) N(0, \infty).
 \end{aligned}$$

Combining (4.8), (4.9), (4.13) and (4.15), we have (4.4). By (3.7) with  $(p, q) = (4, 1)$  and (4.12),

$$\begin{aligned}
 (4.16) \quad \|II(t)\|_{L_\infty} &\leq C \int_0^{t-1} (t-s)^{-3/2} \left( \|F(U)(s)\|_{L_1} + \|F(U)(s)\|_{W_4^{1,0}} \right) ds \\
 &\leq C \int_0^{t-1} (t-s)^{-3/2} (1+s)^{-1/2} ds (M_2(t) + M_3(t)) N(0, \infty) \\
 &\leq C(1+t)^{-1/2} (M_2(t) + M_3(t)) N(0, \infty).
 \end{aligned}$$

By (3.3), (4.12) and the Sobolev inequality

$$\begin{aligned}
 (4.17) \quad \|I(t)\|_{L_\infty} &\leq C \|I(t)\|_{W_4^1} \\
 &\leq C \int_{t-1}^t (t-s)^{-1/2} \|F(U)(s)\|_{W_4^{1,0}} ds \\
 &\leq C \int_{t-1}^t (t-s)^{-1/2} (1+s)^{-1/2} ds (M_2(t) + M_3(t)) N(0, \infty) \\
 &\leq C(1+t)^{-1/2} (M_2(t) + M_3(t)) N(0, \infty).
 \end{aligned}$$

Combining (4.8), (4.11), (4.16) and (4.17), we have (4.6). Next, we shall show that

$$(4.18) \quad \sup_{0 \leq s \leq t} (1+s)^{1/2} \|\partial_x^2 U(t)\|_{L_2} \leq C((M_2(t) + M_3(t))N(0, \infty) + \|U_0\|_{H^2}).$$

In order to prove (4.18), we shall use the following Proposition (see [8, Proposition A, p. 3])

**Proposition 4.1.** *Let  $b$  be an arbitrary number such that  $B_{b-3} = \{x \in \mathbf{R}^3 \mid |x| \leq b-3\} \supset \partial\Omega$ . Let  $1 < p < \infty$  and  $m$  be an integer  $\geq 0$ . Suppose that  $u = (u_1, u_2, u_3) \in W_p^{m+2}(\Omega)$  and  $f = (f_1, f_2, f_3) \in W_p^m(\Omega)$  satisfy the equation:*

$$-\alpha \Delta u - \beta \nabla(\nabla \cdot u) = f \quad \text{in } \Omega \quad \text{and} \quad u|_{\partial\Omega} = 0,$$

where  $\alpha > 0$  and  $\alpha + \beta > 0$ . Then, the following estimate holds:

$$\|\partial_x^{m+2}u\|_{L_p} \leq C_{m,p} \left( \|f\|_{W_p^m} + \|u\|_{W_p^1(\Omega \cap B_b)} \right).$$

Applying Proposition 4.1 to the second equation of (4.1), we have

$$\begin{aligned} & \|\partial_x^2 \mathbf{P}U(t)\|_{L_2} \\ & \leq C \left( \|\partial_t \mathbf{P}U(t)\|_{L_2} + \|\partial_x(\mathbf{I} - \mathbf{P})U(t)\|_{L_2} + \|\mathbf{P}F(U)(t)\|_{L_2} + \|\mathbf{P}U(t)\|_{H^1(\Omega \cap B_b)} \right), \end{aligned}$$

which together with (4.4), (4.5), (4.6) and (4.12) implies that

$$(4.19) \quad \|\partial_x^2 \mathbf{P}U(t)\|_{L_2} \leq C(1+t)^{-1/2}((M_2(t) + M_3(t))N(0, \infty) + \|U_0\|_{H^2}).$$

Therefore, our task is to estimate  $\|\partial_x^2(\mathbf{I} - \mathbf{P})e^{-tA}U_0\|_{L_2}$ . By (3.8) with  $(p, q) = (2, 2)$ ,

$$(4.20) \quad \|\partial_x^2(\mathbf{I} - \mathbf{P})e^{-tA}U_0\|_{L_2} \leq Ct^{-3/4}\|U_0\|_{H^{2,1}}.$$

By (3.8) with  $(p, q) = (2, 2)$  and (4.12),

$$\begin{aligned} (4.21) \quad & \|\partial_x^2(\mathbf{I} - \mathbf{P})II(t)\|_2 \\ & \leq C \int_0^{t-1} (t-s)^{-3/4} \|F(U)(s)\|_{H^{2,1}} ds \\ & \leq C \int_0^{t-1} (t-s)^{-3/4} (1+s)^{-1/2} \|\partial_x U(s)\|_{H^2} ds (M_2(t) + M_3(t)) \\ & \leq C \left( \int_0^{t-1} (t-s)^{-3/2} (1+s)^{-1} ds \right)^{1/2} (M_2(t) + M_3(t))N(0, \infty) \\ & \leq (1+t)^{-1/2} (M_2(t) + M_3(t))N(0, \infty). \end{aligned}$$

On the other hand, by (3.4) and (4.12)

$$\begin{aligned} (4.22) \quad & \|\partial_x^2(\mathbf{I} - \mathbf{P})I(t)\|_{L_2} \leq C \int_{t-1}^t (t-s)^{-1/2} \|F(U)(s)\|_{H^{2,1}} ds \\ & \leq C \int_{t-1}^t (t-s)^{-1/2} (1+s)^{-1/2} ds (M_2(t) + M_3(t))N(0, \infty) \\ & \leq C(1+t)^{-1/2} (M_2(t) + M_3(t))N(0, \infty). \end{aligned}$$

Combining (4.8), (4.20), (4.21), (4.22) with (4.19), we have (4.18). By (4.4), (4.5), (4.6) and (4.18), we have

$$M_1(t) + M_2(t) + M_3(t) \leq C\|U_0\|_{2,2} + C(M_2(t) + M_3(t))N(0, \infty).$$

If  $CN(0, \infty) < 1$ , then we have Step 1.

*Step 2.* Put

$$M_4(t) = \sup_{0 \leq s \leq t} \frac{(1+s)^{3/4}}{\log(2+s)} \|\partial_x^2(\mathbf{I} - \mathbf{P})U(s)\|_{L_2},$$

$$M_5(t) = \sup_{0 \leq s \leq t} \frac{(1+s)^{3/4}}{\log(2+s)} \|U(s)\|_{L_\infty}.$$

Then, there exists an  $\epsilon' > 0$  such that if  $N(0, \infty) \leq \epsilon'$ , then

$$M_4(t) + M_5(t) \leq C(\|U_0\|_{H^2} + M_2(t)^2).$$

By (4.7), (4.10) and the Hölder's inequality,

(4.23)

$$\|F(U)(s)\|_{H^{2,1}} \leq C(\|U(s)\|_{L_\infty} + \|\partial_x^2 \rho(s)\|_{L_2}) \|\partial_x U(s)\|_{H^2} + C\|\partial_x U(s)\|_{H^1}^2.$$

Therefore, by (3.8) with  $(p, q) = (2, 2)$  and (4.23),

(4.24)

$$\begin{aligned} & \|\partial_x^2(\mathbf{I} - \mathbf{P})II(t)\|_2 \\ & \leq C \int_0^{t-1} (t-s)^{-3/4} \|F(U)(s)\|_{H^{2,1}} ds \\ & \leq C \left( \int_0^{t-1} (t-s)^{-3/2} (1+s)^{-3/2} (\log(1+s))^2 ds \right)^{1/2} (M_4(t) + M_5(t))N(0, \infty) \\ & \quad + C \int_0^{t-1} (t-s)^{-3/4} (1+s)^{-1} ds M_2(t)^2 \\ & \leq C(1+t)^{-3/4} \log(1+t) ((M_4(t) + M_5(t))N(0, \infty) + M_2(t)^2); \end{aligned}$$

and by (3.4) and (4.23),

(4.25)

$$\begin{aligned} & \|\partial_x^2(\mathbf{I} - \mathbf{P})I(t)\|_{L_2} \\ & \leq C \int_{t-1}^t (t-s)^{-1/2} \|F(U)(s)\|_{H^{2,1}} ds \\ & \leq C \int_{t-1}^t (t-s)^{-1/2} (1+s)^{-3/4} \log(1+s) ds (M_4(t) + M_5(t))N(0, \infty) \\ & \quad + C \int_{t-1}^t (t-s)^{-1/2} (1+s)^{-1} ds M_2(t)^2 \\ & \leq C(1+t)^{-3/4} ((M_4(t) + M_5(t))N(0, \infty) + M_2(t)^2). \end{aligned}$$

By (4.7), (4.10) and the Hölder's inequality

$$(4.26) \quad \begin{aligned} \|F(U)(s)\|_{L_2} &\leq C\|U(s)\|_{L_\infty}\|\partial_x U(s)\|_{H^1}, \\ \|F(U)(s)\|_{W_4^{1,0}} &\leq C(\|U(s)\|_{L_\infty} + \|\partial_x^2 \rho(s)\|_{L_2})\|\partial_x U(s)\|_{H^2}. \end{aligned}$$

Therefore, by (3.7) with  $(p, q) = (4, 2)$  and (4.26),

$$(4.27) \quad \begin{aligned} &\|II(t)\|_{L_\infty} \\ &\leq C \int_0^{t-1} (t-s)^{-3/4} \left( \|F(U)(s)\|_2 + \|F(U)(s)\|_{W_4^{1,0}} \right) ds \\ &\leq C \left( \int_0^{t-1} (t-s)^{-3/2} (1+s)^{-3/2} (\log(1+s))^2 ds \right)^{1/2} (M_4(t) + M_5(t))N(0, \infty) \\ &\leq C(1+t)^{-3/4} \log(1+t) (M_4(t) + M_5(t))N(0, \infty); \end{aligned}$$

and by (3.3), (4.26) and the Sobolev inequality

$$(4.28) \quad \begin{aligned} \|I(t)\|_{L_\infty} &\leq C\|I(t)\|_{W_4^1} \\ &\leq C \int_{t-1}^t (t-s)^{-1/2} \|F(U)(s)\|_{W_4^{1,0}} ds \\ &\leq C \int_{t-1}^t (t-s)^{-1/2} (1+s)^{-3/4} \log(1+s) ds (M_4(t) + M_5(t))N(0, \infty) \\ &\leq C(1+t)^{-3/4} \log(1+t) (M_4(t) + M_5(t))N(0, \infty). \end{aligned}$$

Combining (4.11), (4.24), (4.25), (4.11), (4.27), (4.28) with (4.8), we have

$$M_4(t) + M_5(t) \leq C(M_4(t) + M_5(t))N(0, \infty) + CM_2(t)^2 + C\|U_0\|_{H^2},$$

which means the Step 2. By Steps 1 and 2, the proof of Theorem 3 is completed.

### §5. Proof of Corollary 2.1

Let  $\tilde{N}(0, \infty)$  be the quantity defined in Section 2. Put

$$\begin{aligned} M_6(t) &= \sup_{0 \leq s \leq t} (1+s)^{-1/2} \|\partial_t \partial_x U(t)\|_{L_2}, \\ M_7(t) &= \sup_{0 \leq s \leq t} (1+s)^{-1/2} \|\partial_x^3 \mathbf{P}U(t)\|_{L_2}, \\ M_8(t) &= \sup_{0 \leq s \leq t} (1+s)^{-3/4} \|U(t)\|_{L_\infty}. \end{aligned}$$

By (2.2), when  $0 \leq t \leq 2$ , we have

$$(5.1) \quad \sum_{j=1}^8 M_j(t) \leq C\tilde{N}(0, \infty).$$

Therefore we consider the case when  $t \leq 1$ , below. By (3.8) with  $(p, q) = (2, 2)$ ,

$$(5.2) \quad \|\partial_t \partial_x e^{-tA} U_0\|_{L_2} \leq Ct^{-3/4} \|U_0\|_{H^{2,1}}.$$

By (3.8) with  $(p, q) = (2, 2)$  and (4.12),

$$(5.3) \quad \begin{aligned} \|\partial_t \partial_x II(t)\|_{L_2} &\leq C \int_0^{t-1} (t-s)^{-3/4} \|F(U)(s)\|_{H^{2,1}} ds \\ &\leq C \int_0^{t-1} (t-s)^{-3/4} (1+s)^{-1/2} \|\partial_x U(s)\|_{H^2} ds (M_2(t) + M_3(t)) \\ &\leq C \left( \int_0^{t-1} (t-s)^{-3/2} (1+s)^{-1} ds \right)^{1/2} (M_2(t) + M_3(t)) N(0, \infty) \\ &\leq C(1+t)^{-1/2} (M_2(t) + M_3(t)) N(0, \infty). \end{aligned}$$

By the equation

$$\partial_t \partial_x I(t) = \int_{t-1}^t \partial_x e^{-(t-s)A} \partial_s F(U)(s) ds$$

and by (3.2) and (4.12),

$$(5.4) \quad \begin{aligned} \|\partial_t \partial_x I(t)\|_{L_2} &\leq C \int_{t-1}^t (t-s)^{-1/2} \|\partial_s F(U)(s)\|_{H^{1,0}} ds \\ &\leq C \int_{t-1}^t (t-s)^{-1/2} (1+s)^{-1/2} ds (M_2(t) + M_3(t)) \tilde{N}(0, \infty) \\ &\leq C(1+t)^{-1/2} (M_2(t) + M_3(t)) \tilde{N}(0, \infty). \end{aligned}$$

Combining (5.1) through (5.4), we have

$$(5.5) \quad M_6(t) \leq C(M_2(t) + M_3(t))\tilde{N}(0, \infty) + C\|U_0\|_{H^{2,1}}.$$

Applying Proposition 4.1 to the second equation of (4.1), we have

$$\begin{aligned} \|\partial_x^3 \mathbf{P}U(t)\|_{L_2} &\leq C (\|\partial_t \partial_x \mathbf{P}U(t)\|_{L_2} + \|\partial_x^2 (\mathbf{I} - \mathbf{P})U(t)\|_{L_2} \\ &\quad + \|\partial_x \mathbf{P}F(U)(t)\|_{L_2} + \|\mathbf{P}U(t)\|_{H^1(\Omega \cap B_b)}), \end{aligned}$$

which together with Steps 1 and 2 in Section 4, (5.5) and (4.12) implies that

$$(5.6) \quad M_7(t) \leq C(M_2(t) + M_3(t))\tilde{N}(0, \infty) + C\|U_0\|_{H^2}.$$

Finally, we shall estimate  $M_8(t)$ . By (4.7), (4.10) and the Hölder's inequality

$$(5.7) \quad \|F(U)(s)\|_{W_4^{1,0}} \leq C(\|U(s)\|_{L^\infty}\|\partial_x U(s)\|_{H^2} + \|\partial_x^2 \rho(s)\|_{L_2}\|\partial_x v(s)\|_{H^2}).$$

Therefore, by (3.7) with  $(p, q) = (4, 2)$ , (4.26) and (5.7)

$$(5.8) \quad \begin{aligned} \|II(t)\|_{L^\infty} &\leq C \int_0^{t-1} (t-s)^{-3/4} \left( \|F(U)(s)\|_{L_2} + \|F(U)(s)\|_{W_4^{1,0}} \right) ds \\ &\leq C \left( \int_0^{t-1} (t-s)^{-3/2} (1+s)^{-3/2} ds \right)^{1/2} M_8(t)N(0, \infty) \\ &\quad + C \int_0^{t-1} (t-s)^{-3/4} (1+s)^{-5/4} \log(1+s) ds M_4(t)M_7(t) \\ &\leq C(1+t)^{-3/4} (M_8(t)N(0, \infty) + M_4(t)M_7(t)); \end{aligned}$$

and by (3.3), (5.7) and the Sobolev inequality

$$(5.9) \quad \begin{aligned} \|I(t)\|_{L^\infty} &\leq C\|I(t)\|_{W_4^1} \\ &\leq C \int_{t-1}^t (t-s)^{-1/2} \|F(U)(s)\|_{W_4^{1,0}} ds \\ &\leq C \int_{t-1}^t (t-s)^{-1/2} (1+s)^{-3/4} ds M_8(t)N(0, \infty) \\ &\quad + C \int_{t-1}^t (t-s)^{-1/2} (1+s)^{-5/4} \log(1+s) ds M_4(t)M_7(t) \\ &\leq C(1+t)^{-3/4} (M_8(t)N(0, \infty) + M_4(t)M_7(t)). \end{aligned}$$

Combining (5.1), (4.11), (5.8) and (5.9), we have

$$M_8(t) \leq C(M_8(t)N(0, \infty) + M_4(t)M_7(t)) + C\|U_0\|_{H^2}.$$

If  $CN(0, \infty) < 1$ , then we have

$$M_8(t) \leq CM_4(t)M_7(t) + C\|U_0\|_{H^2},$$

which completes the proof.

## References

- [1] Deckelnick, K., Decay estimates for the compressible Navier-Stokes equations in unbounded domain, *Math. Z.*, **209** (1992), 115–130.
- [2] ———,  $L^2$ -decay for the compressible Navier-Stokes equations in unbounded domains, *Comm. Partial Differential Equations*, **18** (1993), 1445–1476.
- [3] Hoff, D. and Zumbrun, K., Multi-dimensional diffusion waves for the Navier-Stokes equations of compressible flow, *Indiana Univ. Math. J.*, **44** (1995), 604–676.
- [4] ———, Pointwise decay estimates for multidimensional Navier-Stokes diffusion waves, *Z. Angew. Math. Phys.*, **48** (1997), 597–614.
- [5] Kobayashi, T., On a local energy decay of solutions for the equations of motion of viscous and heat-conductive gases in an exterior domain in  $\mathbf{R}^3$ , *Tsukuba. J. Math.*, **21** (1997), 629–670.
- [6] ———, On the local energy decay of higher derivatives of solutions for the equations of motion of viscous and heat-conductive gases in an exterior domain in  $\mathbf{R}^3$ , *Proc. Japan Acad. Ser. A*, **73** (A) (1997), 126–129.
- [7] ———, Some estimates of solutions for the equations of motion of compressible viscous fluid in an exterior domain in  $\mathbf{R}^3$ , *Preprint*.
- [8] Kobayashi, T. and Shibata, Y., Decay estimates of solutions for the equations of motion of compressible viscous and heat-conductive gases in an exterior domain in  $\mathbf{R}^3$ , *Comm. Math. Phys.*, **200** (1999), 621–659.
- [9] Liu, Tai-P. and Wang, W., The pointwise estimates of diffusion wave for the Navier-Stokes Systems in odd multi-dimensions, *Comm. Math. Phys.*, **196** (1998), 145–173.
- [10] Matsumura, A., An energy method for the equations of motion of compressible viscous and heat-conductive fluids, *University of Wisconsin-Madison, MRC Technical Summary Report*, **2194** (1981), 1–16.
- [11] Matsumura, A. and Nishida, T., The initial value problem for the equations of motion of compressible viscous and heat-conductive fluids, *Proc. Japan Acad. Ser. A*, **55** (1979), 337–342.
- [12] ———, The initial value problems for the equations of motion of viscous and heat-conductive gases, *J. Math. Kyoto Univ.*, **20-1** (1980), 67–104.
- [13] ———, Initial boundary value problems for the equations of motion of general fluids, *Comput. Methods Appl. Sci. Engrg.*, **V** (1981), 389–406
- [14] ———, Initial boundary value problems for the equations of compressible viscous and heat-conductive fluid, *Lecture Notes Numer. Appl. Anal.*, **5** (1982), 153–170.
- [15] ———, Initial boundary value problems for the equations of motion of compressible viscous and heat-conductive fluids, *Comm. Math. Phys.*, **89** (1983), 445–464.
- [16] Ponce, G., Global existence of small solutions to a class of nonlinear evolution equations, *Nonlinear Anal.*, **9** (1989), 339–418.
- [17] Weike, W., Large time behavior of solutions for general Navier-Stokes Systems in Multi-dimension, *Wuhan Univ. J. Nat. Sci.*, **2** (1997), 385–393.