

Fine properties of solutions to the Cauchy problem for a fast diffusion equation with Caffarelli–Kohn–Nirenberg weights

Matteo Bonforte and Nikita Simonov

Abstract. We investigate fine global properties of nonnegative, integrable solutions to the Cauchy problem for the fast diffusion equation with weights (WFDE) $u_t = |x|^\gamma \operatorname{div}(|x|^{-\beta} \nabla u^m)$ posed on $(0, +\infty) \times \mathbb{R}^d$, with $d \geq 3$, in the so-called *good fast diffusion range* $m_c < m < 1$, within the range of parameters γ, β which is optimal for the validity of the so-called Caffarelli–Kohn–Nirenberg inequalities.

It is natural to ask in which sense such solutions behave like the Barenblatt \mathfrak{B} (fundamental solution): for instance, asymptotic convergence, i.e. $\|u(t) - \mathfrak{B}(t)\|_{L^p(\mathbb{R}^d)} \xrightarrow{t \rightarrow \infty} 0$, is well known for all $1 \leq p \leq \infty$, while only a few partial results tackle a finer analysis of the tail behaviour. We characterize the maximal set of data $\mathcal{X} \subset L^1_+(\mathbb{R}^d)$ that produces solutions which are pointwise trapped between two Barenblatt (global Harnack principle), and uniformly converge in relative error (UREC), i.e. $d_\infty(u(t)) = \|u(t)/\mathfrak{B}(t) - 1\|_{L^\infty(\mathbb{R}^d)} \xrightarrow{t \rightarrow \infty} 0$. Such a characterization is in terms of an integral condition on $u(t = 0)$.

To the best of our knowledge, analogous issues for the linear heat equation, $m = 1$, do not possess such clear answers, but only partial results. Our characterization is also new for the classical, nonweighted FDE. We are able to provide *minimal rates of convergence* to \mathfrak{B} in different norms. Such rates are almost optimal in the nonweighted case, and become optimal for radial solutions. To complete the panorama, we show that solutions with data in $L^1_+(\mathbb{R}^d) \setminus \mathcal{X}$, preserve the same “fat” spatial tail for all times, hence UREC fails and $d_\infty(u(t)) = \infty$, even if $\|u(t) - \mathfrak{B}(t)\|_{L^1(\mathbb{R}^d)} \xrightarrow{t \rightarrow \infty} 0$.

Contents

1. Introduction and main results	2
2. Initial data in $L^1_{\gamma,+}(\mathbb{R}^d)$	15
3. Initial data in \mathcal{X} . Global Harnack principle and uniform convergence in relative error	23
4. Counterexamples and generalized global Harnack principle	36
5. On the fast diffusion flow in \mathcal{X}	42
A. Regularity results and some technical details	48
References	53

2020 *Mathematics Subject Classification.* Primary 35K55; Secondary 35B40, 35B45, 35K67, 35K65.

Keywords. Fast diffusion equation, Caffarelli–Kohn–Nirenberg weights, global Harnack inequalities, tail behaviour, asymptotic behaviour.

1. Introduction and main results

The purpose of this paper is to investigate fine decay properties of solutions to the Cauchy problem for the fast diffusion equation posed on $(0, +\infty) \times \mathbb{R}^d$,

$$\partial_t u = \Delta u^m \quad \text{with } m \in (0, 1),$$

which is a classical prototype of singular nonlinear diffusion in homogeneous media; see [92, 93]. Indeed, our results also cover a more general case, in which Caffarelli–Kohn–Nirenberg-type weights are allowed, also including the case of inhomogeneous media; for more details see [66, 86] and Section 1.1. Our goal is to understand the quantitative tail behaviour of nonnegative solutions, depending on the behaviour at infinity of the initial datum. In order to focus on the main questions and answers, we explain here a simplified version of our results. Some of them, in the nonweighted case, were already known (mostly in a nonsharp form) and in this case we provide a new proof; see for instance [17, 19, 21, 22, 24, 35, 68, 72, 91, 92] and also [48, 57].

It is well known that nonnegative and integrable solutions tend to behave like the Barenblatt profile \mathfrak{B} with the same mass [91–93]. However, the issue of making a precise and quantitative statement about such “similar behaviour” presents serious difficulties. Nonnegative integrable solutions (namely, solutions in $L^1_+(\mathbb{R}^d)$) converge to \mathfrak{B} in different norms, $\|u - \mathfrak{B}\| \xrightarrow{t \rightarrow \infty} 0$. However, none of these convergences provide enough information about the tails. Here we explore finer properties of such solutions to give an answer to the following question:

Q_1 : Do solutions in $L^1_+(\mathbb{R}^d)$ have the same tail as the Barenblatt (fundamental solution)?

When the answer to this question is positive we can ask for a finer convergence, namely whether or not the quotient $u/\mathfrak{B} \xrightarrow{t \rightarrow \infty} 1$ uniformly in \mathbb{R}^d . We call this uniform convergence in relative error (UREC), and we can state the main question in this direction as follows:

Q_2 : Do solutions in $L^1_+(\mathbb{R}^d)$ behave asymptotically as the Barenblatt, uniformly in relative error?

In this paper, we completely answer both Q_1 and Q_2 in the so-called good fast diffusion range, when $m \in (\frac{d-2}{d}, 1)$. This is the natural range of parameters to deal with these questions, since the Barenblatt profile \mathfrak{B} essentially represents the asymptotic behaviour of all integrable solutions and the mass is preserved along the flow. In order to provide answers to the above questions, we split the cone of nonnegative integrable initial data $L^1_+(\mathbb{R}^d)$ in two disjoint subspaces $\mathcal{X} \sqcup \mathcal{X}^c = L^1_+(\mathbb{R}^d)$, where \mathcal{X} is a set of functions satisfying a suitable tail condition specified below.

The main results of this paper can be roughly explained as a quantitative version of the following fact. On one hand, the answers to Q_1 and Q_2 are both affirmative *if and only if*

$u_0 \in \mathcal{X}$ (Theorems 1.1, 1.3 and Section 3). On the other hand, the answers to Q_1 and Q_2 are both no *if and only if* $u_0 \in \mathcal{X}^c$ (Theorem 1.5 and Section 4).

It is remarkable that such a complete answer can be given for a nonlinear equation, while, in the – a priori simpler – linear heat equation $m = 1$, things are not so clear. Partial – nonsharp – answers to Q_1 can be deduced from the representation formula, an extremely useful tool that we do not have at our disposal in the nonlinear case. As for Q_2 , the question seems to be completely open: to the best of our knowledge, there is no characterization of the class of initial data for which the corresponding solution converges to the Gaussian (with the same mass) uniformly in relative error. Some examples, in the negative direction, are shown in [95].

Our results are sharp and turn into an explicit characterization of the “tail condition” that the initial datum has to satisfy to be in \mathcal{X} (hence answering yes to Q_1 and Q_2), which amounts to requiring that

$$\sup_{R>0} R^{\frac{2}{1-m}-d} \int_{B_R^c(0)} |f(x)| \, dx < \infty,$$

or equivalently

$$\int_{B_{|x|/2}(x)} |f(y)| \, dy = O(|x|^{d-\frac{2}{1-m}}).$$

The proof of the equivalence of the above two conditions is not trivial, and indeed it requires one of our main results, Theorem 1.1; see Section 5.1. The latter condition was introduced by Vázquez in 2003 [91] to give a positive answer to Q_2 , and we show here that a posteriori it was a sharp one. Notice that this condition allows for a wider class of data than the (nonsharp) pointwise condition used in [17, 92], namely $u_0(x) \lesssim |x|^{-\frac{2}{1-m}}$; see Section 5.2. We also show that the above condition, when fulfilled by the data, is enough to prove polynomial rates of convergence in several norms. In the radial case, we deduce sharp rates of convergence in uniform relative error and we provide an answer to a question left open by Carrillo and Vázquez in [24]; see Remark 1.4.

Concerning initial data in \mathcal{X}^c , we show the existence of a class of solutions which exhibits, for all times, a *fat tail* (bigger than the Barenblatt’s). This is done by constructing explicit sub- and supersolutions. Such a class provides a negative answer to both Q_1 and Q_2 . Furthermore, we show that in \mathcal{X}^c , no (power-like) rate of convergence to the Barenblatt profile is possible; see Theorem 4.6, Section 4.3.

The results of this paper have important consequences. For instance, when $\gamma = \beta = 0$, they allow us to prove uniform Hölder continuity estimates for solutions to (CP), which, together with the characterization of GHP and UREC of Theorems 1.1 and 1.3 respectively, represent key tools in the proof of constructive and quantitative stability estimates for Gagliardo–Nirenberg–Sobolev inequalities: more details can be found in [11, Chapter 4]. Analogous considerations apply to the weighted case.

In the rest of this section we set the problem in its whole generality, also including equations with Caffarelli–Kohn–Nirenberg-type weights, and give precise statements of our results.

1.1. The setup of the problem and precise statement of the main results

In this paper we study the following Cauchy problem for the weighted fast diffusion equation (WFDE):

$$\begin{cases} \partial_t u = |x|^\gamma \nabla \cdot (|x|^{-\beta} \nabla u^m) & \text{in } (0, \infty) \times \mathbb{R}^d, \\ u(0, x) = u_0(x) & \text{in } \mathbb{R}^d, \end{cases} \quad (\text{CP})$$

where the parameters d, γ, β satisfy

$$d \geq 3, \quad \gamma < d, \quad \text{and} \quad \gamma - 2 < \beta \leq \gamma(d - 2)/d.$$

This is a natural restriction since it represents the optimal domain of validity of the so-called Caffarelli–Kohn–Nirenberg inequalities; see [14, 18]. The exponent m is in the so-called *good fast diffusive range*, namely

$$m \in (m_c, 1), \quad \text{where } m_c := \frac{d - 2 - \beta}{d - \gamma}. \quad (1.1)$$

From now on we will fix the parameters d, m, γ, β as above (unless explicitly stated).

Modelling and related results. Problem (CP) was introduced in the 1980s by Kamin and Roseau to model singular/degenerate diffusion in inhomogeneous media; see [65, 66, 86]. Since then, there has been a systematic study of similar equations, mostly in the case $m \geq 1$ and/or with only one weight; see [1, 7, 28, 45, 46, 53, 54, 58, 63, 64, 74, 80–83]. Recently, (CP) has proven to be an essential tool in the study of symmetry/symmetry breaking phenomena in Caffarelli–Kohn–Nirenberg inequalities; see [9, 10, 25, 39–43, 47]. Several intriguing connections between nonlinear diffusion on Riemannian manifolds and weighted parabolic equations were explored in [12, 13, 52, 59–61, 94].

Existence, uniqueness, comparison and mass conservation. The basic theory is well established: existence, uniqueness and comparison for nonnegative and bounded integrable data is well known; see [9, Section 2.2], where a suitable definition of weak solutions can also be found (cf. also [14, Definition 1.1]). In view of the smoothing effects of [14], it is straightforward to extend those results to weak solutions corresponding to merely integrable (and possibly unbounded) data

$$u_0 \in L_{\gamma,+}^1(\mathbb{R}^d) = \{u_0: \mathbb{R}^d \rightarrow \mathbb{R} : u_0 \geq 0, \int_{\mathbb{R}^d} u_0 |x|^{-\gamma} dx < \infty\}. \quad (1.2)$$

We refrain from giving further details that would involve weighted Sobolev spaces, which we never use in this paper and we choose not to define here. What we want to emphasize is that data in $L_{\gamma,+}^1$ produce solutions that turn out to be bounded, positive and regular (at least Hölder continuous; see also Appendix A.3 and [14]) and that solutions considered in this paper possess enough regularity to guarantee the validity of all the calculations performed here. We also recall that in the good fast diffusive range, nonnegative integrable solutions conserve mass along the flow,

$$\int_{\mathbb{R}^d} u(t, x) |x|^{-\gamma} dx = \int_{\mathbb{R}^d} u_0(x) |x|^{-\gamma} dx =: M \quad \text{for any } t \geq 0.$$

For a proof, see [10, Section 2.2] and [14, Proposition 2.4 and the Remark following it].

The *fundamental solution* is of *self-similar type* and it is often called a Barenblatt solution:

$$\begin{aligned} \mathfrak{B}(t + T, x; M) &= \frac{\zeta^{d-\gamma}}{R_\star(t + T)^{d-\gamma}} \mathfrak{B}_M\left(\frac{\zeta x}{R_\star(t + T)}\right) \\ &= \frac{(t + T)^{\frac{1}{\vartheta}}}{\left[b_0 \frac{(t+T)^{\sigma\vartheta}}{M^{\sigma\vartheta(1-m)}} + b_1 |x|^\sigma\right]^{\frac{1}{1-m}}}, \end{aligned} \quad (\text{B})$$

where

$$\sigma := 2 + \beta - \gamma, \quad \frac{1}{\vartheta} = (d - \gamma)(m - m_c), \quad \zeta^{\frac{1}{\vartheta}} = \frac{1 - m}{\sigma m} \quad \text{and} \quad R_\star(t) = \left(\frac{t}{\vartheta}\right)^\vartheta, \quad (1.3)$$

the parameter M is the *mass* of the solution, T is a free parameter and b_0, b_1 are constants which depend on m, d, γ, β . The profile \mathfrak{B}_M is given by

$$\mathfrak{B}_M = (C(M) + |x|^\sigma)^{\frac{1}{m-1}}, \quad (1.4)$$

where $C(M)$ depends on M, d, m, γ, β , and has an explicit expression; see Appendix A.1. In what follows we will frequently use the solution (B) with the parameter $T = 0$. Recall that by the very definition of fundamental solutions we have $\mathfrak{B}(0, x; M) = M\delta_0$, in the sense of measures: the computation goes as for the standard FDE, and the extra weight $|x|^{-\gamma}$ does not cause any problem. Also, we will sometimes drop the dependence on the x variable and write $\mathfrak{B}(t; M)$ or $\mathfrak{B}(t, \cdot; M)$ when no confusion arises.

The tail condition. We say that $f \in L_\gamma^1(\mathbb{R}^d)$ satisfies the *tail condition* – or equivalently that $f \in \mathcal{X}$ – if

$$|f|_{\mathcal{X}} := \sup_{R>0} R^{\frac{2+\beta-\gamma}{1-m} - (d-\gamma)} \int_{B_R^c(0)} |f(x)| |x|^{-\gamma} dx < \infty. \quad (\text{TC})$$

Recall that since $m \in (m_c, 1)$ we have $\frac{2+\beta-\gamma}{1-m} - (d-\gamma) > 0$. We emphasize that \mathcal{X} *does* depend on the parameters m, γ and β ; however, in order to make the notation easier to read, we have dropped such dependencies. It is easily seen that $|\cdot|_{\mathcal{X}}$ is a norm. Intuitively, the quantity $|f|_{\mathcal{X}}$ measures *how fast* the function f decays at ∞ relative to the decay of the Barenblatt profile \mathfrak{B}_M . We now introduce a subspace of $L_\gamma^1(\mathbb{R}^d)$ of functions that satisfy the tail condition (TC), which will play a key role in the rest of the paper:

$$\mathcal{X} := \{u \in L_\gamma^1(\mathbb{R}^d) : |u|_{\mathcal{X}} < +\infty\}. \quad (1.5)$$

We adapt to our setting an alternative tail condition proposed by Vázquez [91]: we say that $f \in L_\gamma^1(\mathbb{R}^d)$ satisfies (TC') if

$$\int_{B_{\frac{|x|}{2}}(x)} |f(y)| |y|^{-\gamma} dy = O(|x|^{d-\gamma-\frac{2+\beta-\gamma}{1-m}}). \quad (\text{TC}')$$

We will show in Section 5.1 that (TC) and (TC') are indeed equivalent.

We will provide now a precise and sharp answer to questions Q_1 and Q_2 , in the form of our main results.

The space \mathcal{X} : Affirmative answers to Q_1 and Q_2 , and a characterization. As we explained in the introduction, the answers to both Q_1 and Q_2 are affirmative *if and only if* the initial data is in \mathcal{X} . The main tool in providing such answers is the so-called *global Harnack principle* (GHP): a lower and upper bound in terms of Barenblatt profiles; see Theorem 1.1 below. The GHP provides a complete answer both to Q_1 and, surprisingly, also to Q_2 , as we will see later. In the nonweighted case, the GHP in the form of Theorem 1.1 was introduced in [17] (under the stronger pointwise assumption $u_0(x) \lesssim |x|^{-\frac{2}{1-m}}$) and was inspired by the pioneering results of [91], in which condition (TC') was introduced.

Our main contribution in this case consists of the *characterization* of the maximal set \mathcal{X} of initial data that generate solutions satisfying the GHP.

We will see in Section 5.3 that the space \mathcal{X} is *invariant* under WFDE flow: indeed, $u(t) \in \mathcal{X}$ if and only if $u(0) \in \mathcal{X}$, and the same holds for \mathcal{X}^c ; see Proposition 5.3 and Theorems 1.1 and 1.5.

Theorem 1.1 (Characterization of the GHP). *Let $m \in (m_c, 1)$ and let u be a solution to (CP) with $u_0 \in L^1_{\gamma,+}(\mathbb{R}^d)$. Then, for any $t_0 > 0$, there exist $\bar{\tau}, \underline{\tau} > 0$ and $\bar{M}, \underline{M} > 0$ such that*

$$\mathfrak{B}(t - \underline{\tau}, x; \underline{M}) \leq u(t, x) \leq \mathfrak{B}(t + \bar{\tau}, x; \bar{M}) \quad \text{for any } x \in \mathbb{R}^d \text{ and } t \geq t_0 \quad (1.6)$$

if and only if

$$u_0 \in \mathcal{X} \setminus \{0\},$$

where $L^1_{\gamma,+}(\mathbb{R}^d)$ is as in (1.2), m_c is as in (1.1), $\mathfrak{B}(t, x; M)$ is as in (B) and \mathcal{X} is as in (1.5). The values $\underline{M}, \underline{\tau}$ are as in (2.3) while \bar{M} and $\bar{\tau}$ can be found in (3.10).

Remark 1.2. The proof of the above result is quantitative and provides explicit expressions for $\underline{\tau}, \bar{\tau}, \underline{M}, \bar{M}$. It follows by combining the upper bound of Theorem 3.1 with the lower bound of Theorem 2.1. For the upper bound the hypothesis $0 \leq u_0 \in \mathcal{X}$ is strictly necessary. Indeed, for data $u_0 \notin \mathcal{X}$ we are able to construct explicit (sub)solutions that provide precise lower bounds that clearly contradict the upper bound of formula (1.6). More precisely, for any $t > 0$ and for any $x \in \mathbb{R}^d$ we have

$$u(t, x) \geq \frac{1}{(D(t) + |x|^\sigma)^{\frac{1}{1-m} - \varepsilon}} \geq \mathfrak{B}(t, x; M),$$

where $\varepsilon > 0$ is small, and $D(t) \sim t^{\frac{2}{\varepsilon(1-m)}}$.

On the other hand, such a hypothesis is not necessary for the lower bound of formula (1.6): indeed, lower bounds hold for any data $0 \leq u_0 \in L^1_{\gamma,\text{loc}}(\mathbb{R}^d)$, and produce a *minimal lower tail*; see Theorem 2.1 and Corollary 2.2. This provides a partial answer to Q_1 .

Let us turn our attention to question Q_2 . The convergence of the solution to the Barenblatt profile has been studied by many researchers, under different sets of assumptions, especially in the nonweighted case $\gamma = \beta = 0$; see for instance [6, 13, 19, 24, 48, 68, 72, 91] and references therein. We will discuss now some of the existing results which are strictly

related to ours, but all in the nonweighted case. To the best of our knowledge, no results for the weighted case are present in the literature, except for some partial results in [9, 10, 14]. In [48] the authors proved uniform convergence on expanding sets of the form $|x| \leq Ct^\vartheta$, namely

$$\lim_{t \rightarrow 0} \sup_{x \in \{|x| \leq Ct^\vartheta\}} \left| \frac{u(t, x) - \mathfrak{B}(t, x; M)}{\mathfrak{B}(t, x; M)} \right| = 0, \quad (1.7)$$

under the condition $u_0 \in L^1(\mathbb{R}^d) \cap L^2(\mathbb{R}^d)$. We will prove an analogous result in the weighted case; see Theorem 2.4. Recently, Vázquez [91] has completed the proof of the previous result for the whole class of positive initial data which belong to $L^1(\mathbb{R}^d)$; he also shows uniform convergence in $L^\infty(\mathbb{R}^d)$ and in $L^1(\mathbb{R}^d)$. In [91] Vázquez proved that UREC takes place for all data which satisfy the pointwise condition $u_0(x) \lesssim |x|^{-\frac{2}{1-m}}$, and he also introduces (TC'). In 2003, Carrillo and Vázquez [24] obtained the estimates

$$\sup_{x \in \mathbb{R}^d} \left\| \frac{u(t, x)}{\mathfrak{B}(t, x; M)} - 1 \right\|_{L^\infty(\mathbb{R}^d)} \leq \frac{C(u_0)}{t} \quad (1.8)$$

for radial initial data which satisfy the pointwise condition $u_0(x) \lesssim |x|^{-\frac{2}{1-m}}$. An intriguing open question was left on [24, p. 1027]: to extend the validity of (1.8) to a larger class of initial data. The question was partially answered in [6, 10, 68] in some nonoptimal classes of data, possibly nonradial.

Our main contribution in this paper is to *characterize* the maximal set \mathcal{X} of initial data whose solution converges to the Barenblatt profile *uniformly in relative error*.

Theorem 1.3 (Characterization of the UREC). *Let $m \in (m_c, 1)$ and let u be a solution to (CP) with initial data $u_0 \in L^1_{\gamma,+}(\mathbb{R}^d)$ and $M = \|u_0\|_{L^1_{\gamma,+}(\mathbb{R}^d)}$. Then*

$$\lim_{t \rightarrow \infty} \left\| \frac{u(t, x)}{\mathfrak{B}(t, x; M)} - 1 \right\|_{L^\infty(\mathbb{R}^d)} = 0 \quad (1.9)$$

if and only if

$$u_0 \in \mathcal{X} \setminus \{0\}.$$

Remark 1.4 (Sharp convergence rates for radial solutions). We notice here that if $f \in \mathcal{X}_{\text{rad}}$, the class of radial functions in \mathcal{X} , then it does not necessarily satisfy the pointwise condition $f \lesssim |x|^{-\frac{2}{1-m}}$; see Section 5.2. Hence, our Theorem 3.8, which shows the validity of (1.8) for any $u_0 \in \mathcal{X}_{\text{rad}}$, provides a sharp answer to the question raised by Carrillo and Vázquez. The maximality of \mathcal{X} is guaranteed by Theorem 1.3. Indeed, if $u_0 \notin \mathcal{X}$ then the limit (1.9) is infinite; see Proposition 5.3.

In Theorem 3.6 we provide *almost optimal rates* of convergence for all data in \mathcal{X} in the nonweighted case, valid also for nonradial solutions. Analogously, Theorem 3.7 shows minimal rates in the weighted case. Sharp rates of convergence can be obtained under more restrictive assumptions (but for the whole range $m < 1$): this happens if the initial datum is trapped between two Barenblatt solutions with exactly the same tail (which is stronger than the GHP of Theorem 1.1). We refer to [6, 8, 9, 13] and references therein for an overview of previous results; see also [8] for a brief historical overview.

The space \mathcal{X}^c : Negative answer to Q_1 and Q_2 . Counterexamples. In order to complete the panorama, we still have to answer the next natural question: *What happens to the solutions with data in \mathcal{X}^c ?* On one hand, the space \mathcal{X}^c is also *invariant* under WFDE flow: indeed, $u(t) \in \mathcal{X}^c$ if and only if $u_0 \in \mathcal{X}^c$. Moreover, the uniform relative error (1.10) is always infinite; see Proposition 5.3. As a consequence, the answers to Q_1 and Q_2 are definitively negative in \mathcal{X}^c . On the other hand, we will show that – somehow stable – *anomalous tail behaviour* can happen in this case. Let us begin with an illuminating example in the simplest possible case, when $\gamma = \beta = 0$. Let $m > \frac{d}{d+2}$; consider the solution $w(t, x)$ with initial data

$$w_0 = \frac{1}{(1 + |x|^2)^{\frac{m}{1-m}}}.$$

It is clear that w_0 does not satisfy the assumption of Theorem 3.1 and, for $|x|$ large enough, we have that $w_0(x) > \mathfrak{B}(t_0, x; M)$ for any $t_0, M > 0$. However, $w_0 \in L^1(\mathbb{R}^d)$ whenever $m > \frac{d}{d+2}$. The tail behaviour of $w(t, x)$ is strongly different from the Barenblatt profiles, and this can be better appreciated in logarithmic scale (see for instance Figures 1 and 2); indeed for all $t > 0$,

$$\frac{1}{((t+1)^{\frac{1}{1-m}} + |x|^2)^{\frac{m}{1-m}}} \lesssim w(t, x) \lesssim \frac{(1+t)^{\frac{m}{1-m}}}{(1+t + |x|^2)^{\frac{m}{1-m}}} \quad \text{for all } x \in \mathbb{R}^d.$$

The above inequality gives us remarkable insights into the long time behaviour of the solution $w(t, x)$. First, for any time $t > 0$, $w(t, x)$ has power-like behaviour at infinity different from the Barenblatt one, namely as $|x| \rightarrow \infty$ we have $w(t, x) \sim |x|^{\frac{-2m}{1-m}}$ versus $\mathfrak{B}(t, x; M) \sim |x|^{\frac{-2}{1-m}}$. The upper part of GHP fails outside a space-time region that we explicitly identify, as a consequence of this anomalous tail behaviour. Indeed,

$$\sup_{x \in \mathbb{R}^d} \left| \frac{w(t, x)}{\mathfrak{B}(t, x; M)} - 1 \right| = \infty, \quad (1.10)$$

where $\mathfrak{B}(t, x; M)$ has the same mass as $w(t, x)$. The same considerations apply by replacing $\mathfrak{B}(t, x; M)$ with any other Barenblatt solution. Obviously, *uniform convergence in relative error (UREC) fails*.

The anomalous behaviour found in this particular example is indeed shared by an entire class of solutions. We prove here a *generalized* version of the GHP, valid for initial data decaying slower than the Barenblatt profile. The proof is based on the construction of two families of sub- and supersolutions. We will cover all admissible γ, β and $m \in (m_c, 1)$, extending the above considerations to the weighted case, as in the following theorem.

Theorem 1.5 (Generalized global Harnack principle). *Let $m \in (m_c, 1)$, $\varepsilon \in (0, \frac{2}{1-m} - \frac{2}{\sigma}(d - \gamma))$ and $\alpha = \frac{1}{1-m} - \frac{\varepsilon}{2} > 0$. Assume that the initial data u_0 satisfies*

$$\frac{A}{(\underline{t}^{1-\alpha(1-m)} + B|x|^\sigma)^\alpha} \leq u_0(x) \leq \frac{E\bar{t}^\sigma}{(\bar{t} + F|x|^\sigma)^\alpha},$$

for some $A, B, E, F, \underline{t}, \bar{t} > 0$. Then for any $t > 0$ we have

$$\underline{V}(t, x) := \frac{A}{(D(t) + B|x|^\sigma)^\alpha} \leq u(t, x) \leq \frac{EG(t)^\alpha}{(G(t) + F|x|^\sigma)^\alpha} =: \bar{V}(t, x),$$

where

$$D(t) := (\sigma A^{m-1} m B (d - \gamma) (1 - \alpha(1 - m)) t + \underline{t})^{\frac{1}{1 - \alpha(1 - m)}} \quad \text{and} \quad G(t) := \bar{t} + Ht,$$

where $H \geq m\sigma F^2 E^{m-1} (2 + \beta - d + \sigma\alpha m)$.

The proof of the above theorem is just the combination of the results of Propositions 4.1 and 4.5, where explicit expressions for all the constants are given.

Remark. The above theorem shows that an initial “fat-tail” is preserved for all times. This marks a clear difference between the “good” range $(m_c, 1)$ and the very fast diffusive range $(0, m_c)$: in the latter case there can be solutions with power-like tails which change with time; see for instance [34].

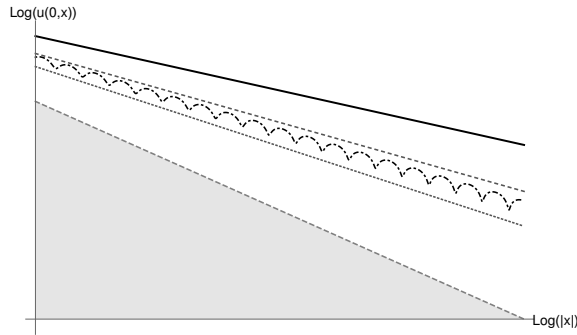


Figure 1. The picture represents different possible power-like tail behaviour of an initial datum in dimension $d = 3$, $\gamma = \beta = 0$ and $m = 2/3$. Every line represents a power-like behaviour. The dashed “sinusoidal” curve represents a function whose behaviour is trapped between two different tail powers. The thick line represents the power $|x|^{-3}$, the dashed line $|x|^{-3.8}$, the dotted line $|x|^{-4.4}$ and the grey region represents any decay below the Barenblatt decay ($|x|^{-6}$). In both plots $|x| \in [10^3, 10^8]$.

The space $L_{\gamma,+}^1$. General picture. We provide here a general picture for solutions in $L_{\gamma,+}^1$, depending on the spatial decay of the initial data. This is better understood in the $\log(u(x))$ – $\log(x)$ plot in Figure 1, where different kinds of possible tail behaviours are represented. For instance, the Barenblatt profile \mathcal{B} , marked in dashed-grey, corresponds to the curve $\log(\mathcal{B}(x)) = -\frac{2+\beta-\gamma}{1-m} \log|x| + o(\log(|x|))$. The different lines represent other possible power-like tail behaviours; the thick line is the natural barrier for solutions to $L_{\gamma,+}^1(\mathbb{R}^d)$ since it corresponds to the case $|x|^{-(d-\gamma)}$.

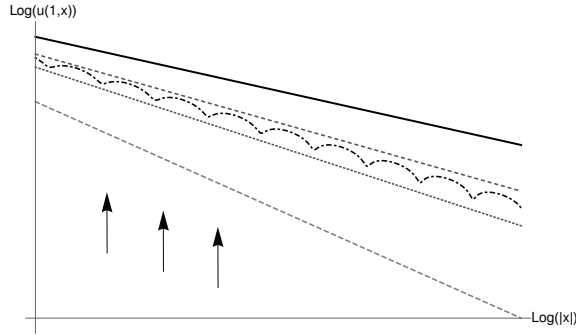


Figure 2. General panorama for $L^1_{\gamma,+}$ solutions. The picture represents different possible power-like tail behaviours in dimension $d = 3$, $\gamma = \beta = 0$ and $m = 2/3$ for a solution $u(t, x)$ at time $t = 1$. We can appreciate that (i) there are no solutions below the lowest line $|x|^{-6}$, the one corresponding to the Barenblatt behaviour; (ii) if the tail of the initial data is a line above the Barenblatt then the solution preserves the same tail behaviour for all times; (iii) the dashed curve remains trapped between the same initial power-like behaviours.

Let us begin our analysis. As we have already explained, every nonnegative solution to (CP) develops a minimal power-like tail, at least $|x|^{-\frac{\sigma}{1-m}}$. Therefore, in Figure 2 there are no solutions below the dashed-grey line. Initial data in \mathcal{X} develop exactly the Barenblatt tail (see Theorem 1.1), hence, roughly speaking “they live on the dashed-grey line”, the \mathcal{X} -curve. Things are different for initial data in \mathcal{X}^c . We only analyze power-like behaviours at infinity. Roughly speaking, the generalized GHP, Theorem 1.5, says that all solutions must live between the dotted and dashed lines (recall that the plain line corresponds to the case $|x|^{-(d-\gamma)}$). More precisely, the generalized GHP tells us that solutions with data decaying like $|x|^{-\alpha}$, with $(d - \gamma) < \alpha < \frac{\sigma}{1-m}$, will have the same decay $|x|^{-\alpha}$. Indeed, Theorem 1.5 tells us more: any initial datum in \mathcal{X}^c with tail behaviour trapped between two different lines (maybe oscillating between two power tails at infinity) produces a solution trapped between the same lines. For instance, if the datum is between the dotted and dashed lines, then the solution is trapped between those barriers for all times.

1.2. A dynamical system interpretation

The aim of this section is to describe a global picture of the fine behaviour of the solutions to the WFDE, in terms of convergence to equilibrium states of an (infinite-dimensional) dynamical system. It is convenient to pass to self-similar variables in order to make stationary the “asymptotic” Barenblatt solution.

Self similar variables. Nonlinear Fokker–Plank equation. Let $u(t, x)$ be a solution to (CP) with initial data u_0 , and consider $R(t) = R_\star(t + 1)$. The *self-similar change of*

variables

$$v(\tau, y) = \frac{R(t)^{d-\gamma}}{\zeta^{d-\gamma}} u(t, x), \quad \text{where } \tau = \frac{1}{\sigma} \log \frac{R(t)}{R(0)}, \quad y = \frac{\zeta x}{R(t)}, \quad (1.11)$$

transforms $u(t, x)$ into a solution to the nonlinear *Fokker–Planck*-type equation

$$\frac{\partial v}{\partial \tau} + |x|^\gamma \operatorname{div}(|x|^{-\beta} v \nabla v^{m-1}) = |x|^\gamma \operatorname{div}(|x|^{-\beta} v \nabla |x|^\sigma), \quad (\text{NLWFP})$$

with initial data $v_0(y) = \frac{\zeta^{d-\gamma}}{R(0)^{d-\gamma}} u_0(\frac{\zeta x}{R(0)})$, with the same mass. Also notice that among all the Barenblatt profiles $\mathfrak{B}(t + \tau, x; M)$, only the one with $\tau = 1$ becomes stationary, and we call it Barenblatt profile $\mathcal{B}_M(y)$: this is the unique attractor or the unique equilibrium (asymptotically stable).

In what follows we will assume that solutions to (CP) with initial data $u_0 \in L^1_{\gamma,+}(\mathbb{R}^d)$ will converge to the Barenblatt solution in the following sense (recall that the mass is preserved along the flow):

$$\|u(t) - \mathfrak{B}(t; M)\|_{L^1_{\gamma}(\mathbb{R}^d)} \rightarrow 0 \quad \text{as } t \rightarrow \infty.$$

In *self-similar* variables the previous result can be restated as

$$\|v(\tau) - \mathcal{B}_M\|_{L^1_{\gamma}(\mathbb{R}^d)} \rightarrow 0 \quad \text{as } \tau \rightarrow \infty.$$

Dynamical system approach: “Infinite-dimensional phase plane analysis” and the space \mathfrak{X} . It is well known that (NLWFP) can be seen as the gradient flow of an entropy functional; cf. [71, 76]. It can be shown that solutions corresponding to nonnegative initial data will converge to a stationary solution with the same mass. To be more precise, let us define the ω -limit of the (NLWFP) as the one-dimensional manifold of the so-called Barenblatt solutions:

$$\mathcal{M}_{\mathcal{B}} := \{\mathcal{B}_M : M > 0\}, \quad (1.12)$$

and the distance

$$d_1(f) := \inf_{\mathcal{B}_M \in \mathcal{M}_{\mathcal{B}}} \|f - \mathcal{B}_M\|_{L^1_{\gamma}(\mathbb{R}^d)}.$$

It has been proven in [48, 91] (for the case $\gamma = \beta = 0$) and in [9, 10] (for the weighted case) that for any $u_0 \in L^1_{\gamma,+}(\mathbb{R}^d)$ there exists a unique \mathcal{B}_{M_0} (M_0 being the mass of u_0) such that

$$d_1(v(t)) \leq \|v(t) - \mathcal{B}_{M_0}\|_{L^1_{\gamma}(\mathbb{R}^d)} \xrightarrow{t \rightarrow \infty} 0.$$

Hence, solutions of (NLWFP) can be seen as a continuous path with respect to the $L^1_{\gamma,+}$ topology, that will eventually converge to a point of the manifold $\mathcal{M}_{\mathcal{B}}$. See Figure 3. This fact can be rephrased as follows: *the basin of attraction of $\mathcal{M}_{\mathcal{B}}$ in the $L^1_{\gamma,+}$ -topology is the whole space $L^1_{\gamma,+}$.*

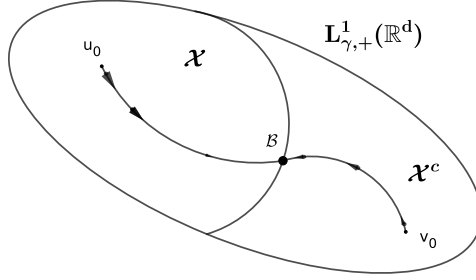


Figure 3. We represent two possible paths in $L^1_{\gamma,+}$. Since \mathcal{X} and \mathcal{X}^c are invariant sets for the flow there are no crossing lines between them. We notice that the manifold $\mathcal{M}_{\mathcal{B}}$ is contained in the topological boundary (with respect to the $L^1_{\gamma,+}$ topology) of \mathcal{X} , $\mathcal{M}_{\mathcal{B}} \subset \partial_{L^1_{\gamma,+}} \mathcal{X}$.

We can ask a similar question for a stronger convergence that allows us to have better asymptotic knowledge of the tails, the *uniform convergence in relative error* (UREC), properly measured by the following distance from $\mathcal{M}_{\mathcal{B}}$:

$$d_{\infty}(f) := \inf_{\mathcal{B}_M \in \mathcal{M}_{\mathcal{B}}} \left\| \frac{f}{\mathcal{B}_M} - 1 \right\|_{L^{\infty}(\mathbb{R}^d)}.$$

The above distance is induced by the norm

$$\|f\|_{m,\gamma,\beta} := \sup_{x \in \mathbb{R}^d} |f(1 + |x|^{2+\beta-\gamma})^{\frac{1}{1-m}}|,$$

hence the topology is equivalent and we will refer to it as relative error topology. As a consequence of L^1_{γ} convergence, $v(\tau, x) \rightarrow \mathcal{B}_M(x)$ as $\tau \rightarrow \infty$ a.e., hence also pointwise, i.e. $v(\tau, x)\mathcal{B}_M^{-1}(x) \rightarrow 1$ as $\tau \rightarrow \infty$. However the *uniform convergence in relative error* (UREC) may fail. The main results of this paper solve this issue: \mathcal{X} is the *basin of attraction of the manifold $\mathcal{M}_{\mathcal{B}}$ in the relative error topology*. Notice that \mathcal{X} is defined in terms of a practical – easy-to-check – condition on the initial datum, which a priori does not have any relation with the asymptotic behaviour. In what follows we explain our main results in terms of assumptions on the initial datum. Our main contribution in this direction is that we show that *only three things can happen*: (i) $v_0 \in \mathcal{X}^c$, (ii) $v_0 \in \mathcal{X}$ and $d_{\infty}(v_0) < \infty$, (iii) $v_0 \in \mathcal{X}$ and $d_{\infty}(v_0) = \infty$. We analyze each case separately. See also Figure 4.

- (i) *If $v_0 \in \mathcal{X}^c$.* Roughly speaking, in this case we show that if the initial datum v_0 has a tail strictly above the Barenblatt one, then that “fat tail” is preserved in time. More precisely, Proposition 5.3 implies that

$$v_0 \in \mathcal{X}^c \Rightarrow \left\| \frac{v(\tau)}{\mathcal{B}_M} - 1 \right\|_{L^{\infty}(\mathbb{R}^d)} = +\infty \Rightarrow v(\tau) \in \mathcal{X}^c \text{ and } d_{\infty}(v(\tau)) = \infty$$

for all $\tau > 0$. In particular, since $d_{\infty}(v_0) = \infty$, this reveals that it is impossible to have bounds of the form $d_{\infty}(v(\tau)) \leq d_1(v(\tau))$ if we do not have them already

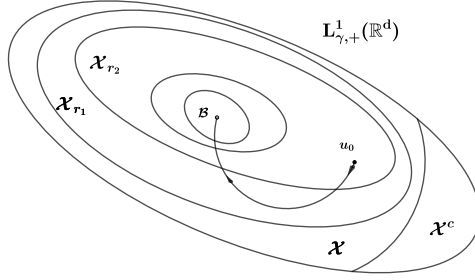


Figure 4. Illustration of the stability of the sets \mathcal{X}_r : if the solution starts from one of those sets, it will forever stay in one of those sets. Indeed, if $u_0 \in \mathcal{X}_r$ there exists a maximal $\mathcal{X}_{\bar{r}}$ which is invariant under the flow, i.e. such that $u(t) \in \mathcal{X}_{\bar{r}}$ for all $t \geq 0$.

(at least) for the initial datum. We can appreciate here the strong difference along the flow between the $L^1_{\gamma,+}(\mathbb{R}^d)$ and the $\|\cdot\|_{m,\gamma,\beta}$ topologies. On one hand, any $v_0 \in L^1_{\gamma,+}$ is sent by the flow to a unique element of $\mathcal{M}_{\mathcal{B}}$ in the d_1 distance. On the other hand, this is not true in the d_∞ distance, in which case the flow stays always at an infinite d_∞ -distance from $\mathcal{M}_{\mathcal{B}}$.

- (ii) If $v_0 \in \mathcal{X}$ and $d_\infty(v_0) < \infty$. This is the stable case: if the initial datum is close to the manifold $\mathcal{M}_{\mathcal{B}}$, then the flow will stay close to it and eventually d_∞ -converge to a unique element of $\mathcal{M}_{\mathcal{B}}$. More precisely, recall that $\mathcal{M}_{\mathcal{B}} \subset \mathcal{X}$ is the ω -limit set, made of stationary solutions or equilibria of our dynamical system. The GHP of Theorem 1.1, together with the UREC of Theorem 1.3, implies

$$\begin{aligned} v_0 \in \mathcal{X} \Rightarrow \left\| \frac{v(\tau)}{\mathcal{B}_M} - 1 \right\|_{L^\infty(\mathbb{R}^d)} < \infty \quad \text{for all } \tau > 0 \Rightarrow v(\tau) \in \mathcal{X} \\ \Rightarrow d_\infty(v(\tau)) \xrightarrow{\tau \rightarrow \infty} 0. \end{aligned}$$

Indeed the GHP tells us a global stability result for the flow, since it can be rewritten as $d_\infty(v(t)) \leq F(\|v_0\|_{\mathcal{X}})$, for some locally bounded function F . A finer analysis is performed below.

- (iii) If $v_0 \in \mathcal{X}$ and $d_\infty(v_0) = \infty$. We show that even if the initial datum is at infinite distance from the manifold $\mathcal{M}_{\mathcal{B}}$, but still in \mathcal{X} , the solution will eventually d_∞ -converge to $\mathcal{M}_{\mathcal{B}}$. Indeed, the GHP of Theorem 1.1 only needs the assumption $v_0 \in \mathcal{X}$, regardless of $d_\infty(v_0) = \infty$. Hence the same argument as case (ii) applies.

Finer analysis in \mathcal{X} . It is possible to show that

$$\mathcal{X} = \{d_\infty = \infty\} \cup \bigcup_{r>0} \mathcal{X}_r^\infty = \{d_\infty = \infty\} \cup \bigcup_{r>0} \{f \in \mathcal{X} : d_\infty(f) < r\}.$$

The GHP of Theorem 1.1 reveals an important stability property of fast diffusion flows: for any $v_0 \in \mathcal{X}$ there exists $r_0, \tau_0 > 0$ such that $d_\infty(v(\tau)) < r_0$ for all $\tau \geq \tau_0$, hence the flow

never exits from a certain $\mathcal{X}_{r_0}^\infty$. Indeed, we show more: $d_\infty(v(\tau)) \rightarrow 0$ as $\tau \rightarrow \infty$, which means that the flow always leaves the manifolds $d_\infty(v(\tau)) = r$ (level sets of distance from $\mathcal{M}_{\mathcal{B}}$) to enter one at a lower level, say $d_\infty(v(\tau)) = r - \varepsilon$.

This can be summarized as follows: we show that the solution map immediately sends \mathcal{X} to a more regular subspace $\mathcal{X}_{r_0}^\infty$,

$$T_\tau : \mathcal{X} \rightarrow \bigcup_{r>0} \mathcal{X}_r^\infty, \quad \text{indeed, there exists } r_0 > 0 : T_\tau : \mathcal{X} \rightarrow \mathcal{X}_{r_0}^\infty.$$

On one hand, in the relative error topology we have a dichotomy: $\lim_{\tau \rightarrow \infty} T_\tau(\mathcal{X} \setminus \{0\}) = \mathcal{M}_{\mathcal{B}}$ and $d_\infty(T_\tau(\mathcal{X}^c)) = \infty$ for all $\tau > 0$. On the other hand, in the d_1 -topology we always have $\lim_{\tau \rightarrow \infty} T_\tau(L_{\gamma,+}^1(\mathbb{R}^d) \setminus \{0\}) = \mathcal{M}_{\mathcal{B}}$.

1.3. Others ranges of m and generalizations

In the study of (CP) a dramatic change occurs when we consider the exponent $m \leq m_c$ or $m \geq 1$. When $m = 1$, it is known that solutions to the Cauchy problem for the classical heat equation eventually develop Gaussian tails whenever the initial data is compactly supported. Surprisingly enough, no sharp results are known in this context, to the best of our knowledge. Uniform convergence in relative error is (in general) false [95] and several tail behaviours are indeed possible [56]. The porous medium case $m > 1$ has been widely investigated; see the monograph [93]. This is the slow diffusive case, with finite speed of propagation, and it is hard to properly define the relative error. To our knowledge, the best result in this direction has been obtained in [67], by means of a careful study of the evolution of the support of the solution. Let us now comment on the very fast diffusion case $m \leq m_c$, where several new difficulties arise such as, for instance, loss of mass and extinction in finite time. In the case $m = m_c$, conservation of mass still holds but nevertheless the asymptotic behaviour becomes quite involved; see [51, 62]. Below the critical exponent m_c , fewer results are known and some considerations are in order. The fundamental solution does not exist anymore [92]. A large class of solutions vanish in finite time with different possible behaviour near the extinction time. The vanishing profile of a suitable class of initial data is represented by the so-called *pseudo*-Barenblatt solutions; see [6]. The extinction behaviour of bounded and integrable solutions for $0 < m < m_c$ is only known in the radial case [50, 92]. In the Yamabe flow case, $m = \frac{d-2}{d+2}$, finer results are known; see [30, 33, 34, 36]. The situation is completely different in the so-called ultrafast diffusion case, $m \leq 0$. Indeed, for the Cauchy problem, nonnegative L^1 data do not produce solutions [84, 90]. As a consequence, there are no solutions for the homogeneous Dirichlet problem, while a special class of solutions can be found for the Neumann problem [87, 88]. Sharp existence and nonexistence conditions for the Cauchy problem appear in [29, 31, 32] and [85].

Possible generalizations. The global Harnack principle (Theorem 1.1) can be generalized to solutions to equations of the form

$$u_t = \operatorname{div}(A(t, x, u, \nabla u^m)), \quad (1.13)$$

where $A(t, x, u, \eta)$ satisfies suitable structural conditions, such as those in [77, 79]. Even if the fundamental solution exists [77, 79], it is not clear whether or not (and in which sense) it represents the large time behaviour of nonnegative, integrable solutions to (1.13).

As already mentioned, the proofs of our results are new and allow us to improve on previous results in the nonweighted case [17], but we would like to emphasize that the strategy of the proof is robust, in the sense that it can also be adapted to prove analogous sharp characterizations of GHP and UREC in the case of the p -Laplacian evolution equation; see [15].

In the case of other equations, the validity of UREC and GHP is not always clear. In the recent papers [96, 97], a form of GHP for the *fractional p -Laplacian* evolution equation is proven to hold for a suitable class of initial data compactly supported. As for other nonlinear equations, in [20] the authors prove UREC for a Newtonian vortex equation. It is quite surprising that when the diffusion is driven by the *fractional Laplacian* $(-\Delta)^s$, with $s \in (0, 1)$, both GHP and UREC hold [16], while the same result is false for $s = 1$, as previously mentioned. However, to the best of our knowledge, no results are available for general linear nonlocal equations.

Organization of the paper. In Section 2 we collect some results that hold for all nonnegative integrable solutions: we find universal lower bounds in terms of Barenblatt profiles, which allow us to identify the minimal tail of all nonnegative solutions, Corollary 2.2. We prove that the answer to Q_2 can be yes for all $u_0 \in L^1_+(\mathbb{R}^d)$, but only on suitable parabolic cones, which represent the optimal domain of validity for such results. In Section 3 we analyze the behaviour of solutions whose initial data are in \mathcal{X} and provide a positive answer to both Q_1 and Q_2 . We prove the upper part of inequality (1.6) (Theorem 1.1) and Theorem 1.3. As a consequence, we obtain rates of convergence to the Barenblatt profile in several norms. In Section 4 we construct sub/supersolutions with the anomalous tail behaviour analyzed above. We also show, by means of counterexamples, that the power-like rates obtained in Section 3 are not possible for data outside \mathcal{X} . In Section 5 we show the equivalence between (TC) and (TC') and we give an example of a function in \mathcal{X} which does not satisfy the pointwise condition $u_0(x) \lesssim |x|^{-\frac{\sigma}{1-m}}$. We also give more details about the natural topology of \mathcal{X} and analyze stability properties of WFDE flow as a curve in \mathcal{X} .

Notation. We will systematically use ∞ to indicate $+\infty$. We will use the following notation throughout the paper: $a \wedge b := \min\{a, b\}$, $a \vee b := \max\{a, b\}$ and $a \asymp b$ means that there exist constants $c_1, c_2 > 0$ such that $c_1 a \leq b \leq c_2 a$; similarly we write $a \lesssim b$ whenever there exists $c > 0$ such that $a \leq cb$. Also, given $B \subset \mathbb{R}^d$ we define χ_B as the characteristic function of B , namely $\chi_B(x) = 1$ if $x \in B$, while $\chi_B(x) = 0$ if $x \notin B$.

2. Initial data in $L^1_{\gamma,+}(\mathbb{R}^d)$

In this section we show the results that hold for all data in $L^1_{\gamma,+}(\mathbb{R}^d)$, namely, we show that the lower part of the GHP estimates holds true (Theorem 2.1) for (just) locally

integrable data: this allows us to measure the (infinite) speed of propagation as “fatness of the tails”. On the other hand, on the whole space it is not possible to match the lower bounds with similar upper bounds for all initial data in $L^1_{\gamma,+}(\mathbb{R}^d)$: we will provide explicit counterexamples and improved lower bounds in Section 4. This latter phenomenon, an anomalous tail behaviour, can only happen if a control of the tail of the initial datum is missing: we will show that the sharp tail condition is encoded in the space \mathcal{X} thoroughly analyzed in Section 3. As a consequence of the estimates of this section, we also show uniform convergence in relative error towards equilibrium on compact sets and even on parabolic cones; see Theorem 2.4. All of these results are sharp, as shown in Section 4 by means of suitable counterexamples.

2.1. A universal global lower bound: Measuring the speed of propagation

We now state the main result of this section, which holds for nonnegative initial data which are merely locally integrable. We recall here a useful quantity t_* that will appear frequently throughout this section:

$$t_* = t_*(u_0, R_0) = \kappa_* \|u_0\|_{L^1_{\gamma}(B_{R_0}(0))}^{1-m} R_0^{\frac{1}{\beta}}, \quad (2.1)$$

where $\kappa_* > 0$ depends on d, m, γ, β . We are now ready to state the main result of this section.

Theorem 2.1. *Let u be a solution to (CP) with initial data $0 \leq u_0 \in L^1_{\gamma,\text{loc}}(\mathbb{R}^d)$ and let $t_0, R_0 > 0$ be such that $\|u_0\|_{L^1_{\gamma}(B_{R_0}(0))} > 0$. Then there exists $\underline{\tau} > 0$ and $\underline{M} > 0$ such that*

$$u(t, x) \geq \mathfrak{B}(t - \underline{\tau}, x; \underline{M}) \quad \text{for all } x \in \mathbb{R}^d \text{ and } t \geq t_0, \quad (2.2)$$

where

$$\underline{\tau} = \frac{1}{2}(t_* \wedge t_0) \quad \text{and} \quad \underline{M} = b \|u_0\|_{L^1_{\gamma}(B_{R_0}(0))} \left(1 \wedge \frac{t_0}{t_*}\right)^{\frac{1}{1-m}}. \quad (2.3)$$

The constant $b > 0$ depends only on d, m, γ, β and has an explicit expression given in the proofs, while t_* is as in (2.1).

Measuring the speed of propagation. The above theorem partially answers Q_1 and reveals a remarkable property of solutions to the WFDE: the positivity spreads immediately for every nonnegative initial datum, showing infinite speed of propagation. A delicate issue is how to discriminate in a quantitative way among two infinite speeds of propagation. Our theorem shows that we can put a (delayed) fundamental solution as a lower barrier for any data: this is how the WFDE immediately creates a fat tail (inverse power), which is clearly bigger than the “standard Gaussian tail” (exponentially decaying) created by the linear heat equation. This can be expressed as follows:

Corollary 2.2 (Minimal tails). *Under the assumption of Theorem 2.1 we have that for any $t > 0$,*

$$\liminf_{|x| \rightarrow \infty} u(t, x) |x|^{\frac{\sigma}{1-m}} \geq b_1 t^{\frac{1}{1-m}}, \quad (2.4)$$

where the constant b_1 depends only on m, d, γ, β and is achieved by the Barenblatt solutions.

We will often call $|x|^{-\sigma/(1-m)}$ a *minimal tail* or a *Barenblatt tail*. Finding matching upper bounds is simply not possible in such generality; we will need to ask for the tail condition (TC) on u_0 .

Proof of Theorem 2.1. Let us first state an inequality proven in [14, Theorem 1.4], a sharp local lower bound (half-Harnack inequality), essential to this proof. We do not use the Aleksandrov principle here, as in [17], nor other moving planes arguments. Under the running assumption we have

$$\inf_{x \in B_{2R}(0)} u(t_*, x) \geq \kappa_1 \frac{\|u_0\|_{L^1_\gamma(B_R(0))}}{R^{d-\gamma}}, \quad (2.5)$$

where κ_1 depends only on d, m, γ, β , and has an explicit expression given in [14]. Let us now explain the strategy of the proof. The quantities $\underline{\tau}$ and \underline{M} take different forms depending on whether or not $t_* \leq t_0$. We will assume first that $t_0 \geq t_*$, and then we will discuss the case $0 < t_0 < t_*$ at the end of the proof.

Let $M_{R_0} = \|u_0\|_{L^1_\gamma(B_{R_0}(0))}$, $\underline{\tau} = at_*$ and $\underline{M} = bM_{R_0}$, where $a \in (0, 1)$ and $b > 0$ will be explicitly chosen later. Without loss of generality, we prove inequality (2.2) only at $t = t_*$, namely

$$u(t_*, x) \geq \mathfrak{B}((1-a)t_*, x; \underline{M}). \quad (2.6)$$

Once proven at $t = t_*$, the case $t \geq t_*$ will follow by comparison. To prove (2.6) we need to determine the values of a, b . We need to separate two cases, namely inside a ball and outside a ball, obtaining different conditions on a, b , respectively conditions (2.8) and (2.12). Finally we check the compatibility of such conditions and choose a, b explicitly as in (2.13).

Condition on a, b inside a ball. We want to find a condition on a, b such that the following inequality holds:

$$\kappa_1 \frac{M_{R_0}}{R_0^{d-\gamma}} \geq \frac{b^{\sigma\vartheta} M_{R_0}}{b_0^{\frac{1}{1-m}} (1-a)^{(d-\gamma)\vartheta} \kappa_*^{(d-\gamma)\vartheta} R_0^{d-\gamma}} = \sup_{x \in B_{2R_0}(0)} \mathfrak{B}(t_* - \underline{\tau}, x; \underline{M}), \quad (2.7)$$

where κ_1 is as in (2.5). It is easily seen that the former is implied by the following condition on a and b :

$$b^{\sigma\vartheta} \leq \kappa_*^{(d-\gamma)\vartheta} \kappa_1 b_0^{\frac{1}{1-m}} (1-a)^{(d-\gamma)\vartheta}. \quad (2.8)$$

Note that by inequality (2.5), the first term in (2.7) is bounded above by $\inf_{x \in B_{2R_0}} u(t_*, x)$; therefore we obtain

$$\inf_{x \in B_{2R_0}} u(t_*, x) \geq \sup_{x \in B_{2R_0}(0)} \mathfrak{B}(t_* - \underline{\tau}, x; \underline{M}),$$

and inequality (2.6) is then proved for any $|x| \leq 2R_0$.

Condition on a, b outside a ball. We want to find suitable conditions on a, b such that (2.6) holds in the outer region $|x| > R_0$. Such an inequality will be deduced by applying the comparison on the parabolic boundary of $Q = (\underline{\tau}, t_*) \times B_{R_0}^c(0)$, namely $\partial_p Q = \{\{\underline{\tau}\} \times B_{R_0}^c(0)\} \cup \{\{\underline{\tau}, t_*\} \times \{x \in \mathbb{R}^d : |x| = R_0\}\}$; see for instance [57, Lemma 3.4].

It is clear that $u(\underline{\tau}, x) \geq \mathfrak{B}(0, x; \underline{M}) = \delta_0(x)$ for any $|x| \geq R_0$, hence we just need to prove that

$$u(t, x) \geq \mathfrak{B}(t - \underline{\tau}, x; \underline{M}) \quad \text{for any } |x| = R_0, t \in (\underline{\tau}, t_*). \quad (2.9)$$

A straightforward computation shows that, under the running assumption, for $|x| = R_0$ we have

$$\sup_{t \geq \underline{\tau}} \mathfrak{B}(t - \underline{\tau}, x; \underline{M}) = \left(\frac{b_1}{b_0 \vartheta} \right)^{\frac{1}{\sigma \vartheta(1-m)}} \frac{[(d - \gamma)(1 - m)]^{\frac{d-\gamma}{\sigma}}}{[\kappa_* \sigma]^{\frac{1}{1-m}}} \left(\frac{t_*}{b_0 R_0^\sigma} \right)^{\frac{1}{1-m}} b. \quad (2.10)$$

The inequality

$$\underline{\kappa} \left(\frac{at_*}{R_0^\sigma} \right)^{\frac{1}{1-m}} \geq \left(\frac{b_1}{b_0 \vartheta} \right)^{\frac{1}{\sigma \vartheta(1-m)}} \frac{[(d - \gamma)(1 - m)]^{\frac{d-\gamma}{\sigma}}}{[\kappa_* \sigma]^{\frac{1}{1-m}}} \left(\frac{t_*}{b_0 R_0^\sigma} \right)^{\frac{1}{1-m}} b \quad (2.11)$$

implies that inequality (2.9) holds; indeed for any $|x| = R_0$ and $t \in (\underline{\tau}, t_*)$ we have

$$u(t, x) \geq \inf_{\substack{t \in (at_*, t_*) \\ x \in B_{2R_0}(0)}} u(t, x) \geq \underline{\kappa} \left(\frac{at_*}{R_0^\sigma} \right)^{\frac{1}{1-m}} \geq \sup_{t \geq \underline{\tau}} \mathfrak{B}(t - \underline{\tau}, x; \underline{M}) \geq \mathfrak{B}(t - \underline{\tau}, x; \underline{M}).$$

It is easy to show that inequality (2.11) is equivalent to

$$b^{\sigma \vartheta} \leq b_0^{\frac{1}{1-m}} a^{\frac{\sigma \vartheta}{1-m}} \underline{\kappa}^{\sigma \vartheta} \left(\frac{\vartheta b_0^{\sigma \vartheta}}{b_1} \right)^{\frac{1}{1-m}} \frac{[\kappa_* \sigma]^{\frac{\sigma \vartheta}{1-m}}}{[(d - \gamma)(1 - m)]^{(d-\gamma)\vartheta}}, \quad (2.12)$$

which is the condition we were looking for.

Compatibility of conditions (2.8) and (2.12). Both conditions are satisfied by the choices

$$\begin{aligned} a &= \frac{1}{2}, \\ b^{\sigma \vartheta} &= b_0^{\frac{1}{1-m}} \left[\left(\frac{\vartheta b_0^{\sigma \vartheta}}{2^{\sigma \vartheta} b_1} \right)^{\frac{1}{1-m}} \frac{\underline{\kappa}^{\sigma \vartheta} [\kappa_* \sigma]^{\frac{\sigma \vartheta}{1-m}}}{[(d - \gamma)(1 - m)]^{(d-\gamma)\vartheta}} \wedge \frac{\kappa_*^{(d-\gamma)\vartheta} \kappa_1}{2^{(d-\gamma)\vartheta}} \right]. \end{aligned} \quad (2.13)$$

This concludes the proof of (2.6) in the case when $t_0 \geq t_*$. It only remains to analyze the case when $t_0 < t_*$.

Case $0 < t_0 < t_$.* Without loss of generality, we only need to prove inequality (2.2) at time $t = t_0$, and the full result will then follow by comparison. Recall the Bénilan–Crandall-type estimate [4]

$$u(t_0, x) \geq u(t_*, x) \left(\frac{t_0}{t_*} \right)^{\frac{1}{1-m}} \quad \text{for all } 0 < t_0 < t_*. \quad (2.14)$$

Now we recall that inequality (2.6) holds under the choices of a, b as in (2.13). Using inequalities (2.6) and (2.14) we get

$$\begin{aligned} u(t_0, x) &\geq u(t_*, x) \left(\frac{t_0}{t_*}\right)^{\frac{1}{1-m}} \geq \frac{2^{-\frac{1}{1-m}} t_*^{\frac{1}{1-m}}}{\left[b_0 \frac{2^{-\sigma\vartheta} t_*^{\sigma\vartheta}}{\underline{M}^{\sigma\vartheta(1-m)}} + b_1 |x|^\sigma\right]^{\frac{1}{1-m}}} \left(\frac{t_0}{t_*}\right)^{\frac{1}{1-m}} \\ &= \frac{2^{-\frac{1}{1-m}} t_0^{\frac{1}{1-m}}}{\left[b_0 \frac{2^{-\sigma\vartheta} t_0^{\sigma\vartheta}}{\underline{M}^{\sigma\vartheta(1-m)} \left[\left(\frac{t_0}{t_*}\right)^{\frac{1}{1-m}}\right]^{\sigma\vartheta}} + b_1 |x|^\sigma\right]^{\frac{1}{1-m}}} \\ &= \mathfrak{B}\left(t_0 - \frac{t_0}{2}, x; \left(\frac{t_0}{t_*}\right)^{\frac{1}{1-m}} \underline{M}\right). \end{aligned}$$

Recalling that in this case $\underline{\tau} = t_0/2$, the proof is concluded. \blacksquare

We can now give the proof of Corollary 2.2.

Proof of Corollary 2.2. Let R_0 be such that $\|u_0\|_{L^1_+(B_{R_0}(0))} > 0$, $t > 0$, and $0 < \varepsilon < t$. By applying Theorem 2.1 at time $t_0 = t - \varepsilon$ and radius R_0 we get the inequality

$$u(t, x) \geq \mathfrak{B}(t - \underline{\tau}, x; \underline{M}).$$

As a consequence we obtain

$$\liminf_{|x| \rightarrow \infty} u(t, x) |x|^{\frac{\sigma}{1-m}} \geq b_1 \left[t - \frac{1}{2}(t_* \wedge t_0)\right]^{\frac{1}{1-m}},$$

from which (2.4) follows just by taking the limit for $\varepsilon \rightarrow t$. Notice that in such a limit $t_0 \rightarrow 0$. \blacksquare

2.2. Harnack inequality in parabolic cones

We showed in [14] that nonnegative local solutions to the WFDE satisfy Harnack inequalities of various kinds: an *elliptic* form (in which the supremum and the infimum are taken at the same time), a *forward* in time (the supremum is taken at a smaller time than the infimum) and a *backward* in time (the supremum is taken at a bigger time than the infimum). We remark that for solutions to the *heat equations* in general only the *forward* Harnack inequality holds. Here we will prove an *elliptic* form of a Harnack inequality on conical space-time domains of the form

$$K(t) = K_M(t) = \{|x| \leq t^\vartheta M^{(m-1)\vartheta}\} \quad (2.15)$$

for some fixed $M > 0$. We will call these sets “parabolic cones”, with a slight abuse of language. Indeed, for $\vartheta = 1$, $K(t)$ are really cones in space-time domains of the form $\mathbb{R}_+ \times \mathbb{R}^N$. A similar inequality on balls has been proven in [17, Theorem 1.4].

Theorem 2.3 (Harnack inequality in parabolic cones). *Let u be a solution to (CP) with initial data $0 \leq u_0 \in L^1_\gamma(\mathbb{R}^d)$. Let $M = \|u_0\|_{L^1_\gamma(\mathbb{R}^d)}$ and $R_0 > 0$ be such that $\|u_0\|_{L^1_\gamma(B_{R_0}(0))} = M/2$, and let $t_* = \kappa_* R_0^{\frac{1}{\beta}} (M/2)^{\frac{1}{1-m}}$. Then there exists a positive constant \mathcal{H} such that*

$$\sup_{x \in K(t)} \frac{u(t, x)}{\mathfrak{B}(t, x; M)} \leq \mathcal{H} \inf_{x \in K(t)} \frac{u(t, x)}{\mathfrak{B}(t, x; M)} \quad \text{for any } t \geq 3t_*, \quad (2.16)$$

where the constant \mathcal{H} depends only on m, d, γ, β and $K(t)$ depends on M as in (2.15)

Proof. In the proof we will make use of the smoothing effect for solutions to (CP), namely the following inequality, which holds for any $t > 0$:

$$\|u(t)\|_{L^\infty(\mathbb{R}^d)} \leq \frac{\bar{\kappa}_1}{t^{(d-\gamma)\vartheta}} \left[\int_{\mathbb{R}^d} |u_0(y)| |y|^{-\gamma} dy \right]^{\sigma\vartheta}, \quad (2.17)$$

where ϑ and σ are defined in (1.3). The constant $\bar{\kappa}_1$ has an explicit form, cf. [14]; it depends only on d, m, γ and β and it is the same constant $\bar{\kappa}_1$ appearing in formula (3.3). Inequality (2.17) has been obtained in [14] and can be easily deduced by taking $x_0 = 0$ and letting $R_0 \rightarrow \infty$ in estimate (3.3).

Let us now begin the proof. By applying Theorem 2.1 we deduce that $u(t, x) \geq \mathfrak{B}(t - \underline{\tau}, x; \underline{M})$ with $\underline{\tau} = \frac{t_*}{2} = \frac{\kappa_*}{2} R_0^{\frac{1}{\beta}} \left(\frac{M}{2}\right)^{\frac{1}{1-m}}$ and $\underline{M} = bM/2$, where b is as in (2.12). In view of the smoothing effects (2.17) and of inequality (2.2), it is enough to prove that there exists \mathcal{H} such that

$$\bar{\kappa}_1 (b_0 + b_1)^{\frac{1}{1-m}} \leq \mathcal{H} b_0^{\frac{1}{1-m}} \frac{t^{(d-\gamma)\vartheta}}{M^{\sigma\vartheta}} \inf_{x \in K(t)} \mathfrak{B}(t - \underline{\tau}, x; \underline{M}).$$

This amounts to proving that the following quotient is uniformly bounded by \mathcal{H} for $t \geq 3t_*$:

$$\bar{\kappa}_1 \left(1 + \frac{b_1}{b_0}\right)^{\frac{1}{1-m}} \frac{M^{\sigma\vartheta}}{t^{(d-\gamma)\vartheta}} \frac{\left[\frac{b_0(t-\underline{\tau})^{\sigma\vartheta}}{M^{2\vartheta(1-m)}} + \frac{b_1 t^{\sigma\vartheta}}{M^{\sigma\vartheta(1-m)}}\right]^{\frac{1}{1-m}}}{(t - \underline{\tau})^{\frac{1}{1-m}}} \leq \mathcal{H}.$$

Since $\underline{\tau} = t_*/2$ we easily conclude that \mathcal{H} can be taken as

$$\mathcal{H} = \bar{\kappa}_1 \left(1 + \frac{b_1}{b_0}\right)^{\frac{1}{1-m}} 5^{\frac{1}{1-m}} \left[b_0 \left(\frac{2}{b}\right)^{\sigma\vartheta} + b_1\right]^{\frac{1}{1-m}}. \quad \blacksquare$$

2.3. Uniform convergence in relative error in parabolic cones

In this section we will prove that solutions to (CP) with initial data $u_0 \in L^1_{\gamma,+}(\mathbb{R}^d)$ converge to the Barenblatt profile $\mathfrak{B}(t, x; M)$ in *relative error* uniformly in parabolic cones, and as a consequence uniformly on compact subsets of \mathbb{R}^d . To obtain such a result we will use the convergence to the Barenblatt profile in $L^1_\gamma(\mathbb{R}^d)$, namely

$$\|u(t) - \mathfrak{B}(t; M)\|_{L^1_\gamma(\mathbb{R}^d)} \rightarrow 0 \quad \text{as } t \rightarrow \infty \quad (2.18)$$

or equivalently, in *self-similar variables*,

$$\|v(\tau) - \mathfrak{B}_M\|_{L^1_\gamma(\mathbb{R}^d)} \rightarrow 0 \quad \text{as } \tau \rightarrow \infty, \quad (2.19)$$

where $v(\tau, y)$ is defined in (1.11) and it is a solution to (NLWFP). The proof of (2.18) can be done by a straightforward adaptation to our setting of the so-called “4 step method”, carefully explained in [91, Theorem 1.1]. We leave the details to the interested reader, just noting that the proof contained in [91] deals with the case $m > 1$, and uses compactly supported initial data, hence compactly supported solutions (when $m > 1$ there is finite speed of propagation). In the present setting, the very same proof works, just by replacing the compactly supported solutions by those which satisfy the GHP.

Theorem 2.4 (Uniform convergence in *relative error* on parabolic cones). *Let $m \in (m_c, 1)$ and let u be a solution to (CP) with initial data $0 \leq u_0 \in L^1_\gamma(\mathbb{R}^d)$ and let $M = \|u_0\|_{L^1_\gamma(\mathbb{R}^d)}$. Then for any $\Upsilon > 0$ we have*

$$\lim_{t \rightarrow \infty} \sup_{x \in \{|x| \leq \Upsilon t^\vartheta\}} \left| \frac{u(t, x) - \mathfrak{B}(t, x; M)}{\mathfrak{B}(t, x; M)} \right| = 0. \quad (2.20)$$

Remark. As an easy corollary of the previous theorem we obtain

$$\left\| \frac{u(t, x) - \mathfrak{B}(t, x; M)}{\mathfrak{B}(t, x; M)} \right\|_{L^\infty(K)} \xrightarrow{t \rightarrow \infty} 0 \quad \text{for any compact set } K \subset \mathbb{R}^d.$$

This follows from inequality (2.20) just by observing that $K \subset \{|x| \leq \Upsilon t^\vartheta\}$ for some $t_0 > 0$.

Proof of Theorem 2.4. We split the proof into several steps. First we prove a uniform pointwise estimate on the solution $u(t, x)$ in domains of the form $\{|x| \leq CR(t)\}$, where $R(t)$ is as in (1.11) and $C > 0$. We remark that for any $t > 0$ we have $\{|x| \leq Ct^\vartheta\} \subset \{|x| \leq CR(t)\}$. As a second step we will rescale $u(t, x)$ to self-similar variables (we recall that domains of type $\{|x| \leq CR(t)\}$ are transformed into $B_\rho(0)$, where $\rho = \zeta C$) and, using the estimates obtained before, we estimate $\lfloor v(\tau, \cdot) - \mathfrak{B}_M \rfloor_{C^v(B_{3r})}$ uniformly in time. Finally, by applying a clever interpolation, Lemma A.1, we prove that $\|v(\tau, \cdot) - \mathfrak{B}_M\|_{L^\infty(B_r)} \rightarrow 0$ as $\tau \rightarrow \infty$, and finally (2.20) follows.

Uniform estimate on $u(t, x)$ in $\{|x| \leq 3\Upsilon R(t)\}$. Let $\rho > 0$ be such that $\int_{B_\rho} u_0(x) |x|^{-\gamma} dx = \frac{M}{2}$ and define $t_\star = \kappa_\star \rho^{\frac{1}{\vartheta}} (\frac{M}{2})^{1-m}$, where κ_\star is as in (2.1). By applying Theorem 2.1 and the global smoothing effect, inequality (2.17), we obtain that for any $t \geq t_\star$,

$$\mathfrak{B}(t - \underline{t}, x; \underline{M}) \leq u(t, x) \leq \bar{\kappa}_1 \frac{M^{\sigma\vartheta}}{t^{(d-\gamma)\vartheta}},$$

where $\underline{t} = \frac{t_\star}{2}$ and $\underline{M} = \frac{b}{2} M$. By the above inequality, we can deduce the following matching lower bound, by means of straightforward estimates relying on the explicit expression for \mathfrak{B} : there exists a constant $\kappa_1 > 0$ which depends on $d, m, \gamma, \beta, \Upsilon$ and \underline{M} such that

$$\kappa_1 \frac{M^{\sigma\vartheta}}{t^{(d-\gamma)\vartheta}} \leq u(t, x) \leq \bar{\kappa}_1 \frac{M^{\sigma\vartheta}}{t^{(d-\gamma)\vartheta}} \quad \text{for any } t \geq t_\star \text{ and any } x \in \{|x| \leq 3\Upsilon R(t)\}. \quad (2.21)$$

Uniform and Hölder estimates in self-similar variables. We first rescale u in self-similar variables, according to (1.11), and get $v(\tau, y)$. Analogously, the domain $\{|x| \leq 3\Upsilon R(t)\}$ is transformed into $B_{3r}(0)$ where $r = \Upsilon\zeta$. Inequality (2.21) reads, in rescaled variables,

$$\frac{\bar{\kappa}_1}{\zeta^{d-\gamma}} \vartheta^\vartheta M^{\sigma\vartheta} \leq v(\tau, y) \leq 2 \frac{\bar{\kappa}_1}{\zeta^{d-\gamma}} \vartheta^\vartheta M^{\sigma\vartheta} \quad \text{for any } \tau \geq \frac{1}{\sigma} \log \frac{R(t_\star \vee 1)}{R(0)}$$

and any $y \in B_{3r}(0)$. (2.22)

By applying Lemma A.3 we deduce that there exist $\nu > 0$, $\bar{\kappa} > 0$ such that for any $\tau \geq \frac{1}{\sigma} \log \frac{R(t_\star \vee 1)}{R(0)} + 1$ we have

$$[v(\tau, \cdot)]_{C^\nu(B_{\frac{3}{2}r}(0))} \leq \bar{\kappa} 2 \frac{\bar{\kappa}_1}{\zeta^{d-\gamma}} \vartheta^\vartheta M^{\sigma\vartheta}.$$

Using the subadditivity of $[\cdot]_{C^\nu(B_r(0))}$ and the fact that the above estimates can also be applied to the Barenblatt profile $\mathcal{B}_M(y)$, we conclude that

$$[v(\tau, \cdot) - \mathcal{B}_M]_{C^\nu(B_{\frac{3}{2}r}(0))} \leq 4\bar{\kappa} \frac{\bar{\kappa}_1}{\zeta^{d-\gamma}} \vartheta^\vartheta M^{\sigma\vartheta} \quad \text{for any } \tau \geq \frac{1}{\sigma} \log \frac{R(t_\star \vee 1)}{R(0)} + 1. \quad (2.23)$$

Convergence in L^∞ norm. We only prove the case $0 < \gamma < d$, which is the most delicate, the case $\gamma \leq 0$ being simpler. In what follows it is convenient to assume that $r \geq 2$, namely that $\Upsilon \geq \frac{2}{\zeta}$, but we will overcome this technical assumption at the end of the proof. From the convergence in L_γ^1 , formula (2.19), we deduce that there exists τ_\star such that for any $\tau \geq \tau_\star$ we have $\|v(\tau, \cdot) - \mathcal{B}_M\|_{L_\gamma^1(B_{\frac{3}{2}r}(0))} \leq \frac{|\gamma|}{d}$. We are in a position to apply inequality (A.4) of Lemma A.1 to $v(\tau, \cdot) - \mathcal{B}_M$ and get that for any $\tau \geq \tau_\star \vee \frac{1}{\sigma} \log \frac{R(t_\star \vee 1)}{R(0)} + 1$,

$$\|v(\tau, \cdot) - \mathcal{B}_M\|_{L^\infty(B_r(0))} \leq C_{d,\gamma,\nu,p} (1+r)^\gamma \left(1 + 4\bar{\kappa} \frac{\bar{\kappa}_1}{\zeta^{d-\gamma}} \vartheta^\vartheta M^{\sigma\vartheta}\right)^{\frac{d}{d+\nu}} \times \|v(\tau, \cdot) - \mathcal{B}_M\|_{L_\gamma^1(B_{\frac{3}{2}r}(0))}^{\frac{\nu}{d+\nu}}, \quad (2.24)$$

where we have used (2.23). Since $\mathfrak{B} \geq (C(M) + r^\sigma)^{\frac{1}{1-m}}$ on $B_r(0)$, it follows that

$$\sup_{y \in B_r(0)} \left| \frac{v(\tau, y) - \mathcal{B}_M(y)}{\mathfrak{B}_M(y)} \right| \leq (C(M) + r^\sigma)^{\frac{1}{1-m}} \|v(\tau, \cdot) - \mathcal{B}_M\|_{L^\infty(B_r(0))},$$

which, combined with (2.24) and the convergence in L_γ^1 , formula (2.19), shows that the *relative error* approaches zero as $\tau \rightarrow \infty$. Rescaling back, we finally obtain (2.20), recalling that $\{|x| \leq \Upsilon t^\vartheta\} \subset \{|x| \leq \Upsilon R(t)\}$.

It only remains to overcome the technical assumption $\Upsilon \geq \frac{2}{\zeta}$. If $\Upsilon \leq \frac{2}{\zeta}$ we can repeat the same argument for $\Upsilon = \frac{2}{\zeta}$. Next we conclude that (2.20) takes place for any $\Upsilon' \leq \frac{2}{\zeta}$ using that $\{|x| \leq \Upsilon' R(t)\} \subset \{|x| \leq \Upsilon R(t)\}$ whenever $\Upsilon' < \Upsilon$. The proof is now concluded. \blacksquare

3. Initial data in \mathcal{X} . Global Harnack principle and uniform convergence in relative error

The space \mathcal{X} is naturally invariant under the fast diffusion flow as explained in the introduction; see also Proposition 5.3. As a consequence, solutions belonging to this space possess some extra properties, that we summarize here:

- The tail is essentially the same as the Barenblatt solution, the GHP holds; see Section 1.1.
- Uniform convergence in relative error (UREC) takes place; see Section 3.2. Moreover, we also provide *almost optimal rates of convergence* in Section 3.5, which turn out to be sharp in some cases.
- Boundary Harnack-type inequalities hold true, and the behaviour at infinity of solutions does not depend on the mass; see Section 3.4.

In Section 4 we will show that the above properties are false if $u_0 \notin \mathcal{X}$.

3.1. Upper bound and proof of Theorem 1.1

As already observed in the introduction, Theorem 1.1 is divided into two parts: the upper bound and the lower bound of inequality (1.6). In this section we are going to discuss the upper bound, and the main result of this section is the following theorem.

Theorem 3.1. *Let u be the solution to (CP) corresponding to the initial data $0 \leq u_0 \in L^1_{\gamma,+}(\mathbb{R}^d)$. Then, for any $t_0 > 0$ there exist $\bar{\tau}, \bar{M} > 0$, explicitly given in (3.10), such that*

$$u(t, x) \leq \mathfrak{B}(t + \bar{\tau}, x; \bar{M}) \quad \text{for any } x \in \mathbb{R}^d \text{ and any } t > t_0, \quad (3.1)$$

if and only if

$$u_0 \text{ satisfies (TC), i.e. } u_0 \in \mathcal{X}.$$

The proof of inequality (3.1) is constructive and we are able to give values of $\bar{\tau}$ and \bar{M} , see formulae (3.10) at the end of the proof. Here we just point out that they depend on $M, A, d, m, \gamma, \beta$ and t_0 .

Proof of Theorem 1.1. The proof is a simple combination of Theorems 3.1 and 2.1. ■

Remark 3.2. We easily deduce from the above upper bound that

$$\limsup_{|x| \rightarrow \infty} u(t, x) |x|^{\frac{\sigma}{1-m}} \leq b_1 (t + \bar{\tau})^{\frac{1}{1-m}}. \quad (3.2)$$

Equality is achieved by the Barenblatt solution translated in time by $\bar{\tau}$. Notice that this maximal tail behaviour only holds for $u_0 \in \mathcal{X}$, in which case it matches the optimal minimal behaviour given in Corollary 2.2. These two pieces of information combine well and allow us to deduce the sharp behaviour at infinity; see Section 3.4, Corollary 3.21.

Proof of Theorem 3.1. Let us first explain the strategy of the proof. We will prove inequality (3.1) only at time t_0 and then, by comparison (see for instance [10, Corollary 9]), it will hold for any $t \geq t_0$. The proof is divided into several steps: first, we estimate the solution $u(t_0, x)$ in two different regions (on $B_{R_1}(0)$ and on $B_{R_1}(0)^c$, with R_1 to be chosen later), and then we find conditions on $\bar{\tau}$ and \bar{M} necessary for inequality (3.1) to hold. Finally, we show that such conditions can be fulfilled, providing an explicit expression of $\bar{\tau}$ and \bar{M} in terms of t_0 , M and A .

In this proof we will make use of the following estimate [14, Theorem 1.2]: there exist $\bar{\kappa}_1, \bar{\kappa}_2 > 0$ such that for any $t > 0$, $x_0 \in \mathbb{R}^d$ and any $R_0 \in [|x_0|/16, |x_0|/32]$ (any $R_0 > 0$ if $x_0 = 0$) we have

$$\sup_{y \in B_{R_0}(x_0)} u(t, y) \leq \frac{\bar{\kappa}_1}{t^{(d-\gamma)\vartheta}} \left[\int_{B_{2R_0}(x_0)} |u_0(y)| |y|^{-\gamma} dy \right]^{\sigma\vartheta} + \bar{\kappa}_2 \left[\frac{t}{R_0^\sigma} \right]^{\frac{1}{1-m}}, \quad (3.3)$$

where ϑ and σ are defined in (1.3). The constants $\bar{\kappa}_1, \bar{\kappa}_2$ are explicit and depend only on d, m, γ and β . The constant $\bar{\kappa}_1$ is the same one that appears in the smoothing effect given in inequality (2.17).

Estimate inside a ball. We want to find suitable conditions on $\bar{M}, \bar{\tau}$ and R_1 such that

$$\begin{aligned} u(t_0, x) &\leq \mathfrak{B}(t_0 + \bar{\tau}, x; \bar{M}) \\ &= \frac{(t_0 + \bar{\tau})^{\frac{1}{1-m}}}{\left[b_0 \frac{(t_0 + \bar{\tau})^{\sigma\vartheta}}{M^{\sigma\vartheta(1-m)}} + b_1 |x|^\sigma \right]^{\frac{1}{1-m}}} \quad \text{holds for all } |x| \leq R_1. \end{aligned} \quad (3.4)$$

Recall that $M = \int_{\mathbb{R}^d} u_0 |x|^{-\gamma} dx$. Inequality (3.3) implies that

$$u(t_0, x) \leq \bar{\kappa}_1 t_0^{-(d-\gamma)\vartheta} M^{\sigma\vartheta} \quad \text{for any } x \in \mathbb{R}^d \text{ and } t_0 > 0.$$

To deduce the above from (3.3) it suffices to take $x_0 = 0$ and let $R_0 \rightarrow \infty$. In view of the above inequality, to prove (3.4) it is enough to find suitable $\bar{M}, \bar{\tau}$ and R_1 such that

$$\bar{\kappa}_1 \frac{M^{\sigma\vartheta}}{t_0^{(d-\gamma)\vartheta}} \leq \frac{(t_0 + \bar{\tau})^{\frac{1}{1-m}}}{\left[b_0 \frac{(t_0 + \bar{\tau})^{\sigma\vartheta}}{M^{\sigma\vartheta(1-m)}} + b_1 |x|^\sigma \right]^{\frac{1}{1-m}}} \quad \text{for any } |x| \leq R_1.$$

Since the right-hand side is decreasing in $|x|$ it suffices to have the previous inequality at $|x| = R_1$, i.e.

$$b_0 \frac{(t_0 + \bar{\tau})^{\sigma\vartheta}}{M^{\sigma\vartheta(1-m)}} + b_1 R_1^\sigma \leq \frac{(t_0 + \bar{\tau}) t_0^{(d-\gamma)\vartheta(1-m)}}{\bar{\kappa}_1^{1-m} M^{\sigma\vartheta(1-m)}}. \quad (3.5)$$

Inequality (3.5) is nothing but a first condition on $\bar{M}, \bar{\tau}$ and R_1 in order to guarantee the validity of (3.4).

Estimate outside a ball. The goal of this step is to extend inequality (3.4) outside a ball, namely for all $|x| \geq R_1$. This will end up with conditions on \bar{M} , $\bar{\tau}$ and R_1 different from (3.5). In the next step we will take care of checking the compatibility of the two conditions.

We first prove that for any fixed $t_0 > 0$ there exists $C_1 = C_1(t_0, A) > 0$ such that

$$u(t_0, x) \leq \frac{C_1}{|x|^{\frac{\sigma}{1-m}}} \quad \text{for any } |x| > R_1. \quad (3.6)$$

Let $x \in \mathbb{R}^d$, $|x| \geq R_1$ and let R be such that $B_{2R}(x) \subset B_{2R}(0)^c$, for instance $R = |x|/16$. Applying inequality (3.3) to $u(t_0, x)$ in the ball $B_R(x)$, we get

$$\begin{aligned} u(t_0, x) &\leq \frac{\bar{k}_1}{t_0^{(d-\gamma)\vartheta}} \left[\int_{B_{2R}^c(0)} u_0(y) |y|^{-\gamma} dy \right]^{\sigma\vartheta} + \bar{k}_2 (16)^{-\frac{\sigma}{1-m}} \left(\frac{t_0}{|x|^\sigma} \right)^{\frac{1}{1-m}} \\ &\leq \frac{\bar{k}_1 8^{\frac{\sigma}{1-m}}}{t_0^{(d-\gamma)\vartheta}} \frac{A^{\sigma\vartheta}}{|x|^{\frac{\sigma}{1-m}}} + \frac{\bar{k}_2}{16^{\frac{\sigma}{1-m}}} \left(\frac{t_0}{|x|^\sigma} \right)^{\frac{1}{1-m}} \leq \frac{C_1}{|x|^{\frac{\sigma}{1-m}}}, \end{aligned}$$

where in the second inequality we have used that $\int_{B_R^c(0)} u_0 |x|^{-\gamma} dx \leq AR^{(d-\gamma)-\frac{2+\beta-\gamma}{1-m}}$ with $R = |x|/16$ and that $C_1 = C_1(t_0, A)$ is given by

$$C_1 = 8^{\frac{\sigma}{1-m}} \frac{\bar{k}_1}{t_0^{(d-\gamma)\vartheta}} A^{\sigma\vartheta} + \frac{\bar{k}_2}{16^{\frac{\sigma}{1-m}}} t_0^{\frac{1}{1-m}}.$$

Hence inequality (3.6) holds. It only remains to show that

$$\frac{C_1}{|x|^{\sigma/(1-m)}} \leq \frac{(t_0 + \bar{\tau})^{\frac{1}{1-m}}}{\left[b_0 \frac{(t_0 + \bar{\tau})^{\sigma\vartheta}}{\bar{M}^{\sigma\vartheta(1-m)}} + b_1 |x|^\sigma \right]^{\frac{1}{1-m}}} \quad \text{for any } |x| \geq R_1. \quad (3.7)$$

This will give a condition on $\bar{\tau}$, \bar{M} and R_1 , as we explain next. Indeed, the above inequality is equivalent to

$$b_1 C_1^{1-m} + b_0 \frac{(t_0 + \bar{\tau})^{\sigma\vartheta}}{|x|^\sigma \bar{M}^{\sigma\vartheta(1-m)}} \leq t_0 + \bar{\tau}.$$

It is indeed enough to choose $R_1 > 0$ such that

$$b_1 C_1^{1-m} + b_0 \frac{(t_0 + \bar{\tau})^{\sigma\vartheta}}{R_1^\sigma \bar{M}^{\sigma\vartheta(1-m)}} \leq t_0 + \bar{\tau}, \quad (3.8)$$

since the second term on left-hand side is decreasing in $|x|$. We conclude that inequality (3.4) holds for any $|x| \geq R_1$ whenever $\bar{\tau}$, \bar{M} and R_1 satisfy condition (3.8).

Compatibility between conditions (3.5) and (3.8). We only need to show the compatibility of the conditions that imply the main estimates of the previous steps, i.e. that inequality (3.4) holds for all $x \in \mathbb{R}^d$. The two conditions (3.5) and (3.8) correspond to the following

system of inequalities:

$$(A) = \begin{cases} b_1 C_1^{1-m} R_1^\sigma + b_0 \frac{(t_0 + \bar{\tau})^{\sigma\vartheta}}{\bar{M}^{\sigma\vartheta(1-m)}} \leq R_1^\sigma (t_0 + \bar{\tau}), \\ b_1 R_1^\sigma + b_0 \frac{(t_0 + \bar{\tau})^{\sigma\vartheta}}{\bar{M}^{\sigma\vartheta(1-m)}} \leq \frac{(t_0 + \bar{\tau}) t_0^{(d-\gamma)\vartheta(1-m)}}{\bar{\kappa}_1^{1-m} M^{\sigma\vartheta(1-m)}}. \end{cases}$$

It is convenient to simplify the above system in order to be able to make explicit choices of $\bar{\tau}$, \bar{M} and R_1 . The first simplification is

$$(B) = \begin{cases} b_1 (1 \vee C_1)^{1-m} R_1^\sigma \leq \frac{t_0 + \bar{\tau}}{2} \left[R_1^\sigma \wedge \frac{t_0^{(d-\gamma)\vartheta(1-m)}}{\bar{\kappa}_1^{1-m} M^{\sigma\vartheta(1-m)}} \right], \\ b_0 \frac{(t_0 + \bar{\tau})^{\sigma\vartheta}}{\bar{M}^{\sigma\vartheta(1-m)}} \leq \frac{t_0 + \bar{\tau}}{2} \left[R_1^\sigma \wedge \frac{t_0^{(d-\gamma)\vartheta(1-m)}}{\bar{\kappa}_1^{1-m} M^{\sigma\vartheta(1-m)}} \right]. \end{cases} \quad (3.9)$$

It is clear that any choice of $\bar{\tau}$, \bar{M} and R_1 that satisfies **(B)** also satisfies **(A)**. We need a further simplification, but this time we will choose $R_1 = R_1(R_0, t_0, M)$ in the particular way that

$$R_1 := \left(\frac{t_0^{(d-\gamma)\vartheta}}{\bar{\kappa}_1 M^{\sigma\vartheta}} \right)^{\frac{1-m}{\sigma}} \quad \text{so that} \quad R_1^\sigma = \frac{t_0^{(d-\gamma)\vartheta(1-m)}}{\bar{\kappa}_1^{1-m} M^{\sigma\vartheta(1-m)}},$$

and system **(B)** simplifies to

$$(B') = \begin{cases} b_1 (1 \vee C_1)^{1-m} \leq \frac{(t_0 + \bar{\tau})}{2}, \\ b_0 \frac{(t_0 + \bar{\tau})^{\sigma\vartheta}}{\bar{M}^{\sigma\vartheta(1-m)}} \leq \frac{(t_0 + \bar{\tau})}{2} \frac{t_0^{(d-\gamma)\vartheta(1-m)}}{\bar{\kappa}_1^{1-m} M^{\sigma\vartheta(1-m)}}, \\ \bar{\tau} \geq [2b_1(1 \vee C_1)^{1-m} - t_0], \\ \Rightarrow \begin{cases} \bar{M} \geq (2b_0 \bar{\kappa}_1^{1-m})^{\frac{1}{\sigma\vartheta(1-m)}} \left(\frac{t_0 + \bar{\tau}}{t_0} \right)^{\frac{d-\gamma}{\sigma}} M, \end{cases} \end{cases}$$

It is now clear that choosing $\bar{\tau} = \bar{\tau}(t_0, M_\infty, C_1, R_1)$ and $M = M(\bar{\tau}, t_0, M)$ of the form

$$\bar{\tau} := 0 \vee [2b_1(1 \vee C_1)^{1-m} - t_0] \quad \text{and} \quad \bar{M} := (2b_0 \bar{\kappa}_1^{1-m})^{\frac{1}{\sigma\vartheta(1-m)}} \left(\frac{t_0 + \bar{\tau}}{t_0} \right)^{\frac{d-\gamma}{\sigma}} M$$

implies the validity of the two inequalities of system **(B')**, hence of system **(B)**, and finally of **(A)**.

Values of the constants. Letting $A := |u_0|_{\mathcal{X}}$ and $M := \|u_0\|_{L^1_\gamma(\mathbb{R}^d)}$ we have

$$\begin{aligned} \bar{\tau} &:= 0 \vee \left\{ 2b_1 \left[1 \vee \left(8^{\frac{\sigma}{1-m}} \frac{\bar{\kappa}_1 A^{\sigma\vartheta}}{t_0^{(d-\gamma)\vartheta}} + \frac{\bar{\kappa}_2 t_0^{\frac{1}{1-m}}}{16^{\frac{\sigma}{1-m}}} \right) \right]^{1-m} - t_0 \right\}, \\ \bar{M} &:= (2b_0 \bar{\kappa}_1^{1-m})^{\frac{1}{\sigma\vartheta(1-m)}} \left(\frac{t_0 + \bar{\tau}}{t_0} \right)^{\frac{d-\gamma}{\sigma}} M, \end{aligned} \quad (3.10)$$

where $\bar{\kappa}_1, \bar{\kappa}_2 > 0$ depend on d, m, γ, β , and they have an explicit expression given at the end of the proof of [14, Theorem 1.2]. The proof is concluded. ■

3.2. Proof of Theorem 1.3. Uniform relative error convergence

In this subsection we prove the sufficiency part of Theorem 1.3. The converse implication will be proven in the next section. We just recall that partial results in the nonweighted case, $\gamma = \beta = 0$, have been proven in [6, 8, 13, 17, 24, 68, 91, 92]. For the weighted case, see [9, 10].

Theorem 3.3 (UREC). *Let $m \in (m_c, 1)$ and let u be a solution to (CP) with initial data $0 \leq u_0 \in \mathcal{X} \setminus \{0\}$. Then we have*

$$\lim_{t \rightarrow 0} \left\| \frac{u(t, x) - \mathfrak{B}(t, x; M)}{\mathfrak{B}(t, x; M)} \right\|_{L^\infty(\mathbb{R}^d)} = 0, \quad \text{where } M = \|u_0\|_{L^1_\gamma(\mathbb{R}^d)}. \quad (3.11)$$

Proof. It is convenient to work in self-similar variables: we transform $u(t, x)$ into $v(\tau, y)$ according to formula (1.11). We will prove that for any $\varepsilon > 0$ there exists $\tau_\varepsilon > 0$ such that

$$\left\| \frac{v(\tau, y) - \mathfrak{B}_M(y)}{\mathfrak{B}_M(y)} \right\|_{L^\infty(\mathbb{R}^d)} < 2\varepsilon \quad \text{for any } \tau \geq \tau_\varepsilon. \quad (3.12)$$

We argue that we only need to prove the following claim.

Claim. *For any $1 > \varepsilon > 0$ there exists $\rho_\varepsilon > 0$ and $\bar{\tau}_\varepsilon > 0$ such that*

$$\sup_{|y| \geq \rho_\varepsilon} \left| \frac{v(\tau, y) - \mathfrak{B}_M(y)}{\mathfrak{B}_M(y)} \right| < \varepsilon \quad \text{for any } \tau \geq \bar{\tau}_\varepsilon. \quad (3.13)$$

Indeed, once the claim is proven, we just combine it with the convergence inside parabolic cones, i.e. the main result of Theorem 2.4 and obtain inequality (3.12) as follows:

$$\left\| \frac{v(\tau, y)}{\mathfrak{B}_M(y)} - 1 \right\|_{L^\infty(\mathbb{R}^d)} \leq \left\| \frac{v(\tau, y)}{\mathfrak{B}_M(y)} - 1 \right\|_{L^\infty(\{|y| \leq \Upsilon\})} + \left\| \frac{v(\tau, y)}{\mathfrak{B}_M(y)} - 1 \right\|_{L^\infty(\{|y| \geq \Upsilon\})} \leq 2\varepsilon.$$

Recall that the change of variables (1.11) transforms the parabolic cones $\{|x| \leq \Upsilon R(t)\}$ into balls $\{|y| \leq \Upsilon\}$.

Proof of the claim. Let $t_0, R_0 > 0$ be such that $\|u_0\|_{L^1_\gamma(B_{R_0}(0))} > 0$. We know by Theorem 1.1 that

$$\mathfrak{B}(t - \underline{t}, x; \underline{M}) \leq u(t, x) \leq \mathfrak{B}(t + \bar{t}, x; \bar{M})$$

for suitable $\bar{t}, \underline{t} > 0$ and $\bar{M}, \underline{M} > 0$. As a consequence, recalling the change of variables (1.11), we get

$$a(t)^{d-\gamma} (1 \wedge a(t))^{\frac{-\sigma}{1-m}} \mathfrak{B}_{\underline{M}}(y) \leq v(\tau, y) \leq b(t)^{d-\gamma} (b(t) \vee 1)^{\frac{-\sigma}{1-m}} \mathfrak{B}_{\bar{M}}(y), \quad (3.14)$$

where $R(t) = R_\star(t + 1)$ and

$$\tau = \frac{1}{\sigma} \log \frac{R(t)}{R(0)}, \quad a(t) = \frac{R_\star(t + 1)}{R_\star(t + \underline{t})} \quad \text{and} \quad b(t) = \frac{R_\star(t + 1)}{R_\star(t + \bar{t})}.$$

Since $a(t), b(t) \rightarrow 1$ as $t \rightarrow \infty$ we deduce that there exists $\tau_\varepsilon > 0$ such that

$$\left(1 - \frac{\varepsilon}{3}\right) \mathcal{B}_{\underline{M}}(y) \leq v(\tau, y) \leq \left(1 + \frac{\varepsilon}{3}\right) \mathcal{B}_{\bar{M}}(y) \quad \text{for every } \tau > \tau_\varepsilon. \quad (3.15)$$

Recall that all the Barenblatt solutions \mathcal{B}_M have the same behaviour at infinity, which is independent of the mass M , namely $\lim_{|y| \rightarrow \infty} \mathcal{B}_{M_1}(y) / \mathcal{B}_{M_2}(y) = 1$ for any $M_1, M_2 > 0$. Hence, there exists $\rho_\varepsilon = \rho_\varepsilon(\underline{M}, \bar{M}) > 0$ such that

$$1 - \frac{\varepsilon}{3} \leq \frac{\mathcal{B}_{\underline{M}}(y)}{\mathcal{B}_M(y)} \quad \text{and} \quad \frac{\mathcal{B}_{\bar{M}}(y)}{\mathcal{B}_M(y)} \leq 1 + \frac{\varepsilon}{3} \quad \text{for any } |y| \geq \rho_\varepsilon.$$

Combining the above inequality with (3.15) we obtain the proof of the claim. The proof is concluded. \blacksquare

3.3. Proof of Theorem 1.3. The necessary part

We have already shown that the tail condition (TC) implies uniform convergence in relative error: this result is contained in Theorem 3.3. As a consequence, the sufficiency part of Theorem 1.3 has already been proven, and we only need to prove the converse implication: if a solution converges uniformly in relative error, then the initial datum u_0 satisfies the tail condition (TC) or equivalently $u_0 \in \mathcal{X}$. The proof of the necessary part is based on the following lemma, which bears some similarity to a result by Herrero and Pierre [57, Lemma 3.1], valid in the nonweighted case, and then generalized by the authors in [14] to the present weighted case. Notice that we need to integrate over the complementary set of balls, instead of balls as in [14, 57]. The proof is quite similar hence we only sketch it.

Lemma 3.4. *Let $m \in (m_c, 1)$ and let u be a solution to (CP) with initial data $0 \leq u_0 \in L^1_\gamma(\mathbb{R}^d)$. Then for any $R > 0$ and for any $t, s \geq 0$ there exist constants $C_1, C_2 > 0$ which depend on m, d, γ, β such that*

$$\int_{B_{2R}^c(0)} u(t, x) |x|^{-\gamma} dx \leq C_1 \int_{B_R^c(0)} u(s, x) |x|^{-\gamma} dx + C_2 |t - s|^{\frac{1}{1-m}} R^{(d-\gamma) - \frac{\sigma}{1-m}}. \quad (3.16)$$

Proof. Let $R > 0$ and define $A(R) = B_{2R}(0) \setminus B_R(0)$. Let $0 \leq \psi \in C^\infty(\mathbb{R}^d)$ be such that $\psi = 1$ in $B_{2R}(0)^c$ and $\psi = 0$ in $B_R(0)$. Let us sketch the proof. For any $t > 0$, let us formally compute

$$\begin{aligned} & \left| \frac{d}{dt} \int_{\mathbb{R}^d} u(t, x) \psi(x) \frac{dx}{|x|^\gamma} \right| \\ &= \left| \int_{A(R)} u^m(t, x) |x|^\gamma \operatorname{div}(|x|^{-\beta} \nabla \psi) \frac{dx}{|x|^\gamma} \right| \\ &\leq \int_{A(R)} u^m(t, x) \psi^m(x) \psi^{-m}(x) |x|^\gamma \operatorname{div}(|x|^{-\beta} \nabla \psi) \Big| \frac{dx}{|x|^\gamma} \end{aligned}$$

$$\begin{aligned}
&\leq \left(\int_{A(R)} u(t, x) \psi(x) \frac{dx}{|x|^\gamma} \right)^m \left(\int_{A(R)} \psi^{\frac{-m}{1-m}} |x|^\gamma \operatorname{div}(|x|^{-\beta} \nabla \psi) \Big| \frac{1}{1-m} \frac{dx}{|x|^\gamma} \right)^{1-m} \\
&\leq C \left(\int_{\mathbb{R}^d} u(t, x) \psi(x) \frac{dx}{|x|^\gamma} \right)^m R^{(d-\gamma)(1-m)-(2+\beta-\gamma)}, \tag{3.17}
\end{aligned}$$

where we have used Hölder and the fact that $\psi^{\frac{-m}{1-m}} |x|^\gamma \operatorname{div}(|x|^{-\beta} \nabla \psi) \Big| \frac{1}{1-m} \leq c R^{\frac{-2-\beta+\gamma}{1-m}}$, which can be easily derived following the lines of the proof of [14, Lemma 5.2]. Recall that the integration by parts done in the first line of (3.17) presents no difficulties, since $\psi = 0$ in a neighbourhood of the origin, where the weight $|x|^{-\gamma}$ can be singular or degenerate. Integrating the differential inequality (3.17), as done in [57, Lemma 3.1], one obtains (3.16). The proof is concluded. A rigorous proof starts from the integrated version of (3.17) (which follows by definition of weak solutions) and then follows by a Gronwall-type argument. ■

Proof of the necessary part of Theorem 1.3. The proof is based on Lemma 3.4 proven in the Appendix, which we restate here for reader’s convenience in the form that we need. Let u be a solution to (CP) with initial data $0 \leq u_0 \in L_\gamma^1(\mathbb{R}^d)$. Then there exist constants $C_1, C_2 > 0$ which depend on m, d, γ, β such that for any $R > 0$ and for any $t \geq 0$,

$$\int_{B_{2R}^c(0)} u_0(x) |x|^{-\gamma} dx \leq C_1 \int_{B_R^c(0)} u(t, x) |x|^{-\gamma} dx + C_2 t^{\frac{1}{1-m}} R^{(d-\gamma)-\frac{\sigma}{1-m}}. \tag{3.18}$$

Let us proceed with the rest of the proof. Assume now that (1.9) holds, hence there exists a time $\bar{t} > 0$ such that for any $x \in \mathbb{R}^d$,

$$\left| \frac{u(\bar{t}, x) - \mathfrak{B}(\bar{t}, x; M)}{\mathfrak{B}(\bar{t}, x; M)} \right| < 1, \quad \text{hence } u(\bar{t}, x) \leq 2\mathfrak{B}(\bar{t}, x; M).$$

Integrating the latter inequality over B_R^c we get that there exists a constant $\kappa > 0$ which depends on m, d, γ, β and on \bar{t} such that

$$\int_{B_R^c(0)} u(\bar{t}, x) |x|^{-\gamma} dx \leq \kappa R^{d-\frac{\sigma}{1-m}} \quad \text{for all } R > 0.$$

Combining (3.18) at time $t = \bar{t}$ with the above estimate, we conclude that for any $R > 0$,

$$(2R)^{\frac{\sigma}{1-m}-d} \int_{B_{2R}^c(0)} u_0(x) |x|^{-\gamma} dx \leq C \kappa^{\frac{1}{1-m}} + C_{m,d,\gamma,\beta}^{\frac{1}{1-m}} |\bar{t}|^{\frac{1}{1-m}}.$$

As a consequence, the initial data u_0 satisfies the tail condition (TC) and the proof is concluded. ■

3.4. Harnack inequalities for quotients and sharp behaviour at infinity

In this subsection we show a result which can be interpreted as a boundary Harnack inequality, since it extends to the whole space the Harnack inequalities on parabolic cones

of Theorem 2.3. More precisely, there exist a constant $H > 0$ such that for any $t > 0$ (large enough) we have

$$\frac{u(t, x)}{u(t, y)} \leq H \frac{\mathfrak{B}(t, x; M)}{\mathfrak{B}(t, y; M)} \quad \text{for all } x, y \in \mathbb{R}^d. \quad (3.19)$$

The above inequality is equivalent to (3.20) and provides interesting information about the behaviour at infinity of solutions to (CP) with data $u_0 \in \mathcal{X}$. Somehow, the behaviour at infinity does not depend on the mass: indeed, we can show that for all $\bar{M}, \underline{M} > 0$, there exists $\bar{\tau} > 0$ such that

$$1 \leq \frac{\limsup_{|x| \rightarrow \infty} u(t, x) \mathfrak{B}^{-1}(t, x; \bar{M})}{\liminf_{|x| \rightarrow \infty} u(t, x) \mathfrak{B}^{-1}(t, x; \underline{M})} \leq \left(1 + \frac{\bar{\tau}}{t}\right)^{\frac{1}{1-m}} \quad \text{if and only if } u_0 \in \mathcal{X} \setminus \{0\}.$$

The above inequalities are sharp, and are an equivalent statement to (3.21). It is remarkable that equality is attained by Barenblatt profiles, possibly with different mass.

Theorem 3.5. *Let u be a solution to (CP) with $0 \leq u_0 \in \mathcal{X}$ such that $\|u_0\|_{m, \gamma, \beta} = A$ and $\|u_0\|_{L^1_\gamma(\mathbb{R}^d)} = M$ and let $R_0 > 0$ be such that $\|u_0\|_{L^1_\gamma(B_{R_0}(0))} = M/2$. Then there exists a constant $H > 0$, which depends only on m, d, γ, β , such that*

$$\sup_{x \in \mathbb{R}^d} \frac{u(t, x)}{\mathfrak{B}(t, x; M)} \leq H \inf_{x \in \mathbb{R}^d} \frac{u(t, x)}{\mathfrak{B}(t, x; M)} \quad \text{for any } t \geq \bar{\tau}, \quad (3.20)$$

where

$$\bar{\tau} = 3 \max \left\{ A^{1-m} \left(\frac{\bar{\kappa}_1}{\bar{\kappa}_2} \right)^{\frac{1-m}{\sigma\vartheta}} 2^{\frac{7}{\vartheta}}, \kappa_* R_0^{\frac{1}{\vartheta}} (M/2)^{\frac{1}{1-m}} \right\}.$$

The constants $\bar{\kappa}_1, \bar{\kappa}_2$ and κ_* are as in (3.3) and as in (2.1) respectively.

Moreover, we obtain a characterization of the sharp behaviour at infinity, namely we have

$$1 \leq \frac{\limsup_{|x| \rightarrow \infty} u(t, x) |x|^{\frac{\sigma}{1-m}}}{\liminf_{|x| \rightarrow \infty} u(t, x) |x|^{\frac{\sigma}{1-m}}} \leq \left(1 + \frac{\bar{\tau}}{t}\right)^{\frac{1}{1-m}} \quad \text{if and only if } u_0 \in \mathcal{X} \setminus \{0\}. \quad (3.21)$$

Here, $\bar{\tau}$ depends on the initial data and is as in Theorem 3.1.

Proof. We begin by proving inequality (3.20). In what follows we will assume without loss of generality that

$$\kappa_2^{\frac{1}{1-m}} \geq b_1^{-1} 2^{4\sigma+m-2}, \quad (3.22)$$

indeed, since κ_2 comes from the upper bound (3.3) we can choose it as large as needed. By applying Theorem 3.1 at time $t_0 = A^{1-m} (\bar{\kappa}_1/\bar{\kappa}_2)^{\frac{1-m}{\sigma\vartheta}} 2^{\frac{7}{\vartheta}}$ and Theorem 2.1 at time $t_1 = \kappa_* R_0^{\frac{1}{\vartheta}} (M/2)^{\frac{1}{1-m}}$ we obtain that for any $t \geq \bar{\tau}$ the following inequality holds:

$$\mathfrak{B}(t - \underline{\tau}, x; \underline{M}) \leq u(t, x) \leq \mathfrak{B}(t + \bar{\tau}, x; \bar{M}), \quad (3.23)$$

where

$$\underline{t} = t_1/2 = \left(\frac{\kappa_\star}{2}\right) R_0^{\frac{1}{\beta}} \left(\frac{M}{2}\right)^{\frac{1}{1-m}}, \quad \underline{M} = b \frac{M}{2}$$

and

$$\bar{t} = (b_1 2^{2-m-4\sigma} (\bar{\kappa}_2)^{1-m} - 1) t_0, \quad \bar{M} = (2b_0 \bar{\kappa}_1)^{1-m} (b_1 2^{2-m-4\sigma} (\bar{\kappa}_2)^{1-m})^{\frac{d-\gamma}{\sigma}} M.$$

Here is the point where assumption (3.22) enters the game, since it implies that $\bar{t} \geq 0$. By inequality (3.23) it is enough to show that there exists a constant H such that for any $t \geq \bar{t}$,

$$\sup_{x \in \mathbb{R}^d} \frac{\mathfrak{B}(t + \bar{t}, x; \bar{M})}{\mathfrak{B}(t, x; M)} \leq H \inf_{x \in \mathbb{R}^d} \frac{\mathfrak{B}(t - \underline{t}, x; \underline{M})}{\mathfrak{B}(t, x; M)}.$$

A simple computation, which is left to the interested reader, shows that the previous inequality holds with a constant which depends only on m, d, γ, β and not on the mass M nor on the parameter A . The proof of (3.20) is then concluded. ■

Proof of (3.21). We just combine inequality (2.4) of Corollary 2.2 with inequality (3.2) of Remark 3.2. ■

3.5. Rates of convergence in \mathfrak{X}

As we mentioned in Section 1.2, we know that solutions starting from $0 \leq u_0 \in \mathfrak{X}$, will eventually converge to a Barenblatt profile \mathcal{B}_M (with the same mass as u_0), i.e. an element of the manifold $\mathcal{M}_{\mathcal{B}}$. The natural question that we address here is whether there are “universal rates” of convergence towards $\mathcal{M}_{\mathcal{B}}$. More precisely,

in self-similar variables, can we find a speed of convergence to the stationary profile valid for all solutions starting from data in \mathfrak{X} ?

The answer to this question is delicate and cannot be easily given for all $m \in (0, 1)$, nor for all $m \in (m_c, 1)$. Some preliminary remarks are in order. In the case $\gamma = \beta = 0$ the question has a long story (see [37]). When $\frac{d-2}{d} < m < 1$ it has been proven in [21–23, 35, 38, 70, 72, 76], that, under suitable assumptions, there exist (sharp) rates of convergence in different topologies, the most common being d_1 (see Section 1.2). These results were proven by means of relative entropy functionals introduced in [73, 78], or by means of the so-called Bakry–Émery method [3]. The rate t^{-1} of uniform convergence in relative error in the whole range $\frac{d-2}{d} < m < 1$ was first computed in [24] for radial data, and later extended to a larger class of data in [68]. In a series of papers, similar results were obtained in the whole range $m < 1$, cf. [5, 6, 8, 13]; notice that in the range $m < \frac{d-2}{d}$ there is a dramatic change in the behaviour of solutions since mass is not preserved and they can extinguish in finite time; see [92, 93]. In the general case $\gamma \neq 0, \beta \neq 0$, rates of convergence were studied in [9, 10].

In what follows we will show how we can combine the techniques of this paper with those used in [6, 8, 13], to obtain rates of convergence to the Barenblatt profile with an

(almost) uniform rate in the whole of \mathcal{X} . For reasons that are not entirely clear up to now, we need to restrict ourselves to the range $\frac{d-1}{d} = m_1 < m < 1$ in the case $\gamma = \beta = 0$, and to the range $\frac{2d-2-\beta-\gamma}{2(d-\gamma)} < m < 1$ for the general case; see [9, 10] for further remarks. The latter restriction is somehow natural, since, at least when $\gamma = \beta = 0$, we have that the FDE is a gradient flow of a displacement convex functional (the relative entropy) with respect to the so-called Wasserstein distance; see [71, 72, 76]. The displacement convexity is lost below m_1 . The main result reads as follows:

Theorem 3.6 (Almost optimal rates of convergence in the nonweighted case). *Let u be the solution to (CP) corresponding to the initial data $0 \leq u_0 \in \mathcal{X} \setminus \{0\}$, $\int_{\mathbb{R}^d} x u_0(x) dx = 0$ and assume that $\beta = \gamma = 0$ and $m \in (\frac{d-1}{d}, 1)$. Then, for every $\delta \in (0, 1)$, there exist $t_\delta, c_\delta > 0$ (which may also depend on u_0) such that for all $t > t_\delta$,*

$$\|u(t) - \mathfrak{B}(t; M)\|_{L^1(\mathbb{R}^d)} \leq \frac{c_\delta}{t^{1-\delta}} \quad \text{and} \quad t^{d\vartheta} \|u(t) - \mathfrak{B}(t; M)\|_{L^\infty(\mathbb{R}^d)} \leq \frac{c_\delta}{t^{1-\delta}}, \quad (3.24)$$

where $M = \|u_0\|_{L^1(\mathbb{R}^d)}$.

Remark. Notice that the above result is new for the whole space \mathcal{X} even if we are dealing with the case $\gamma = \beta = 0$. Indeed, all the previous results deal with more restrictive assumptions such as radial data, a very precise control for $|x| \rightarrow \infty$ or being sandwiched between two Barenblatt profiles.

When dealing with CKN weights, the result is a bit weaker, because of the possible lack of C^k regularity at the origin and reads as follows:

Theorem 3.7 (Minimal rates of convergence in the weighted case). *Assume $\gamma < 0$ and let u be the solution to (CP) corresponding to the initial data $0 \leq u_0 \in \mathcal{X} \setminus \{0\}$ and let $\frac{2d-2-\beta-\gamma}{2(d-\gamma)} < m < 1$. Then there exists a $\delta_* \in (0, 1)$ such that for every $\delta \in (0, \delta_*)$ there exist $t_\delta, c_\delta > 0$ (which may also depend on u_0) such that for all $t > t_\delta$,*

$$\begin{aligned} \|u(t) - \mathfrak{B}(t; M)\|_{L^1_\gamma(\mathbb{R}^d)} &\leq \frac{c_\delta}{t^{1-\delta}}, \\ t^{(d-\gamma)\vartheta} \|u(t) - \mathfrak{B}(t; M)\|_{L^\infty(\mathbb{R}^d)} &\leq \frac{c_\delta}{t^{1-\delta}}, \end{aligned} \quad (3.25)$$

where $M = \|u_0\|_{L^1_\gamma(\mathbb{R}^d)}$.

If we consider radial initial data in \mathcal{X} we can provide a *universal* rate of convergence, very much in the spirit of [24] or [68].

Theorem 3.8 (Sharp universal rates for radial data). *Let u be the solution to (CP) corresponding to the radial initial data $0 \leq u_0 \in \mathcal{X} \setminus \{0\}$ and assume that $\gamma = \beta = 0$ and $m \in (\frac{d-2}{d}, 1)$. Then there exist $t_0, c_0 > 0$ (which may also depend on u_0) such that for all $t > t_0$,*

$$\left\| \frac{u(t)}{\mathfrak{B}(t; M)} - 1 \right\|_{L^\infty(\mathbb{R}^d)} \leq \frac{c_0}{t}, \quad (3.26)$$

where $M = \|u_0\|_{L^1(\mathbb{R}^d)}$.

Remark 3.9. As an immediate consequence of (3.26) we obtain that for all $t \geq t_0$,

$$\begin{aligned} \|u(t) - \mathfrak{B}(t; M)\|_{L^1(\mathbb{R}^d)} &\leq \frac{c_0}{t}, \\ t^{(d-\gamma)\vartheta} \|u(t) - \mathfrak{B}(t; M)\|_{L^\infty(\mathbb{R}^d)} &\leq \frac{c_0}{t}. \end{aligned} \quad (3.27)$$

The above theorem solves a problem left open in [24], i.e. identifying the largest class of nonnegative *radial* L^1 data for which the above rate of convergence holds. Such rates are proven to be sharp, since they are fulfilled by two time-shifted Barenblatt, with the same mass, see [24, 68]. Finally, we observe that, even if we restrict the analysis to radial data, the class \mathcal{X} is much larger than those considered up to now in the literature: we refer to Section 5 for examples of functions in \mathcal{X} with substantially different behaviour from the Barenblatt profile.

Finally, let us give the proof of the above statements.

Proof of Theorem 3.6. Here we exploit the techniques introduced in [6, 8, 13]. Let us rescale $u(t, x)$ to $v(\tau, y)$ according to the change of variables (1.11) and define $w := \frac{v(\tau, y)}{\mathfrak{B}_M(y)}$ where $M = \|u_0\|_{L^1(\mathbb{R}^d)}$. Let us define the *free energy* or *relative entropy* $\mathcal{F}[w]$ and the *Fisher information* $\mathcal{I}[w]$ as

$$\begin{aligned} \mathcal{F}[w(\tau)] &:= \frac{m}{m-1} \int_{\mathbb{R}^d} \left[\frac{w^m - 1}{m} - (w-1) \right] \mathfrak{B}_M^m dy, \\ \mathcal{I}[w] &:= \frac{m}{1-m} \int_{\mathbb{R}^d} w \mathfrak{B}_M |\nabla[(w^{m-1} - 1)\mathfrak{B}_M^{m-1}]|^2 dy. \end{aligned} \quad (3.28)$$

The Fisher information is related to the relative entropy by the time derivative along the flow,

$$\frac{d}{d\tau} \mathcal{F}[w] = -\mathcal{I}[w]. \quad (3.29)$$

It is well known that the relative entropy controls the L^1 distance between the solution $v(\tau, y)$ and the Barenblatt profile \mathfrak{B}_M , via the celebrated Csiszár–Kullback inequality (see e.g. [24, 27, 44, 69, 76]), more precisely,

$$\|v(\tau) - \mathfrak{B}_M\|_{L^1(\mathbb{R}^d)} \leq \left(\frac{8}{m} \|\mathfrak{B}_M^{2-m}\|_{L^1(\mathbb{R}^d)} \right)^{\frac{1}{2}} \sqrt{\mathcal{F}[w]}. \quad (3.30)$$

Therefore the decay of relative entropy implies the same decay as $\|u(t) - \mathfrak{B}_M(t)\|_{L^1(\mathbb{R}^d)}$.

For any $m \in (\frac{d-1}{d}, 1)$ the *entropy–entropy production* inequality reads

$$4\mathcal{F}[w] \leq \mathcal{I}[w], \quad (3.31)$$

and it is well known to be equivalent to a member of a suitable family of (optimal) Gagliardo–Nirenberg inequalities; more details can be found in the pioneering work of

Del Pino and Dolbeault [35]. The best constant in (3.31) is 4, and combining it with (3.29), we obtain the (sharp) exponential decay of the relative entropy along the flow, namely

$$\mathcal{F}[w] \leq \mathcal{F}[w_0]e^{-4\tau}.$$

The strategy in [6, 8, 13] consists in proving a faster decay of the entropy along the flow using an improved (with a larger constant) entropy–entropy production inequality along the flow. Such an improved inequality is obtained by means of Hardy–Poincaré-type inequalities (with improved constants), and by means of quantitative inequalities that compare the linear(ized) entropy and Fisher information with their nonlinear counterparts. Notice that we are in position to apply the results of [6, 13], since the running assumption guarantees that GHP (Theorem 1.1) holds, and implies the validity of assumption (H1)'' in those papers. Combining [6, Lemma 3, Theorem 7] with [8, Lemma 1], we can prove the following claim.

Claim. *For any $0 < \delta < 4\frac{1-\vartheta}{\vartheta}$ there exists a time $\tau_\delta > 0$ such that*

$$\left(\frac{4}{\vartheta} - \delta\right) \mathcal{F}[w(\tau)] \leq \mathcal{I}[w(\tau)] \quad \text{for any } \tau \geq \tau_\delta, \quad (3.32)$$

where ϑ is as in (1.3).

Sketch of the proof of the claim. We will not provide lengthy details of the proof of the above claim, but we will just explain how to deduce it as a straightforward combination of already published results, adapting them to the current notation. The claim follows by [8, formula (11)], which in the current notation takes the form (at least for sufficiently large times)

$$\frac{2[\Lambda_{\alpha,d} - d(1-m)((1+\varepsilon)^{4(2-m)} - 1)]}{(1+\varepsilon)^{7-3m}} \mathcal{F}[w(\tau)] \leq \mathcal{I}[w(\tau)], \quad (3.33)$$

where ε is (roughly speaking) the size of the relative error $|w - 1| \sim \varepsilon$, which we need to be small in order to guarantee the validity of the result (note that in [8, formula (11)] $h = \max\{\sup_{\mathbb{R}^d} |w|, 1 - \inf_{\mathbb{R}^d} |w|\} \sim 1 + \varepsilon$). Notice that everything is quantified explicitly in terms of ε in [8] which also relies on precise results of [6, 13]. The smallness of ε for sufficiently large times follows by our Theorem 1.1, the global Harnack principle, together with the uniform convergence in relative norm, Theorem 3.3. Now recalling [13, Lemma 1], we get the expression for $\Lambda_{\alpha,d} = -4\alpha - 2d$, which in our notation becomes $\Lambda_{\alpha,d} = \frac{2}{\vartheta(1-m)}$. Note that we need to assume that the first moment is fixed, but this is well known to be true along the nonlinear flow as well; see [13]. This concludes the proof of the claim. ■

As a consequence of inequality (3.32), we obtain a faster decay of the relative entropy and conclude that

$$\|v(\tau) - \mathcal{B}_M\|_{L^1(\mathbb{R}^d)} \leq C_\delta e^{-\left(\frac{2}{\vartheta} - \frac{\delta}{2}\right)\tau} \quad \text{for any } \tau \geq \tau_\delta. \quad (3.34)$$

By rescaling back to original variables and observing that $e^{2\tau} = R(t) \sim t^\delta$ one concludes that

$$\|u(t, \cdot) - \mathfrak{B}(t, \cdot; M)\|_{L^1(\mathbb{R}^d)} \leq C_\delta t^{-1 + \frac{\delta}{4}\vartheta}.$$

Notice that $\delta > 0$ is arbitrary, but C_δ may diverge as $\delta \rightarrow 0^+$. This proves the left inequality in (3.24).

It only remains to prove the second inequality in (3.24). To do so we need to invoke the following interpolation lemma which goes back to Gagliardo (see [49]) and Nirenberg (see [75, p. 126]): let $f \in C^k(\mathbb{R}^d) \cap L^1(\mathbb{R}^d)$ for some $p \geq 1$ and k be a positive integer; then

$$\|f\|_{L^\infty(\mathbb{R}^d)} \leq C_{p,k,d} \|f\|_{C^k(\mathbb{R}^d)}^{\frac{d}{d+k}} \|f\|_{L^1(\mathbb{R}^d)}^{\frac{k}{d+k}}, \quad (3.35)$$

where $\|\cdot\|_{C^k(\mathbb{R}^d)}$ is given by

$$\|f\|_{C^k(\mathbb{R}^d)} := \max_{|\eta|=k} \sup_{z \in \mathbb{R}^d} |\partial^\eta f(z)|,$$

where $|\eta| = \eta_1 + \dots + \eta_d$ is the length of the multi-index $\eta = (\eta_1, \dots, \eta_d) \in \mathbb{Z}^d$. We recall that in the case $\gamma = \beta = 0$, solutions to (CP) are $C^\infty(\mathbb{R}^d)$ and for any $k \geq 1$ we have

$$\sup_{\tau \geq \tau_0} \|v(\tau) - \mathcal{B}_M\|_{C^k(\mathbb{R}^d)} < \infty.$$

For a proof of the above inequality see [6, Theorems 2 and 4]. Fix $k \geq 1$ to be chosen later. Combining the above interpolation inequality (3.35) with the decay of the L^1 norm given in (3.34) one obtains

$$\|v(\tau) - \mathcal{B}_M\|_{L^\infty} \lesssim e^{-(\frac{2}{\vartheta} - \frac{\delta}{2})(\frac{k}{d+k})},$$

and rescaling back to original variables we easily find that

$$t^{d\vartheta} \|u(t, \cdot) - \mathfrak{B}(t, \cdot; M)\|_{L^\infty} \leq C t^{-(1 - \frac{\delta\vartheta}{4} - \frac{d}{d+k} + \frac{d\delta\vartheta}{4(d+k)})},$$

since both k and δ are arbitrary we conclude that the second inequality in (3.24) holds by choosing k sufficiently large, for instance $k \geq \delta^{-4}$. This concludes the proof. ■

Proof of Theorem 3.7. The proof is very similar to that of Theorem 3.6; here we only explain the main differences. We cannot reach the rate $t^{-1+\delta}$ for two reasons. First, we can obtain an inequality like (3.32), although the constant is smaller than $(\frac{4}{\vartheta} - \delta)$; see [9, 10]. Second, we need to assume $\gamma < 0$ to obtain an inequality similar to (3.35); see Lemma A.1. Finally, solutions to (CP) do not enjoy C^∞ regularity, indeed they can be only C^α at the origin (cf. [14]), hence inequality (3.35) does not apply as in the proof of Theorem 3.6: we can only interpolate with C^α norms, and we cannot choose k arbitrarily large. This concludes our considerations. ■

Proof of Theorem 3.8. In [24], the authors proved Theorem 3.8 under the assumption that the initial data u_0 is bounded, radially symmetric and satisfies $u_0 = O(|x|^{-\frac{2}{1-m}})$. It is only needed to show that radial data in \mathcal{X} produce solutions that satisfy the decay assumption above for any time $t > t_0$ for some given t_0 . This is exactly the statement of the GHP, Theorem 3.1. The proof is concluded. \blacksquare

4. Counterexamples and generalized global Harnack principle

In this section we carefully construct the family of sub/supersolutions presented in the introduction. Such a phenomenon is possible only in \mathcal{X}^c . We also show examples of an anomalous “fat-tail” behaviour for both integrable and non-integrable solutions. At the end of this section we show that in \mathcal{X}^c the convergence toward the Barenblatt is slower due to the different tail behaviour of solutions.

4.1. Construction of a family of subsolutions and anomalous tail behaviour

In the following proposition we construct an explicit family of subsolutions parameterized by the powers of their decay at infinity. Every subsolution decays in space more slowly than the Barenblatt profile.

Proposition 4.1 (Family of L^1_γ -subsolutions). *Let $m \in (m_c, 1)$, $\varepsilon \in (0, \frac{2}{1-m} - \frac{2}{\sigma}(d - \gamma))$, $A, B > 0$ and*

$$\alpha = \frac{1}{1-m} - \frac{\varepsilon}{2} > 0.$$

Define for some $t_0 \in \mathbb{R}$ the function

$$D(t) := (\sigma A^{m-1} m B (d - \gamma) (1 - \alpha(1 - m)) t + t_0)^{\frac{1}{1-\alpha(1-m)}}. \quad (4.1)$$

Then, for all $t > 0$, the $L^1_\gamma(\mathbb{R}^d)$ function

$$\underline{V}(t, x) = \frac{A}{(D(t) + B|x|^\sigma)^\alpha} \quad (4.2)$$

is a subsolution to (CP). If $m \in (\frac{d-\gamma}{(d-\gamma+\sigma)}, 1)$ and $\varepsilon \in (0, \frac{2}{1-m} - \frac{2(d-\gamma)}{\sigma} - 2)$ then $|x|^\sigma \underline{V}(t, x) \in L^1_\gamma(\mathbb{R}^d)$.

Remark 4.2. We notice that $\|V(t, \cdot)\|_{L^\infty(\mathbb{R}^d)} \asymp t^{-\frac{\alpha}{1-\alpha(1-m)}}$ as $t \rightarrow \infty$. This is not in contrast with the smoothing effect (inequality (2.17)) which implies that any solution $u(t, x)$ to (CP) decays in time less than $t^{-(d-\gamma)\vartheta}$: a simple computation shows that the condition $\varepsilon \in (0, \frac{2}{1-m} - 2\frac{d-\gamma}{\sigma})$ implies that $t^{-\frac{\alpha}{1-\alpha(1-m)}} < t^{-(d-\gamma)\vartheta}$.

However, as $|x| \rightarrow \infty$, $V(t, x)$ exhibits quite interesting behaviour, namely $V(t, x) \asymp |x|^{-\sigma\alpha}$. The power $-\sigma\alpha$ does not match that of the fundamental solution: indeed, we have $\mathfrak{B}(t, x; M) \asymp |x|^{-\frac{\sigma}{1-m}}$ as $|x| \rightarrow \infty$. This proves that, for any choice of the parameters A ,

B , t_0 and for any choice of the mass M , the inequality $V(t, x) > \mathfrak{B}(t, x; M)$ holds for $|x|$ large enough.

As a final remark, we can define another family of subsolutions. Indeed, for some choice of the parameters B' , F' and T the function W defined as

$$W(t, x) = \frac{(T - t)^{\frac{1}{1-m}}}{(B' + F'|x|^\sigma)^\alpha}$$

is a subsolution to (CP) which has the same qualitative behaviour as $|x| \rightarrow \infty$. The drawback is that this is meaningful only on a finite time interval, hence we prefer to use \underline{V} .

Proof of Proposition 4.1. We just need to verify that the function $\underline{V}(t, x)$ defined in (4.2) satisfies, in a pointwise sense, the inequality

$$\partial_t \underline{V}(t, x) \leq |x|^\gamma \operatorname{div}(|x|^{-\beta} \nabla \underline{V}^m). \quad (4.3)$$

Let $r = |x|$, and write $-$ with a small abuse of notation $-\underline{V}(t, r)$ instead of $\underline{V}(t, x)$. We recall that the operator $\mathcal{L}_{\gamma, \beta} = |x|^\gamma \nabla \cdot (|x|^{-\beta} \nabla f)$ acts on a radial function $f(r)$ in the following way:

$$\mathcal{L}_{\gamma, \beta}(f) = r^{\gamma-\beta} \left(f''(r) + \frac{(d-1-\beta)}{r} f'(r) \right). \quad (4.4)$$

A straightforward computation shows the following identities:

$$\begin{aligned} \partial_t \underline{V}(t, r) &= \frac{-A\alpha \partial_t D(t)}{(D(t) + Fr^\sigma)^{\alpha+1}}, \\ \mathcal{L}_{\gamma, \beta}(\underline{V}(t, r)^m) &= \frac{-(\sigma\alpha m A^m F)}{(D(t) + Fr^\sigma)^{\alpha m+2}} [(d-\gamma)D(t) + Fr^\sigma(-\sigma\alpha m + d - 2 - \beta)]. \end{aligned}$$

As a consequence, the inequality $\partial_t \underline{V}(t, r) \leq \mathcal{L}_{\gamma, \beta}(\underline{V}(t, r)^m)$ is satisfied if and only if

$$\partial_t D(t) \geq \frac{\sigma m F A^{m-1}}{(D(t) + Fr^\sigma)^{\alpha(m-1)+1}} [(d-\gamma)D(t) + Fr^\sigma(-\sigma\alpha m + d - 2 - \beta)]. \quad (4.5)$$

The reader may notice that if $\varepsilon < 2/(1-m) - 2\frac{(d-\gamma)}{\sigma}$ then on the right-hand side of inequality (4.5) the term $Fr^\sigma(-\sigma\alpha m + d - 2 - \beta)$ is negative. A simple computation then shows that the supremum of the right-hand side of inequality (4.5) is achieved at $r = 0$. Hence the inequality $\partial_t \underline{V}(t, r) \leq \mathcal{L}_{\gamma, \beta}(\underline{V}(t, r)^m)$ will follow by asking that

$$\begin{aligned} \partial_t D(t) &\geq \sigma m F A^{m-1} (d-\gamma) D(t)^{\alpha(1-m)} \\ &= \sup_{r \geq 0} \frac{\sigma m F A^{m-1}}{(D(t) + Fr^\sigma)^{\alpha(m-1)+1}} [(d-\gamma)D(t) + Fr^\sigma(-\sigma\alpha m + d - 2 - \beta)]. \end{aligned}$$

We conclude the proof observing that, for any $t_0 \in \mathbb{R}^d$, such an inequality is satisfied by the function $D(t)$ defined in (4.1). \blacksquare

Anomalous tail behaviour. As a corollary of Proposition 4.1 we have the following results about unexpected “fat-tails”, both integrable and nonintegrable.

Corollary 4.3 (Anomalous integrable tail behaviour). *Under the assumptions of Proposition 4.1,*

$$\text{if } u_0(x) \geq \frac{A}{(C + B|x|^\sigma)^\alpha}, \text{ then for any } t > 0 \text{ we have } \liminf_{|x| \rightarrow \infty} |x|^{\sigma\alpha} u(t, x) \geq \frac{A}{B}.$$

Moreover, for any $t > 0$ we have

$$\sup_{x \in \mathbb{R}^d} \left| \frac{u(t, x)}{\mathfrak{B}(t, x; M)} - 1 \right| = \infty. \quad (4.6)$$

Proof. The proof of the first statement of the above corollary is an immediate application of Proposition 4.1. As for formula (4.6), we see that if the initial data u_0 do not satisfy the tail condition (TC) there is no chance to conclude the convergence to the Barenblatt profile in *uniform relative error*. Indeed, since

$$u_0(x) \geq \frac{A}{(C + B|x|^\sigma)^\alpha},$$

by Proposition 4.1 we get

$$u(t, x) \geq \frac{A}{(D(t) + B|x|^\sigma)^\alpha}$$

for any $t > 0$, where $D(t) = (\sigma A^{m-1} m B (d - \gamma) (1 - \alpha(1 - m)) t + C^{1-\alpha(1-m)})^{\frac{1}{1-\alpha(1-m)}}$. A simple computation shows that the quotient

$$\begin{aligned} \frac{1}{\mathfrak{B}(t, x; M)} \frac{A}{(D(t) + B|x|^\sigma)^\alpha} &= \frac{A}{t^{\frac{1}{1-m}}} \frac{[b_0 \frac{t^{\sigma\delta}}{M^{\sigma\delta(1-m)}} + b_1 |x|^\sigma]^{\frac{1}{1-m}}}{(D(t) + B|x|^\sigma)^{\frac{1}{1-m} - \frac{\varepsilon}{1}}} \\ &\sim \frac{A b_1^{\frac{1}{1-m}} |x|^{\frac{\sigma\varepsilon}{2}}}{B^{\frac{1}{1-m} - \frac{\varepsilon}{2}} t^{\frac{1}{1-m}}} \quad \text{as } |x| \rightarrow \infty, \end{aligned}$$

from which we deduce (4.6). The proof is then complete. \blacksquare

A family of subsolutions not in L_γ^1 . A closer inspection of the proof of Proposition 4.1 reveals that the condition on ε for \underline{V} to be a subsolution is actually $m\varepsilon \leq \frac{2}{1-m} - \frac{2}{\sigma}(d - \gamma)$, while for $\varepsilon > \frac{1}{m}(\frac{2}{1-m} - \frac{2}{\sigma}(d - \gamma))$, \underline{V} ceases to be a subsolution. Indeed, when $\varepsilon > \frac{2}{1-m} - \frac{2}{\sigma}(d - \gamma)$, we can construct subsolutions which do not belong to $L_\gamma^1(\mathbb{R}^d)$, so that, in general, initial data $u_0 \notin L_\gamma^1(\mathbb{R}^d)$ do not produce $L_\gamma^1(\mathbb{R}^d)$ solutions for any time $t > 0$. We resume this fact in the following corollary, considering the parameters $m, \varepsilon, A, B, t_0$ in the “nonintegrability range”:

$$\begin{aligned} m \in (m_c, 1), \quad \varepsilon \in \left[\frac{2}{1-m} - \frac{2}{\sigma}(d - \gamma), \frac{1}{m} \left(\frac{2}{1-m} - \frac{2}{\sigma}(d - \gamma) \right) \right], \\ A, B, t_0 > 0, \quad \alpha = \frac{1}{1-m} - \frac{\varepsilon}{2} > 0. \end{aligned} \quad (4.7)$$

Corollary 4.4 (Family of subsolutions not in L_γ^1). *Under assumptions (4.7), the function $\underline{V}(t, x)$ defined in (4.2) is a nonintegrable subsolution to (CP). If $u(t, x)$ is a (super)solution to (CP), we have*

$$u_0 \geq \underline{V}(0, \cdot) \notin L_\gamma^1(\mathbb{R}^d), \quad \text{then } u(t, \cdot) \geq \underline{V}(t, \cdot) \notin L_\gamma^1(\mathbb{R}^d).$$

4.2. Construction of a family of supersolutions

In this section we construct a family of supersolutions which share the same spatial tail behaviour with the subsolutions constructed in the previous section in Proposition 4.1.

Proposition 4.5 (Family of L_γ^1 -supersolutions). *Let $m \in (m_c, 1)$, $\varepsilon \in (0, \frac{2}{1-m} - \frac{2}{\sigma}(d - \gamma))$, $E, F > 0$ and $\alpha = \frac{1}{1-m} - \frac{\varepsilon}{2} > 0$. Define for some $t_0 \in \mathbb{R}$ and $H > 0$ the function*

$$\bar{V}(t, x) = \frac{EG(t)^\alpha}{(G(t) + F|x|^\sigma)^\alpha} \in L_\gamma^1(\mathbb{R}^d), \quad \text{where } G(t) := t_0 + Ht. \quad (4.8)$$

Then \bar{V} is a supersolution for all $t > 0$ whenever H is sufficiently big or, more precisely, for all

$$H \geq m\sigma F^2 E^{m-1} (2 + \beta - d + \sigma\alpha m). \quad (4.9)$$

Proof of Proposition 4.5. We just need to verify that the function $\bar{V}(t, x)$ defined in (4.8) satisfies the inequality

$$\partial_t \bar{V}(t, x) \geq |x|^\gamma \operatorname{div}(|x|^{-\beta} \nabla \bar{V}^m)(t, x) \quad (4.10)$$

under assumption (4.9). Let $r = |x|$, and write $-$ with a small abuse of notation $-\bar{V}(t, r)$ instead of $\bar{V}(t, x)$. We have the following identities (recalling the radial form of $\mathcal{L}_{\gamma, \beta}$, formula (4.4)):

$$\begin{aligned} \partial_t \bar{V}(t, r) &= \frac{\alpha EG(t)^{\alpha-1} H}{(G(t) + Fr^\sigma)^{\alpha+1}} Fr^\sigma, \\ \mathcal{L}_{\gamma, \beta}(\bar{V}^m(t, r)) &= \frac{\sigma\alpha m FE^m G(t)^{\alpha m}}{(G(t) + Fr^\sigma)^{\alpha m+2}} [Fr^\sigma (2 + \beta - d + \sigma\alpha m) - (d - \gamma)G(t)]. \end{aligned}$$

It is straightforward to verify that (4.10) holds at $r = 0$ since for any $t > 0$ the derivative in time $\partial_t \bar{V}(t, 0) = 0$ and $\mathcal{L}_{\gamma, \beta}(\bar{V}^m(t, 0))$ is negative. When $r > 0$, a direct computation shows that (4.10) is equivalent to

$$\begin{aligned} H &\geq \left(\frac{G(t)}{G(t) + Fr^\sigma} \right)^{1-\alpha(1-m)} m\sigma FE^{m-1} \\ &\quad \times \left[F(2 + \beta - d + \sigma\alpha m) - (d - \gamma) \frac{G(t)}{r^\sigma} \right]. \end{aligned} \quad (4.11)$$

Finally, we check that (4.9) implies (4.11), just by using that $(\frac{G(t)}{G(t)+Fr^\sigma})^{1-\alpha(1-m)} < 1$ and $(d - \gamma) \frac{G(t)}{r^\sigma} > 0$. Therefore, $\bar{V}(t, x)$ is a supersolution and the proof is concluded. \blacksquare

A closer look at the above proof reveals that if we allow $\varepsilon < 0$, then $\bar{V}(t, x)$ ceases to be a supersolution. Indeed, if $\varepsilon < 0$ we have $1 - \alpha(1 - m) < 0$ and so in (4.11) we would have

$$\left(\frac{G(t)}{G(t) + Fr^\sigma}\right)^{1-\alpha(1-m)} = \left(1 + \frac{Fr^\sigma}{G(t)}\right)^{\alpha(1-m)-1} \rightarrow \infty \quad \text{as } r \rightarrow \infty,$$

and an inequality such as (4.11) would be impossible.

4.3. Slower convergence rates in \mathcal{X}^c

We showed in Section 3.5 that when $u_0 \in \mathcal{X}$, there are always power-type rates of convergence to a Barenblatt profile and in some cases the sharp decay rate $O(1/t)$ is obtained. In this paragraph we show, by means of an explicit counterexample, that power-like decay rates are simply not possible for general data outside \mathcal{X} . However, we are not able to exclude the possibility of slower decay rates (e.g. log, log log). The latter question is really delicate and deserves a thorough study that goes beyond the scope of this paper.

Theorem 4.6. *For any $\delta > 0$ there exists initial data $u_{0,\delta} \in L^1_\gamma(\mathbb{R}^d)$ such that the corresponding solution $u_\delta(t, x)$ to (CP) satisfies*

$$\begin{aligned} \lim_{t \rightarrow \infty} t^{\delta+(d-\gamma)\vartheta} \|u_\delta(t) - \mathfrak{B}(t; M)\|_{L^\infty(\mathbb{R}^d)} &= \infty, \\ \lim_{t \rightarrow \infty} t^\delta \|u_\delta(t) - \mathfrak{B}(t; M)\|_{L^1_\gamma(\mathbb{R}^d)} &= \infty. \end{aligned} \quad (4.12)$$

Proof. We only give a detailed proof of the first limit in (4.12) in the nonweighted case $\gamma = \beta = 0$, the weighted case being completely analogous. For the same reasons, we will just sketch the proof of the second limit in the nonweighted case.

Proof of the first limit. Fix any $\delta > 0$ and let $\varepsilon \in (0, \frac{2}{1-m} - d)$ be such that

$$\delta > \frac{2\vartheta}{\varepsilon(1-m)} \left(\frac{2}{1-m} - d - \varepsilon \right). \quad (4.13)$$

A simple computation shows that this choice of ε is always possible. Let $u_\delta(t, x)$ be the solution to (CP) corresponding to the initial data

$$u_{0,\delta}(x) = \frac{A}{(1 + B|x|^2)^\alpha}, \quad \text{where } \alpha = \frac{1}{1-m} - \frac{\varepsilon}{2}$$

and $A, B > 0$ are chosen such that $\|u_{0,\delta}\|_{L^1(\mathbb{R}^d)} = 1$.

A simple rescaling shows that for any choice of B , we can always choose A such that $\|u_{0,\delta}\|_{L^1(\mathbb{R}^d)} = 1$, namely

$$A \int_{\mathbb{R}^d} \frac{dy}{(1 + |y|^2)^\alpha} = B^{\frac{d}{2}}. \quad (4.14)$$

Let us consider the subsolution \underline{V} given in Proposition (4.1):

$$\underline{V}(t, x) = \frac{A}{(D(t) + B|x|^2)^\alpha},$$

where $D(t) := (2A^{m-1}mBd(1 - \alpha(1 - m))t + 1)^{\frac{1}{1-\alpha(1-m)}}$. We recall that $u_{0,\delta}(x) = \underline{V}(0, x)$ so that by comparison $u_\delta(t, x) \geq \underline{V}(t, x)$ for all $t \geq 0$. The following claim allows us to conclude the proof:

Claim. *For any B sufficiently large, there exists $t_0 = t_0(B) > 0$ and $\underline{c} > 0$ such that for any $t \geq t_0$ and for any $|x|^2 \in [D(t), 2D(t)]$, we have*

$$\underline{V}(t, x) > \mathfrak{B}(t, x; 1) \quad \text{and} \quad (\underline{V}(t, x) - \mathfrak{B}(t, x; 1)) > \underline{c}t^{-\frac{\alpha}{1-\alpha(1-m)}}. \quad (4.15)$$

Let us assume momentarily the validity of the claim and conclude the proof. The above inequality (4.15) immediately implies that $u_\delta(t, x) \geq \mathfrak{B}(t, x; 1)$ for any $|x|^2 \in [D(t), 2D(t)]$. Now let $|x|^2 = D(t)$; we then have

$$\begin{aligned} t^{d\vartheta+\delta} \|u_\delta(t) - \mathfrak{B}(t; 1)\|_{L^\infty(\mathbb{R}^d)} &\geq t^{d\vartheta+\delta} |u_\delta(t, x) - \mathfrak{B}(t, x; 1)| \\ &\geq t^{d\vartheta+\delta} (u_\delta(t, x) - \mathfrak{B}(t, x; 1)) \\ &\geq t^{d\vartheta+\delta} (\underline{V}(t, x) - \mathfrak{B}(t, x; 1)) \\ &\geq \underline{c}t^{d\vartheta+\delta-\frac{\alpha}{1-\alpha(1-m)}}, \end{aligned} \quad (4.16)$$

where in the last line we have used (4.15) again. Notice that $d\vartheta + \delta - \frac{\alpha}{1-\alpha(1-m)} > 0$, since δ is as in (4.13). Therefore (4.12) follows as a consequence of (4.16). The proof of the first limit is complete.

Proof of the second limit. The proof follows by integrating the last line of inequality (4.16) over the region $|x|^2 \in [D(t), 2D(t)]$, and recalling that $u_\delta(t, x) \geq \underline{V}(t, x)$ for all $t \geq 0$.

It only remains to prove the claim.

Proof of the claim. Let us define $R(t) = \{x \in \mathbb{R}^d : D(t) \leq |x|^2 \leq 2D(t)\}$ for any $t > 0$. To prove the first inequality in (4.15) we need to check that for t large enough we have

$$\inf_{x \in R(t)} \underline{V}(t, x) > \sup_{x \in R(t)} \mathfrak{B}(t, x; 1),$$

which amounts to proving that $\underline{V}(t, 2D(t)) > \mathfrak{B}(t, D(t); 1)$, where in the last inequality there is a small abuse of notation: we write \mathfrak{B} as a radial function $\mathfrak{B}(t, r; 1)$. We rewrite $\underline{V}(t, 2D(t)) > \mathfrak{B}(t, D(t); 1)$ in the equivalent form

$$\frac{A}{(1 + 2B)^\alpha} \frac{D(t)^{\frac{\alpha}{2}}}{t^{\frac{1}{1-m}}} - \left[b_1 + b_0 \frac{t^{2\vartheta}}{D(t)} \right]^{\frac{1}{1-m}} \geq a > 0 \quad (4.17)$$

for some $a > 0$. The proof of (4.17) follows by observing that in the limit $t \rightarrow \infty$, we have that for all B large enough,

$$(1 + 2B)^{\frac{\alpha}{2}} \left(\frac{B}{1 + 2B} \right)^{\frac{1}{1-m}} (md\varepsilon(1 - m)) > b_1^{\frac{1}{m-1}}.$$

Hence, inequality (4.17) holds for B and t large enough, since it is true in the limit $t \rightarrow \infty$ and all the quantities that appear in (4.17) are continuous with respect to t , $B > 0$.

It only remains to prove the last inequality in (4.15): for any $x \in R(t)$ we have

$$\begin{aligned} \underline{V}(t, x) - \mathfrak{B}(t, x; 1) &\geq \underline{V}(t, 2D(t)) - \mathfrak{B}(t, D(t); 1) \\ &= \left(\frac{t}{D(t)}\right)^{\frac{1}{1-m}} \left[\frac{A}{(1+2B)^\alpha} \frac{D(t)^{\frac{\alpha}{2}}}{t^{\frac{1}{1-m}}} - \frac{1}{(b_1 + b_0 \frac{t^{2\beta}}{D(t)})^{\frac{1}{1-m}}} \right] \\ &\geq \underline{c} t^{\frac{-\alpha}{1-\alpha(1-m)}}, \end{aligned}$$

where we have used that $t/D(t) \asymp t^{\frac{-\alpha(1-m)}{1-\alpha(1-m)}}$ and (4.17). The proofs of the claim and of the theorem are now concluded. \blacksquare

5. On the fast diffusion flow in \mathfrak{X}

In this section we analyze some properties of the tail space \mathfrak{X} that play a key role in the proof of the upper estimates of Theorem 3.1 and it is the optimal space for *uniform convergence in relative error*.

5.1. An equivalent tail condition

Here we analyze the tail condition (TC') introduced by Vázquez [91]: we will prove that it is equivalent to (TC), but this fact is not trivial. Its proof needs the GHP of Theorem 1.1, as we will see below.

Proposition 5.1. *Let $d \geq 3$, $\gamma, \beta < d$ be real numbers such that $\gamma - 2 < \beta \leq \gamma(d - 2)/d$ and $m \in (m_c, 1)$. Assume $f \in L^1_\gamma(\mathbb{R}^d)$. Then*

f satisfies (TC) if and only if it satisfies (TC').

Proof. We will first prove that (TC) implies (TC'). Assume that $|f|_{\mathfrak{X}} < \infty$ and let $x \in \mathbb{R}^d$, $x \neq 0$. We have the following chain of inequalities:

$$\begin{aligned} \int_{B_{\frac{|x|}{2}}(x)} |f(y)| \frac{dy}{|y|^\gamma} &\leq \int_{B_{\frac{|x|}{2}}(0)} |f(y)| \frac{dy}{|y|^\gamma} \\ &\leq 2^{\frac{2+\beta-\gamma}{1-m} - (d-\gamma)} |f|_{\mathfrak{X}} |x|^{d-\gamma - \frac{2+\beta-\gamma}{1-m}} = O(|x|^{d-\gamma - \frac{2+\beta-\gamma}{1-m}}), \end{aligned}$$

which is exactly (TC'), where we have used $B_{\frac{|x|}{2}}(x) \subset B_{\frac{|x|}{2}}(0)$ in the line above. Assume now that f satisfies (TC'); without loss of generality we can assume that $f \neq 0$. Let $u(t, x)$ be the solution to (CP) with initial data $u(0, x) = |f(x)|$. As proven by Vázquez [91] and also by Vázquez and one of the authors [17], $u(t, x)$ satisfies inequality (1.6), i.e. the GHP. Therefore, by Theorem 1.1 we have $u_0 = |f| \in \mathfrak{X} \setminus \{0\}$, which means that f satisfies (TC). The proof is concluded. \blacksquare

5.2. A nonequivalent tail condition: An example of a “bad” function in \mathcal{X}

In [17] a (nonsharp) sufficient condition for the validity of the GHP takes the following form: there exist $R > 0$ and $A > 0$ such that

$$|f(x)| \leq \frac{A}{|x|^{\frac{2+\beta-\gamma}{1-m}}} \quad \text{for any } |x| \geq R. \tag{5.1}$$

It is easy to check that the above condition is sufficient to guarantee that $f \in \mathcal{X}$, but not necessary, as we will explain by means of the following example: we construct a function $f \in \mathcal{X}$ which does not satisfy (5.1).

Let $d \geq 3$ and, to fix ideas, let us assume that $\gamma = \beta = 0$. Define the function f to be

$$f(y) = \sum_{N=2}^{\infty} \frac{\chi_{B_{N-2}(x_N)}(y)}{N^{\frac{2}{1-m}-1}},$$

where $N \in [2, \infty) \cap \mathbb{N}$, $x_N = (N, \mathbf{0})$, where $\mathbf{0} \in \mathbb{R}^{d-1}$ is the zero vector. The function f is well defined, since by construction $B_{N-2}(x_N) \cap B_{M-2}(x_M) = \emptyset$ unless $N = M$, and $\frac{2}{1-m} - 1 > 1$ since $\frac{d-2}{d} < m < 1$. We then have

$$\begin{aligned} \int_{\mathbb{R}^d} |x|^{\frac{2}{1-m}} f(x) \, dx &= \sum_{N=2}^{\infty} \int_{B_{\frac{1}{N^2}}(x_N)} |x|^{\frac{2}{1-m}} f(x) \, dx \\ &\leq c \sum_{N=2}^{\infty} \frac{N^{\frac{2}{1-m}-2d}}{N^{\frac{2}{1-m}-1}} \leq c \sum_{N=2}^{\infty} \frac{1}{N^{2d-1}} < \infty, \end{aligned}$$

where $0 < c = c(m, d) < \infty$ and we have used the fact that if $x \in B_{\frac{1}{N^2}}(x_N)$ then $|x| \leq 2N$ and we recall that the last series converges since $d \geq 3$. As a consequence, we have that $f \in \mathcal{X}$: indeed for any $R > 1$,

$$R^{\frac{2}{1-m}-d} \int_{B_R^c(0)} f \, dx \leq R^{\frac{2}{1-m}} \int_{B_R^c(0)} f \, dx \leq \int_{B_R^c(0)} |x|^{\frac{2}{1-m}} f \, dx < \infty.$$

On the other hand, a straightforward computation shows that f does not satisfy the pointwise decay condition in (5.1). As expected, the pointwise condition is more restrictive than the integral one.

To conclude, we also give an example of a *radial function* $h \in \mathcal{X}$ which does not satisfy (5.1). Let

$$h(y) = \sum_{N=2}^{\infty} \frac{\chi_{A_N}(y)}{|N - |y||^\eta},$$

where $N \geq 2$ is an integer, $A_N := \{x \in \mathbb{R}^d : N \leq |x| \leq N + N^{-\alpha}\}$ with $0 < \eta < 1$ and $(1 - \eta)\alpha > 2/(1 - m)$.

5.3. The fast diffusion flow as a curve in \mathcal{X}

In this section we will consider solutions to (CP) as continuous curves in \mathcal{X} . To this end, we provide some details about the natural topology of \mathcal{X} . Indeed, the metric associated to the “natural” norm $|\cdot|_{\mathcal{X}}$ would provide \mathcal{X} with a noncomplete topology and this is a bit unpleasant. To see this, just consider $0 < \varepsilon < \frac{\sigma}{1-m} - (d - \gamma)$ and define the function

$$f(x) = |x|^{-\frac{\sigma}{1-m}}(1 - \chi_{B_1(0)}) + |x|^{-(d-\gamma)-\varepsilon} \chi_{B_1(0)}.$$

The above function f does not belong to $L^1_{\gamma}(\mathbb{R}^d)$ but nevertheless $|f|_{\mathcal{X}} < \infty$, hence $f \in \mathcal{X}$. Unfortunately, f can be approximated in the topology induced by $|\cdot|_{\mathcal{X}}$, by the family $\{f_r(x)\}_{r \in (0,1]} \subset L^1_{\gamma}(\mathbb{R}^d)$, where

$$f_r(x) = |x|^{-\frac{\sigma}{1-m}}(1 - \chi_{B_1(0)}) + |x|^{-(d-\gamma)-\varepsilon} \chi_{B_1(0) \setminus B_r(0)}, \quad 0 < r < 1.$$

We have that $f_r \in \mathcal{X}$ for any $0 < r \leq 1$ and a simple (but lengthy) computation shows that $|f_r - f|_{\mathcal{X}} \rightarrow 0$ as $r \rightarrow 0$. Hence, we prefer to introduce the following norm on \mathcal{X} :

$$\|f\|_{\mathcal{X}} := \sup_{R>0} (1 \vee R)^{\frac{2+\beta-\gamma}{1-m}-(d-\gamma)} \int_{B_R^c(0)} |f(x)| |x|^{-\gamma} dx < \infty. \quad (5.2)$$

The main difference between $|\cdot|_{\mathcal{X}}$ and $\|\cdot\|_{\mathcal{X}}$ is that the latter takes into account the influence of the $L^1_{\gamma}(\mathbb{R}^d)$ norm, and provides a complete topology, as in the following:

Proposition 5.2. *Let $d \geq 3$, $\gamma, \beta < d$ be real numbers such that $\gamma - 2 < \beta \leq \gamma(d - 2)/d$ and $m \in (m_c, 1)$. Then*

(i) *for any $f \in \mathcal{X}$ we have*

$$\|f\|_{\mathcal{X}} = \max\{\|f\|_{L^1_{\gamma}(\mathbb{R}^d)}, |f|_{\mathcal{X}}\}; \quad (5.3)$$

(ii) *\mathcal{X} equipped with the norm $\|\cdot\|_{\mathcal{X}}$, defined in (5.2), is a Banach space;*

(iii) *compactly supported functions are not dense in \mathcal{X} equipped with the norm $\|\cdot\|_{\mathcal{X}}$;*

(iv) *the Barenblatt profile has finite \mathcal{X} norm, indeed,*

$$|\mathcal{B}_M|_{\mathcal{X}} = \lim_{R \rightarrow \infty} R^{\frac{\sigma}{1-m}-(d-\gamma)} \int_{B_R^c(0)} \mathcal{B}_M |x|^{-\gamma} dx = (1 - m)\vartheta \omega_d. \quad (5.4)$$

The proof of the above proposition is long but straightforward, hence we refrain from giving it here; however, it can be found in [89, Chapter 4].

As we previously explained, the space $L^1_{\gamma,+}(\mathbb{R}^d)$ can be split into two disjoint sets \mathcal{X} and \mathcal{X}^c . A remarkable fact is that \mathcal{X} and \mathcal{X}^c are two invariant sets of $L^1_{\gamma,+}(\mathbb{R}^d)$ for the fast diffusion flow, in a sense made precise in the following proposition.

Proposition 5.3 (Stability of \mathcal{X} and \mathcal{X}^c). *Let $u(t)$ be a solution to (CP) with $u_0 \in L^1_{\gamma,+}(\mathbb{R}^d)$. Then*

- (i) \mathcal{X} is invariant under the flow, namely $u_0 \in \mathcal{X}$ if and only if $u(t, \cdot) \in \mathcal{X}$ for all $t > 0$;
- (ii) \mathcal{X}^c is invariant under the flow, namely $u_0 \in \mathcal{X}^c$ if and only if $u(t, \cdot) \in \mathcal{X}^c$ for all $t > 0$;
- (iii) if $u_0 \in \mathcal{X}^c$, then

$$\left\| \frac{u(t)}{\mathcal{B}(t; M)} - 1 \right\|_{L^\infty(\mathbb{R}^d)} = \infty \quad \text{for all } t > 0;$$

- (iv) if $u_0 \in \mathcal{X} \setminus \{0\}$, then the following limit holds:

$$\lim_{t \rightarrow \infty} \frac{|u(t)|_{\mathcal{X}}}{t^{\frac{1}{1-m}}} = \lim_{t \rightarrow \infty} \frac{\|u(t)\|_{\mathcal{X}}}{t^{\frac{1}{1-m}}} = (\sigma m)^{\frac{1}{1-m}} \left(\frac{\vartheta}{1-m} \right)^{\frac{m}{1-m}} \omega_d; \quad (5.5)$$

- (v) if $u_0 \in \mathcal{X} \setminus \{0\}$, then for any $t > 0$ the function $t \mapsto t^{\frac{1}{m-1}} |u(t, \cdot)|_{\mathcal{X}}$ is nonincreasing, and

$$\frac{|u(t-h)|_{\mathcal{X}}}{(t-h)^{\frac{1}{1-m}}} \geq \frac{|u(t)|_{\mathcal{X}}}{t^{\frac{1}{1-m}}} \geq (\sigma m)^{\frac{1}{1-m}} \left(\frac{\vartheta}{1-m} \right)^{\frac{m}{1-m}} \omega_d \quad \text{for all } 0 \leq h \leq t. \quad (5.6)$$

Remark 5.4. Inequality (5.6) is sharp, because equality is achieved by the Barenblatt profile $\mathcal{B}(t, \cdot; M)$ for any $M > 0$.

Proof of Proposition 5.3. First recall Lemma 3.4, whose proof is contained in the appendix: for any $R, t, s \geq 0$ we have

$$\left(\int_{B_R^c(0)} u(t, x) \frac{dx}{|x|^\gamma} \right)^{1-m} \leq \left(\int_{B_{2R}^c(0)} u(s, x) \frac{dx}{|x|^\gamma} \right)^{1-m} + C |t-s| R^{(d-\gamma)(1-m)-\sigma}, \quad (5.7)$$

where C is a positive constant which only depends on d, m, γ, β . By taking the supremum in $R > 0$ in inequality (5.7) we can deduce that there exists $c > 0$ depending only on d, m, γ, β , such that for any $t, s \geq 0$,

$$|u(t, \cdot)|_{\mathcal{X}} \leq c(|u(s, \cdot)|_{\mathcal{X}} + |t-s|^{\frac{1}{1-m}}). \quad (5.8)$$

Let us prove first (i). From (5.8) we deduce that if $u_0 \in \mathcal{X}$ then $u(t) \in \mathcal{X}$ for all $t > 0$ just by letting $s = 0$. The opposite choice leads to the converse implication. Part (ii) follows analogously from (5.8). To prove (iii) we proceed by contradiction. Suppose that there exists \bar{t} such that

$$\left\| \frac{u(\bar{t})}{\mathcal{B}(\bar{t}; M)} - 1 \right\|_{L^\infty(\mathbb{R}^d)} \leq C < \infty.$$

Then, reasoning as in the proof of Theorem 1.3 (the necessary part, Section 3.3), we conclude that $u(\bar{t}) \in \mathcal{X}$ and, by point (i), we would have $u_0 \in \mathcal{X}$, a contradiction.

As for the proof of (iv), we will use the GHP of Theorem 1.1: fix $t_0 > 0$, then there exist $\underline{\tau}, \bar{\tau} > 0$ and $\underline{M}, \bar{M} > 0$ such that for any $t \geq t_0 > 0$ we have

$$\mathfrak{B}(t - \underline{\tau}, x; \underline{M}) \leq u(t, x) \leq \mathfrak{B}(t + \bar{\tau}, x; \bar{M}), \quad \text{for any } x \in \mathbb{R}^d.$$

Then a lengthy but still not difficult computation gives the inequality

$$\left(\frac{R_*(t - \underline{\tau})}{\zeta} \right)^{\frac{1}{(1-m)\vartheta}} \leq \frac{|u(t, \cdot)|_{\mathcal{X}}}{\omega_d(1-m)\vartheta} \leq \left(\frac{R_*(t + \bar{\tau})}{\zeta} \right)^{\frac{1}{(1-m)\vartheta}}. \quad (5.9)$$

Recall that R_* is defined in (1.3), so that inequality (5.5) follows from (5.9), together with

$$\lim_{t \rightarrow \infty} R_*(t + T)^{\frac{1}{(1-m)\vartheta}} t^{-\frac{1}{1-m}} = \vartheta^{-\frac{1}{1-m}} \quad \text{for any } T \in \mathbb{R}.$$

We prove (v) through a smooth approximation by means of auxiliary norms: Let $k > 0$ be a positive integer and $\phi_k(x)$ be such that

$$\begin{aligned} \phi_k(x) &= 1 \quad \text{on } |x| \geq 1 + \frac{1}{k}, \quad \phi_k(x) = 0 \quad \text{on } |x| \leq 1, \\ \phi_k(x) &> 0 \quad \text{on } 1 < |x| < 1 + \frac{1}{k}. \end{aligned}$$

Let us define

$$\|f\|_{k, \mathcal{X}} = \sup_{R > 0} R^{\frac{\sigma}{1-m} - (d-\gamma)} \int_{\mathbb{R}^d} f(x) \phi_k\left(\frac{x}{R}\right) |x|^{-\gamma} dx. \quad (5.10)$$

For any $k \geq 1$ and for any $f \in \mathcal{X}$ we have

$$\left(\frac{k}{k+1} \right)^{\frac{\sigma}{1-m} - (d-\gamma)} |f|_{\mathcal{X}} \leq \|f\|_{k, \mathcal{X}} \leq |f|_{\mathcal{X}}.$$

As a consequence of the above inequality, for any $f \in \mathcal{X}$ the following limit holds:

$$\lim_{k \rightarrow \infty} \|f\|_{k, \mathcal{X}} = |f|_{\mathcal{X}}. \quad (5.11)$$

We take advantage of the auxiliary norms (5.10). Let $k > 0$ be a positive integer and $R > 0$ and define $Y_k(t) = \int_{\mathbb{R}^d} \phi_k\left(\frac{x}{R}\right) u(t, x) \frac{dx}{|x|^\gamma}$. Now using time monotonicity, the so-called Bénilan–Crandall estimates [4], $u_t \leq \frac{u}{(1-m)t}$ valid in the distributional sense, we find that

$$Y_k'(t) \leq \frac{1}{(1-m)t} Y_k(t), \quad \text{so that for all } \tau > s > 0, \quad \frac{Y_k(s)}{s^{\frac{1}{1-m}}} \geq \frac{Y_k(\tau)}{\tau^{\frac{1}{1-m}}}.$$

Multiplying by $R^{\sigma/(1-m) - (d-\gamma)}$ and taking the supremum in $R > 0$ in the above inequality, we get

$$\frac{\|u(s, \cdot)\|_{k, \mathcal{X}}}{s^{\frac{1}{1-m}}} \geq \frac{\|u(t, \cdot)\|_{k, \mathcal{X}}}{t^{\frac{1}{1-m}}}.$$

Then taking the limit as $k \rightarrow \infty$ in the above inequality one gets the monotonicity of $t^{-1/(1-m)}|u(t, \cdot)|_{\mathcal{X}}$. Let $\tau > s > 0$ as before (we can take any $\tau > s$, hence also $\tau \rightarrow \infty$) to obtain inequality (5.6), namely

$$\frac{|u(s, \cdot)|}{s^{\frac{1}{1-m}}} \geq \frac{|u(\tau, \cdot)|}{\tau^{\frac{1}{1-m}}} \geq \lim_{\tau \rightarrow \infty} \frac{|u(\tau, \cdot)|}{\tau^{\frac{1}{1-m}}} = (\sigma m)^{\frac{1}{1-m}} \left(\frac{\vartheta}{1-m} \right)^{\frac{m}{1-m}} \omega_d,$$

where in the last step we have used (5.5). The proof is then concluded. \blacksquare

5.4. Convergence to the Barenblatt in \mathcal{X}

Finally, we address here the question of convergence to the Barenblatt profile of solutions to (CP) in \mathcal{X} with the topology induced by $\|\cdot\|_{\mathcal{X}}$. We find that it is false in general that

$$\|u(t, \cdot) - \mathfrak{B}(t, \cdot; M)\|_{\mathcal{X}} \rightarrow 0 \quad \text{as } t \rightarrow \infty.$$

To see this fact we provide an explicit counterexample. Consider the Barenblatt solution $\mathfrak{B}(t, x; M)$ and its translation in time $\mathfrak{B}(t + \tau, x; M)$. For $R > 0$ large enough we have

$$|\mathfrak{B}(t + \tau, x; M) - \mathfrak{B}(t, x; M)| \geq \frac{1}{2} \left[\left(1 + \frac{\tau}{t} \right)^{\frac{1}{1-m}} - 1 \right] \mathfrak{B}(t, x; M) \quad \text{for any } |x| \geq R,$$

and we therefore conclude, thanks to (5.4) and (5.6), that

$$|\mathfrak{B}(t + \tau, x; M) - \mathfrak{B}(t, x; M)|_{\mathcal{X}} \gtrsim t^{\frac{m}{1-m}}.$$

However, if we suitably renormalize $\|\cdot\|_{\mathcal{X}}$ by the factor $t^{\frac{1}{1-m}}$ we find the following result.

Proposition 5.5. *Under the assumption of Theorem 1.3 we have*

$$\lim_{t \rightarrow \infty} \frac{\|u(t, \cdot) - \mathfrak{B}(t, \cdot; M)\|_{\mathcal{X}}}{t^{\frac{1}{1-m}}} = 0. \quad (5.12)$$

Remark 5.6. It is interesting to stress that the above limit, rewritten in *self-similar* variables (1.11), does not need the renormalization factor $t^{\frac{1}{1-m}}$, namely

$$\lim_{\tau \rightarrow \infty} \|v(\tau, \cdot) - \mathcal{B}_M(\cdot)\|_{\mathcal{X}} = 0.$$

Proof of Proposition 5.5. By Proposition 5.2 we have $\|\cdot\|_{\mathcal{X}} = \max\{\|\cdot\|_{\mathcal{X}}, \|\cdot\|_{L^1_{\gamma}(\mathbb{R}^d)}\}$. To prove (5.5) we need to show

$$\lim_{t \rightarrow \infty} \frac{\|u(t, \cdot) - \mathfrak{B}(t, \cdot; M)\|_{\mathcal{X}}}{t^{\frac{1}{1-m}}} = 0 \quad \text{and} \quad \lim_{t \rightarrow \infty} \frac{\|u(t, \cdot) - \mathfrak{B}(t, \cdot; M)\|_{L^1_{\gamma}(\mathbb{R}^d)}}{t^{\frac{1}{1-m}}} = 0.$$

By conservation of mass the second limit is straightforward, hence we only need to prove the first. Under the running assumption we know by Theorem 1.3 that $u(t, x)$ converges to $\mathfrak{B}(t, x; M)$ uniformly in relative error. We restate this result in the following way: there

exists a positive function $g(t) \rightarrow 0$ as $t \rightarrow \infty$ such that for any $x \in \mathbb{R}^d$ and for any t large enough we have

$$|u(t, x) - \mathfrak{B}(t, x; M)| \leq g(t)\mathfrak{B}(t, x; M).$$

By the above inequality we deduce that

$$\limsup_{t \rightarrow \infty} \frac{|u(t, \cdot) - \mathfrak{B}(t, \cdot; M)|_{\mathcal{X}}}{t^{\frac{1}{1-m}}} \leq \lim_{t \rightarrow \infty} g(t) \frac{|\mathfrak{B}(t, \cdot; M)|_{\mathcal{X}}}{t^{\frac{1}{1-m}}} = 0,$$

where we have used identity (5.5) and the fact that $g(t) \rightarrow 0$ as $t \rightarrow \infty$. The proof is therefore concluded. \blacksquare

A. Regularity results and some technical details

A.1. How to recover the mass of the Barenblatt profile \mathfrak{B}_M

The following identity is a consequence of the integral representation formula of the Euler beta function; see [2, 6.2.1, p. 258]:

$$\bar{M} = \int_{\mathbb{R}^d} (1 + |x|^{2+\beta-\gamma})^{\frac{1}{m-1}} |x|^{-\gamma} dx = |\mathbb{S}^{d-1}| \frac{\Gamma(\frac{d-\gamma}{2+\beta-\gamma})\Gamma(\frac{1}{1-m} - \frac{d-\gamma}{2+\beta-\gamma})}{(2+\beta-\gamma)\Gamma(\frac{1}{1-m})}.$$

By scaling we obtain

$$\begin{aligned} M &= \int_{\mathbb{R}^d} (C(M) + |x|^{2+\beta-\gamma})^{\frac{1}{m-1}} |x|^{-\gamma} dx \\ &= C(M)^{\frac{1}{m-1} - \frac{\gamma}{\sigma}} \int_{\mathbb{R}^d} (1 + |x/C(M)^{1/\sigma}|^{\sigma})^{\frac{1}{m-1}} |x/C(M)^{1/\sigma}|^{-\gamma} dx \\ &= C(M)^{\frac{1}{m-1} + \frac{d-\gamma}{\sigma}} \int_{\mathbb{R}^d} (1 + |x|^{2+\beta-\gamma})^{\frac{1}{m-1}} |x|^{-\gamma} dx \\ &= C(M)^{\frac{1}{m-1} + \frac{d-\gamma}{\sigma}} \bar{M}, \end{aligned}$$

and therefore we have $C(M) = (\frac{\bar{M}}{M})^{\frac{\sigma(1-m)}{\sigma-(d-\gamma)(1-m)}}$.

A.2. Interpolation inequality

Let $\Omega \subset \mathbb{R}^d$ be a bounded domain and $u: \Omega \rightarrow \mathbb{R}$ be a function and define for any $\nu \in (0, 1)$,

$$[u]_{C^\nu(\Omega)} := \sup_{\substack{x, y \in \Omega \\ x \neq y}} \frac{|u(x) - u(y)|}{|x - y|^\nu}. \quad (\text{A.1})$$

We say that $u \in C^\nu(\Omega)$ whenever $[u]_{C^\nu(\Omega)} < \infty$. Notice that $[u]_{C^\nu(\Omega)} = 0$ if and only if u is constant, and since in what follows we need to use strictly positive quantities, we will use the following inequality which holds for $u \in C^\nu(\Omega)$:

$$|u(x) - u(y)| \leq (1 + [u]_{C^\nu(\mathbb{R}^d)})|x - y|^\nu. \quad (\text{A.2})$$

Let $\Omega' \subset \Omega$ be a subdomain; we define the distance between Ω and Ω' as

$$\text{dist}(\Omega, \Omega') = \inf_{\substack{x \in \partial\Omega \\ y \in \partial\Omega'}} |x - y|,$$

where $\partial\Omega$ is the boundary of Ω and $\partial\Omega'$ is the boundary of Ω' . The purpose of this appendix is to prove the following lemma.

Lemma A.1. *Let $p \geq 1$, $v \in (0, 1)$ and $u: \Omega \rightarrow \mathbb{R}$ be a function such that $u \in L^p_\gamma(\Omega) \cap C^v(\Omega)$. Assume $\gamma \leq 0$ and let $\Omega' \subset \Omega$ be such that $\text{dist}(\Omega, \Omega') > 0$. Then there exists a positive constant $C_{d,\gamma,v,p}$, which depends on d, γ, p and v , such that*

$$\begin{aligned} \|u\|_{L^\infty(\Omega')} &\leq C_{d,\gamma,v,p} \left(1 + \frac{\|u\|_{L^p_\gamma(\Omega)}}{(1 + [u]_{C^v(\Omega)}) \text{dist}(\Omega, \Omega')^{\frac{1}{p}}} \right)^{\frac{d-\gamma}{d-\gamma+pv}} \\ &\quad \times (1 + [u]_{C^v(\Omega)})^{\frac{d-\gamma}{d-\gamma+pv}} \|u\|_{L^p_\gamma(\Omega)}^{\frac{pv}{d-\gamma+pv}}. \end{aligned} \quad (\text{A.3})$$

Assume $0 < \gamma < d$, and in addition let $\Omega' \subset \Omega$ be two bounded domains. Then there exists a positive constant $C_{d,\gamma,v,p}$, which depends on d, γ, p and v , such that

$$\begin{aligned} \|u\|_{L^\infty(\Omega')} &\leq C_{d,\gamma,v,p} \left(\text{dist}(\Omega, \Omega') + \sup_{x \in \Omega'} |x| \right)^{\frac{\gamma}{p}} \\ &\quad \times \left(1 + \frac{\|u\|_{L^p_\gamma(\Omega)}}{(1 + [u]_{C^v(\Omega)}) \text{dist}(\Omega, \Omega')^{\frac{1}{p}}} \right)^{\frac{d}{d+pv}} \\ &\quad \times (1 + [u]_{C^v(\Omega)})^{\frac{d}{d+pv}} \|u\|_{L^p_\gamma(\Omega)}^{\frac{pv}{d+pv}}. \end{aligned} \quad (\text{A.4})$$

Proof. For any $x, y \in \Omega'$ we have, by the triangle inequality,

$$|u(x)|^p \leq (|u(x) - u(y)| + |u(y)|)^p \leq 2^p (|u(x) - u(y)|^p + |u(y)|^p).$$

Let $x \in \Omega'$, $0 \leq R < \text{dist}(\Omega, \Omega')$; averaging on a ball $B_R(x) \subset \Omega$ we have

$$\begin{aligned} |u(x)|^p &\leq \frac{2^p}{\mu_\gamma(B_R(x))} \int_{B_R(x)} |u(x) - u(y)|^p |y|^{-\gamma} dy \\ &\quad + \frac{2^p}{\mu_\gamma(B_R(x))} \int_{B_R(x)} |u(y)|^p |y|^{-\gamma} dy \\ &\leq 2^p R^{pv} (1 + [u]_{C^v(\Omega)})^p + 2^p \frac{\|u\|_{L^p_\gamma(\Omega)}^p}{\mu_\gamma(B_R(x))}, \end{aligned} \quad (\text{A.5})$$

where in the last step we have used (A.2) and that $\int_{B_R(x)} |u(y)|^p |y|^{-\gamma} dy \leq \|u\|_{L^p_\gamma(\Omega)}^p$. We claim that, for any $x_0 \in \mathbb{R}^d$ and for any $R \geq 0$, there exist positive constants $c_{\gamma,d}$, $C_{\gamma,d}$ such that

$$c_{\gamma,d} R^d \left(|x_0| \vee \frac{R}{2} \right)^{-\gamma} \leq \mu_\gamma(B_R(x_0)) \leq C_{\gamma,d} R^d \left(|x_0| \vee \frac{R}{2} \right)^{-\gamma}. \quad (\text{A.6})$$

The above inequality can be proven using the techniques developed in [14, Lemma 5.2, Appendix B], but we will not include the proof here. Now we consider two cases, namely $\gamma \leq 0$ and $\gamma > 0$.

Assume that $\gamma \leq 0$. Plugging the lower bound of (A.6) into (A.5) we deduce

$$|u(x)|^p \leq C(R^{p\nu}(1 + [u]_{C^\nu(\Omega)})^p + R^{\gamma-d} \|u\|_{L^\gamma(\Omega)}^p). \quad (\text{A.7})$$

Then inequality (A.3) follows by letting

$$2R = \left(\frac{(d - \gamma) \|u\|_{L^\gamma(\Omega)}^p}{p\nu(1 + [u]_{C^\nu(\Omega)})^p} \right)^{\frac{1}{d-\gamma+p\nu}} \wedge \text{dist}(\Omega, \Omega').$$

Assume that $0 < \gamma < d$. Then using (A.6) in (A.5) we have, for any $0 \leq R < \text{dist}(\Omega, \Omega') := D$,

$$\begin{aligned} |u(x)|^p &\leq C \left(R^{p\nu}(1 + [u]_{C^\nu(\Omega)})^p + \frac{(|x_0| \vee \frac{R}{2})^\gamma}{R^d} \|u\|_{L^\gamma(\Omega)}^p \right) \\ &\leq C(D + |x_0|)^\gamma (R^{p\nu}(1 + [u]_{C^\nu(\Omega)})^p + R^{-d} \|u\|_{L^\gamma(\Omega)}^p), \end{aligned} \quad (\text{A.8})$$

where in the last step we have used that $(|x_0| \vee \frac{R}{2})^\gamma \leq (D + |x_0|)^\gamma$. Letting

$$2R = \left(\frac{d \|u\|_{L^\gamma(\Omega)}^p}{p\nu(1 + [u]_{C^\nu(\Omega)})^p} \right)^{\frac{1}{d+p\nu}} \wedge \text{dist}(\Omega, \Omega'),$$

taking the supremum in $x_0 \in \Omega'$ we get (A.4). The proof is then complete. \blacksquare

A.3. Hölder continuity of solutions to weighted equations

We present here some regularity results for nonnegative local weak solutions to both linear and nonlinear parabolic equations with weights. The results contained in this section are mainly contained in [14, Part III] and references therein. We provide similar results here, adapted to the present setting, and assumptions, for convenience of the reader. Consider local weak solutions in the cylinder $Q := (0, T) \times \Omega$ to the equation

$$v_t = w_\gamma \sum_{i,j=1}^N \partial_i (A_{i,j}(t, x) \partial_j v), \quad (\text{A.9})$$

where $A_{i,j} = A_{j,i}$ and there exist constants $0 < \lambda_0 \leq \lambda_1 < +\infty$ such that for some $\gamma, \beta < N$, satisfying $\gamma - 2 < \beta \leq (\frac{N-2}{N})\gamma$, we have for any $\xi \in \mathbb{R}^d$ and any $x \in \mathbb{R}^d$,

$$w_\gamma \asymp |x|^\gamma \quad \text{and} \quad 0 < \lambda_0 |x|^{-\beta} |\xi|^2 \leq \sum_{i,j=1}^N A_{i,j}(t, x) \xi_i \xi_j \leq \lambda_1 |x|^{-\beta} |\xi|^2. \quad (\text{A.10})$$

We will restrict ourselves to the class of bounded, nonnegative, local weak solutions to equation (A.9), precisely defined in [14, 26, 55]. Notice that this class of solutions is large enough for our purposes.

It is convenient to introduce the notion of distance between nested cylinders of the form $Q = (0, T) \times \Omega$. Let $Q' = (T_1, T_2) \times \Omega' \subset Q$. We define

$$d_{\gamma, \beta}(Q, Q') := \inf_{\substack{(t, x) \in \{[0, T] \times \partial\Omega\} \cup \{0\} \times \Omega \\ (s, y) \in Q'}} |x - y| \vee (\rho_y^{\gamma, \beta})^{-1}(|t - s|), \quad (\text{A.11})$$

where γ, β are as above and $(\rho_y^{\gamma, \beta})^{-1}$ is the inverse of $\rho_y^{\gamma, \beta}$ (well defined for any $y \in \mathbb{R}^d$) defined as

$$\rho_y^{\gamma, \beta}(R) := \left(\int_{B_R(y)} |x|^{(\beta - \gamma)\frac{N}{2}} dx \right)^{\frac{2}{N}}.$$

Finally, we introduce the notion of a C^α norm which takes into account the presence of the weights. With the above notation we define

$$[u]_{C_{\gamma, \beta}^\alpha(Q)} := \sup_{\substack{(t, x), (\tau, y) \in Q' \\ (t, x) \neq (\tau, y)}} \frac{|v(t, x) - v(\tau, y)|}{(|x - y| + |t - \tau|^{\frac{1}{2\gamma\sigma}})^\alpha}. \quad (\text{A.12})$$

The proof of the following result can be found in [14, Proposition 4.2, Corollary 4.3].

Notice that the following results involve both the “parabolic” Hölder norm $[u]_{C_{\gamma, \beta}^\alpha(Q)}$ defined in (A.13), and the “elliptic” one, $[u]_{C^v(\Omega)}$, defined in (A.1).

Proposition A.2 (Hölder continuity for linear equations with weights). *Let v be a non-negative bounded local weak solution to equation (A.9) on $Q := (0, T) \times \Omega$, under the assumption (A.10). Let $Q' := (T_1, T_2) \times \Omega' \subset Q$. Then there exist $\alpha \in (0, 1)$ and $\bar{\kappa}_\alpha > 0$, such that for all $(t, x), (s, y) \in Q'$,*

$$[u]_{C_{\gamma, \beta}^\alpha(Q')} \leq \frac{\bar{\kappa}_\alpha}{d_{\gamma, \beta}(Q, Q')^\alpha} \|v\|_{L^\infty(Q)}, \quad (\text{A.13})$$

where $\bar{\kappa}_\alpha > 0$ is given by

$$\bar{\kappa}_\alpha = \bar{\kappa}'_\alpha \begin{cases} 1 & \text{if } \sigma \geq 2, \\ (T^{\frac{1}{\sigma}} \vee \sup_{x_0 \in \Omega} |x_0|)^{\frac{\gamma - \beta}{2}} & \text{if } 0 < \sigma < 2. \end{cases} \quad (\text{A.14})$$

The constants $\alpha, \bar{\kappa}'_\alpha$ depend only on $N, \gamma, \beta, \lambda_0, \lambda_1$.

Proposition A.13 can be fruitfully used to deduce regularity results for the nonlinear parabolic equation also: for example, we can consider nonnegative bounded solutions to $u_t = |x|^\gamma \nabla(|x|^{-\beta} \nabla u^m)$ as solutions to the linear equation $u_t = |x|^\gamma \nabla(|x|^{-\beta} a(t, x) \nabla u)$, where $a(t, x) = mu(t, x)^{m-1}$. Indeed, the same can be done for solutions to the Fokker–Planck-type equation (NLWFP) as follows.

Lemma A.3. *Let $\rho, \tau_0 > 0$, $0 < \lambda_0 \leq \lambda_1 < \infty$, $m \in (0, 1)$ and let $v(\tau, y) : (0, \infty) \times \mathbb{R}^d \rightarrow \mathbb{R}$ be a nonnegative bounded solution to (NLWFP). Assume that*

$$\lambda_0 \leq mv^{m-1}(\tau, y) \leq \lambda_1 \quad \text{for any } \tau \geq \tau_0 \text{ and } |y| \leq \rho.$$

Then there exist $\nu > 0$ and $\bar{\kappa} > 0$ such that if $\tau_1 > \tau_0$ and $\tau \in [\tau_1 + \frac{1}{\sigma} \log R_\star(2), \tau_1 + \frac{1}{\sigma} \log R_\star(3)]$ then

$$[v(\tau, \cdot)]_{C^\nu(B_{\rho/2}(0))} \leq \bar{\kappa} \|v\|_{L^\infty([\tau_1, \tau_1 + \frac{1}{\sigma} \log R_\star(4)] \times B_\rho(0))}. \quad (\text{A.15})$$

The constants $\nu, \bar{\kappa}$ depend on $m, d, \gamma, \beta, \lambda_0, \lambda_1$; $\bar{\kappa}$ also depends on ρ .

Proof. The proof is divided into several steps. It is convenient to consider a time-shifted solution:

$$\bar{v}(\tau, y) := v(\tau + \tau_1) \quad \text{for any } \tau \geq 0.$$

Rescaling to original variables. The rescaled function $\bar{u}(t, x)$ defined by

$$\begin{aligned} \bar{u}(t, x) &:= \frac{\zeta^{d-\gamma}}{R_\star(t+1)^{d-\gamma}} \bar{v}\left(\frac{1}{\sigma} \log \frac{R_\star(t+1)}{R_\star(1)}, \frac{\zeta x}{R_\star(t)}\right) \\ &= \frac{\zeta^{d-\gamma}}{R_\star(t+1)^{d-\gamma}} \bar{v}(\tau, y), \end{aligned} \quad (\text{A.16})$$

satisfies (CP). Define the following domains:

$$\begin{aligned} \bar{Q}_1 &:= \{(t, x) : 0 \leq t \leq 3, |x| \leq \frac{\rho R_\star(t+1)}{\zeta}\}, \\ \bar{Q}_2 &:= \{(t, x) : 1 \leq t \leq 2, |x| \leq \frac{\rho R_\star(t+1)}{2\zeta}\}. \end{aligned}$$

On both \bar{Q}_1 and \bar{Q}_2 the following estimate holds true:

$$\frac{R_\star(1)^{(d-\gamma)(1-m)}}{\zeta^{(d-\gamma)(1-m)}} \lambda_0 \leq m \bar{u}^{m-1}(t, x) \leq \lambda_1 \frac{R_\star(4)^{(d-\gamma)(1-m)}}{\zeta^{(d-\gamma)(1-m)}}.$$

Application of the linear result. We can consider \bar{u} as a bounded solution to the linear equation

$$\bar{u}_t = |x|^\gamma \nabla (|x|^{-\beta} a(t, x) \nabla \bar{u}) \quad \text{where } a(t, x) = m \bar{u}(t, x),$$

on the domain \bar{Q}_1 . From Proposition A.2 we deduce that there exist $\nu > 0$ and $\bar{\kappa}_\nu > 0$ such that

$$\|\bar{u}\|_{C_{\gamma, \beta}^\nu(\bar{Q}_2)} \leq \bar{\kappa}_\nu \frac{\|\bar{u}\|_{L^\infty(\bar{Q}_1)}}{d_{\gamma, \beta}(\bar{Q}_1, \bar{Q}_2)^\nu}. \quad (\text{A.17})$$

The constant ν will depend only on d, m, γ, β and λ_0, λ_1 , since $R_\star(1), R_\star(4), \zeta$ are numerical constants which only depend on d, m, γ, β . However, the constant $\bar{\kappa}_\nu$ will depend on ρ when $0 < \sigma < 2$; see the expression for the constant $\bar{\kappa}_\alpha$ in Proposition A.2. Finally, we notice that $d_{\gamma, \beta}(\bar{Q}_1, \bar{Q}_2)^\nu$ depends on ρ as well. We will now freeze the time variable and consider $\bar{u}(t, x)$ as a function in space only. From (A.17) we deduce that for any $t \in [1, 2]$ we have

$$[u(t, \cdot)]_{C^\nu(B_{\bar{R}(t)}(0))} \leq \bar{\kappa}_\nu \frac{\|\bar{u}\|_{L^\infty(\bar{Q}_1)}}{d_{\gamma, \beta}(\bar{Q}_1, \bar{Q}_2)^\nu}, \quad \text{where } \bar{R}(t) = \frac{\rho R_\star(t+1)}{2\zeta}. \quad (\text{A.18})$$

Rescaling back to self-similar variables. The domains \bar{Q}_1 and \bar{Q}_2 will be back to $[\tau_1, \tau_1 + \frac{1}{\sigma} \log R_\star(4)] \times B_\rho(0)$ and $[\tau_1 + \frac{1}{\sigma} \log R_\star(2), \tau_1 + \frac{1}{\sigma} \log R_\star(3)] \times B_{\rho/2}(0)$ respectively, while (A.18) become (A.15) where

$$\bar{\kappa} = \left(\frac{R_\star(4)}{R_\star(2)} \right)^{d-\gamma} \left(\frac{\zeta}{R_\star(1)} \right)^\nu \frac{\bar{\kappa}_\nu}{d_{\gamma,\beta}(\bar{Q}_1, \bar{Q}_2)^\nu}.$$

The proof is then concluded. ■

Funding. This work was partially funded by Projects MTM2017-85757-P and PID2020-113596GB-I00 (Spain) and by the EU H2020 MSCA programme, grant agreement 777822. M.B. acknowledges financial support from the Spanish Ministry of Science and Innovation, through the “Severo Ochoa Programme for Centres of Excellence in R&D” (CEX2019-000904-S). N.S. was partially funded by the FPI-grant BES-2015-072962, associated to the project MTM2014-52240-P (Spain) and also by Regione Ile-de-France. This work has been partially supported by the Project EFI ANR-17-CE40-0030 of the French National Research Agency.

References

- [1] B. Abdellaoui and I. P. Alonso, Hölder regularity and Harnack inequality for degenerate parabolic equations related to Caffarelli-Kohn-Nirenberg inequalities. *Nonlinear Anal.* **57** (2004), no. 7-8, 971–1003 Zbl [1053.35070](#) MR [2070620](#)
- [2] M. Abramowitz and I. A. Stegun, *Handbook of mathematical functions with formulas, graphs, and mathematical tables*. National Bureau of Standards Appl. Math. Ser. 55, US Government Printing Office, Washington, DC, 1964 Zbl [0171.38503](#) MR [0167642](#)
- [3] D. Bakry and M. Émery, Diffusions hypercontractives. In *Séminaire de probabilités, XIX, 1983/84*, pp. 177–206, Lecture Notes in Math. 1123, Springer, Berlin, 1985 Zbl [0561.60080](#) MR [889476](#)
- [4] P. Bénilan and M. G. Crandall, Regularizing effects of homogeneous evolution equations. In *Contributions to analysis and geometry (Baltimore, MD, 1980)*, pp. 23–39, Johns Hopkins University Press, Baltimore, MD, 1981 Zbl [0556.35067](#) MR [648452](#)
- [5] A. Blanchet, M. Bonforte, J. Dolbeault, G. Grillo, and J.-L. Vázquez, Hardy-Poincaré inequalities and applications to nonlinear diffusions. *C. R. Math. Acad. Sci. Paris* **344** (2007), no. 7, 431–436 Zbl [1190.35119](#) MR [2320246](#)
- [6] A. Blanchet, M. Bonforte, J. Dolbeault, G. Grillo, and J. L. Vázquez, Asymptotics of the fast diffusion equation via entropy estimates. *Arch. Ration. Mech. Anal.* **191** (2009), no. 2, 347–385 Zbl [1178.35214](#) MR [2481073](#)
- [7] S. Bonafede and I. I. Skrypnik, On Hölder continuity of solutions of doubly nonlinear parabolic equations with weight. *Ukrain. Mat. Zh.* **51** (1999), no. 7, 890–903 Zbl [0937.35021](#) MR [1727693](#)
- [8] M. Bonforte, J. Dolbeault, G. Grillo, and J. L. Vázquez, Sharp rates of decay of solutions to the nonlinear fast diffusion equation via functional inequalities. *Proc. Natl. Acad. Sci. USA* **107** (2010), no. 38, 16459–16464 Zbl [1256.35026](#) MR [2726546](#)

- [9] M. Bonforte, J. Dolbeault, M. Muratori, and B. Nazaret, Weighted fast diffusion equations (Part I): Sharp asymptotic rates without symmetry and symmetry breaking in Caffarelli-Kohn-Nirenberg inequalities. *Kinet. Relat. Models* **10** (2017), no. 1, 33–59 Zbl [1361.35075](#) MR [3579563](#)
- [10] M. Bonforte, J. Dolbeault, M. Muratori, and B. Nazaret, Weighted fast diffusion equations (Part II): Sharp asymptotic rates of convergence in relative error by entropy methods. *Kinet. Relat. Models* **10** (2017), no. 1, 61–91 Zbl [1361.35076](#) MR [3579564](#)
- [11] M. Bonforte, J. Dolbeault, B. Nazaret, and N. Simonov, Stability in Gagliardo-Nirenberg-Sobolev inequalities: Flows, regularity and the entropy method. 2022, arXiv:[2007.03674](#). To appear in *Mem. Amer. Math. Soc.*
- [12] M. Bonforte, G. Grillo, and J. L. Vázquez, Fast diffusion flow on manifolds of nonpositive curvature. *J. Evol. Equ.* **8** (2008), no. 1, 99–128 Zbl [1139.35065](#) MR [2383484](#)
- [13] M. Bonforte, G. Grillo, and J. L. Vázquez, Special fast diffusion with slow asymptotics: entropy method and flow on a Riemann manifold. *Arch. Ration. Mech. Anal.* **196** (2010), no. 2, 631–680 Zbl [1209.35069](#) MR [2609957](#)
- [14] M. Bonforte and N. Simonov, Quantitative a priori estimates for fast diffusion equations with Caffarelli-Kohn-Nirenberg weights. Harnack inequalities and Hölder continuity. *Adv. Math.* **345** (2019), 1075–1161 Zbl [1408.35073](#) MR [3903914](#)
- [15] M. Bonforte, N. Simonov, and D. Stan, The Cauchy problem for the fast p -Laplacian evolution equation. Characterization of the global Harnack principle and fine asymptotic behaviour. *J. Math. Pures Appl. (9)* **163** (2022), 83–131 Zbl [07541864](#) MR [4438897](#)
- [16] M. Bonforte, Y. Sire, and J. L. Vázquez, Optimal existence and uniqueness theory for the fractional heat equation. *Nonlinear Anal.* **153** (2017), 142–168 Zbl [1364.35416](#) MR [3614666](#)
- [17] M. Bonforte and J. L. Vázquez, Global positivity estimates and Harnack inequalities for the fast diffusion equation. *J. Funct. Anal.* **240** (2006), no. 2, 399–428 Zbl [1107.35063](#) MR [2261689](#)
- [18] L. Caffarelli, R. Kohn, and L. Nirenberg, First order interpolation inequalities with weights. *Compositio Math.* **53** (1984), no. 3, 259–275 Zbl [0563.46024](#) MR [768824](#)
- [19] J. A. Carrillo, M. Di Francesco, and G. Toscani, Intermediate asymptotics beyond homogeneity and self-similarity: long time behavior for $u_t = \Delta\phi(u)$. *Arch. Ration. Mech. Anal.* **180** (2006), no. 1, 127–149 Zbl [1096.35015](#) MR [2211709](#)
- [20] J. A. Carrillo, D. Gómez-Castro, and J. L. Vázquez, A fast regularisation of a Newtonian vortex equation. *Ann. Inst. H. Poincaré C Anal. Non Linéaire* **39** (2022), no. 3, 705–747 MR [4412079](#)
- [21] J. A. Carrillo, A. Jüngel, P. A. Markowich, G. Toscani, and A. Unterreiter, Entropy dissipation methods for degenerate parabolic problems and generalized Sobolev inequalities. *Monatsh. Math.* **133** (2001), no. 1, 1–82 Zbl [0984.35027](#) MR [1853037](#)
- [22] J. A. Carrillo, C. Lederman, P. A. Markowich, and G. Toscani, Poincaré inequalities for linearizations of very fast diffusion equations. *Nonlinearity* **15** (2002), no. 3, 565–580 Zbl [1011.35025](#) MR [1901093](#)
- [23] J. A. Carrillo and G. Toscani, Asymptotic L^1 -decay of solutions of the porous medium equation to self-similarity. *Indiana Univ. Math. J.* **49** (2000), no. 1, 113–142 Zbl [0963.35098](#) MR [1777035](#)
- [24] J. A. Carrillo and J. L. Vázquez, Fine asymptotics for fast diffusion equations. *Comm. Partial Differential Equations* **28** (2003), no. 5-6, 1023–1056 Zbl [1036.35100](#) MR [1986060](#)

- [25] F. Catrina and Z.-Q. Wang, On the Caffarelli-Kohn-Nirenberg inequalities: sharp constants, existence (and nonexistence), and symmetry of extremal functions. *Comm. Pure Appl. Math.* **54** (2001), no. 2, 229–258 Zbl [1072.35506](#) MR [1794994](#)
- [26] F. M. Chiarenza and R. P. Serapioni, A Harnack inequality for degenerate parabolic equations. *Comm. Partial Differential Equations* **9** (1984), no. 8, 719–749 Zbl [0546.35035](#) MR [748366](#)
- [27] I. Csiszár, Information-type measures of difference of probability and indirect observations. *Studia Sci. Math. Hungar.* **2** (1967), 299–318 Zbl [0157.25802](#) MR [219345](#)
- [28] A. Dall’Aglio, D. Giachetti, and I. Peral, Results on parabolic equations related to some Caffarelli-Kohn-Nirenberg inequalities. *SIAM J. Math. Anal.* **36** (2004/05), no. 3, 691–716 Zbl [1078.35056](#) MR [2111912](#)
- [29] P. Daskalopoulos and M. Del Pino, On nonlinear parabolic equations of very fast diffusion. *Arch. Rational Mech. Anal.* **137** (1997), no. 4, 363–380 Zbl [0886.35081](#) MR [1463800](#)
- [30] P. Daskalopoulos, M. del Pino, and N. Sesum, Type II ancient compact solutions to the Yamabe flow. *J. Reine Angew. Math.* **738** (2018), 1–71 Zbl [1472.53102](#) MR [3794888](#)
- [31] P. Daskalopoulos and M. A. del Pino, On fast diffusion nonlinear heat equations and a related singular elliptic problem. *Indiana Univ. Math. J.* **43** (1994), no. 2, 703–728 Zbl [0806.35086](#) MR [1291536](#)
- [32] P. Daskalopoulos and M. A. del Pino, On a singular diffusion equation. *Comm. Anal. Geom.* **3** (1995), no. 3-4, 523–542 Zbl [0851.35072](#) MR [1371208](#)
- [33] P. Daskalopoulos, J. King, and N. Sesum, Extinction profile of complete non-compact solutions to the Yamabe flow. *Comm. Anal. Geom.* **27** (2019), no. 8, 1757–1798 Zbl [1462.53093](#)
- [34] P. Daskalopoulos and N. Sesum, On the extinction profile of solutions to fast diffusion. *J. Reine Angew. Math.* **622** (2008), 95–119 Zbl [1173.35074](#) MR [2433613](#)
- [35] M. Del Pino and J. Dolbeault, Best constants for Gagliardo-Nirenberg inequalities and applications to nonlinear diffusions. *J. Math. Pures Appl. (9)* **81** (2002), no. 9, 847–875 Zbl [1112.35310](#) MR [1940370](#)
- [36] M. del Pino and M. Sáez, On the extinction profile for solutions of $u_t = \Delta u^{(N-2)/(N+2)}$. *Indiana Univ. Math. J.* **50** (2001), no. 1, 611–628 Zbl [0991.35011](#) MR [1857048](#)
- [37] J. Denzler, H. Koch, and R. J. McCann, Higher-order time asymptotics of fast diffusion in Euclidean space: a dynamical systems approach. *Mem. Amer. Math. Soc.* **234** (2015), no. 1101 Zbl [1315.35004](#) MR [3307161](#)
- [38] J. Denzler and R. J. McCann, Fast diffusion to self-similarity: complete spectrum, long-time asymptotics, and numerology. *Arch. Ration. Mech. Anal.* **175** (2005), no. 3, 301–342 Zbl [1083.35074](#) MR [2126633](#)
- [39] J. Dolbeault, M. J. Esteban, S. Filippas, and A. Tertikas, Rigidity results with applications to best constants and symmetry of Caffarelli-Kohn-Nirenberg and logarithmic Hardy inequalities. *Calc. Var. Partial Differential Equations* **54** (2015), no. 3, 2465–2481 Zbl [1336.26023](#) MR [3412380](#)
- [40] J. Dolbeault, M. J. Esteban, and M. Loss, Rigidity versus symmetry breaking via nonlinear flows on cylinders and Euclidean spaces. *Invent. Math.* **206** (2016), no. 2, 397–440 Zbl [1379.49021](#) MR [3570296](#)
- [41] J. Dolbeault, M. J. Esteban, M. Loss, and M. Muratori, Symmetry for extremal functions in subcritical Caffarelli-Kohn-Nirenberg inequalities. *C. R. Math. Acad. Sci. Paris* **355** (2017), no. 2, 133–154 Zbl [1379.49009](#) MR [3612700](#)

- [42] J. Dolbeault, M. J. Esteban, G. Tarantello, and A. Tertikas, Radial symmetry and symmetry breaking for some interpolation inequalities. *Calc. Var. Partial Differential Equations* **42** (2011), no. 3-4, 461–485 Zbl [1246.26014](#) MR [2846263](#)
- [43] J. Dolbeault, M. Muratori, and B. Nazaret, Weighted interpolation inequalities: a perturbation approach. *Math. Ann.* **369** (2017), no. 3-4, 1237–1270 Zbl [1390.49005](#) MR [3713540](#)
- [44] J. Dolbeault and G. Toscani, Improved interpolation inequalities, relative entropy and fast diffusion equations. *Ann. Inst. H. Poincaré C Anal. Non Linéaire* **30** (2013), no. 5, 917–934 Zbl [1450.46020](#) MR [3103175](#)
- [45] D. Eidus, The Cauchy problem for the nonlinear filtration equation in an inhomogeneous medium. *J. Differential Equations* **84** (1990), no. 2, 309–318 Zbl [0707.35074](#) MR [1047572](#)
- [46] D. Eidus and S. Kamin, The filtration equation in a class of functions decreasing at infinity. *Proc. Amer. Math. Soc.* **120** (1994), no. 3, 825–830 Zbl [0791.35065](#) MR [1169025](#)
- [47] V. Felli and M. Schneider, Perturbation results of critical elliptic equations of Caffarelli-Kohn-Nirenberg type. *J. Differential Equations* **191** (2003), no. 1, 121–142 Zbl [1088.35023](#) MR [1973285](#)
- [48] A. Friedman and S. Kamin, The asymptotic behavior of gas in an n -dimensional porous medium. *Trans. Amer. Math. Soc.* **262** (1980), no. 2, 551–563 Zbl [0447.76076](#) MR [586735](#)
- [49] E. Gagliardo, Proprietà di alcune classi di funzioni in più variabili. *Ricerche Mat.* **7** (1958), 102–137 Zbl [0089.09401](#) MR [102740](#)
- [50] V. A. Galaktionov and L. A. Peletier, Asymptotic behaviour near finite-time extinction for the fast diffusion equation. *Arch. Rational Mech. Anal.* **139** (1997), no. 1, 83–98 Zbl [0885.35058](#) MR [1475779](#)
- [51] V. A. Galaktionov, L. A. Peletier, and J. L. Vázquez, Asymptotics of the fast-diffusion equation with critical exponent. *SIAM J. Math. Anal.* **31** (2000), no. 5, 1157–1174 Zbl [0954.35094](#) MR [1759202](#)
- [52] G. Grillo and M. Muratori, Radial fast diffusion on the hyperbolic space. *Proc. Lond. Math. Soc. (3)* **109** (2014), no. 2, 283–317 Zbl [1336.35201](#) MR [3254926](#)
- [53] G. Grillo, M. Muratori, and M. M. Porzio, Porous media equations with two weights: smoothing and decay properties of energy solutions via Poincaré inequalities. *Discrete Contin. Dyn. Syst.* **33** (2013), no. 8, 3599–3640 Zbl [1277.35217](#) MR [3021373](#)
- [54] G. Grillo, M. Muratori, and F. Punzo, On the asymptotic behaviour of solutions to the fractional porous medium equation with variable density. *Discrete Contin. Dyn. Syst.* **35** (2015), no. 12, 5927–5962 Zbl [1336.35057](#) MR [3393261](#)
- [55] C. E. Gutiérrez and R. L. Wheeden, Harnack’s inequality for degenerate parabolic equations. *Comm. Partial Differential Equations* **16** (1991), no. 4-5, 745–770 Zbl [0746.35007](#) MR [1113105](#)
- [56] L. Herreraiz, Asymptotic behaviour of solutions of some semilinear parabolic problems. *Ann. Inst. H. Poincaré C Anal. Non Linéaire* **16** (1999), no. 1, 49–105 Zbl [0918.35025](#) MR [1668560](#)
- [57] M. A. Herrero and M. Pierre, The Cauchy problem for $u_t = \Delta u^m$ when $0 < m < 1$. *Trans. Amer. Math. Soc.* **291** (1985), no. 1, 145–158 Zbl [0583.35052](#) MR [797051](#)
- [58] R. G. Iagar and A. Sánchez, Large time behavior for a porous medium equation in a nonhomogeneous medium with critical density. *Nonlinear Anal.* **102** (2014), 226–241 Zbl [1284.35059](#) MR [3182811](#)
- [59] K. Ishige, On the behavior of the solutions of degenerate parabolic equations. *Nagoya Math. J.* **155** (1999), 1–26 Zbl [0932.35131](#) MR [1711391](#)

- [60] K. Ishige and M. Murata, An intrinsic metric approach to uniqueness of the positive Cauchy problem for parabolic equations. *Math. Z.* **227** (1998), no. 2, 313–335 Zbl [0893.35042](#) MR [1609065](#)
- [61] K. Ishige and M. Murata, Uniqueness of nonnegative solutions of the Cauchy problem for parabolic equations on manifolds or domains. *Ann. Scuola Norm. Sup. Pisa Cl. Sci. (4)* **30** (2001), no. 1, 171–223 Zbl [1024.35010](#) MR [1882029](#)
- [62] A. S. Kalashnikov, Some problems of the qualitative theory of second-order nonlinear degenerate parabolic equations. *Uspekhi Mat. Nauk* **42** (1987), no. 2(254), 135–176, 287 Zbl [0642.35047](#) MR [898624](#)
- [63] S. Kamin, R. Kersner, and A. Tesei, On the Cauchy problem for a class of parabolic equations with variable density. *Atti Accad. Naz. Lincei Cl. Sci. Fis. Mat. Natur. Rend. Lincei (9) Mat. Appl.* **9** (1998), no. 4, 279–298 (1999) Zbl [0926.35045](#) MR [1722787](#)
- [64] S. Kamin, G. Reyes, and J. L. Vázquez, Long time behavior for the inhomogeneous PME in a medium with rapidly decaying density. *Discrete Contin. Dyn. Syst.* **26** (2010), no. 2, 521–549 Zbl [1196.35052](#) MR [2556497](#)
- [65] S. Kamin and P. Rosenau, Propagation of thermal waves in an inhomogeneous medium. *Comm. Pure Appl. Math.* **34** (1981), no. 6, 831–852 Zbl [0458.35042](#) MR [634287](#)
- [66] S. Kamin and P. Rosenau, Nonlinear thermal evolution in an inhomogeneous medium. *J. Math. Phys.* **23** (1982), no. 7, 1385–1390 Zbl [0499.76111](#) MR [666194](#)
- [67] C. Kienzler, H. Koch, and J. L. Vázquez, Flatness implies smoothness for solutions of the porous medium equation. *Calc. Var. Partial Differential Equations* **57** (2018), no. 1, Paper No. 18 Zbl [1403.35065](#) MR [3740398](#)
- [68] Y. J. Kim and R. J. McCann, Potential theory and optimal convergence rates in fast nonlinear diffusion. *J. Math. Pures Appl. (9)* **86** (2006), no. 1, 42–67 Zbl [1112.35025](#) MR [2246356](#)
- [69] S. Kullback, A lower bound for discrimination information in terms of variation (Corresp.). *IEEE Trans. Inform. Theory* **13** (1967), no. 1, 126–127
- [70] C. Lederman and P. A. Markowich, On fast-diffusion equations with infinite equilibrium entropy and finite equilibrium mass. *Comm. Partial Differential Equations* **28** (2003), no. 1-2, 301–332 Zbl [1024.35040](#) MR [1974458](#)
- [71] R. J. McCann, A convexity principle for interacting gases. *Adv. Math.* **128** (1997), no. 1, 153–179 Zbl [0901.49012](#) MR [1451422](#)
- [72] R. J. McCann and D. Slepčev, Second-order asymptotics for the fast-diffusion equation. *Int. Math. Res. Not.* (2006), Art. ID 24947 Zbl [1130.35080](#) MR [2211152](#)
- [73] W. I. Newman, A Lyapunov functional for the evolution of solutions to the porous medium equation to self-similarity. I. *J. Math. Phys.* **25** (1984), no. 10, 3120–3123 Zbl [0583.76114](#) MR [760591](#)
- [74] S. Nieto and G. Reyes, Asymptotic behavior of the solutions of the inhomogeneous porous medium equation with critical vanishing density. *Commun. Pure Appl. Anal.* **12** (2013), no. 2, 1123–1139 Zbl [1267.35037](#) MR [2982811](#)
- [75] L. Nirenberg, On elliptic partial differential equations. *Ann. Scuola Norm. Sup. Pisa Cl. Sci. (3)* **13** (1959), 115–162 Zbl [0088.07601](#) MR [109940](#)
- [76] F. Otto, The geometry of dissipative evolution equations: the porous medium equation. *Comm. Partial Differential Equations* **26** (2001), no. 1-2, 101–174 Zbl [0984.35089](#) MR [1842429](#)
- [77] F. Ragnedda, S. Vernier Piro, and V. Vespi, Pointwise estimates for the fundamental solutions of a class of singular parabolic problems. *J. Anal. Math.* **121** (2013), 235–253 Zbl [1282.35106](#) MR [3127384](#)

- [78] J. Ralston, A Lyapunov functional for the evolution of solutions to the porous medium equation to self-similarity. II. *J. Math. Phys.* **25** (1984), no. 10, 3124–3127 Zbl [0583.76115](#) MR [760592](#)
- [79] M. V. C. Recalde and V. Vespri, Backward estimates for nonnegative solutions to a class of singular parabolic equations. *Nonlinear Anal.* **144** (2016), 194–203 Zbl [1346.35118](#) MR [3534102](#)
- [80] G. Reyes and J. L. Vázquez, The Cauchy problem for the inhomogeneous porous medium equation. *Netw. Heterog. Media* **1** (2006), no. 2, 337–351 Zbl [1124.35035](#) MR [2223075](#)
- [81] G. Reyes and J. L. Vázquez, A weighted symmetrization for nonlinear elliptic and parabolic equations in inhomogeneous media. *J. Eur. Math. Soc. (JEMS)* **8** (2006), no. 3, 531–554 Zbl [1162.35033](#) MR [2250171](#)
- [82] G. Reyes and J. L. Vázquez, The inhomogeneous PME in several space dimensions. Existence and uniqueness of finite energy solutions. *Commun. Pure Appl. Anal.* **7** (2008), no. 6, 1275–1294 Zbl [1157.35412](#) MR [2425009](#)
- [83] G. Reyes and J. L. Vázquez, Long time behavior for the inhomogeneous PME in a medium with slowly decaying density. *Commun. Pure Appl. Anal.* **8** (2009), no. 2, 493–508 Zbl [1169.35313](#) MR [2461559](#)
- [84] A. Rodríguez and J. L. Vázquez, A well-posed problem in singular Fickian diffusion. *Arch. Rational Mech. Anal.* **110** (1990), no. 2, 141–163 Zbl [0695.76043](#) MR [1037346](#)
- [85] A. Rodríguez, J. L. Vázquez, and J. R. Esteban, The maximal solution of the logarithmic fast diffusion equation in two space dimensions. *Adv. Differential Equations* **2** (1997), no. 6, 867–894 Zbl [1023.35515](#) MR [1606339](#)
- [86] P. Rosenau and S. Kamin, Nonlinear diffusion in a finite mass medium. *Comm. Pure Appl. Math.* **35** (1982), no. 1, 113–127 Zbl [0469.35060](#) MR [637497](#)
- [87] G. Schimperna, A. Segatti, and S. Zelik, Asymptotic uniform boundedness of energy solutions to the Penrose-Fife model. *J. Evol. Equ.* **12** (2012), no. 4, 863–890 Zbl [1270.35151](#) MR [3000460](#)
- [88] G. Schimperna, A. Segatti, and S. Zelik, On a singular heat equation with dynamic boundary conditions. *Asymptot. Anal.* **97** (2016), no. 1-2, 27–59 Zbl [1344.35146](#) MR [3475117](#)
- [89] N. Simonov, *Fast diffusion equations with Caffarelli-Kohn-Nirenberg weights: Regularity and asymptotics*. Ph.D. thesis, Universidad Autónoma de Madrid, 2020
- [90] J. L. Vázquez, Nonexistence of solutions for nonlinear heat equations of fast-diffusion type. *J. Math. Pures Appl. (9)* **71** (1992), no. 6, 503–526 Zbl [0694.35088](#) MR [1193606](#)
- [91] J. L. Vázquez, Asymptotic behaviour for the porous medium equation posed in the whole space. pp. 67–118, 3, 2003 *J. Evol. Equ.* **3** (2003), no. 1, 67–118 Zbl [1036.35108](#) MR [1977429](#)
- [92] J. L. Vázquez, *Smoothing and decay estimates for nonlinear diffusion equations*. Oxford Lecture Ser. Math. Appl. 33, Oxford University Press, Oxford, 2006 Zbl [1113.35004](#) MR [2282669](#)
- [93] J. L. Vázquez, *The porous medium equation*. Oxford Math. Monogr., The Clarendon Press, Oxford University Press, Oxford, 2007 Zbl [1107.35003](#) MR [2286292](#)
- [94] J. L. Vázquez, Fundamental solution and long time behavior of the porous medium equation in hyperbolic space. *J. Math. Pures Appl. (9)* **104** (2015), no. 3, 454–484 Zbl [1327.35213](#) MR [3383174](#)
- [95] J. L. Vázquez, Asymptotic behaviour methods for the heat equation. Convergence to the Gaussian. 2018, arXiv:[1706.10034](#)

- [96] J. L. Vázquez, The evolution fractional p -Laplacian equation in \mathbb{R}^N . Fundamental solution and asymptotic behaviour. *Nonlinear Anal.* **119** (2020), Article ID 112034 Zbl [1447.35205](#) MR [4114983](#)
- [97] J. L. Vázquez, The fractional p -Laplacian evolution equation in \mathbb{R}^N in the sublinear case. *Calc. Var. Partial Differ. Equ.* **60** (2021), no. 4, Paper No. 140 Zbl [1471.35312](#) MR [4280519](#)

Received 10 June 2021; revised 24 November 2021; accepted 26 November 2021.

Matteo Bonforte

Departamento de Matemáticas, Universidad Autónoma de Madrid, ICMAT – Instituto de Ciencias Matemáticas, CSIC-UAM-UC3M-UCM, Calle Nicolás Cabrera 13-15, Campus de Cantoblanco, 28049 Madrid, Spain; matteo.bonforte@uam.es

Nikita Simonov

Université Paris-Saclay, CNRS, and Laboratoire de Mathématiques et Modélisation d'Évry, Université d'Évry, 91037 Évry-Courcouronnes, France; simonov.nikita@gmail.com