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On the structure of Hardy–Sobolev–Maz’ya inequalities

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Abstract. We establish new improvements of the optimal Hardy inequality in the half-space. We first add all possible linear combinations of Hardy type terms, thus revealing the structure of this type of inequalities and obtaining best constants. We then add the critical Sobolev term and obtain necessary and sufficient conditions for the validity of Hardy–Sobolev–Maz’ya type inequalities.

1. Introduction

One version of the Hardy inequality states that for convex domains $\Omega \subset \mathbb{R}^n$ the following estimate holds:

$$\int_{\Omega} |\nabla u|^2 dx \geq \frac{1}{4} \int_{\Omega} \frac{|u|^2}{d(x)^2} dx, \quad u \in C_0^\infty(\Omega),$$

where $d(x) = \text{dist}(x, \partial\Omega)$ and the constant $1/4$ is best possible. This result has been improved and generalized in many different ways: see for example [1], [2], [4]–[9], [12], [13].

One pioneering result due to Brezis and Marcus [4] is the following improved Hardy inequality:

$$\int_{\Omega} |\nabla u|^2 dx \geq \frac{1}{4} \int_{\Omega} \frac{|u|^2}{d(x)^2} dx + C_2(\Omega) \int_{\Omega} |u|^2 dx, \quad u \in C_0^\infty(\Omega), \quad (1.1)$$

valid for any convex domain $\Omega \subset \mathbb{R}^n$. This estimate has recently been extended in [7]:

$$\int_{\Omega} |\nabla u|^2 dx \geq \frac{1}{4} \int_{\Omega} \frac{|u|^2}{d(x)^2} dx + C_q(\Omega) \left(\int_{\Omega} |u|^q dx \right)^{2/q}, \quad u \in C_0^\infty(\Omega). \quad (1.2)$$

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Moreover, it is shown in [7] that there exist constants c_1 and c_2 only depending on q and the dimension n of Ω such that the best constant $C_q(\Omega)$ satisfies

$$c_1 D^{n-2-2n/q} \geq C_q(\Omega) \geq c_2 D^{n-2-2n/q},$$

where $D = \sup_{x \in \Omega} d(x) < \infty$ and $2 \leq q < 2n/(n-2)$. We note that the critical Sobolev exponent $q = 2^* := 2n/(n-2)$ is not included in the above theorem. For results in the critical case we refer to [8].

Let $S_n = \pi n(n-2)(\Gamma(n/2)/\Gamma(n))^{2/n}$, $n \geq 3$, denote the best constant in the Sobolev inequality

$$\int_{\Omega} |\nabla u|^2 dx \geq S_n \left(\int_{\Omega} |u|^{2^*} dx \right)^{2/2^*}, \quad u \in C_0^\infty(\Omega).$$

The first inequality that combines both the critical Sobolev exponent term and the Hardy term, the latter with best constant, is due to Maz'ya [10], and is the following Hardy–Sobolev–Maz'ya inequality:

$$\int_{\mathbb{R}_+^n} |\nabla u|^2 dx \geq \frac{1}{4} \int_{\mathbb{R}_+^n} \frac{|u|^2}{x_1^2} dx + C_n \left(\int_{\mathbb{R}_+^n} |u|^{2^*} dx \right)^{2/2^*}, \quad u \in C_0^\infty(\mathbb{R}_+^n), \quad (1.3)$$

where $\mathbb{R}_+^n = \{(x_1, \dots, x_n) : x_1 > 0\}$ denotes the upper half-space, C_n is a positive constant and $2^* = 2n/(n-2)$, $n \geq 3$. Recently, it was shown in [3] that in the 3-dimensional case, the best constant C_3 coincides with the best Sobolev constant S_3 ! On the other hand, when $n \geq 4$ one has $C_n < S_n$ (see [11]).

We next mention an improvement of Hardy's inequality that involves two distance functions:

$$\int_{\mathbb{R}_+^n} |\nabla u|^2 dx \geq \frac{1}{4} \int_{\mathbb{R}_+^n} \frac{|u|^2}{x_1^2} dx + C(\tau) \int_{\mathbb{R}_+^n} \frac{|u|^2}{x_1^{2-\tau} (x_1^2 + x_2^2)^{\tau/2}} dx, \quad u \in C_0^\infty(\mathbb{R}_+^n),$$

where $0 < \tau \leq 1$. This is a special case of a more general inequality proved in [13].

In this work we study improvements of Hardy's inequality that involve various distance functions. Working in the upper half-space \mathbb{R}_+^n , we obtain Hardy type inequalities that involve constant multiples of the inverse square of the distance to linear submanifolds of different codimensions of the boundary $\partial\mathbb{R}_+^n$. Actually, we are able to give a complete description of the structure of this kind of improved Hardy inequalities. In particular, we have a lot of freedom in choosing these constants and we will show that all our configurations of constants are, in a natural sense, optimal. More precisely, our first result reads:

Theorem A (Improved Hardy inequality).

(i) Let $\alpha_1, \dots, \alpha_n$ be arbitrary real numbers and

$$\beta_1 = -\alpha_1^2 + 1/4, \quad \beta_m = -\alpha_m^2 + (\alpha_{m-1} - 1/2)^2, \quad m = 2, \dots, n.$$

Then for any $u \in C_0^\infty(\mathbb{R}_+^n)$,

$$\int_{\mathbb{R}_+^n} |\nabla u|^2 dx \geq \int_{\mathbb{R}_+^n} \left(\frac{\beta_1}{x_1^2} + \frac{\beta_2}{x_1^2 + x_2^2} + \dots + \frac{\beta_n}{x_1^2 + \dots + x_n^2} \right) u^2 dx.$$

(ii) Suppose that for some real numbers β_1, \dots, β_n the inequality

$$\int_{\mathbb{R}_+^n} |\nabla u|^2 dx \geq \int_{\mathbb{R}_+^n} \left(\frac{\beta_1}{x_1^2} + \frac{\beta_2}{x_1^2 + x_2^2} + \dots + \frac{\beta_n}{x_1^2 + \dots + x_n^2} \right) u^2 dx$$

holds for any $u \in C_0^\infty(\mathbb{R}_+^n)$. Then there exist nonpositive constants $\alpha_1, \dots, \alpha_n$ such that

$$\beta_1 = -\alpha_1^2 + 1/4, \quad \beta_m = -\alpha_m^2 + (\alpha_{m-1} - 1/2)^2, \quad m = 2, \dots, n.$$

We next investigate the possibility of adding Sobolev type remainder terms. It turns out that almost every choice of the constants in Theorem A allows one to add a positive Sobolev term as well. The details are in our second main theorem.

Theorem B (Improved Hardy–Sobolev–Maz’ya inequality). *Let $\alpha_1, \dots, \alpha_n$ be arbitrary nonpositive real numbers and*

$$\beta_1 = -\alpha_1^2 + 1/4, \quad \beta_m = -\alpha_m^2 + (\alpha_{m-1} - 1/2)^2, \quad m = 2, \dots, n.$$

If $\alpha_n < 0$ then there exists a positive constant C such that for any $u \in C_0^\infty(\mathbb{R}_+^n)$,

$$\begin{aligned} \int_{\mathbb{R}_+^n} |\nabla u|^2 dx &\geq \int_{\mathbb{R}_+^n} \left(\frac{\beta_1}{x_1^2} + \frac{\beta_2}{x_1^2 + x_2^2} + \dots + \frac{\beta_n}{x_1^2 + \dots + x_n^2} \right) u^2 dx \\ &\quad + C \left(\int_{\mathbb{R}_+^n} |u|^{2^*} dx \right)^{2/2^*}. \end{aligned} \quad (1.4)$$

If $\alpha_n = 0$ then there is no positive constant C such that (1.4) holds.

It is interesting to note that the Sobolev term vanishes precisely when the constant β_n , in front of the Hardy-type term containing the point singularity, is chosen optimal. It is a bit curious that the size of the other constants, $\beta_1, \dots, \beta_{n-1}$, does not matter at all for this question. Only the relative size of β_n compared to the other constants matters.

Our results depend heavily on the Gagliardo–Nirenberg–Sobolev inequality and also on an interesting relation between the existence of an L^1 Hardy inequality and the possibility of adding a Sobolev type remainder term to the corresponding L^2 inequality. The precise result reads:

Theorem C. *Let $\Omega \subset \mathbb{R}^n$, $n \geq 3$, be a smooth domain. Assume that $\phi \in C^2(\Omega)$, $\phi > 0$, and the following weighted L^1 inequality holds:*

$$\int_{\Omega} \phi^{\frac{2(n-1)}{n-2}} |\nabla v| dx \geq C \int_{\Omega} \phi^{\frac{n}{n-2}} |\nabla \phi| |v| dx, \quad v \in C_0^\infty(\Omega). \quad (1.5)$$

Then there exists $c > 0$ such that

$$\int_{\Omega} |\nabla u|^2 dx \geq - \int_{\Omega} \frac{\Delta \phi}{\phi} |u|^2 dx + c \left(\int_{\Omega} |u|^{2^*} dx \right)^{2/2^*}, \quad u \in C_0^\infty(\Omega). \quad (1.6)$$

The regularity assumptions on ϕ can be weakened, but for our purposes it is enough to restrict ourselves to $\phi \in C^2(\Omega)$. We note that under the sole assumption $\phi > 0$ and $\phi \in C^2(\Omega)$ the inequality

$$\int_{\Omega} |\nabla u|^2 dx \geq - \int_{\Omega} \frac{\Delta \phi}{\phi} |u|^2 dx, \quad u \in C_0^\infty(\Omega), \tag{1.7}$$

is always true; see Lemma 2.1. It is the validity of (1.5) that makes possible the addition of the Sobolev term in (1.7). An easy example where both (1.5) and (1.6) fail is the case where ϕ is taken to be the first Dirichlet eigenfunction of the Laplacian of Ω , for Ω bounded.

Our methods are not restricted to the case $\Omega = \mathbb{R}_+^n$. In the last section of the paper we give an example of how to apply the method to get some results for the quarter-space. Moreover, as one can easily check, our results remain valid even for complex-valued functions.

The paper is organized as follows. In Section 2 we give the proof of Theorem A. In Section 3 we prove Theorems B and C. Finally, in the last section we obtain some results for the quarter-space.

2. Improved Hardy inequalities in the half-space

The half-space \mathbb{R}_+^n has some nice features that are not present for an arbitrary convex domain. The fact that the boundary has zero curvature is very useful when one tries to prove certain sorts of inequalities, as we shall see below.

We start with a general auxiliary lemma.

Lemma 2.1. (i) *Let $\mathbf{F} \in C^1(\Omega)$. Then*

$$\int_{\Omega} |\nabla u|^2 dx = \int_{\Omega} (\operatorname{div} \mathbf{F} - |\mathbf{F}|^2) |u|^2 dx + \int_{\Omega} |\nabla u + \mathbf{F}u|^2 dx, \quad u \in C_0^\infty(\Omega). \tag{2.1}$$

(ii) *Let $\phi > 0$, $\phi \in C^2(\Omega)$ and $u = \phi v$. Then*

$$\int_{\Omega} |\nabla u|^2 dx = - \int_{\Omega} \frac{\Delta \phi}{\phi} u^2 dx + \int_{\Omega} \phi^2 |\nabla v|^2 dx, \quad u \in C_0^\infty(\Omega). \tag{2.2}$$

Proof. By expanding the square we have

$$\int_{\Omega} |\nabla u + \mathbf{F}u|^2 dx = \int_{\Omega} |\nabla u|^2 dx + \int_{\Omega} |\mathbf{F}|^2 u^2 dx + \int_{\Omega} \mathbf{F} \cdot \nabla u^2 dx.$$

Identity (2.1) now follows by integrating by parts the last term.

To prove (2.2) we apply (2.1) to $\mathbf{F} = -\nabla \phi / \phi$. Elementary calculations now yield the result. □

We especially want to study inequalities of the type

$$\int_{\mathbb{R}_+^n} |\nabla u|^2 dx \geq \int_{\mathbb{R}_+^n} \left(\frac{\beta_1}{x_1^2} + \frac{\beta_2}{x_1^2 + x_2^2} + \cdots + \frac{\beta_n}{x_1^2 + x_2^2 + \cdots + x_n^2} \right) |u|^2 dx, \quad u \in C_0^\infty(\mathbb{R}_+^n),$$

where $\beta = (\beta_1, \dots, \beta_n)$ is a vector of nonnegative constants. The case when $\beta_1 = 1/4$ is especially interesting since it corresponds to the term in the standard Hardy inequality. So every legitimate choice of β with $\beta_1 = 1/4$ corresponds to an improved Hardy inequality. Let us introduce some notation. Let

$$\mathbf{X}_k := (x_1, \dots, x_k, 0, \dots, 0) \quad \text{so that} \quad |\mathbf{X}_k|^2 = x_1^2 + \cdots + x_k^2.$$

We now give the proof of the first part of Theorem A:

Proof of Theorem A(i). Let $\gamma_1, \dots, \gamma_n$ be arbitrary real numbers and set

$$\phi := |\mathbf{X}_1|^{-\gamma_1} \cdots |\mathbf{X}_n|^{-\gamma_n}, \quad \mathbf{F} := -\frac{\nabla \phi}{\phi}.$$

An easy calculation shows that

$$\mathbf{F} = \sum_{m=1}^n \gamma_m \frac{\mathbf{X}_m}{|\mathbf{X}_m|^2}.$$

With this choice of \mathbf{F} , we get

$$\operatorname{div} \mathbf{F} = \sum_{m=1}^n \gamma_m \frac{m-2}{|\mathbf{X}_m|^2}$$

and

$$|\mathbf{F}|^2 = \sum_{m=1}^n \frac{\gamma_m^2}{|\mathbf{X}_m|^2} + 2 \sum_{m=1}^n \sum_{j=1}^{m-1} \gamma_m \gamma_j \frac{\mathbf{X}_m}{|\mathbf{X}_m|^2} \frac{\mathbf{X}_j}{|\mathbf{X}_j|^2} = \sum_{m=1}^n \frac{\gamma_m^2}{|\mathbf{X}_m|^2} + 2 \sum_{m=1}^n \sum_{j=1}^{m-1} \frac{\gamma_m \gamma_j}{|\mathbf{X}_j|^2}.$$

We then get

$$-\frac{\Delta \phi}{\phi} = \operatorname{div} \mathbf{F} - |\mathbf{F}|^2 = \sum_{m=1}^n \frac{\beta_m}{|\mathbf{X}_m|^2}, \quad (2.3)$$

where

$$\beta_1 = -\gamma_1(\gamma_1 + 1), \quad \beta_m = -\gamma_m \left(2 - m + \gamma_m + 2 \sum_{j=1}^{m-1} \gamma_j \right), \quad m = 2, \dots, n.$$

We next set

$$\gamma_1 = \alpha_1 - 1/2, \quad \gamma_m = \alpha_m - \alpha_{m-1} + 1/2, \quad m = 2, \dots, n.$$

With this choice of γ 's the β 's are given as in the statement of the theorem.

As a consequence of Lemma 2.1 we have

$$\int_{\mathbb{R}_+^n} |\nabla u|^2 dx \geq \int_{\mathbb{R}_+^n} (\operatorname{div} \mathbf{F} - |\mathbf{F}|^2) u^2 dx. \quad (2.4)$$

The result then follows from (2.3) and (2.4). \square

Remark. It is easy to check that for any choice of n real numbers $\alpha_1, \dots, \alpha_n$, we can find n nonpositive real numbers $\alpha'_1, \dots, \alpha'_n$ that give the same constants β_1, \dots, β_n . Consequently, without loss of generality, we may assume that $\alpha_1, \dots, \alpha_n$ are nonpositive.

In the above theorem we have a lot of freedom. We can choose the γ 's in many different ways, each choice giving a different inequality. We may, for instance, first maximize β_1 and then β_2 and so on. More generally, we might try to make the first $m - 1$ β_m 's equal to zero and then maximize the β_m 's in increasing order.

In fact, we have the following corollary:

Corollary 2.2. *Let $k = 1, \dots, n$. Then*

$$\int_{\mathbb{R}_+^n} |\nabla u|^2 dx \geq \int_{\mathbb{R}_+^n} \left(\frac{k^2}{4} \frac{1}{x_1^2 + \dots + x_k^2} + \frac{1}{4} \frac{1}{x_1^2 + \dots + x_{k+1}^2} + \dots \right. \\ \left. + \frac{1}{4} \frac{1}{x_1^2 + \dots + x_n^2} \right) u^2 dx, \quad u \in C_0^\infty(\mathbb{R}_+^n).$$

Proof. In the case $k = 1$ we choose $\alpha_1 = \dots = \alpha_n = 0$. In this case all β_k 's are equal to $1/4$.

In the general case $k > 1$ we choose $\alpha_m = -m/2$ when $m = 1, \dots, k - 1$, and $\alpha_m = 0$ when $m = k, \dots, n$. \square

We next prove the second part of Theorem A:

Proof of Theorem A(ii). We will first prove that $\beta_1 \leq 1/4$, therefore $\beta_1 = -\alpha_1^2 + 1/4$ for suitable $\alpha_1 \leq 0$. Then, for this β_1 , we will prove that $\beta_2 \leq (\alpha_1 - 1/2)^2$, and therefore $\beta_2 = -\alpha_2^2 + (\alpha_1 - 1/2)^2$ for suitable $\alpha_2 \leq 0$, and so on.

Step 1. Let us first prove the estimate for β_1 . To this end we set

$$Q_1[u] := \frac{\int_{\mathbb{R}_+^n} |\nabla u|^2 dx - \sum_{i=2}^n \beta_i \int_{\mathbb{R}_+^n} \frac{u^2}{(x_1^2 + \dots + x_i^2)} dx}{\int_{\mathbb{R}_+^n} \frac{u^2}{x_1^2} dx}. \quad (2.5)$$

We clearly have $\beta_1 \leq \inf_{u \in C_0^\infty(\mathbb{R}_+^n)} Q_1[u]$. We will show that

$$\inf_{u \in C_0^\infty(\mathbb{R}_+^n)} Q_1[u] \leq 1/4, \quad (2.6)$$

whence $\beta_1 \leq 1/4$.

At this point we introduce a family of cutoff functions for later use. For $j = 1, \dots, n$ and $k_j > 0$ we set

$$\phi_j(t) = \begin{cases} 0, & t < 1/k_j^2, \\ 1 + \frac{\ln k_j t}{\ln k_j}, & 1/k_j^2 \leq t < 1/k_j, \\ 1, & t \geq 1/k_j, \end{cases}$$

and

$$h_{k_j}(x) := \phi_j(r_j) \quad \text{where} \quad r_j := |\mathbf{X}_j| = (x_1^2 + \dots + x_j^2)^{1/2}.$$

Note that

$$|\nabla h_{k_j}(x)|^2 = \begin{cases} \frac{1}{\ln^2 k_j} \frac{1}{r_j^2}, & 1/k_j^2 \leq r_j \leq 1/k_j, \\ 0, & \text{otherwise.} \end{cases}$$

We also denote by $\phi(x)$ a radially symmetric $C_0^\infty(\mathbb{R}^n)$ function such that $\phi = 1$ for $|x| < 1/2$ and $\phi = 0$ for $|x| > 1$.

To prove (2.6) we consider the family of functions

$$u_{k_1}(x) = x_1^{1/2} h_{k_1}(x) \phi(x). \tag{2.7}$$

We will show that as $k_1 \rightarrow \infty$,

$$\frac{\int_{\mathbb{R}_+^n} |\nabla u_{k_1}|^2 dx - \sum_{i=2}^n \beta_i \int_{\mathbb{R}_+^n} \frac{u_{k_1}^2}{(x_1^2 + \dots + x_i^2)} dx}{\int_{\mathbb{R}_+^n} \frac{u_{k_1}^2}{x_1^2} dx} = \frac{\int_{\mathbb{R}_+^n} |\nabla u_{k_1}|^2 dx}{\int_{\mathbb{R}_+^n} \frac{u_{k_1}^2}{x_1^2} dx} + o(1). \tag{2.8}$$

To see this, let us first examine the behavior of the denominator. For k_1 large we easily compute

$$\int_{\mathbb{R}_+^n} \frac{u_{k_1}^2}{x_1^2} dx = \int_{\mathbb{R}_+^n} x_1^{-1} h_{k_1}^2 \phi^2 dx > C \int_{1/k_1}^{1/2} x_1^{-1} dx > C \ln k_1. \tag{2.9}$$

On the other hand, the terms $\sum_{i=2}^n \beta_i \int_{\mathbb{R}_+^n} \frac{u_{k_1}^2}{(x_1^2 + \dots + x_i^2)} dx$ are easily seen to be bounded as $k_1 \rightarrow \infty$, by the Lebesgue dominated convergence theorem. From this and (2.9) we deduce (2.8).

We now estimate the gradient term in (2.8):

$$\begin{aligned} \int_{\mathbb{R}_+^n} |\nabla u_{k_1}|^2 dx &= \frac{1}{4} \int_{\mathbb{R}_+^n} x_1^{-1} h_{k_1}^2 \phi^2 dx + \int_{\mathbb{R}_+^n} x_1 |\nabla h_{k_1}|^2 \phi^2 \\ &\quad + \int_{\mathbb{R}_+^n} x_1 h_{k_1}^2 |\nabla \phi|^2 + \text{mixed terms.} \end{aligned} \tag{2.10}$$

The first integral of the right hand side behaves exactly as the denominator (cf. (2.9)), that is, it goes to infinity like $O(\ln k_1)$. The last integral is easily seen to be bounded as $k_1 \rightarrow \infty$. For the middle integral we have

$$\int_{\mathbb{R}_+^n} x_1 |\nabla h_{k_1}|^2 \phi^2 \leq \frac{C}{\ln^2 k_1} \int_{1/k_1^2 \leq x_1 \leq 1/k_1} x_1^{-1} dx_1 \leq \frac{C}{\ln k_1}.$$

As a consequence of these estimates, we easily see that the mixed terms in (2.10) are of the order $o(\ln k_1)$ as $k_1 \rightarrow \infty$. Hence, as $k_1 \rightarrow \infty$,

$$\int_{\mathbb{R}_+^n} |\nabla u_{k_1}|^2 dx = \frac{1}{4} \int_{\mathbb{R}_+^n} x_1^{-1} h_{k_1}^2 \phi^2 dx + o(\ln k_1). \tag{2.11}$$

From (2.8), (2.9) and (2.11) we conclude that as $k_1 \rightarrow \infty$,

$$Q_1[u_{k_1}] = 1/4 + o(1),$$

hence $\inf_{u \in C_0^\infty(\mathbb{R}_+^n)} Q_1[u] \leq 1/4$ and consequently $\beta_1 \leq 1/4$. Therefore for a suitable nonnegative constant α_1 we have $\beta_1 = -\alpha_1^2 + 1/4$. We also set

$$\gamma_1 := \alpha_1 - 1/2. \tag{2.12}$$

Step 2. We will next show that $\beta_2 \leq (\alpha_1 - 1/2)^2$. To this end, setting

$$Q_2[u] := \frac{\int_{\mathbb{R}_+^n} |\nabla u|^2 dx - (1/4 - \alpha_1^2) \int_{\mathbb{R}_+^n} \frac{u^2}{x_1^2} dx - \sum_{i=3}^n \beta_i \int_{\mathbb{R}_+^n} \frac{u^2}{|\mathbf{x}_i|^2} dx}{\int_{\mathbb{R}_+^n} \frac{u^2}{|\mathbf{x}_2|^2} dx}, \tag{2.13}$$

we will prove that

$$\inf_{u \in C_0^\infty(\mathbb{R}_+^n)} Q_2[u] \leq (\alpha_1 - 1/2)^2.$$

We now consider the family of functions

$$u_{k_1, k_2}(x) := x_1^{-\gamma_1} |\mathbf{x}_2|^{\alpha_1 - 1/2} h_{k_1}(x) h_{k_2}(x) \phi(x) =: x_1^{-\gamma_1} v_{k_1, k_2}(x). \tag{2.14}$$

An easy calculation shows that

$$Q_2[u_{k_1, k_2}] = \frac{\int_{\mathbb{R}_+^n} x_1^{-2\gamma_1} |\nabla v_{k_1, k_2}|^2 dx - \sum_{i=3}^n \beta_i \int_{\mathbb{R}_+^n} x_1^{-2\gamma_1} |\mathbf{x}_i|^{-2} v_{k_1, k_2}^2 dx}{\int_{\mathbb{R}_+^n} x_1^{-2\gamma_1} |\mathbf{x}_2|^{-2} v_{k_1, k_2}^2 dx}. \tag{2.15}$$

We next use the precise form of $v_{k_1, k_2}(x)$. Concerning the denominator of $Q_2[u_{k_1, k_2}]$ we have

$$\int_{\mathbb{R}_+^n} x_1^{-2\gamma_1} |\mathbf{x}_2|^{-2} v_{k_1, k_2}^2 dx = \int_{\mathbb{R}_+^n} x_1^{1-2\alpha_1} (x_1^2 + x_2^2)^{\alpha_1 - 3/2} h_{k_1}^2 h_{k_2}^2 \phi^2 dx.$$

Sending k_1 to infinity, using the structure of the cutoff functions and then introducing polar coordinates we get

$$\begin{aligned} \int_{\mathbb{R}_+^n} x_1^{-2\gamma_1} |\mathbf{X}_2|^{-2} v_{\infty, k_2}^2 dx &= \int_{\mathbb{R}_+^n} x_1^{1-2\alpha_1} (x_1^2 + x_2^2)^{\alpha_1-3/2} h_{k_2}^2 \phi^2 dx \\ &\geq C \int_{1/k_2 < x_1^2 + x_2^2 < 1/2} x_1^{1-2\alpha_1} (x_1^2 + x_2^2)^{\alpha_1-3/2} dx_1 dx_2 \\ &\geq C \int_0^\pi \int_{1/k_2}^{1/2} r^{-1} (\sin \theta)^{1-2\alpha_1} dr d\theta \geq C \ln k_2. \end{aligned} \quad (2.16)$$

The terms in the numerator that are multiplied by the β_i 's stay bounded as k_1 or k_2 go to infinity (cf. the estimates related to (2.29) in Step 3). Next,

$$\begin{aligned} \int_{\mathbb{R}_+^n} x_1^{-2\gamma_1} |\nabla v_{k_1, k_2}|^2 dx &= (\alpha_1 - 1/2)^2 \int_{\mathbb{R}_+^n} x_1^{-2\gamma_1} |\mathbf{X}_2|^{2\alpha_1-3} h_{k_1}^2 h_{k_2}^2 \phi^2 dx \\ &\quad + \int_{\mathbb{R}_+^n} x_1^{-2\gamma_1} |\mathbf{X}_2|^{2\alpha_1-1} |\nabla(h_{k_1} h_{k_2})|^2 \phi^2 dx \\ &\quad + \int_{\mathbb{R}_+^n} x_1^{-2\gamma_1} |\mathbf{X}_2|^{2\alpha_1-1} h_{k_1}^2 h_{k_2}^2 |\nabla \phi|^2 dx \\ &\quad + \text{mixed terms.} \end{aligned} \quad (2.17)$$

The first integral on the right hand side above is the same as the denominator of Q_2 , and therefore is finite as $k_1 \rightarrow \infty$ and increases like $\ln k_2$ as $k_2 \rightarrow \infty$ (cf. (2.16)). The last integral is bounded, no matter how big the k_1 and k_2 are. Concerning the middle term we have

$$\begin{aligned} M[v_{k_1, k_2}] &:= \int_{\mathbb{R}_+^n} x_1^{-2\gamma_1} |\mathbf{X}_2|^{2\alpha_1-1} |\nabla(h_{k_1} h_{k_2})|^2 \phi^2 dx \\ &= \int_{\mathbb{R}_+^n} x_1^{-2\gamma_1} |\mathbf{X}_2|^{2\alpha_1-1} |\nabla h_{k_1}|^2 h_{k_2}^2 \phi^2 dx \\ &\quad + \int_{\mathbb{R}_+^n} x_1^{-2\gamma_1} |\mathbf{X}_2|^{2\alpha_1-1} h_{k_1}^2 |\nabla h_{k_2}|^2 \phi^2 dx + \text{mixed terms} \\ &=: I_1 + I_2 + \text{mixed terms.} \end{aligned} \quad (2.18)$$

Since

$$|\mathbf{X}_2|^{2\alpha_1-1} h_{k_2}^2 = r_2^{2\alpha_1-1} \phi_2(r_2) \leq C k_2, \quad 0 < r_2 < 1,$$

we easily get

$$I_1 \leq \frac{C}{(\ln k_1)^2} \int_{1/k_1^2}^{1/k_1} x_1^{-1-2\alpha_1} dx_1,$$

and therefore, since $\alpha_1 \leq 0$,

$$I_1 \leq \frac{C}{\ln k_1}, \quad k_1 \rightarrow \infty. \quad (2.19)$$

Also, since $h_{k_1}^2 \leq 1$, we similarly get (for any k_1)

$$I_2 \leq \frac{C}{(\ln k_2)^2} \int_{1/k_2}^{1/k_1} r_2^{-1} dr_2 \leq \frac{C}{\ln k_2}, \quad k_2 \rightarrow \infty. \tag{2.20}$$

From (2.18)–(2.20) we infer that as $k_2 \rightarrow \infty$,

$$M[v_{\infty, k_2}] = o(1).$$

Returning to (2.17) we see that as $k_2 \rightarrow \infty$,

$$\int_{\mathbb{R}_+^n} x_1^{-2\gamma_1} |\nabla v_{\infty, k_2}|^2 dx = (\alpha_1 - 1/2)^2 \int_{\mathbb{R}_+^n} x_1^{-2\gamma_1} |\mathbf{X}_2|^{-2} v_{\infty, k_2}^2 dx + o(\ln k_2). \tag{2.21}$$

We then have as $k_2 \rightarrow \infty$,

$$Q_2[u_{\infty, k_2}] = (\alpha_1 - 1/2)^2 + o(1); \tag{2.22}$$

consequently, $\beta_2 \leq (\alpha_1 - 1/2)^2$, and therefore $\beta_2 = -\alpha_2^2 + (\alpha_1 - 1/2)^2$ for suitable $\alpha_2 \leq 0$. We also set

$$\gamma_2 = \alpha_2 - \alpha_1 + 1/2.$$

Step 3. The general case. At the $(q - 1)$ th step, $1 \leq q \leq n$, we have already established that

$$\beta_1 = -\alpha_1^2 + 1/4, \quad \beta_m = -\alpha_m^2 + (\alpha_{m-1} - 1/2)^2, \quad m = 2, \dots, q - 1,$$

for suitable nonpositive constants α_i . Also, we have defined

$$\gamma_1 = \alpha_1 - 1/2, \quad \gamma_m = \alpha_m - \alpha_{m-1} + 1/2, \quad m = 2, \dots, q - 1.$$

Our goal for the rest of the proof is to show that $\beta_q \leq (\alpha_{q-1} - 1/2)^2$. To this end we consider the quotient

$$Q_q[u] := \frac{\int_{\mathbb{R}_+^n} |\nabla u|^2 dx - \sum_{q \neq i=1}^n \beta_i \int_{\mathbb{R}_+^n} \frac{u^2}{|\mathbf{X}_i|^2} dx}{\int_{\mathbb{R}_+^n} \frac{u^2}{|\mathbf{X}_q|^2} dx}. \tag{2.23}$$

The test function is now given by

$$\begin{aligned} u_{k_1, k_q}(x) &:= x_1^{-\gamma_1} |\mathbf{X}_2|^{-\gamma_2} \dots |\mathbf{X}_{q-1}|^{-\gamma_{q-1}} |\mathbf{X}_q|^{\alpha_{q-1}-1/2} h_{k_1}(x) h_{k_q}(x) \phi(x) \\ &=: x_1^{-\gamma_1} |\mathbf{X}_2|^{-\gamma_2} \dots |\mathbf{X}_{q-1}|^{-\gamma_{q-1}} v_{k_q}(x). \end{aligned} \tag{2.24}$$

A straightforward calculation shows that

$$\begin{aligned} &Q_q[u_{k_1, k_q}] \\ &= \frac{\int_{\mathbb{R}_+^n} \prod_{j=1}^{q-1} |\mathbf{X}_j|^{-2\gamma_j} |\nabla v_{k_1, k_q}|^2 dx - \sum_{i=q+1}^n \beta_i \int_{\mathbb{R}_+^n} \prod_{j=1}^{q-1} |\mathbf{X}_j|^{-2\gamma_j} |\mathbf{X}_i|^{-2} v_{k_1, k_q}^2 dx}{\int_{\mathbb{R}_+^n} \prod_{j=1}^{q-1} |\mathbf{X}_j|^{-2\gamma_j} |\mathbf{X}_q|^{-2} v_{k_1, k_q}^2 dx}. \end{aligned} \tag{2.25}$$

Let us first look at the denominator,

$$D_q[u_{k_1, k_q}] := \int_{\mathbb{R}_+^n} \prod_{j=1}^{q-1} |\mathbf{X}_j|^{-2\gamma_j} |\mathbf{X}_q|^{2\alpha_{q-1}-3} h_{k_1}(x) h_{k_q}(x) \phi(x) dx.$$

Sending $k_1 \rightarrow \infty$, we find that $h_{k_1} \rightarrow 1$ and therefore

$$D_q[u_{\infty, k_q}] = \int_{\mathbb{R}_+^n} \prod_{j=1}^{q-1} |\mathbf{X}_j|^{-2\gamma_j} |\mathbf{X}_q|^{2\alpha_{q-1}-3} h_{k_q}(x) \phi(x) dx.$$

To see that this is finite we note that with $B_R^+ := \{x \in \mathbb{R}^n : |x| < R, x_1 \geq 0\}$,

$$\begin{aligned} D_q[u_{\infty, k_q}] &\leq \int_{B_1^+ \cap \{1/k_q^2 \leq r_q \leq 1/k_q\}} \prod_{j=1}^{q-1} |\mathbf{X}_j|^{-2\gamma_j} |\mathbf{X}_q|^{2\alpha_{q-1}-3} dx \\ &\leq C \int_{\{1/k_q^2 \leq r_q \leq 1/k_q\}} \prod_{j=1}^{q-1} |\mathbf{X}_j|^{-2\gamma_j} |\mathbf{X}_q|^{2\alpha_{q-1}-3} dx_1 \dots dx_q. \end{aligned} \quad (2.26)$$

To estimate this, we introduce polar coordinates $(x_1, \dots, x_q) \rightarrow (r_q, \theta_1, \dots, \theta_{q-1})$,

$$\begin{aligned} x_1 &= r_q \sin \theta_{q-1} \sin \theta_{q-2} \dots \sin \theta_2 \sin \theta_1 \\ x_2 &= r_q \sin \theta_{q-1} \sin \theta_{q-2} \dots \sin \theta_2 \cos \theta_1 \\ x_3 &= r_q \sin \theta_{q-1} \sin \theta_{q-2} \dots \cos \theta_2 \\ &\vdots \\ x_q &= r_q \cos \theta_{q-1}, \end{aligned}$$

where $0 \leq \theta_1 < 2\pi$ and $0 \leq \theta_m < \pi$ for $m = 2, \dots, q-1$. The surface measure on the unit sphere S^{q-1} then becomes

$$C(\sin \theta_{q-1})^{q-2} (\sin \theta_{q-2})^{q-3} \dots \sin \theta_2 d\theta_1 \dots d\theta_{q-1}.$$

Also, $r_q = |\mathbf{X}_q|$ and for $1 \leq m \leq q-1$,

$$r_m = |\mathbf{X}_m| = (x_1^2 + \dots + x_m^2)^{1/2} = r_q \sin \theta_{q-1} \sin \theta_{q-2} \dots \sin \theta_m.$$

We then have

$$\begin{aligned} &\int_{\{1/k_q^2 \leq r_q \leq 1/k_q\}} \prod_{j=1}^{q-1} |\mathbf{X}_j|^{-2\gamma_j} |\mathbf{X}_q|^{2\alpha_{q-1}-3} dx_1 \dots dx_q \\ &= C \int_{\{1/k_q^2 \leq r_q \leq 1/k_q\}} r_q^{-1} \prod_{j=1}^{q-1} (\sin \theta_j)^{1-2\alpha_j} d\theta_1 \dots d\theta_{q-1} dr_q \leq C \ln k_q. \end{aligned} \quad (2.27)$$

On the other hand, since

$$D_q[u_{\infty, k_q}] \geq \int_{B_{1/2}^+ \cap \{1/k_q \leq r_q \leq 1/2\}} \prod_{j=1}^{q-1} |\mathbf{X}_j|^{-2\gamma_j} |\mathbf{X}_q|^{2\alpha_{q-1}-3} dx,$$

by practically the same argument we find that as $k_q \rightarrow \infty$,

$$D_q[u_{\infty, k_q}] \geq C \ln k_q. \tag{2.28}$$

For $i = q + 1, \dots, n$, we consider the terms

$$\begin{aligned} \int_{\mathbb{R}_+^n} \prod_{j=1}^{q-1} |\mathbf{X}_j|^{-2\gamma_j} |\mathbf{X}_i|^{-2} v_{k_1, k_q}^2 dx &= \int_{\mathbb{R}_+^n} \prod_{j=1}^{q-1} |\mathbf{X}_j|^{-2\gamma_j} |\mathbf{X}_q|^{2\alpha_{q-1}-1} |\mathbf{X}_i|^{-2} h_{k_1}^2 h_{k_q}^2 \phi(x)^2 dx \\ &\leq \int_{\mathbb{R}_+^n} \prod_{j=1}^{q-1} |\mathbf{X}_j|^{-2\gamma_j} |\mathbf{X}_q|^{2\alpha_{q-1}-1} |\mathbf{X}_{q+1}|^{-2} h_{k_1}^2 h_{k_q}^2 \phi(x)^2 dx. \end{aligned} \tag{2.29}$$

If we let first $k_1 \rightarrow \infty$ and then $k_q \rightarrow \infty$, the above integral converges to

$$I_q := \int_{\mathbb{R}_+^n} \prod_{j=1}^{q-1} |\mathbf{X}_j|^{-2\gamma_j} |\mathbf{X}_q|^{2\alpha_{q-1}-1} |\mathbf{X}_{q+1}|^{-2} \phi(x)^2 dx.$$

To see this is finite we introduce polar coordinates $(x_1, \dots, x_{q+1}) \rightarrow (r_{q+1}, \theta_1, \dots, \theta_q)$ and use elementary estimates to get

$$I_q \leq C \int_{B_1^+} \sin \theta_q \prod_{j=1}^q (\sin \theta_j)^{1-2\alpha_j} d\theta_1 \dots d\theta_q dr_{q+1} < \infty.$$

We next consider the gradient term

$$\begin{aligned} &\int_{\mathbb{R}_+^n} \prod_{j=1}^{q-1} |\mathbf{X}_j|^{-2\gamma_j} |\nabla v_{k_1, k_q}|^2 dx \\ &= (\alpha_{q-1} - 1/2)^2 \int_{\mathbb{R}_+^n} \prod_{j=1}^{q-1} |\mathbf{X}_j|^{-2\gamma_j} |\mathbf{X}_q|^{2\alpha_{q-1}-3} h_{k_1}^2 h_{k_q}^2 \phi^2 dx \\ &\quad + \int_{\mathbb{R}_+^n} \prod_{j=1}^{q-1} |\mathbf{X}_j|^{-2\gamma_j} |\mathbf{X}_q|^{2\alpha_{q-1}-1} |\nabla(h_{k_1} h_{k_q})|^2 \phi^2 dx \\ &\quad + \int_{\mathbb{R}_+^n} \prod_{j=1}^{q-1} |\mathbf{X}_j|^{-2\gamma_j} |\mathbf{X}_q|^{2\alpha_{q-1}-1} h_{k_1}^2 h_{k_q}^2 |\nabla \phi|^2 dx + \text{mixed terms}. \end{aligned} \tag{2.30}$$

The first term of the right hand side is the same as the denominator. Using polar coordinates and arguments similar to the ones used in estimating the gradient term in (2.17), all other terms of (2.30) are bounded as $k_1 \rightarrow \infty$ and $k_q \rightarrow \infty$. In particular, we end up with

$$\int_{\mathbb{R}_+^n} \prod_{j=1}^{q-1} |\mathbf{X}_j|^{-2\gamma_j} |\nabla v_{\infty, k_q}|^2 dx = (\alpha_{q-1} - 1/2)^2 D_q[u_{\infty, k_q}] + o(\ln k_q), \quad k_q \rightarrow \infty.$$

Putting things together we obtain

$$Q_q[u_{\infty, k_q}] = (\alpha_{q-1} - 1/2)^2 + o(1), \quad k_q \rightarrow \infty,$$

from which it follows that $\beta_q \leq (\alpha_{q-1} - 1/2)^2$. This completes the proof of the theorem. \square

The previous analysis can also lead to the following result:

Theorem 2.3. *Let $\alpha_1, \dots, \alpha_k$, $1 \leq k \leq n - 1$, be nonpositive constants and*

$$\beta_1 = -\alpha_1^2 + 1/4, \quad \beta_m = -\alpha_m^2 + (\alpha_{m-1} - 1/2)^2, \quad m = 2, \dots, k.$$

Suppose that there exists a constant β_{k+1} such that the inequality

$$\int_{\mathbb{R}_+^n} |\nabla u|^2 dx \geq \int_{\mathbb{R}_+^n} \left(\frac{\beta_1}{x_1^2} + \frac{\beta_2}{x_1^2 + x_2^2} + \dots + \frac{\beta_{k+1}}{x_1^2 + \dots + x_{k+1}^2} \right) u^2 dx \quad (2.31)$$

holds for any $u \in C_0^\infty(\mathbb{R}_+^n)$. Then

$$\beta_{k+1} \leq (\alpha_k - 1/2)^2. \quad (2.32)$$

Moreover,

$$\inf_{u \in C_0^\infty(\mathbb{R}_+^n)} \frac{\int_{\mathbb{R}_+^n} |\nabla u|^2 dx - \beta_1 \int_{\mathbb{R}_+^n} \frac{|u|^2}{x_1^2} dx - \dots - \beta_k \int_{\mathbb{R}_+^n} \frac{|u|^2}{x_1^2 + \dots + x_k^2} dx}{\int_{\mathbb{R}_+^n} \frac{|u|^2}{x_1^2 + \dots + x_{k+1}^2} dx} = (\alpha_k - 1/2)^2. \quad (2.33)$$

Proof. The proof of the first part, that is, of estimate (2.32), is contained in the proof of Theorem A(ii).

To establish (2.33), we first use (2.32) to deduce that the infimum in (2.33) is less than or equal to $(\alpha_k - 1/2)^2$. To obtain the reverse inequality we use Theorem A(i) with $\alpha_{k+l} = -(l - 1)/2$, $l = 1, \dots, n - k$. For this choice we have $\beta_{k+2} = \dots = \beta_n = 0$. \square

The following is an interesting consequence of the previous theorem.

Corollary 2.4. For $1 \leq k \leq n$,

$$\inf_{u \in C_0^\infty(\mathbb{R}_+^n)} \frac{\int_{\mathbb{R}_+^n} |\nabla u|^2 dx}{\int_{\mathbb{R}_+^n} \frac{|u|^2}{x_1^2 + \dots + x_k^2} dx} = \frac{k^2}{4}, \tag{2.34}$$

and for $k \leq m < n$,

$$\inf_{u \in C_0^\infty(\mathbb{R}_+^n)} \frac{\int_{\mathbb{R}_+^n} |\nabla u|^2 dx - \frac{k^2}{4} \int_{\mathbb{R}_+^n} \frac{|u|^2}{x_1^2 + \dots + x_k^2} dx - \frac{1}{4} \int_{\mathbb{R}_+^n} \frac{|u|^2}{x_1^2 + \dots + x_{k+1}^2} dx - \dots - \frac{1}{4} \int_{\mathbb{R}_+^n} \frac{|u|^2}{x_1^2 + \dots + x_m^2} dx}{\int_{\mathbb{R}_+^n} \frac{|u|^2}{x_1^2 + \dots + x_{m+1}^2} dx} = \frac{1}{4}. \tag{2.35}$$

Proof. To establish (2.34) we use (2.33) with $\alpha_l = -l/2, l = 1, \dots, k - 1$.

To establish (2.35) we again use (2.33) with $\alpha_l = -l/2, l = 1, \dots, k - 1$, and $\alpha_l = 0, k \leq l \leq m$. With this choice we have $\beta_1 = \dots = \beta_{k-1} = 0, \beta_k = k^2/4$ and $\beta_l = 1/4, l = k + 1, \dots, m$. \square

3. Hardy–Sobolev–Maz’ya inequalities

We begin by proving Theorem C.

Proof of Theorem C. Our starting point is the Gagliardo–Nirenberg–Sobolev inequality

$$C_n \int_{\Omega} |f|^{\frac{n}{n-1}} dx \leq \left(\int_{\Omega} |\nabla f| dx \right)^{\frac{n}{n-1}}, \quad f \in C_0^\infty(\Omega). \tag{3.1}$$

Let $f = \phi^\alpha w$, where $\alpha = \frac{2(n-1)}{n-2}$. This leads to

$$C_n \int_{\Omega} \phi^{\frac{\alpha n}{n-1}} |w|^{\frac{n}{n-1}} dx \leq \left(\int_{\Omega} \alpha \phi^{\alpha-1} |\nabla \phi| |w| + \phi^\alpha |\nabla w| dx \right)^{\frac{n}{n-1}}, \quad w \in C_0^\infty(\Omega).$$

We now estimate the first term in the integral according to inequality (1.5) and let $w = |v|^\theta$. Then we get

$$\begin{aligned} C \left(\int_{\Omega} \phi^{\alpha n/n-1} |v|^{\theta n/n-1} dx \right)^{\frac{n-1}{n}} &\leq \int_{\Omega} \phi^\alpha |v|^{\theta-1} |\nabla v| dx \\ &\leq \left(\int_{\Omega} \phi^{2\alpha-2} |v|^{2\theta-2} dx \right)^{1/2} \left(\int_{\Omega} \phi^2 |\nabla v|^2 dx \right)^{1/2}. \end{aligned}$$

The choice

$$\theta = \alpha = \frac{2(n-1)}{n-2}$$

gives us the inequality

$$C \left(\int_{\Omega} \phi^{\frac{2n}{n-2}} |v|^{\frac{2n}{n-2}} dx \right)^{\frac{n-2}{n}} \leq \int_{\Omega} \phi^2 |\nabla v|^2 dx. \quad (3.2)$$

Let $u = \phi v$. By Lemma 2.1 we have

$$\int_{\Omega} |\nabla u|^2 dx = - \int_{\Omega} \frac{\Delta \phi}{\phi} u^2 dx + \int_{\Omega} \phi^2 |\nabla v|^2 dx.$$

We conclude the proof by combining this result with (3.2). \square

Condition (1.5) might seem to be unnatural and not easily checked. However, it will be very natural and is easily verified for our choices of ϕ .

To produce Hardy inequalities in the half-space with remainder terms also including the Sobolev term, we will need a weighted version of the Sobolev inequality.

Theorem 3.1. *Let $\sigma_1, \dots, \sigma_k$ be real numbers for some k with $1 \leq k \leq n$. Set $c_l := |\sigma_1 + \dots + \sigma_l + l - 1|$ for $1 \leq l \leq k$. Assume that*

$$c_l \neq 0 \quad \text{whenever} \quad \sigma_l \neq 0.$$

Then there exists a positive constant C such that for any $w \in C_0^\infty(\mathbb{R}_+^n)$,

$$\int_{\mathbb{R}_+^n} x_1^{\sigma_1} |\mathbf{X}_2|^{\sigma_2} \dots |\mathbf{X}_k|^{\sigma_k} |\nabla w| dx \geq C \left(\int_{\mathbb{R}_+^n} \left(x_1^{\sigma_1} |\mathbf{X}_2|^{\sigma_2} \dots |\mathbf{X}_k|^{\sigma_k} |w| \right)^{\frac{n}{n-1}} dx \right)^{\frac{n-1}{n}}, \quad (3.3)$$

and

$$\begin{aligned} & \int_{\mathbb{R}_+^n} x_1^{\frac{\sigma_1(n-2)}{(n-1)}} |\mathbf{X}_2|^{\frac{\sigma_2(n-2)}{(n-1)}} \dots |\mathbf{X}_k|^{\frac{\sigma_k(n-2)}{(n-1)}} |\nabla w|^2 dx \\ & \geq C \left(\int_{\mathbb{R}_+^n} \left(x_1^{\frac{\sigma_1(n-2)}{2(n-1)}} |\mathbf{X}_2|^{\frac{\sigma_2(n-2)}{2(n-1)}} \dots |\mathbf{X}_k|^{\frac{\sigma_k(n-2)}{2(n-1)}} |w| \right)^{\frac{2n}{n-2}} dx \right)^{\frac{n-2}{n}}. \end{aligned} \quad (3.4)$$

Proof. For $\Omega = \mathbb{R}_+^n$ we let $u = x_1^{\sigma_1} v$ in the Sobolev inequality (3.1) to get

$$C_n \int_{\mathbb{R}_+^n} x_1^{\frac{n\sigma_1}{n-1}} |v|^{\frac{n}{n-1}} dx \leq \left(\int_{\mathbb{R}_+^n} (|\sigma_1| x_1^{\sigma_1-1} |v| + x_1^{\sigma_1} |\nabla v|) dx \right)^{\frac{n}{n-1}}, \quad v \in C_0^\infty(\mathbb{R}_+^n).$$

Using the inequality

$$\left| \int_{\mathbb{R}_+^n} \operatorname{div} \mathbf{F} |v| dx \right| \leq \int_{\mathbb{R}_+^n} |\mathbf{F}| |\nabla v| dx \quad (3.5)$$

with the vector field $\mathbf{F} = (x_1^{\sigma_1}, 0, \dots, 0)$ one obtains

$$|\sigma_1| \int_{\mathbb{R}_+^n} x_1^{\sigma_1-1} |v| dx \leq \int_{\mathbb{R}_+^n} x_1^{\sigma_1} |\nabla v| dx$$

and hence

$$C_n \int_{\mathbb{R}_+^n} x_1^{\frac{n\sigma_1}{n-1}} |v|^{\frac{n}{n-1}} dx \leq \left(\int_{\mathbb{R}_+^n} x_1^{\sigma_1} |\nabla v| dx \right)^{\frac{n}{n-1}}, \quad v \in C_0^\infty(\mathbb{R}_+^n).$$

Now let $v = |\mathbf{X}_2|^{\sigma_2} w = (x_1^2 + x_2^2)^{\sigma_2/2} w$ in the above inequality. This gives

$$\begin{aligned} C_n \int_{\mathbb{R}_+^n} x_1^{\frac{n\sigma_1}{n-1}} (x_1^2 + x_2^2)^{\frac{n\sigma_2}{2(n-1)}} |w|^{\frac{n}{n-1}} dx \\ \leq \left(\int_{\mathbb{R}_+^n} x_1^{\sigma_1} (x_1^2 + x_2^2)^{\sigma_2/2} |\nabla w| dx + \int_{\mathbb{R}_+^n} |\sigma_2| x_1^{\sigma_1} (x_1^2 + x_2^2)^{\sigma_2/2-1/2} |w| dx \right)^{\frac{n}{n-1}}. \end{aligned}$$

Letting $\mathbf{F} = x_1^{\sigma_1} (x_1^2 + x_2^2)^{\sigma_2/2-1/2} \mathbf{X}_2$ in (3.5), we get

$$|\sigma_1 + \sigma_2 + 1| \int_{\mathbb{R}_+^n} x_1^{\sigma_1} (x_1^2 + x_2^2)^{\sigma_2/2-1/2} |w| dx \leq \int_{\mathbb{R}_+^n} x_1^{\sigma_1} (x_1^2 + x_2^2)^{\sigma_2/2} |\nabla w| dx. \quad (3.6)$$

Combining the previous two estimates we conclude that

$$c \int_{\mathbb{R}_+^n} x_1^{\frac{n\sigma_1}{n-1}} (x_1^2 + x_2^2)^{\frac{n\sigma_2}{2(n-1)}} |w|^{\frac{n}{n-1}} dx \leq \left(\int_{\mathbb{R}_+^n} x_1^{\sigma_1} (x_1^2 + x_2^2)^{\sigma_2/2} |\nabla w| dx \right)^{\frac{n}{n-1}}.$$

Note that, in case $\sigma_2 = 0$, we have the desired result immediately and we do not have to check whether the constant $\sigma_1 + \sigma_2 + 1$ is zero or not. We may repeat this procedure iteratively. In the l -th step we need the analogue of (3.6) which is

$$\begin{aligned} c_l \int_{\mathbb{R}_+^n} x_1^{\sigma_1} (x_1^2 + x_2^2)^{\sigma_2/2} \dots (x_1^2 + \dots + x_l^2)^{(\sigma_l-1)/2} |w| dx \\ \leq \int_{\mathbb{R}_+^n} x_1^{\sigma_1} (x_1^2 + x_2^2)^{\sigma_2/2} \dots (x_1^2 + \dots + x_l^2)^{\sigma_l/2} |\nabla w| dx \end{aligned}$$

for some positive constant c_l . This follows from (3.5) with

$$\mathbf{F} = x_1^{\sigma_1} (x_1^2 + x_2^2)^{\sigma_2/2} \dots (x_1^2 + \dots + x_l^2)^{(\sigma_l-1)/2} \mathbf{X}_l.$$

For this choice we get

$$c_l = |\sigma_1 + \dots + \sigma_l + (l - 1)|.$$

So our procedure works nicely in case $c_l \neq 0$ for those l such that $\sigma_l \neq 0$. This proves (3.3).

To show (3.4) we apply (3.3) to the function $w = |v|^\theta$. Trivial estimates give

$$\begin{aligned} C \int_{\Omega} x_1^{\frac{n\sigma_1}{n-1}} (x_1^2 + x_2^2)^{\frac{n\sigma_2}{2(n-1)}} \dots (x_1^2 + \dots + x_k^2)^{\frac{n\sigma_k}{2(n-1)}} |v|^{\frac{n\theta}{n-1}} dx \\ \leq \left(\theta \int_{\Omega} x_1^{\sigma_1} (x_1^2 + x_2^2)^{\sigma_2/2} \dots (x_1^2 + \dots + x_k^2)^{\sigma_k/2} |v|^{\theta-1} |\nabla v| dx \right)^{\frac{n}{n-1}}. \end{aligned}$$

We will now apply Hölder’s inequality to the right hand side. We want to do it in such a way that one of the factors becomes identical to the left hand side raised to some power. Therefore we need to choose θ so that

$$\frac{n\theta}{n-1} = 2\theta - 2 \Leftrightarrow \theta = \frac{2(n-1)}{n-2}.$$

Hölder’s inequality then immediately gives the result. \square

We are now ready to give the proof of Theorem B:

Proof of Theorem B. For $\phi > 0$ and $u = \phi v$, Lemma 2.1 gives us the inequality

$$\int_{\mathbb{R}_+^n} |\nabla u|^2 dx + \int_{\mathbb{R}_+^n} \frac{\Delta\phi}{\phi} |u|^2 dx \geq \int_{\mathbb{R}_+^n} \phi^2 |\nabla v|^2 dx. \quad (3.7)$$

We will choose

$$\begin{aligned} \phi(x) &= (x_1^{\sigma_1} \cdot (x_1^2 + x_2^2)^{\sigma_2/2} \dots (x_1^2 + \dots + x_n^2)^{\sigma_n/2})^{\frac{n-2}{2(n-1)}} \\ &= |\mathbf{X}_1|^{-\gamma_1} \dots |\mathbf{X}_n|^{-\gamma_n}, \end{aligned} \quad (3.8)$$

where

$$\gamma_1 = \alpha_1 - 1/2, \quad \gamma_m = \alpha_m - \alpha_{m-1} + 1/2, \quad m = 2, \dots, n,$$

and

$$\sigma_m = -\frac{2(n-1)}{n-2} \gamma_m, \quad m = 1, \dots, n.$$

We now apply (3.4) of Theorem 3.1 to obtain

$$\int_{\mathbb{R}_+^n} \phi^2 |\nabla v|^2 dx \geq C \left(\int_{\mathbb{R}_+^n} |\phi v|^{\frac{2n}{n-2}} dx \right)^{\frac{n-2}{n}},$$

provided that

$$c_l := |\sigma_1 + \dots + \sigma_l + l - 1| \neq 0 \quad \text{whenever} \quad \sigma_l \neq 0, \quad (3.9)$$

for $1 \leq l \leq n$. Combining this with (3.7) we get

$$\int_{\mathbb{R}_+^n} |\nabla u|^2 dx + \int_{\mathbb{R}_+^n} \frac{\Delta\phi}{\phi} |u|^2 dx \geq C \left(\int_{\mathbb{R}_+^n} |u|^{\frac{2n}{n-2}} dx \right)^{\frac{n-2}{n}}.$$

On the other hand, by Theorem A(i),

$$-\frac{\Delta\phi}{\phi} = \frac{\beta_1}{x_1^2} + \frac{\beta_2}{x_1^2 + x_2^2} + \dots + \frac{\beta_n}{x_1^2 + \dots + x_n^2},$$

and the desired inequality follows. It remains to check condition (3.9). After some elementary calculations we see that

$$c_l = \frac{2(n-1)}{n-2} \left| \alpha_l - \frac{n-l}{2(n-1)} \right|, \quad l = 1, \dots, n.$$

Since $\alpha_l \leq 0$ we clearly have $c_l \neq 0$ for $l = 1, \dots, n - 1$. Moreover, $c_n \neq 0$ when $\alpha_n < 0$. This completes the proof of (1.4).

In the rest of the proof we will show that (1.4) fails in case $\alpha_n = 0$. To this end we will establish that

$$\inf_{u \in C_0^\infty(\mathbb{R}_+^n)} \frac{\int_{\mathbb{R}_+^n} |\nabla u|^2 dx - \beta_1 \int_{\mathbb{R}_+^n} \frac{|u|^2}{x_1^2} dx - \dots - \beta_n \int_{\mathbb{R}_+^n} \frac{|u|^2}{x_1^2 + \dots + x_n^2} dx}{\left(\int_{\mathbb{R}_+^n} |u|^{\frac{2n}{n-2}} dx\right)^{\frac{n-2}{n}}} = 0, \quad (3.10)$$

where $\beta_n = (\alpha_{n-1} - 1/2)^2$. Let

$$u(x) = x_1^{-\gamma_1} |\mathbf{X}_2|^{-\gamma_2} \dots |\mathbf{X}_{n-1}|^{-\gamma_{n-1}} v(x).$$

A straightforward calculation, quite similar to the one leading to (2.15), shows that the infimum in (3.10) is the same as

$$\inf_{v \in C_0^\infty(\mathbb{R}_+^n)} \frac{\int_{\mathbb{R}_+^n} \prod_{j=1}^{n-1} |\mathbf{X}_j|^{-2\gamma_j} |\nabla v|^2 dx - \beta_n \int_{\mathbb{R}_+^n} \prod_{j=1}^{n-1} |\mathbf{X}_j|^{-2\gamma_j} |\mathbf{X}_n|^{-2} v^2 dx}{\left(\int_{\mathbb{R}_+^n} \left(\prod_{j=1}^{n-1} |\mathbf{X}_j|^{-\gamma_j}\right)^{\frac{2n}{n-2}} |v|^{\frac{2n}{n-2}} dx\right)^{\frac{n-2}{n}}}. \quad (3.11)$$

We now choose the test functions

$$v_{k_1, \varepsilon} = |\mathbf{X}_n|^{-\gamma_n + \varepsilon} h_{k_1}(x) \phi(x), \quad \varepsilon > 0, \quad (3.12)$$

where $h_{k_1}(x)$ and $\phi(x)$ are as in the first step of the proof of Theorem A(ii). For this choice, after straightforward calculations, quite similar to the ones used in the proof of Theorem A(ii), we obtain the following estimate for the numerator N in (3.11):

$$\begin{aligned} N[v_{\infty, \varepsilon}] &= ((\alpha_{n-1} - 1/2 + \varepsilon)^2 - (\alpha_{n-1} - 1/2)^2) \\ &\quad \times \int_{\mathbb{R}_+^n} \prod_{j=1}^{n-1} |\mathbf{X}_j|^{-2\gamma_j} |\mathbf{X}_n|^{-2\gamma_n + 2\varepsilon} \phi(x)^2 dx + O_\varepsilon(1), \\ &= C\varepsilon \int_{\mathbb{R}_+^n} r^{-1+2\varepsilon} \prod_{j=1}^n (\sin \theta_j)^{1-2\alpha_j} \phi(r)^2 d\theta_1 \dots d\theta_{n-1} dr + O_\varepsilon(1) \\ &= C\varepsilon \int_0^1 r^{-1+\varepsilon} dr + O_\varepsilon(1). \end{aligned}$$

In the above calculations we have taken the limit $k_1 \rightarrow \infty$ and we have used polar coordinates $(x_1, \dots, x_n) \rightarrow (\theta_1, \dots, \theta_{n-1}, r)$. We then conclude that

$$N[v_{\infty, \varepsilon}] < C \quad \text{as } \varepsilon \rightarrow 0. \quad (3.13)$$

Similar calculations for the denominator D in (3.11) reveal that

$$\begin{aligned} D[v_{\infty, \varepsilon}] &= C \left(\int_{\mathbb{R}_+^n} r^{-1 + \frac{2\varepsilon n}{n-2}} \prod_{j=1}^{n-1} (\sin \theta_j)^{\frac{n-j}{n-2} - \frac{2n\alpha_j}{n-2} - 1} \phi^{\frac{2n}{n-2}} d\theta_1 \dots d\theta_{n-1} dr \right)^{\frac{n-2}{n}} \\ &\geq C \left(\int_0^{1/2} r^{-1 + \frac{2\varepsilon n}{n-2}} dr \right)^{\frac{n-2}{n}} = C\varepsilon^{-\frac{n-2}{n}}. \end{aligned}$$

We then have

$$\frac{N[v_{\infty,\varepsilon}]}{D[v_{\infty,\varepsilon}]} \rightarrow 0 \quad \text{as } \varepsilon \rightarrow 0,$$

and therefore the infimum in (3.11) or (3.10) is equal to zero. This completes the proof of the theorem. \square

Here is a consequence of Theorem B.

Corollary 3.2. *Let $1 \leq k < n$. For any $\beta_n < 1/4$, there exists a positive constant C such that for all $u \in C_0^\infty(\mathbb{R}_+^n)$,*

$$\int_{\mathbb{R}_+^n} |\nabla u|^2 dx \geq \int_{\mathbb{R}_+^n} \left(\frac{k^2}{4} \frac{1}{x_1^2 + \dots + x_k^2} + \frac{1}{4} \frac{1}{x_1^2 + \dots + x_{k+1}^2} + \dots \right. \\ \left. + \frac{1}{4} \frac{1}{x_1^2 + \dots + x_{n-1}^2} + \frac{\beta_n}{x_1^2 + \dots + x_n^2} \right) |u|^2 dx + C \left(\int_{\mathbb{R}_+^n} |u|^{2^*} dx \right)^{2/2^*}.$$

If $\beta_n = 1/4$ the previous inequality fails.

In case $k = n$ for any $\beta_n < n^2/4$, there exists a positive constant C such that for all $u \in C_0^\infty(\mathbb{R}_+^n)$,

$$\int_{\mathbb{R}_+^n} |\nabla u|^2 dx \geq \beta_n \int_{\mathbb{R}_+^n} \frac{|u|^2}{x_1^2 + \dots + x_n^2} dx + C \left(\int_{\mathbb{R}_+^n} |u|^{2^*} dx \right)^{2/2^*}.$$

The above inequality fails for $\beta_n = n^2/4$.

Proof. In Theorem B we make the following choices: In the case $k = 1$ we choose $\alpha_1 = \dots = \alpha_{n-1} = 0$. In this case $\beta_k = 1/4$, $k = 1, \dots, n-1$. The condition $\alpha_n < 0$ is equivalent to $\beta_n < 1/4$.

In the case $1 < k \leq n-1$ we choose $\alpha_m = -m/2$ when $m = 1, \dots, k-1$, and $\alpha_m = 0$ when $m = k, \dots, n-1$. Finally, in case $k = n$ we choose $\alpha_m = -m/2$ for $m = 1, \dots, n-1$. \square

4. Further generalizations

The techniques used in the previous sections can be generalized to other situations as well. For example, consider the subset of \mathbb{R}^n where $x_1, \dots, x_k > 0$. We denote this domain by \mathbb{R}_{k+}^n . Then we can easily prove the Hardy-Sobolev inequality

Theorem 4.1. *There exists a positive constant C such that for any $u \in C_0^\infty(\mathbb{R}_{k+}^n)$*

$$\int_{\mathbb{R}_{k+}^n} |\nabla u|^2 dx \geq \frac{1}{4} \int_{\mathbb{R}_{k+}^n} \left(\frac{1}{x_1^2} + \dots + \frac{1}{x_k^2} \right) |u|^2 dx + C \left(\int_{\mathbb{R}_{k+}^n} |u|^{2^*} dx \right)^{2/2^*}.$$

Proof. Let $\phi = \sqrt{x_1 \dots x_k}$. For $u = \phi w$ we calculate

$$\begin{aligned} \int_{\mathbb{R}_{k+}^n} |\nabla u|^2 dx &= \int_{\mathbb{R}_{k+}^n} \left| \sqrt{x_1 \dots x_k} \cdot \nabla w + \frac{1}{2} \sqrt{x_1 \dots x_k} \left(\frac{1}{x_1}, \dots, \frac{1}{x_k} \right) w \right|^2 dx \\ &= \int_{\mathbb{R}_{k+}^n} x_1 \dots x_k |\nabla w|^2 dx + \frac{1}{4} \int_{\mathbb{R}_{k+}^n} x_1 \dots x_k \left(\frac{1}{x_1^2} + \dots + \frac{1}{x_k^2} \right) |w|^2 dx \\ &\quad + \frac{1}{2} \int_{\mathbb{R}_{k+}^n} x_1 \dots x_k \left(\frac{1}{x_1}, \dots, \frac{1}{x_k} \right) \nabla w^2 dx. \end{aligned}$$

By partial integration, we see that the last term is zero. If the second term is expressed in terms of u , we see that it is equal to the Hardy term

$$\frac{1}{4} \int_{\mathbb{R}_{k+}^n} \left(\frac{1}{x_1^2} + \dots + \frac{1}{x_k^2} \right) |u|^2 dx.$$

By Theorem C, the first term may be estimated from below by the Sobolev term provided that we can prove the following L^1 Hardy inequality:

$$C \int_{\mathbb{R}_{k+}^n} (x_1 \dots x_k)^{\frac{n-1}{n-2}} \left(\frac{1}{x_1^2} + \dots + \frac{1}{x_k^2} \right)^{1/2} |v| dx \leq \int_{\mathbb{R}_{k+}^n} (x_1 \dots x_k)^{\frac{n-1}{n-2}} |\nabla v| dx.$$

To do this we work as in the previous section, using the inequality

$$\left| \int_{\mathbb{R}_{k+}^n} \operatorname{div} \mathbf{F} |v| dx \right| \leq \int_{\mathbb{R}_{k+}^n} |\mathbf{F}| |\nabla v| dx,$$

with the proper choice of vector field, which turns out to be

$$\mathbf{F} = (x_1 \dots x_k)^\tau \left(\frac{1}{x_1^2} + \dots + \frac{1}{x_k^2} \right)^\beta \left(\frac{1}{x_1}, \dots, \frac{1}{x_k} \right),$$

where

$$\tau = \frac{n-1}{n-2} \quad \text{and} \quad \beta = -\frac{1}{2}.$$

We immediately see that $|\mathbf{F}| = \phi^{2\tau} = (x_1 \dots x_k)^{\frac{n-1}{n-2}}$. Also,

$$\begin{aligned} \operatorname{div} \mathbf{F} &= -(x_1 \dots x_k)^\tau \left(\frac{1}{x_1^2} + \dots + \frac{1}{x_k^2} \right)^{\beta+1} \\ &\quad + \tau (x_1 \dots x_k)^\tau \left(\frac{1}{x_1^2} + \dots + \frac{1}{x_k^2} \right)^{\beta+1} \\ &\quad - 2\beta \left(\frac{1}{x_1^4} + \dots + \frac{1}{x_k^4} \right) \cdot (x_1 \dots x_k)^\tau \left(\frac{1}{x_1^2} + \dots + \frac{1}{x_k^2} \right)^{\beta-1}. \end{aligned}$$

Since $\tau - 1 > 0$ and the last term is positive, we get the result. □

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