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# Bistable traveling waves for monotone semiflows with applications

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Abstract. This paper is devoted to the study of traveling waves for monotone evolution systems of bistable type. In an abstract setting, we establish the existence of traveling waves for discrete time and continuous-time monotone semiflows in homogeneous and periodic habitats. The results are then extended to monotone semiflows with weak compactness. We also apply the theory to four classes of evolution systems.

Keywords. Monotone semiflows, traveling waves, bistable dynamics, periodic habitat

## 1. Introduction

In this paper, we study traveling waves for monotone (i.e., order preserving) semiflows  ${Q_t}_{t \in \mathcal{T}}$  with the bistability structure on some subsets of the space  $\mathcal{C} := C(\mathcal{H}, \mathcal{X})$  consisting of all continuous functions from the habitat  $\mathcal{H} (= \mathbb{R}$  or  $\mathbb{Z})$  to the Banach lattice X, where  $\mathcal{T} = \mathbb{Z}^+$  or  $\mathbb{R}^+$  is the set of evolution times. Here the bistability structure is generalized from a number of studies for various evolution equations. It means that the restricted semiflow on  $X$  admits two ordered stable equilibria, between which all others are unstable. We focus on the existence of traveling waves connecting these two stable equilibria, which we call bistable (traveling) waves. This setting allows us to study not only autonomous and time-periodic evolution systems in a homogeneous habitat (medium), but also those in a periodic habitat. Moreover, the results obtained can be extended to semiflows with weak compactness on some subsets of the space  $M$  of all monotone functions from  $\mathbb R$  to  $\mathcal X$ .

To explain the concept of bistability structure, we recall some related works on typical evolution equations. Fife and McLeod [\[19,](#page-44-0) [20\]](#page-44-1) proved the existence and global asymptotic stability of monotone traveling waves for the reaction-diffusion equation

<span id="page-0-1"></span>
$$
u_t = u_{xx} + u(1 - u)(u - a), \quad x \in \mathbb{R}, t > 0,
$$
\n(1.1)

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where  $a \in (0, 1)$ . Clearly, the restriction of system  $(1.1)$  to  $\mathcal{X} = \mathbb{R}$  is the ordinary differential equation  $u' = u(1-u)(u-a)$ , which admits a unique unstable equilibrium between two ordered and stable ones. The same property is shared by the nonlocal dispersal equation in [\[4,](#page-43-0) [16,](#page-44-2) [46\]](#page-45-1) and the lattice equations in [\[3,](#page-43-1) [49,](#page-45-2) [50\]](#page-45-3). Chen [\[13\]](#page-44-3) studied a general nonlocal evolution equation  $u_t = A(u(\cdot, t))$ , which also possesses the above bistability structure. Some related investigations on discrete-time equations can be found in [\[27,](#page-44-4) [15\]](#page-44-5). For the time-periodic reaction-diffusion equation  $u_t = u_{xx} + f(t, u)$ , the spatially homogeneous system is a time-periodic ordinary differential equation. In this case, the equilibrium in the bistability structure should be understood as the time-periodic solution. Under such bistability assumption, Alikakos, Bates and Chen [\[1\]](#page-43-2) obtained the existence of bistable time-periodic traveling waves. Recently, Yagisita [\[46\]](#page-45-1) studied bistable traveling waves for discrete-time and continuous-time semiflows on the space consisting of all left-continuous and nondecreasing functions from  $\mathbb R$  to  $\mathcal X = \mathbb R$  under the assumption that there is exactly one intermediate unstable equilibrium. It should be mentioned that the result in [\[46\]](#page-45-1) for continuous-time semiflows requires an assumption on the existence of a pair of upper and lower solutions.

Note that the restrictions to  $\mathcal{X} = \mathbb{R}$  of the aforementioned systems are all scalar equations, and hence there is only one unstable equilibrium in between two stable ones. But in the case where  $\mathcal{X} = \mathbb{R}^n$ , there may be multiple unstable equilibria. This is one of the main reasons why some ideas and techniques developed for scalar equations cannot be easily extended to higher dimensional systems. Volpert [\[36\]](#page-45-4) established the existence and stability of traveling waves for the bistable reaction-diffusion system  $u_t = D\Delta u + f(u)$ by using topological methods, where  $D$  is a positive definite diagonal matrix. Fang and Zhao  $[18]$  further extended these results to the case where D is semi-positive definite via the vanishing viscosity approach.

Consider the following parabolic equation in a cylindrical domain  $\Sigma = \mathbb{R} \times \Omega$ :

$$
\begin{cases} u_t = \Delta u + \alpha(y)u_x + f(u), & x \in \mathbb{R}, y = (y_1, \dots, y_{n-1}) \in \Omega, t > 0, \\ \partial u/\partial v = 0 & \text{on } \mathbb{R} \times \partial \Omega \times (0, \infty), \end{cases}
$$
(1.2)

where f is of the same type as the nonlinearity in  $(1.1)$  and  $\Omega$  is a bounded domain with smooth boundary in  $\mathbb{R}^{n-1}$ . Obviously, the restriction of the solution semiflow of [\(1.2\)](#page-1-0) to  $\mathcal{X} = C(\overline{\Omega}, \mathbb{R})$  gives rise to the following x-independent system:

<span id="page-1-1"></span><span id="page-1-0"></span>
$$
\begin{cases} u_t = \Delta_y u + f(u), & y \in \Omega, \ t > 0, \\ \frac{\partial u}{\partial v} = 0 & \text{on } \frac{\partial \Omega \times (0, \infty)}{\partial u}. \end{cases}
$$
(1.3)

One can see from Matano [\[30\]](#page-44-7) (or Casten and Holland [\[12\]](#page-43-3)) that any nonconstant steady state of [\(1.3\)](#page-1-1) is linearly unstable when the domain  $\Omega$  is convex. It follows that if  $\Omega$ is convex, then  $(1.2)$  admits the bistability structure: its x-independent system has two (constant) linearly stable steady states, between which all others are linearly unstable. In such a case, Berestycki and Nirenberg [\[11\]](#page-43-4) obtained the existence and uniqueness of bistable traveling waves. In the case where  $\Omega$  is an appropriate dumbbell-shaped domain, Matano [\[30\]](#page-44-7) constructed a counterexample to show that [\(1.3\)](#page-1-1) has stable nonconstant steady states, and Berestycki and Hamel [\[6\]](#page-43-5) also proved the nonexistence of traveling waves connecting two stable constant steady states. For bistable traveling waves in timedelayed reaction-diffusion equations, we refer to [\[31,](#page-44-8) [35,](#page-45-5) [29,](#page-44-9) [37\]](#page-45-6). For such an equation with time delay  $\tau > 0$ , one can choose  $\mathcal{X} = C([- \tau, 0], \mathbb{R})$  so that its solution semiflow has the bistability structure.

Recently, there is an increasing interest in reaction-diffusion equations in periodic habitats. A typical example is

<span id="page-2-0"></span>
$$
u_t = (du_x)_x + f(u), \quad x \in \mathbb{R}, \ t > 0,
$$
\n(1.4)

where  $d \in C^1(\mathbb{R}, \mathbb{R})$  is a positive periodic function with period  $r > 0$ . Define  $\mathcal{Y} :=$  $C([0, r], \mathbb{R})$  and  $C_{per}(\mathbb{R}, \mathbb{R}) := \{g \in C(\mathbb{R}, \mathbb{R}) : g(x+r) = g(x), \forall x \in \mathbb{R}\}.$  It is easy to see that

$$
C(\mathbb{R}, \mathbb{R}) = \{ g \in C(r\mathbb{Z}, \mathcal{Y}) : g(ri)(r) = g(r(i+1))(0), \forall i \in \mathbb{Z} \} =: \mathcal{K},
$$

and that any element in  $C_{\text{per}}(\mathbb{R}, \mathbb{R})$  is a constant function in K. Thus, the solution semi-flow of [\(1.4\)](#page-2-0) on  $C(\mathbb{R}, \mathbb{R})$  can be regarded as a conjugate semiflow on K, and hence the bistability structure should be understood as: the restriction of the solution semiflow of [\(1.4\)](#page-2-0) on  $C_{\text{per}}(\mathbb{R}, \mathbb{R})$  has two ordered *r*-periodic steady states, between which all others are unstable. Assuming that the function  $f$  is of bistable type, Xin [\[43\]](#page-45-7) obtained the existence of a spatially periodic (pulsating) traveling wave as long as  $d$  is sufficiently close to a positive constant in a certain sense (see also  $[44, 42]$  $[44, 42]$  $[44, 42]$ ). However, whether the solution semiflow of  $(1.4)$  admits the bistability structure has remained an open problem. We will give an affirmative answer in Section 6.3 and further improve Xin's existence result. Meanwhile, a counterexample will be constructed to show that the solution semiflow of [\(1.4\)](#page-2-0) has no bistability structure in the general case of varying  $d(x)$ . More recently, Chen, Guo and Wu [\[14\]](#page-44-10) proved the existence, uniqueness and stability of spatially periodic traveling waves for one-dimensional lattice equations in a periodic habitat under the bistability assumption. There are also other types of bistable waves (see, e.g., [\[7,](#page-43-6) [33\]](#page-45-10)). For monostable systems in periodic habitats, we refer to [\[5,](#page-43-7) [8,](#page-43-8) [9,](#page-43-9) [22,](#page-44-11) [23,](#page-44-12) [26\]](#page-44-13) and references therein.

In general, there are multiple intermediate unstable equilibria in between two stable ones in the case where the space  $\mathcal X$  is higher dimensional. Meanwhile, it is possible for the given autonomous system to have intermediate unstable time-periodic orbits in  $\mathcal{X}$ . These make the study of bistable semiflows more difficult than that of monostable ones, whose restrictions to  $X$  have only one unstable and one stable equilibrium. To overcome these difficulties, we will show that all these unstable equilibria and all points in these periodic orbits are unordered in  $\mathcal X$  under some appropriate assumptions. With this in mind, a bistable system can be regarded as the union of two monostable systems, although such a union is not unique. From this point of view, we establish a link between monostable subsystems and the bistable system itself, which plays a vital role in the propagation of bistable traveling waves. This link is stated in terms of spreading speeds of monostable subsystems (see assumption (A6)). For spreading speeds of various monostable evolution systems, we refer to [\[2,](#page-43-10) [8,](#page-43-8) [22,](#page-44-11) [25,](#page-44-14) [26,](#page-44-13) [28,](#page-44-15) [38,](#page-45-11) [39,](#page-45-12) [48\]](#page-45-13) and references therein.

Next we use a diagram to outline the proofs. For case (I), we combine the above observations for general bistable semiflows and Yagisita's [\[46\]](#page-45-1) perturbation idea to prove



Fig. 1. Scheme of the proof.

the existence of traveling waves. For case (III), we use the candidates  $\phi_{\pm}(x + c_{\pm, s}n)$ for bistable waves of the discrete-time semiflows  ${((Q_s)^n)_{n \geq 0}}$  to approximate that of the continuous-time semiflow  ${Q_t}_{t\geq0}$ . This new approach heavily relies on an estimation of the boundedness of  $(1/s)c_{\pm,s}$  as  $s \to 0$ , which is proved surprisingly by using the bistability structure of the semiflow (see inequalities  $(3.9)$  and  $(3.10)$ ). It turns out that our result does not require the assumption on the existence of a pair of upper and lower solu-tions as in [\[46\]](#page-45-1). In case (II), both the evolution time  $\mathcal T$  and the habitat  $\mathcal H$  are discrete, a traveling wave  $\psi(i + cn)$  of  ${Q<sup>n</sup>}_{n \geq 0}$  cannot be well defined in the usual way because the wave speed c, and hence the domain of  $\psi$ , is unknown. So we define it to be a traveling wave on R of an associated map  $\ddot{Q}$ . However,  $\ddot{Q}$  has much weaker compactness than Q. To overcome this difficulty, we establish a variant of Helly's theorem for monotone functions from  $\mathbb R$  to  $\mathcal X$  in the Appendix, which is of independent interest. This discovery also enables us to study monotone semiflows in a periodic habitat and with weak compactness. Further, we can deal with case (IV) by a similar idea to that in case (III) because now traveling waves in case  $(II)$  are defined on  $\mathbb R$ . Traveling waves for a time-periodic system can be obtained with the help of the discrete-time semiflow generated by the associated Poincaré map. Motivated by the discussions in  $[26, Section 5]$  $[26, Section 5]$  $[26, Section 5]$ , we can regard a semiflow in a periodic habitat as a conjugate semiflow in a homogeneous discrete habitat, and hence we can employ the arguments for cases (II) and (IV) to establish the existence of spatially periodic bistable traveling waves.

The rest of this paper is organized as follows. In Section 2, we present our main assumptions as well as their explanations. In Section 3, we deal with discrete-time, continuous-time, and time-periodic compact semiflows on some subsets of  $C$ . In Section 4, we extend our results to compact semiflows in a periodic habitat. In Section 5, we further investigate semiflows with weak compactness. In Section 6, we apply the abstract results to four classes of evolution systems: a time-periodic reaction-diffusion system, a parabolic system in a cylinder, a parabolic equation with periodic diffusion, and a time-delayed reaction-diffusion equation. Finally, a short appendix section completes the paper.

#### 2. Notation and assumptions

Throughout this paper, we assume that X is an ordered Banach space with the norm  $\|\cdot\|_{\mathcal{X}}$ and the cone  $\mathcal{X}^+$ . Further, we assume that  $\mathcal X$  is also a vector lattice with the following

monotonicity condition:

$$
|x|\chi \le |y|\chi \implies ||x||\chi \le ||y||\chi,
$$

where  $|z|\chi := \sup\{z, -z\}$  denotes the least upper bound of z and  $-z$ . Such a Banach space is often called a *Banach lattice*. We use  $C(M, \mathbb{R}^d)$  to denote the set of all continuous functions from the compact metric space M to the d-dimensional Euclidean space  $\mathbb{R}^d$ . We equip  $C(M, \mathbb{R}^d)$  with the maximum norm and the standard positive cone consisting of all nonnegative functions. Then  $C(M, \mathbb{R}^d)$  is a special Banach lattice, which will be used in this paper. For more general information about Banach lattices, we refer to the book [\[32\]](#page-44-16).

Let the spatial habitat  $H$  be the real line  $\mathbb R$  or the lattice

$$
r\mathbb{Z}:=\{\ldots,-2r,-r,0,r,2r,\ldots\}
$$

for some positive number r. For simplicity, we let  $r = 1$ . We say a function  $\phi : \mathcal{H} \to \mathcal{X}$ is bounded if the set  $\{\|\phi(x)\|_{\mathcal{X}} : x \in \mathcal{H}\}$  is bounded. Throughout this paper, we always use B to denote the set of all bounded functions from  $\mathbb R$  to X, and C to denote the set of all bounded and continuous functions from H to X. Moreover, any element in X can be regarded as a constant function in  $\beta$  and  $\beta$ .

In this paper, we equip C with the compact-open topology, that is, a sequence  $\phi_n$ converges to  $\phi$  in C if and only if  $\phi_n(x)$  converges to  $\phi(x)$  in X uniformly for x in any compact subset of  $H$ . The following norm on  $C$  induces this topology:

<span id="page-4-2"></span>
$$
\|\phi\|_{\mathcal{C}} = \sum_{k=1}^{\infty} \frac{\max_{|x| \le k} \|\phi(x)\|_{\mathcal{X}}}{2^k}, \quad \forall \phi \in \mathcal{C}.
$$
 (2.1)

Clearly, if  $\mathcal{H} = \mathbb{Z}$ , then  $\phi_n \to \phi$  with respect to the compact-open topology if and only if  $\phi_n(x) \to \phi(x)$  pointwise in  $x \in \mathbb{Z}$ .

We assume that Int( $\mathcal{X}^+$ ) is not empty. For any  $u, v \in \mathcal{X}$ , we write  $u \geq v$  provided that  $u-v \in \mathcal{X}^+, u > v$  provided  $u \geq v$  but  $u \neq v$ , and  $u \gg v$  provided that  $u-v \in \text{Int}(\mathcal{X}^+).$ A subset in  $X$  is said to be *totally unordered* if no two elements are ordered. For any  $\phi, \psi \in \mathcal{C}, \phi \geq \psi$  provided that  $\phi(x) \geq \psi(x)$  for all  $x \in \mathcal{H}, \phi > \psi$  provided that  $\phi \geq \psi$ but  $\phi \neq \psi$ , and  $\phi \gg \psi$  provided that  $\phi(x) \gg \psi(x)$  for all  $x \in \mathcal{H}$ . For any  $\gamma \in \mathcal{X}$ with  $\gamma > 0$ , we define  $\mathcal{X}_{\gamma} := \{u \in \mathcal{X} : \gamma \ge u \ge 0\}$ ,  $\mathcal{C}_{\gamma} := \{\phi \in \mathcal{C} : \gamma \ge \phi \ge 0\}$ and  $\mathcal{B}_{\gamma} := \{\phi \in \mathcal{B} : \gamma \ge \phi \ge 0\}$ . For any  $\phi, \psi \in \mathcal{C}$ , the interval  $[\phi, \psi]_{\mathcal{C}}$  is the set  ${w \in \mathcal{C} : \phi \leq w \leq \psi}, [[\phi, \psi]]_C$  is the set  ${w \in \mathcal{C} : \phi \ll w \ll \psi},$  and similarly for  $[\phi, \psi]]_C$  and  $[[\phi, \psi]_C$ . For any  $u \le v$  in X, we define  $[u, v]_{\mathcal{X}}$ ,  $[[u, v]_{\mathcal{X}}$ ,  $[[u, v]_{\mathcal{X}}$  and  $[u, v]]_{\mathcal{X}}$  in a similar way.

Let  $\beta \in \text{Int}(\mathcal{X}^+)$  and Q be a map from  $\mathcal{C}_{\beta}$  to  $\mathcal{C}_{\beta}$ . Let E be the set of all fixed points of Q restricted to  $\mathcal{X}_{\beta}$ .

<span id="page-4-1"></span>**Definition 2.1.** For the map  $Q : \mathcal{X}_{\beta} \to \mathcal{X}_{\beta}$ , a fixed point  $\alpha \in E$  is said to be *strongly stable from below* if there exist a number  $\delta > 0$  and a unit vector  $e \in Int(\mathcal{X}^+)$  such that

<span id="page-4-0"></span>
$$
Q[\alpha - \eta e] \gg \alpha - \eta e \quad \text{for any } \eta \in (0, \delta]. \tag{2.2}
$$

Strong instability from below is defined by reversing the inequality in [\(2.2\)](#page-4-0). Similarly, we can define strong stability (and instability) from above.

Given  $y \in H$ , define the translation operator  $T_y$  on B by  $T_y[\phi](x) = \phi(x - y)$ . Assume that 0 and  $\beta$  are in E. We impose the following hypotheses on Q:

- (A1) (*Translation invariance*)  $T_v \circ Q[\phi] = Q \circ T_v[\phi]$  for all  $\phi \in C_\beta$  and  $y \in \mathcal{H}$ .
- (A2) (*Continuity*)  $Q : C_\beta \to C_\beta$  is continuous in the compact-open topology.
- (A3) (*Monotonicity*) Q is order preserving in the sense that  $Q[\phi] \ge Q[\psi]$  whenever  $\phi \geq \psi$  in  $\mathcal{C}_{\beta}$ .
- (A4) (*Compactness*)  $Q : C_\beta \to C_\beta$  is compact in the compact-open topology.
- (A5) (*Bistability*) The fixed points 0 and  $\beta$  are strongly stable from above and below, respectively, for the map  $Q : \mathcal{X}_{\beta} \to \mathcal{X}_{\beta}$ , and the set  $E \setminus \{0, \beta\}$  is totally unordered.

Note that the above bistability assumption is imposed on the spatially homogeneous map  $Q: \mathcal{X}_{\beta} \to \mathcal{X}_{\beta}$ . We allow the existence of other fixed points on the boundary of  $\mathcal{X}_{\beta}$  so that the theory is applicable to competitive evolution models. The unordering property of  $E \setminus \{0, \beta\}$  can be obtained by the strong instability of all fixed points in this set if the semiflow is eventually strongly monotone. More precisely, a sufficient condition for hypothesis  $(A5)$  to hold is:

(A5<sup>'</sup>) (*Bistability*)  $Q: \mathcal{X}_{\beta} \to \mathcal{X}_{\beta}$  is eventually strongly monotone in the sense that there exists  $m_1 \in \mathbb{Z}_+$  such that  $Q^m[u] \gg Q^m[v]$  for all  $m \ge m_1$  whenever  $u > v$  in  $\mathcal{X}_{\beta}$ . Further, for the map  $Q : \mathcal{X}_{\beta} \to \mathcal{X}_{\beta}$ , the fixed points 0 and  $\beta$  are strongly stable from above and below, respectively, and each  $\alpha \in E \setminus \{0, \beta\}$  is strongly unstable from both below and above.

Figure [2](#page-5-0) illustrates the bistability structures in  $(A5)$  and  $(A5')$ . Next we show that  $(A5')$ implies (A5). In applications, however, one may find other sufficient conditions for (A5) to hold, weaker than  $(A5')$ .



<span id="page-5-0"></span>Fig. 2. Left: the set E satisfying (A5). Right: the set E satisfying  $(A5')$ .

<span id="page-5-1"></span>**Proposition 2.1.** *If* (A5') *holds, then for any*  $\alpha_1, \alpha_2 \in E \setminus \{0, \beta\}$ *, we have*  $\alpha_1 \nless \alpha_2$  *and*  $\alpha_2 \neq \alpha_1$ .

*Proof.* Without loss of generality, we only show  $\alpha_1 \nless \alpha_2$ . Assume for contradiction that  $\alpha_1 < \alpha_2$ . Then  $\alpha_1 = Q^{m_1}[\alpha_1] \ll Q^{m_1}[\alpha_2] = \alpha_2$ . Since  $\alpha_1$  is strongly unstable from above, there exist  $\delta_{\alpha_1} > 0$  and  $e_{\alpha_1} \in \text{Int}(\mathcal{X}^+)$  such that  $u_0 := \alpha_1 + \delta_{\alpha_1} e_{\alpha_1} \in [[\alpha_1, \alpha_2]]$ and  $Q[u_0] \gg u_0$ . Define the recursion  $u_{n+1} = Q[u_n]$ ,  $n \ge 0$ . Then  $u_n$  is convergent to some  $\alpha \in \mathcal{X}$  with  $\alpha_1 \ll \alpha \le \alpha_2$  due to hypothesis (A4). By the eventual strong monotonicity of  $Q$ , we see that

$$
u_n = Q^{m_1}[u_{n-m_1}] \ll Q^{m_1}[u_{n+1-m_1}] = u_{n+1} \ll Q^{m_1}[\alpha] = \alpha, \quad \forall n \ge m_1.
$$

Since  $\alpha$  is strongly unstable from below, we can find  $\delta_{\alpha} > 0$  and  $e_{\alpha} \in Int(\mathcal{X}^+)$  such that  $Q[\alpha - \delta e_{\alpha}] \ll \alpha - \delta e_{\alpha}$  for all  $\delta \in (0, \delta_{\alpha}]$ . Choose  $n_1 \geq m_1$  such that  $u_{n_1} \geq \alpha - \delta_{\alpha} e_{\alpha}$ . Define  $\eta := \sup \{ \delta \in (0, \delta_{\alpha}] : u_{n_1} \leq \alpha - \delta e_{\alpha} \}.$  Thus,  $u_{n_1} \nless \alpha - \eta \delta_{\alpha}$ . On the other hand,

$$
u_{n_1}\ll u_{n_1+1}=Q[u_{n_1}]\leq Q[\alpha-\eta e_\alpha]\ll \alpha-\eta e_\alpha,
$$

a contradiction.  $\Box$ 

Due to assumption (A5), for any given  $\alpha \in E \setminus \{0, \beta\}$ , we have two monostable subsystems:  ${Q^n}_{n \geq 0}$  restricted to  $[0, \alpha]_C$  and  $[\alpha, \beta]_C$ , respectively. With this in mind, we next construct an initial function  $\phi_{\alpha}^-$  so that we can define the leftward asymptotic speed of propagation of  $\phi_{\alpha}^-$ , and then present our last assumption.

Note that in (A5) we do not require  $\alpha \gg 0$  or  $\alpha \ll \beta$  in X. But assumption (A5) is sufficient to guarantee that  $\alpha$  and  $\beta$  can be separated by two neighborhoods in [ $\alpha$ ,  $\beta$ ] $\chi$ , and a similar claim is valid for 0 and  $\alpha$  (see Lemma 3.1). In view of assumption (A5), we can find  $\delta_{\beta} > 0$  and a unit vector  $e_{\beta} \in \text{Int}(\mathcal{X}^+)$  such that

$$
Q[\beta - \eta e_{\beta}] \gg \beta - \eta e_{\beta}, \quad \forall \eta \in (0, \delta_{\beta}].
$$
 (2.3)

<span id="page-6-0"></span>Define

<span id="page-6-2"></span><span id="page-6-1"></span>
$$
v_{\alpha}^- := \sup{\alpha, \beta - \delta_{\beta}e_{\beta}}.
$$
 (2.4)

Hence,  $Q[v_\alpha^-] \geq Q[\beta - \delta_\beta e_\beta] \gg \beta - \delta_\beta e_\beta$ . This, together with the definition of  $v_\alpha^-$  in [\(2.4\)](#page-6-0), implies that there exists a neighborhood  $\mathcal N$  of  $Q[v_\alpha^-]$  in  $[\alpha, \beta]_{\mathcal X}$  such that  $v_\alpha^- \leq \gamma$ for all  $\gamma \in \mathcal{N}$ . Choose a nondecreasing initial function  $\phi_{\alpha}^{-} \in C_{\beta}$  with

$$
\phi_{\alpha}^{-}(x) = \alpha, \ \forall x \le -1, \quad \text{and} \quad \phi_{\alpha}^{-}(x) = v_{\alpha}^{-}, \ \forall x \ge 0. \tag{2.5}
$$

It then follows from assumptions  $(A1)$ ,  $(A2)$  and  $(A5)$  that

$$
\lim_{x \to \infty} Q[\phi_{\alpha}^{-}](x) = Q[\phi_{\alpha}^{-}(\infty)](0) = Q[v_{\alpha}^{-}],
$$

and hence there exists  $\sigma > 0$  such that

$$
Q[\phi_{\alpha}^{-}](x) \gg \beta - \delta_{\beta}e_{\beta}
$$
 and  $Q[\phi_{\alpha}^{-}](x) \ge v_{\alpha}^{-}$ ,  $\forall x \ge \sigma - 1$ .

Define a sequence  $a_{n,\sigma}$  of points in X as follows:

$$
a_{n,\sigma}=Q^n[\phi_\alpha^-](\sigma n), \quad n\geq 1.
$$

Then

$$
a_{2,\sigma} = Q^2[\phi_\alpha^-](2\sigma) = Q[Q[\phi_\alpha^-](\cdot + \sigma)](\sigma) \ge Q[\phi_\alpha^-](\sigma) = a_{1,\sigma}.
$$

By induction, we see that  $a_{n,\sigma}$  is nondecreasing in n. Thus, assumption (A4) implies that  $a_{n,\sigma}$  tends to a fixed point e with  $e \ge a_{1,\sigma} \gg \beta - \delta_\beta e_\beta$ . Therefore,  $e = \beta$ .

By the above observation, we have

$$
\beta \geq \lim_{n \to \infty, x \geq \sigma_n} Q^n[\phi_\alpha^{-1}(x) \geq \lim_{n \to \infty} Q^n[\phi_\alpha^{-1}(\sigma n) = \lim_{n \to \infty} a_{n,\sigma} = \beta,
$$

and hence

$$
(-\infty, -\sigma] \subset \Lambda(\phi_{\alpha}^{-}):=\Big\{c\in\mathbb{R}:\lim_{n\to\infty, x\geq -cn}\mathcal{Q}^{n}[\phi_{\alpha}^{-}](x)=\beta\Big\}.
$$

<span id="page-7-1"></span>Define

$$
c_{-}^{*}(\alpha, \beta) := \sup \Lambda(\phi_{\alpha}^{-}).
$$
\n(2.6)

Clearly,  $c^*_{-}(\alpha, \beta) \in [-\sigma, \infty]$  and  $(-\infty, c^*_{-}(\alpha, \beta)) \subset \Lambda(\phi_{\alpha}^{-})$ . We further claim that  $c_{-}^{*}(\alpha, \beta)$  is independent of the choice of  $\phi_{\alpha}^{-}$  as long as  $\phi_{\alpha}^{-}$  has the property [\(2.5\)](#page-6-1). Indeed, for any given  $\phi$  with the property [\(2.5\)](#page-6-1), we have

$$
\phi_{\alpha}^-(x-1) \le \phi(x) \le \phi_{\alpha}^-(x+1), \quad \forall x \in \mathcal{H}.
$$

It then follows that for any  $c \in \Lambda(\phi_{\alpha}^-)$  and  $\epsilon > 0$ ,

$$
\beta = \lim_{n \to \infty, x \ge -cn} Q^n[\phi_\alpha^{-}](x) = \lim_{n \to \infty, x \ge - (c - \epsilon)n} Q^n[\phi_\alpha^{-}](x - 1)
$$
  
\n
$$
\le \lim_{n \to \infty, x \ge - (c - \epsilon)n} Q^n[\phi](x) \le \lim_{n \to \infty, x \ge - (c - \epsilon)n} Q^n[\phi_\alpha^{-}](x + 1)
$$
  
\n
$$
= \lim_{n \to \infty, x \ge -cn} Q^n[\phi_\alpha^{-}](x) = \beta,
$$

which implies that  $c - \epsilon \in \Lambda(\phi)$ , and hence sup  $\Lambda(\phi_{\alpha}^{-}) = \sup \Lambda(\phi)$ . For convenience, we may call  $c^*_{-}(\alpha, \beta)$  the *leftward asymptotic speed of propagation* of  $\phi_{\alpha}^-$ .

Following the above procedure, we can find  $\delta_0 > 0$  and  $e_0 \in Int(\mathcal{X}^+)$  such that

$$
Q[\eta e_0] \ll \eta e_0, \quad \forall \eta \in (0, \delta_0]. \tag{2.7}
$$

Here we emphasize that  $\delta_0$ ,  $e_0$  above and  $\delta_\beta$ ,  $e_\beta$  will play a vital role in the whole paper because they describe the local stability of the fixed points 0 and  $\beta$ . Similarly, we can define

<span id="page-7-0"></span>
$$
v_{\alpha}^+ := \inf{\{\alpha, \delta_0 e_0\}}.
$$

Let  $\phi_{\alpha}^+ \in C_{\beta}$  be a nondecreasing initial function with

<span id="page-7-2"></span>
$$
\phi_{\alpha}^{+}(x) = \alpha, \forall x \ge 1, \text{ and } \phi_{\alpha}^{+}(x) = v_{\alpha}^{+}, \forall x \le 0.
$$

For the same reason, we define

$$
c_{+}^{*}(0,\alpha) := \sup \Biggl\{ c \in \mathbb{R} : \lim_{n \to \infty, x \leq cn} Q^{n} [\phi_{\alpha}^{+}](x) = 0 \Biggr\},
$$
 (2.8)

which is called the *rightward asymptotic speed of propagation* of  $\phi_{\alpha}^{+}$ . As showed above, these two speeds are bounded below, but may be plus infinity. To better understand these two spreading speeds, see Figure [3](#page-8-0) (left).

Now we are ready to state our last assumption on  $Q$ :

(A6) (*Counter-propagation*) For each  $\alpha \in E \setminus \{0, \beta\}$ ,  $c^*_{-}(\alpha, \beta) + c^*_{+}(0, \alpha) > 0$ .



<span id="page-8-0"></span>**Fig. 3.** Left:  $c^*_{-}(\alpha, \beta)$  and  $c^*_{+}(0, \alpha)$ . Right:  $c^*_{+}(\alpha, \beta)$  and  $c^*_{-}(0, \alpha)$ .

Assumption (A6) ensures that the two initial functions in the left part of Figure [3](#page-8-0) will eventually propagate in opposite directions although one of these two speeds may be negative. It is interesting to note that assumption (A6) is nearly necessary for the propagation of a bistable traveling wave. Indeed, if a monotone evolution system admits a bistable traveling wave, then it is usually unique (up to translation) and globally attractive (see, e.g., Remark [6.2\)](#page-40-0). This implies that the solution starting from the initial data  $\frac{1}{2}(\phi_{\alpha}^+ + \phi_{\alpha}^-)$ converges to a phase shift of the bistable wave. If  $c^*_{-}(\alpha, \beta) + c^*_{+}(0, \alpha) < 0$ , then the comparison principle would force the solutions starting from  $\phi_{\alpha}^{\pm}$  to split the bistable wave.

Comparing this with the definition of spreading speeds (short for asymptotic speeds of spread/propagation) for monostable semiflows (see, e.g., [\[2,](#page-43-10) [26\]](#page-44-13)), one can find that the leftward spreading speed of the monostable subsystem  ${Q<sup>n</sup>}_{n\geq 0}$  restricted to  $[\alpha, \beta]_C$  is shared by a large class of initial functions, and in many applications, it equals  $c^*_{-}(\alpha, \beta)$ . A similar observation holds for  $c^*_{+}(0, \alpha)$ . Thus, for a specific bistable system, assumption (A6) can be verified by using the properties of spreading speeds for monostable subsystems.

Remark 2.1. If we consider the nonincreasing traveling waves, then we can similarly define the numbers  $c^*_+(\alpha, \beta)$  and  $c^*(0, \alpha)$  (see Figure [3](#page-8-0) (right)). Then (A6) should be stated as  $c^*_+(\alpha, \beta) + c^*_-(0, \alpha) > 0$ .

#### 3. Semiflows in a homogeneous habitat

We say a habitat is *homogeneous* for the semiflow  $\{Q_t\}_{t\in\mathcal{T}}$  on a metric space  $\mathcal{E} \subset \mathcal{C}$  if

$$
Q_t[\phi](x - y) = Q_t[\phi(\cdot - y)](x), \quad \forall \phi \in \mathcal{E}, x, y \in \mathcal{H}, t \in \mathcal{T}.
$$

In this section, we will establish the existence of bistable traveling waves for the semiflow  ${Q_t}_{t\in\mathcal{T}}$  on  $\mathcal E$  in the following order: discrete-time semiflows in a continuous habitat, discrete-time semiflows in a discrete habitat, time-periodic semiflows, continuous-time semiflows in a continuous habitat, and continuous-time semiflows in a discrete habitat.

## *3.1. Discrete-time semiflows in a continuous habitat*

<span id="page-8-1"></span>In this case, time  $\mathcal T$  is discrete and habitat  $\mathcal H$  is continuous:  $\mathcal T = \mathbb{Z}^+$  and  $\mathcal H = \mathbb{R}$ . For convenience, we use Q to denote  $Q_1$ , and consider the semiflow  $\{Q^n\}_{n\geq 0}$ , where  $Q^n$  is the *n*-th iteration of  $Q$ .

**Definition 3.1.**  $\psi(x + cn)$  with  $\psi \in C$  is said to be a *traveling wave with speed*  $c \in \mathbb{R}$  of the discrete semiflow  $\{Q^n\}_{n\geq 0}$  if  $Q^n[\psi](x) = \psi(x + cn)$  for all  $x \in \mathbb{R}$  and  $n \geq 0$ . We say that  $\psi$  *connects* 0 *to*  $\beta$  if  $\psi(-\infty) := \lim_{x \to -\infty} \psi(x) = 0$  and  $\psi(\infty) :=$  $\lim_{x\to\infty}\psi(x)=\beta.$ 

<span id="page-9-1"></span>We first show that 0 and  $\beta$  are two isolated fixed points of Q in  $\mathcal{X}_{\beta}$  if (A5) holds.

**Lemma 3.1.** *Let*  $\delta_0$ ,  $e_0$  *and*  $\delta_\beta$ ,  $e_\beta$  *be chosen so that* [\(2.7\)](#page-7-0) *and* [\(2.3\)](#page-6-2) *hold, respectively. Then*  $E \cap \mathcal{X}_{\delta_0 e_0} = \{0\}$  *and*  $E \cap [\beta - \delta_\beta e_\beta, \beta]$  $\chi = \{\beta\}.$ 

*Proof.* Assume for contradiction that  $0 \neq \alpha \in E \cap \mathcal{X}_{\delta_0 e_0}$ . Define

$$
\bar{\delta} := \inf \{ \delta \in (0, \delta_0] : \alpha \in [0, \delta e_0]_{\mathcal{X}} \}.
$$

Then  $\alpha \leq \bar{\delta}e_0$  but  $\alpha \notin [0, \bar{\delta}e_0]$ ] $\chi$ . However, by the monotonicity of Q and the fact that 0 is strongly stable,

$$
\alpha = Q[\alpha] \le Q[\bar{\delta}e_0] \ll \bar{\delta}e_0.
$$

This contradicts  $\alpha \notin [0, \bar{\delta}e_0]]_X$ . Hence,  $E \cap \mathcal{X}_{\delta_0e_0} = \{0\}$ . Similarly,  $E \cap [\beta - \delta_\beta e_\beta, \beta]_X$  $= {\beta}.$ 

Choose  $\delta > 0$  such that

<span id="page-9-2"></span>
$$
\delta < \min\{\delta_0, \delta_\beta\} \quad \text{and} \quad \delta e_0 \ll \beta - \delta e_\beta. \tag{3.1}
$$

Assume that  $\psi$  and  $\bar{\psi}$  are nondecreasing functions in  $C(\mathbb{R}, \mathcal{X}_{\beta})$  with

$$
\underline{\psi}(x) = \begin{cases} 0, & x \le 0, \\ \beta - \delta e_{\beta}, & x \ge 1, \end{cases} \text{ and } \overline{\psi}(x) = \begin{cases} \delta e_0, & x \le -1, \\ \beta, & x \ge 0. \end{cases}
$$

Clearly  $\psi \leq \bar{\psi}$ . We have the following observation.

<span id="page-9-0"></span>Lemma 3.2. *Assume that* Q *satisfies* (A1)–(A3) *and* (A5)*. Then there exists a positive rational number*  $\bar{c}$  *such that for any*  $c \geq \bar{c}$ *, we have* 

$$
Q[\psi](x) \ge \psi(x-c)
$$
 and  $Q[\bar{\psi}](x) \le \bar{\psi}(x+c)$  for any  $x \in \mathbb{R}$ .

*Proof.* Assume that  $x_n \to \infty$  is an increasing sequence in R. Then  $\psi_n := \psi(\cdot + x_n)$ converges to  $\beta - \delta e_{\beta}$  in  $C_{\beta}$  since  $\psi(x) = \beta - \delta e_{\beta}$  for all  $x \ge 1$ . It then follows from  $(A1)$ – $(A2)$  and  $(A5)$  that

$$
Q[\psi](\infty) = \lim_{n \to \infty} Q[\psi](x_n) = \lim_{n \to \infty} Q[\psi(\cdot + x_n)](0) = Q[\beta - \delta e_{\beta}] \gg \beta - \delta e_{\beta}.
$$

Therefore, there exists  $y_0 > 0$  such that  $Q[\psi](y_0) \ge \beta - \delta e_\beta$ . Note that  $Q[\psi](x)$  is nondecreasing in x. Then for any  $c \geq y_0$  we have

$$
Q[\psi](x) \ge Q[\psi](y_0) \ge \beta - \delta e_\beta \ge \psi(x - c), \quad \forall x \ge y_0,
$$

and

$$
Q[\psi](x) \ge 0 = \psi(0) \ge \psi(x - y_0) \ge \psi(x - c), \quad \forall x < y_0,
$$

which means  $Q[\psi](x) \ge \psi(x - c)$  for all  $c \ge y_0$ . Similarly,

$$
Q[\bar{\psi}](-\infty) = \lim_{n \to \infty} Q[\psi](-x_n) = \lim_{n \to \infty} Q[\bar{\psi}(\cdot - x_n)](0) = Q[\delta e_0] \ll \delta e_0 = \bar{\psi}(-\infty),
$$

and hence there exists  $z_0 > 0$  such that  $Q[\bar{\psi}](x) \le \bar{\psi}(x + c)$  for all  $c \ge z_0$ . Choosing  $\bar{c} = \max\{y_0, z_0\}$  completes the proof.

Let  $\kappa_n := (n + \bar{c})/n$ . Clearly,  $\kappa_n$  is a rational number for all  $n \ge 1$ . For any  $\xi \in \mathbb{R}$ , define a map  $A_{\xi}: \mathcal{B} \to \mathcal{B}$  by  $A_{\xi}[\phi](x) = \phi(\xi x)$  for all  $x \in \mathbb{R}$ . Define  $\psi_n$ ,  $\bar{\psi}_n \in C_{\beta}$  by

$$
\psi_n(x) = \psi(x - (n + \bar{c})) \quad \text{and} \quad \bar{\psi}_n(x) = \bar{\psi}(x + (n + \bar{c})).
$$

<span id="page-10-1"></span>**Lemma 3.3.** Assume that Q satisfies (A1)–(A5). Then for each  $n \in \mathbb{N}$ ,  $G_n := Q \circ A_{\kappa_n}$ *has a fixed point*  $\phi_n$  *in*  $C_\beta$  *such that*  $\phi_n$  *is nondecreasing and*  $\psi_n \leq \phi_n \leq \bar{\psi}_n$ *.* 

*Proof.* We first show that  $\psi_n \leq G_n[\psi_n]$ . Indeed, when  $x < n$  we have

$$
\psi_n(x+\bar{c}) \le \psi_n(n+\bar{c}) = \psi(0) = 0 \le A_{\kappa_n}[\psi_n](x);
$$

when  $x > n$  we have

$$
A_{\kappa_n}[\underline{\psi}_n](x) = \underline{\psi}_n(\kappa_n x) = \underline{\psi}_n\left(x + \frac{\bar{c}}{n}x\right) \ge \underline{\psi}_n(x + \bar{c}),
$$

and hence  $\psi_n(x + \bar{c}) \leq A_{\kappa_n}[\psi_n](x)$  for all  $x \in \mathbb{R}$ . Consequently, by the monotonicity of Q and  $\bar{\psi}(x) \leq Q[\psi](x + \bar{c})$  (see Lemma [3.2\)](#page-9-0) we obtain

$$
\underline{\psi}_n(x) \le Q[\underline{\psi}_n](x+\overline{c}) \le Q \circ A_{\kappa_n}[\underline{\psi}_n](x) = G_n[\underline{\psi}_n](x).
$$

Similarly,  $\bar{\psi}_n \geq G_n[\bar{\psi}_n]$ . It follows that

<span id="page-10-0"></span>
$$
\Psi_n \le G_n^k[\Psi_n] \le G_n^k[\bar{\Psi}_n] \le \bar{\Psi}_n, \quad \forall k \in \mathbb{N}.
$$
\n(3.2)

For any  $k \geq 1$ , we have

$$
G_n^k[\psi_n] = G_n \circ G_n^{k-1}[\psi_n] \in G_n[\mathcal{C}_\beta].\tag{3.3}
$$

Since  $G_n$  is order preserving and  $\psi_n(x)$  is nondecreasing in x, we know that  $G_n^k[\psi_n](x)$ is nondecreasing in both k and x. Recall that  $G_n$  is compact due to assumption (A4). It follows that  $G_h^k[\psi_n]$  converges in  $\mathcal{C}_\beta$ . Denote the limit by  $\phi_n$ . By [\(3.2\)](#page-10-0), we also get  $\psi_n \leq \phi_n \leq \bar{\psi}_n$ . Moreover,  $\phi_n(x)$  is nondecreasing due to Proposition [7.1\(](#page-41-0)2). Obviously,

<span id="page-10-2"></span>
$$
\phi_n = \lim_{k \to \infty} G_n^{k+1}[\psi_n] = G_n \Big[ \lim_{k \to \infty} G_n^k[\psi_n] \Big] = G_n[\phi_n].
$$

This completes the proof.  $\Box$ 

The following lemma reveals a relation between the wave speeds of monostable traveling waves in sub-monostable systems and the numbers defined in  $(2.6)$  and  $(2.8)$ .

**Lemma 3.4.** *Let*  $c^*_{-}(\alpha, \beta)$  *and*  $c^*_{+}(0, \alpha)$  *be defined as in* [\(2.6\)](#page-7-1) *and* [\(2.8\)](#page-7-2)*. Assume that* Q *satisfies* (A3)*. Then the following statements are valid:*

- (1) *If*  $\psi(x+ct)$  *is a monotone traveling wave connecting*  $\alpha$  *to*  $\beta$  *of the discrete semiflow*  ${Q<sup>n</sup>}_{n \geq 1}$ *, then*  $c \geq c^*_{-}(\alpha, \beta)$ *.*
- (2) If  $\psi(x+ct)$  *is a monotone traveling wave connecting* 0 *to*  $\alpha$  *of the discrete semiflow*  ${Q<sup>n</sup>}_{n \geq 1}$ *, then*  $c \leq -c^*_{+}(0, \alpha)$ *.*

*Proof.* We only prove (1) since the proof of (2) is similar. In view of Lemma [3.1,](#page-9-1) there exists a neighborhood  $\mathcal N$  of  $\beta$  in [ $\alpha$ ,  $\beta$ ] $\chi$  such that

$$
v^-_\alpha<\gamma,\quad \forall \gamma\in\mathcal{N},
$$

where  $v_\alpha^-$  is defined in [\(2.4\)](#page-6-0). Since  $\psi(-\infty) = \alpha$  and  $\psi(\infty) = \beta$ , there must exist a translation of  $\psi$ , still denoted by  $\psi$ , such that  $\phi_{\alpha}^{-} \leq \psi$ . Assume for contradiction that  $c < c<sup>*</sup>_{-}(\alpha, \beta)$ . Choose  $q/p \in (c, c<sup>*</sup>_{-}(\alpha, \beta))$  with  $p, q \in \mathbb{Z}$ . It follows from [\(2.6\)](#page-7-1) that

$$
\beta = \lim_{n \to \infty} Q^{pn} [\phi_{\alpha}^{-}] \left( -\frac{q}{p} \cdot pn \right) \le \lim_{n \to \infty} Q^{pn} [\psi] (-qn)
$$
  
= 
$$
\lim_{n \to \infty} \psi (-qn + cpn) = \psi (-\infty) = \alpha,
$$

a contradiction. Thus,  $c \geq c^*_{-}(\alpha, \beta)$ .

Now we are ready to prove the main result of this subsection.

**Theorem 3.1.** Assume that Q satisfies (A1)–(A6). Then there exists  $c \in \mathbb{R}$  such that the discrete semiflow  ${Q^n}_{n \geq 1}$  admits a nondecreasing traveling wave with speed c and *connecting* 0 *to*  $\beta$ .

*Proof.* The proof is in three steps. Firstly, we construct  $\phi_+$ ,  $\phi_- \in C_\beta$  and  $c_+ \leq c_- \in \mathbb{R}$ such that

$$
Q[\phi_+] (x) = \phi_+ (x + c_+)
$$
 and  $Q[\phi_-] (x) = \phi_- (x + c_-)$ 

with

$$
\phi_{-}(0) \in (0, \delta e_0] \chi
$$
 and  $\phi_{+}(0) \in [\beta - \delta e_{\beta}, \beta] \chi$ .

Indeed, let  $\phi_n$  be as in Lemma [3.3.](#page-10-1) Since  $0 \ll \bar{\psi}(-1) = \delta e_0 \ll \psi(1) = \beta - \delta e_\beta \ll \beta$ and  $\psi_n \leq \phi_n \leq \bar{\psi}_n$ , we have

$$
\bar{\psi}(-1) = \bar{\psi}_n(-1 - (n + \bar{c})) \ge \phi_n(-1 - (n + \bar{c})),
$$
  

$$
\psi(1) = \psi_n(1 + (n + \bar{c})) \le \phi_n(1 + (n + \bar{c})).
$$

Now we define

$$
a_n := \sup_{x \in \mathbb{R}} \{ \phi_n(x) \in [0, \delta e_0]_{\mathcal{X}} \}, \quad b_n := \inf_{x \in \mathbb{R}} \{ \phi_n(x) \in [\beta - \delta e_\beta, \beta]_{\mathcal{X}} \}.
$$

Then

$$
-1 - (n + \bar{c}) \le a_n \le b_n \le 1 + (n + \bar{c}), \quad \phi_n(a_n) \le \delta e_0 \le \beta - \delta e_\beta \le \phi_n(b_n).
$$

<span id="page-11-0"></span>

Define  $\phi_{-,n}(x) := \phi_n(x + a_n)$  and  $\phi_{+,n}(x) := \phi_n(x + b_n)$ . Then

$$
\phi_{-,n} = \phi_n(\cdot + a_n) = G_n[\phi_n](\cdot + a_n) = Q[\phi_n(\kappa_n \cdot)](\cdot + a_n) = Q[\phi_n(\kappa_n(\cdot + a_n))] \in Q[\mathcal{C}_{\beta}].
$$

Similarly,  $\phi_{+,n} = Q[\phi_n(\kappa_n(\cdot + b_n))] \in Q[\mathcal{C}_{\beta}].$  Thus, there exists a subsequence (still indexed by *n*), two nondecreasing functions  $\phi_-, \phi_+ \in C_\beta$  and  $\xi_-, \xi_+ \in [-1, 1]$  with ξ<sup>−</sup>  $\leq$  ξ<sup>+</sup> such that

$$
\lim_{n \to \infty} \frac{a_n}{n} = \xi_-, \quad \lim_{n \to \infty} \frac{b_n}{n} = \xi_+, \quad \lim_{n \to \infty} \phi_{-,n} = \phi_-, \quad \lim_{n \to \infty} \phi_{+,n} = \phi_+.
$$

Obviously,  $\phi_-(0) = \lim_{n \to \infty} \phi_n(a_n)$  and  $\phi_+(0) = \lim_{n \to \infty} \phi_n(b_n)$ . By the definitions of  $a_n$  and  $b_n$ , we immediately have  $\phi_-(0) \neq 0$  and  $\phi_+(0) \neq \beta$ , and hence  $0 < \phi_-(0) \leq$  $\psi(-1) = \delta e_0$  and  $\beta - \delta e_\beta = \psi(1) \leq \phi_+(0) < \beta$ . Define  $c_- := -\bar{c}\xi_-$  and  $c_+ := -\bar{c}\xi_+$ . Obviously,  $c_-\geq c_+$  because  $\xi_-\leq \xi_+$ .

We only prove  $Q[\phi_{-}](x) = \phi_{-}(x + c_{-})$  because the proof of the other identity is similar. Note that the following limit is uniform for x in any bounded subset  $M \subset \mathbb{R}$ :

<span id="page-12-0"></span>
$$
\lim_{n\to\infty}\kappa_n(x+a_n)-a_n=\lim_{n\to\infty}\left(x+\bar{c}\cdot\frac{x+a_n}{n}\right)=x-c_-.
$$

Hence for any  $x \in \mathbb{R}$ ,

$$
\phi_{-}(x+c_{-}) = \lim_{n \to \infty} \phi_{-,n}(x+c_{-}) = \lim_{n \to \infty} \phi_{n}(x+c_{-}+a_{n}) = \lim_{n \to \infty} G_{n}[\phi_{n}](x+c_{-}+a_{n})
$$
  
\n
$$
= \lim_{n \to \infty} Q[\phi_{n}(\kappa_{n} \cdot)](x+c_{-}+a_{n}) = \lim_{n \to \infty} Q[\phi_{n}(\kappa_{n}(\cdot+a_{n}))](x+c_{-})
$$
  
\n
$$
= \lim_{n \to \infty} Q[\phi_{-,n}(\kappa_{n}(\cdot+a_{n})-a_{n})](x+c_{-}) = Q[\phi_{-}](x), \qquad (3.4)
$$

where the last equality is obtained from Proposition  $7.2(2)$  $7.2(2)$  and the continuity of  $Q$ .

Secondly, we prove that  $\phi_{\pm}(x)$  obtained in the first step have the following properties:

(i)  $\phi_{-}(-\infty) = 0$  and  $\phi_{+}(\infty) = \beta$ ;

(ii)  $\phi_-(\infty)$  and  $\phi_+(-\infty)$  are ordered.

Indeed, let  $x_n \to \infty$  be an increasing sequence in R. Note that  $\phi_-(x_n) = Q[\phi_-(\cdot - c_- +$  $(x_n)[(0) \in Q[\mathcal{C}_{\beta}](0)$ , which is precompact in  $\mathcal{X}_{\beta}$ . Therefore, there exists a subsequence  ${n_l}$  and  $v \in \mathcal{X}_{\beta}$  such that  $\lim_{l \to \infty} \phi_-(x_{n_l}) = v$ , which, together with the fact that  $\phi_-$ is nondecreasing and Proposition [7.2\(](#page-41-1)1), implies that  $\phi_-(\infty) := \lim_{x \to \infty} \phi_-(x) = v$ . Moreover, from [\(3.4\)](#page-12-0) we see that  $\phi_+(\infty) \in \mathcal{X}_\beta$  is a fixed point of Q. Similar results hold for  $\phi_-(-\infty)$  and  $\phi_+(\pm\infty)$ . Recall that  $\phi_-(-\infty) \leq \phi_-(0) \leq \delta e_0$  and  $\phi_+(\infty) \geq$  $\phi_{+}(0) \geq \beta - \delta e_{\beta}$ , which, together with the choice of  $\delta$ , implies that  $\phi_{-}(-\infty) = 0$  and  $\phi_+(\infty) = \beta$ . Further, since any two real numbers are ordered, there exist sequences

$$
\{n\}_{n\geq 0}\supset \{n_{1m}\}_{m\geq 1}\supset \{n_{2m}\}_{m\geq 2}\supset \cdots \supset \{n_{km}\}_{m\geq 1}\supset \cdots
$$

such that for each  $k \geq 1$ ,

$$
k + a_{n_{km}} \leq -k + b_{n_{km}}, \forall m \geq 1, \text{ or } k + a_{n_{km}} \geq -k + b_{n_{km}}, \forall m \geq 1.
$$

Define  $\Gamma_1 := \{k \in \mathbb{N} : k + a_{n_{km}} \leq -k + b_{n_{km}}, \forall m \geq 1\}$  and  $\Gamma_2 := \mathbb{N} \setminus \Gamma_1$ . Then either  $\Gamma_1$  or  $\Gamma_2$  has infinitely many elements. If  $\Gamma_1$  does, then

$$
\phi_{-,n_{km}}(k) = \phi_{n_{km}}(k + a_{n_{km}}) \le \phi_{n_{km}}(-k + b_{n_{km}}) = \phi_{+,n_{km}}(-k), \quad \forall k \in \Gamma_1, m \in \mathbb{N}.
$$

This implies that  $\phi_-(k) \leq \phi_+(-k)$  for all  $k \in \Gamma_1$ , and hence  $\phi_-(\infty) \leq \phi_+(-\infty)$ . If  $\Gamma_2$  has infinitely many elements, then  $\phi_-(\infty) \ge \phi_+(-\infty)$  by a similar argument. Thus,  $\phi_-(\infty)$  and  $\phi_+(-\infty)$  must be ordered in  $\mathcal{X}_{\beta}$ .

Finally, we prove that either  $\phi_-\$  or  $\phi_+\$  connects 0 to  $\beta$ . Indeed, we have shown in the second step that  $\phi_-(\infty)$  and  $\phi_+(-\infty)$  are ordered. It then follows from the bistability assumption (A5) that there are only three possibilities:

(i)  $\beta = \phi_-(\infty) \ge \phi_+(-\infty);$ 

(ii)  $\phi_-(\infty) \ge \phi_+(-\infty) = 0;$ 

(iii)  $\phi_{-}(\infty) = \alpha = \phi_{+}(-\infty)$  for some  $\alpha \in E \setminus \{0, \beta\}.$ 

We further claim that the possibility (iii) cannot happen: Otherwise, Lemma [3.4](#page-10-2) implies that  $c_+ \geq c^*_{-}(\alpha, \beta)$  and  $c_- \leq -c^*_{+}(0, \alpha)$ . Since  $c_- \geq c_+$ , it then follows that

$$
0 \ge c_{+} + (-c_{-}) \ge c_{-}^{*}(\alpha, \beta) + c_{+}^{*}(0, \alpha),
$$

which contradicts (A6). Thus, either (i) or (ii) holds, completing the proof.  $\square$ 

## *3.2. Discrete-time semiflows in a discrete habitat*

In this case, both  $\mathcal T$  and  $\mathcal H$  are discrete:  $\mathcal T = \mathbb{Z}^+$  and  $\mathcal H = \mathbb{Z}$ . Without confusion, we consider the semiflow  ${Q<sup>n</sup>}_{n \geq 0}$  in a metric space  $\mathcal{E} \subset \mathcal{C}$ . Since the habitat is discrete, we cannot use the definition of traveling waves with an unknown speed as in Definition [3.1.](#page-8-1) This is because the wave profile  $\psi(x)$  may not be well defined for all  $x \in \mathbb{R}$ . So we modify the definition of traveling waves in a discrete habitat.

<span id="page-13-0"></span>**Definition 3.2.**  $\psi(x + cn)$  with  $\psi \in \mathcal{B}$  is said to be a *traveling wave with speed*  $c \in \mathbb{R}$  of the discrete semiflow  $\{Q^n\}_{n\geq 0}$  if there exists a countable set  $\Gamma \subset \mathbb{R}$  such that  $Q[\psi(\cdot + x)](i) = \psi(i + x + c)$  for all  $i \in \mathbb{Z}$  and  $x \in \mathbb{R} \setminus \Gamma$ .

By Definition [3.2](#page-13-0) and Proposition [7.3,](#page-41-2) there exists  $x_0 \in \mathbb{R}$  such that  $Q^n[\psi(\cdot + x_0)](i) =$  $\psi(i + x_0 + cn)$  for all  $i \in \mathbb{Z}$  and  $n \ge 0$ . Define  $\phi(x) := \psi(x + x_0)$  for all  $x \in \mathbb{R}$ . Then, with a little abuse of notation,  $Q^{n}[\phi](i) = \phi(i + cn)$  for all  $i \in \mathbb{Z}$  and  $n \ge 0$ . Therefore, Definition [3.2](#page-13-0) is a generalized version of the classical one, which is an analogue of Definition [3.1.](#page-8-1)

Let  $\beta \gg 0$  be a fixed point of Q. Define  $\tilde{Q} : \mathcal{B}_{\beta} \to \mathcal{B}_{\beta}$  by

$$
\tilde{Q}[\phi](x) = Q[\phi(\cdot + x)](0), \quad \forall x \in \mathbb{R}.
$$

Then we see from [\[25,](#page-44-14) Lemma 2.1] that  $\tilde{Q}$  satisfies (A1)–(A3) and (A5) with  $Q = \tilde{Q}$  and  $C_\beta = \mathcal{B}_\beta$  if Q itself satisfies (A1)–(A3) and (A5). Further, if Q satisfies (A4), then the set  $\tilde{Q}[\mathcal{B}_{\beta}](x) \subset \mathcal{X}_{\beta}$  is precompact for any  $x \in \mathbb{R}$ .

For  $\tilde{Q}$  :  $\mathcal{B}_{\beta} \rightarrow \mathcal{B}_{\beta}$ , we have similar results to Lemmas [3.2](#page-9-0) and [3.3.](#page-10-1)

Lemma 3.5. *Assume that* Q *satisfies* (A1)–(A3) *and* (A5)*. Then there exists a positive rational number*  $\bar{c}$  *such that for any*  $c \geq \bar{c}$ *, we have* 

$$
\tilde{Q}[\psi](x) \ge \psi(x - c) \quad \text{and} \quad \tilde{Q}[\bar{\psi}](x) \le \bar{\psi}(x + c) \quad \text{for any } x \in \mathbb{R}.
$$

**Lemma 3.6.** Assume that Q satisfies (A1)–(A5). Then for each  $n \in \mathbb{N}$ ,  $\tilde{G}_n := \tilde{Q} \circ A_{\kappa_n}$ *has a fixed point*  $\tilde{\phi}_n$  in  $\mathcal{B}_\beta$  *such that*  $\tilde{\phi}_n$  is nondecreasing and  $\psi_n \leq \tilde{\phi}_n \leq \tilde{\psi}_n$ .

*Proof.* By the same arguments as in the proof of Lemma [3.3,](#page-10-1) we can obtain a similar inequality to [\(3.2\)](#page-10-0):

<span id="page-14-0"></span>
$$
\psi_n \leq \tilde{G}_n^k[\psi_n] \leq \tilde{G}_n^k[\bar{\psi}_n] \leq \bar{\psi}_n, \quad \forall k \in \mathbb{N}.
$$

Define  $w_{n,1} := \psi_n$  and  $w_{n,k+1} := \tilde{G}_n[w_{n,k}], k \ge 1$ . Then

$$
w_{n,k+1}(x) = \tilde{Q} \circ A_{\kappa_n}[w_{n,k}](x) = \tilde{Q}[w_{n,k}(\kappa_n \cdot)](x) = Q[w_{n,k}(\kappa_n(\cdot + x))](0). \quad (3.5)
$$

Note that  $Q[C_\beta]$  is compact and  $w_{n,k}$  is nondecreasing in k. Hence, for any fixed  $x \in \mathbb{R}$ ,  $w_{n,k}(x)$  converges in  $\mathcal{X}_{\beta}$ . Denote the limit by  $\tilde{\phi}_n(x)$ . Then  $\tilde{\phi}_n(x)$  is nondecreasing in  $x \in \mathbb{R}$  and  $\psi_n \leq \tilde{\phi}_n \leq \bar{\psi}_n$ . Letting  $k \to \infty$  in [\(3.5\)](#page-14-0), we arrive at  $\tilde{\phi}_n(x) = Q[\tilde{\phi}_n(\kappa_n(\cdot + x))](0)$ . Consequently,

<span id="page-14-1"></span>
$$
\tilde{\phi}_n = \tilde{Q}[\tilde{\phi}_n(\kappa_n \cdot)] = \tilde{Q} \circ A_{\kappa_n}[\tilde{\phi}_n] = \tilde{G}_n[\tilde{\phi}_n].
$$

This completes the proof.  $\Box$ 

Due to the lack of compactness for  $\tilde{Q}$ , we will use the properties of monotone functions established in the Appendix to show the convergence of a sequence in  $Q[\mathcal{B}_\beta]$ .

**Theorem 3.2.** Assume that  $X = C(M, \mathbb{R}^d)$  and Q satisfies (A1)–(A6). Then there exists  $c \in \mathbb{R}$  such that the semiflow  ${Q^n}_{n \geq 1}$  on  $\mathcal{C}_\beta$  admits a nondecreasing traveling wave ψ(x+cn) *with speed* c *and connecting* 0 *to* β*. Further,* ψ *is either left or right continuous.*

*Proof.* As in the proof of Theorem [3.1,](#page-11-0) we define

$$
\tilde{a}_n := \sup_{x \in \mathbb{R}} \{ \tilde{\phi}_n(x) \in [0, \delta e_0]_{\mathcal{X}} \}, \quad \tilde{b}_n := \inf_{x \in \mathbb{R}} \{ \tilde{\phi}_n(x) \in [\beta - \delta e_\beta, \beta]_{\mathcal{X}} \}.
$$

Then  $-1 - (n + \bar{c}) \leq \tilde{a}_n \leq \tilde{b}_n \leq 1 + (n + \bar{c})$ . Note that for any  $x \in \mathbb{R}$ , we have

$$
\tilde{\phi}_n(x) = \tilde{G}_n[\tilde{\phi}_n](x) = \tilde{Q}[\tilde{\phi}_n(\kappa_n \cdot)](x) = Q[\tilde{\phi}_n(\kappa_n(\cdot + x))](0) \in Q[\mathcal{C}_{\beta}](0).
$$

Since  $Q[\mathcal{C}_{\beta}](0)$  is precompact in  $\mathcal{X}_{\beta}$ , for any  $x \in \mathbb{R}$  the limits  $\tilde{\phi}_n(x^-) := \lim_{y \uparrow x} \tilde{\phi}_n(y)$ and  $\tilde{\phi}_n(x^+) := \lim_{y \downarrow x} \tilde{\phi}_n(y)$  both exist. Hence, by the definitions of  $\tilde{a}_n$  and  $\tilde{b}_n$ ,

$$
\tilde{\phi}_n(\tilde{a}_n^-) \leq \delta e_0 \leq \beta - \delta e_\beta \leq \tilde{\phi}_n(\tilde{b}_n^+),
$$

but

$$
\tilde{\phi}_n(\tilde{a}_n^+) \notin [0, \delta e_0]]_{\mathcal{X}}
$$
 and  $\tilde{\phi}_n(\tilde{b}_n^-) \notin [[\beta - \delta e_\beta, \beta]_{\mathcal{X}}$ .

Define  $\tilde{\phi}_{-,n}(x) := \tilde{\phi}_n(x + \tilde{a}_n)$  and  $\tilde{\phi}_{+,n}(x) := \tilde{\phi}_n(x + \tilde{b}_n)$ . Then

$$
\tilde{\phi}_{-,n}(0^-) \leq \delta e_0 \leq \beta - \delta e_{\beta} \leq \tilde{\phi}_{+,n}(0^+),
$$

but

$$
\tilde{\phi}_{-,n}(0^+) \notin [0, \delta e_0] | \chi
$$
 and  $\tilde{\phi}_{+,n}(0^-) \notin [[\beta - \delta e_\beta, \beta] \chi]$ .

Since  $\tilde{\phi}_n = \tilde{G}_n[\tilde{\phi}_n]$ , we have

$$
\tilde{\phi}_{-,n}(x) = \tilde{G}_n[\tilde{\phi}_n](x + \tilde{a}_n) = Q[\tilde{\phi}_n(\kappa_n(\cdot + \tilde{a}_n + x))](0) \in Q[\mathcal{C}_{\beta}](0).
$$

Similarly,  $\tilde{\phi}_{+,n}(x) = Q[\tilde{\phi}_n(\kappa_n(\cdot + \tilde{b}_n + x))](0) \in Q[\mathcal{C}_{\beta}](0)$ . Let  $\mathbb Q$  be the set of all rational numbers, and  $\{x_l\}_{l\geq 1} \subset \mathbb{Q}$  be an increasing sequence converging to x. Using  $\tilde{\phi}_n = \tilde{G}_n[\tilde{\phi}_n]$  again, we see that for any  $i \in \mathbb{Z}$  and  $l \geq 1$ ,

$$
\tilde{\phi}_n(\kappa_n(i+\tilde{a}_n+x_l))=Q[\tilde{\phi}_n(\kappa_n(\cdot+\kappa_n(i+\tilde{a}_n+x_l)))](0)\in Q[\mathcal{C}_{\beta}](0).
$$

Similarly,  $\tilde{\phi}_n(\kappa_n(i + \tilde{b}_n + x_l)) \in Q[\mathcal{C}_{\beta}](0)$ . Note that  $Q[\mathcal{C}_{\beta}](0)$  is precompact in  $\mathcal{X}_{\beta}$  and Q is countable. Hence there exists a subsequence (still indexed by  $\{n\}$ ) and  $\xi_-\leq \xi_+\in \mathbb{R}$ such that  $\lim_{n\to\infty} \tilde{a}_n/n = \xi_-, \lim_{n\to\infty} \tilde{b}_n/n = \xi_+$  and for any  $x \in \mathbb{Q}, i \in \mathbb{Z}$  and  $i \ge 1$ , the sequences  $\tilde{\phi}_{\pm,n}(x)$ ,  $\tilde{\phi}_n(\kappa_n(i + \tilde{a}_n + x_i))$  and  $\tilde{\phi}_n(\kappa_n(i + \tilde{b}_n + x_i))$  converge in  $\mathcal{X}_{\beta}$ . Hence, the limits

$$
\lim_{l \to \infty} \lim_{n \to \infty} \tilde{\phi}_{-,n}(x_l) = \lim_{l \to \infty} \lim_{n \to \infty} Q[\tilde{\phi}_n(\kappa_n(\cdot + \tilde{a}_n + x_l))](0),
$$
  

$$
\lim_{l \to \infty} \lim_{n \to \infty} \tilde{\phi}_{+,n}(x_l) = \lim_{l \to \infty} \lim_{n \to \infty} Q[\tilde{\phi}_n(\kappa_n(\cdot + \tilde{b}_n + x_l))](0)
$$

both exist. This means the limits

$$
\lim_{y \in \mathbb{Q}, y \uparrow x} \lim_{n \to \infty} \tilde{\phi}_{\pm,n}(y) \quad \text{and} \quad \lim_{y \in \mathbb{Q}, y \downarrow x} \lim_{n \to \infty} \tilde{\phi}_{\pm,n}(y)
$$

exist for all  $x \in \mathbb{R}$ . Define

$$
\hat{\phi}_{-}(x) := \begin{cases}\n\lim_{n \to \infty} \tilde{\phi}_{-,n}(x), & x \in \mathbb{Q}, \\
\lim_{y \in \mathbb{Q}, y \uparrow x} \lim_{n \to \infty} \tilde{\phi}_{-,n}(x), & x \in \mathbb{R} \setminus \mathbb{Q}, \\
\hat{\phi}_{+}(x) := \begin{cases}\n\lim_{n \to \infty} \tilde{\phi}_{+,n}(x), & x \in \mathbb{Q}, \\
\lim_{y \in \mathbb{Q}, y \downarrow x} \lim_{n \to \infty} \tilde{\phi}_{+,n}(x), & x \in \mathbb{R} \setminus \mathbb{Q}.\n\end{cases}\n\end{cases}
$$

Clearly,  $\hat{\phi}_{\pm}$  are nondecreasing functions in  $\mathcal{B}_{\beta}$  and for any  $x \in \mathbb{R} \setminus \mathbb{Q}$ ,  $\hat{\phi}_{\pm}(x^{\pm})$  all exist. Hence, we see from Theorem [7.1](#page-42-0) that there exists a countable subset  $\Gamma_1$  of R such that  $\tilde{\phi}_{\pm n}(x)$  converges to  $\hat{\phi}_{\pm}(x)$  for all  $x \in \mathbb{R} \setminus \Gamma_1$ . Define

$$
\tilde{\phi}_{-}(x) := \lim_{y \in \mathbb{Q}, y \uparrow x} \lim_{n \to \infty} \tilde{\phi}_{-,n}(y), \qquad \tilde{\phi}_{+}(x) := \lim_{y \in \mathbb{Q}, y \downarrow x} \lim_{n \to \infty} \tilde{\phi}_{+,n}(y), \quad \forall x \in \mathbb{R}.
$$

Thus,  $\tilde{\phi}_-(x)$  is left continuous and  $\tilde{\phi}_+(x)$  is right continuous. Note that  $\tilde{\phi}_\pm(x) = \hat{\phi}_\pm(x)$ for all  $x \in \mathbb{R} \setminus \mathbb{Q}$ . It then follows that  $\phi_{\pm,n}(x)$  converges to  $\phi_{\pm}(x)$  for  $x \in \mathbb{R} \setminus \Gamma_2$ , where  $\Gamma_2 := \mathbb{Q} \cup \Gamma_1$  is also countable.

Let  $y_k \in \mathbb{R} \setminus \Gamma_2$  be an increasing sequence converging to 0 and  $z_k \in \mathbb{R} \setminus \Gamma_2$  be an increasing sequence converging to 1. Note that

$$
\tilde{\phi}_{-}(0) = \lim_{k \to \infty} \tilde{\phi}_{-}(y_k) = \lim_{k \to \infty} \lim_{n \to \infty} \tilde{\phi}_{-,n}(y_k) \le \delta e_0,
$$
  

$$
\tilde{\phi}_{-}(1) = \lim_{k \to \infty} \tilde{\phi}_{-}(z_k) = \lim_{k \to \infty} \lim_{n \to \infty} \tilde{\phi}_{-,n}(z_k) \notin [0, \delta e_0]] \chi.
$$

Similarly,  $\tilde{\phi}_+(0) \ge \beta - \delta e_\beta$  but  $\tilde{\phi}_+(-1) \notin [[\beta - \delta e_\beta, \beta]_{\mathcal{X}}]$ . Define  $c_- := -\bar{c}\xi_-$  and  $c_+ := -\bar{c}\xi_+$ . Obviously,  $c_- \geq c_+$  since  $\xi_- \leq \xi_+$ .

Now we want to prove  $Q[\tilde{\phi}_-(\cdot+x)](0) = \tilde{\phi}_-(x+c_-)$  for all  $x \in \mathbb{R} \setminus \Gamma_2$ . Note that

$$
\lim_{n \to \infty} \kappa_n(x + \tilde{a}_n) - \tilde{a}_n = \lim_{n \to \infty} \left( x + \bar{c} \cdot \frac{x + \tilde{a}_n}{n} \right) = x - c_-.
$$

It follows that

$$
\tilde{\phi}_{-}(x+c_{-}) = \lim_{n \to \infty} \tilde{\phi}_{-,n}(x+c_{-}) = \lim_{n \to \infty} \tilde{\phi}_{n}(x+c_{-}+\tilde{a}_{n}) = \lim_{n \to \infty} \tilde{G}_{n}[\tilde{\phi}_{n}](x+c_{-}+\tilde{a}_{n})
$$
\n
$$
= \lim_{n \to \infty} \tilde{Q}[\tilde{\phi}_{n}(\kappa_{n} \cdot)](x+c_{-}+\tilde{a}_{n}) = \lim_{n \to \infty} \tilde{Q}[\tilde{\phi}_{n}(\kappa_{n}(\cdot+a_{n}))](x+c_{-})
$$
\n
$$
= \lim_{n \to \infty} \tilde{Q}[\tilde{\phi}_{-,n}(\kappa_{n}(\cdot+\tilde{a}_{n})-\tilde{a}_{n})](x+c_{-})
$$
\n
$$
= \lim_{n \to \infty} Q[\tilde{\phi}_{-,n}(\kappa_{n}(\cdot+x+c_{-}+\tilde{a}_{n})-\tilde{a}_{n})](0).
$$

In view of Proposition [7.5,](#page-41-3) we obtain  $\tilde{\phi}_-(x + c_-) = Q[\tilde{\phi}_-(x + x)](0)$  for all  $x \in \mathbb{R} \setminus \Gamma_2$ . A similar result holds for  $\phi_+$ .

Now, the same argument as in the proof of Theorem [3.1](#page-11-0) completes the proof.  $\Box$ 

#### *3.3. Time-periodic semiflows*

Let  $\omega \in \mathcal{T}$ , where  $\mathcal{T} = \mathbb{R}^+$  or  $\mathbb{Z}^+$ . Recall that a family  $\{Q_t\}_{t \in \mathcal{T}}$  of mappings is said to be an  $\omega$ -*time periodic semiflow* on a metric space  $\mathcal{E} \subset \mathcal{C}$  provided that:

- (i)  $Q_0[\phi] = \phi$  for all  $\phi \in \mathcal{E}$ .
- (ii)  $Q_t \circ Q_{\omega}[\phi] = Q_{t+\omega}[\phi]$  for all  $t \ge 0$  and  $\phi \in \mathcal{E}$ .
- (iii)  $Q_t[\phi]$  is jointly continuous in  $(t, \phi)$  on  $[0, \infty) \times \mathcal{E}$ .

The mapping  $Q_{\omega}$  is called the *Poincaré map* associated with this periodic semiflow.

- **Definition 3.3.** (i) In the case where  $\mathcal{H} = \mathbb{R}$ ,  $U(t, x + ct)$  is said to be an  $\omega$ -time *periodic traveling wave with speed c* of the semiflow  ${Q_t}_{t\in\mathcal{T}}$  if  $Q_t[U(0, \cdot)](x) =$  $U(t, x + ct)$  and  $U(t, x) = U(t + \omega, x)$  for all  $t \in \mathcal{T}$  and  $x \in \mathbb{R}$ .
- <span id="page-16-0"></span>(ii) In the case where  $\mathcal{H} = \mathbb{Z}$ ,  $U(t, x + ct)$  is said to be an  $\omega$ -time periodic traveling *wave with speed* c of  $\{Q_t\}_{t\in\mathcal{T}}$  if there exists a countable subset  $\Gamma \subset \mathbb{R}$  such that  $Q_t[U(0, \cdot+x)](0) = U(t, x+ct)$  for all  $t \in \mathcal{T}$  and  $x \in \mathbb{R}$ , and  $U(t, x) = U(t+\omega, x)$ for all  $t \in \mathcal{T}$  and  $x \in \mathbb{R} \setminus \Gamma$ .

**Theorem 3.3.** Let  $\beta(t)$  be a strongly positive  $\omega$ -time periodic orbit of  $\{Q_t\}_{t\in\mathcal{T}}$  restricted *to* X. Assume that  $Q := Q_\omega$  satisfies hypotheses (A1)–(A6) with  $\beta = \beta(0)$ . Then  ${Q_t}_{t\in\mathcal{T}}$  *admits a traveling wave*  $U(t, x + ct)$  *with*  $U(t, -\infty) = 0$  *and*  $U(t, \infty) = \beta(t)$ *uniformly for*  $t \in \mathcal{T}$ *. Furthermore,*  $U(t, x)$  *is nondecreasing in*  $x \in \mathbb{R}$ *.* 

*Proof.* Assume that  $\mathcal{H} = \mathbb{R}$ . Since  $Q_{\omega}$  satisfies (A1)–(A6), there exist  $c \in \mathbb{R}$  and a nondecreasing function  $\phi \in \mathcal{C}$  connecting 0 to  $\beta(0)$  such that  $Q_{\omega}[\phi](x) = \phi(x + c\omega)$ . Clearly,  $T_{c\omega}Q_{\omega}[\phi] = \phi$ . Define  $U(t, x) := T_{ct}Q_t[\phi](x)$ . Then  $U(t, x + ct) = Q_t[\phi](x) =$  $Q_t[U(0, \cdot)](x)$ , and

$$
U(t+\omega,x)=T_{ct+\omega}Q_{t+\omega}[\phi](x)=T_{ct}Q_tT_{\omega}Q_{\omega}[\phi](x)=T_{ct}Q_t[\phi](x)=U(t,x).
$$

Note that  $Q_t[\beta(0)] = \beta(t)$  and  $\phi$  is nondecreasing and connects 0 to  $\beta(0)$ . It follows that  $U(t, -\infty) = 0$  and  $U(t, \infty) = \beta(t)$ .

Assume now  $\mathcal{H} = \mathbb{Z}$ . Since  $Q_{\omega}$  satisfies (A1)–(A6), there exists  $c \in \mathbb{R}$ , a countable subset  $\Gamma \subset \mathbb{R}$  and a nondecreasing function  $\phi \in \mathcal{B}$  connecting 0 to  $\beta(0)$  such that  $\overline{Q}_{\omega}[\phi](x) = \phi(x + c\omega)$  for all  $x \in \mathbb{R} \backslash \Gamma$ . Clearly,  $T_{c\omega}\overline{Q}_{\omega}[\phi](x) = \phi(x)$  for all  $x \in \mathbb{R} \backslash \Gamma$ . Define  $U(t, x) := T_{ct} \tilde{Q}_t[\phi](x)$ . Thus,

$$
U(t, x + ct) = \tilde{Q}_t[\phi](x) = \tilde{Q}_t[U(0, \cdot)](x) = Q_t[U(0, \cdot + x)](0), \quad \forall x \in \mathbb{R},
$$

and

$$
U(t + \omega, x) = T_{ct + \omega} \tilde{Q}_{t + \omega}[\phi](x) = T_{ct} \tilde{Q}_t T_{c\omega} \tilde{Q}_{\omega}[\phi](x) = T_{ct} \tilde{Q}_t[\phi](x) = U(t, x)
$$

for all  $x \in \mathbb{R} \setminus \Gamma$ . Note that  $O_t[\beta(0)] = \beta(t)$  and  $\phi$  is nondecreasing and connects 0 to  $\beta(0)$ . Hence  $U(t, -\infty) = 0$  and  $U(t, \infty) = \beta(t)$ .

#### *3.4. Continuous-time semiflows in a continuous habitat*

In this subsection, we consider continuous-time semiflows in the continuous habitat  $\mathcal{H} = \mathbb{R}$ . Recall that a family  $\{Q_t\}_{t\geq 0}$  of mappings  $Q_t : \mathcal{E} \to \mathcal{E}$  is said to be a *semiflow* on a metric space  $\mathcal{E} \subset \mathcal{C}$  provided that

(1)  $Q_0[\phi] = \phi$  for all  $\phi \in \mathcal{E}$ .

(2)  $Q_t \circ Q_s[\phi] = Q_{t+s}[\phi]$  for all  $t, s \ge 0$  and  $\phi \in \mathcal{E}$ .

(3)  $Q_t[\phi]$  is jointly continuous in  $(t, \phi)$  on  $[0, \infty) \times \mathcal{E}$ .

Before moving to the study of traveling waves of  ${Q_t}_{t\ge0}$ , we first investigate the spatially homogeneous system, that is, the system restricted to X. Let  $\beta \gg 0$  be an equilibrium in X. For each  $t > 0$ , we use  $\Sigma_t$  to denote the set of all fixed points of  $Q_t$  restricted to  $\mathcal{X}_{\beta}$ . Clearly, the equilibrium set of the semiflow is  $\Sigma := \bigcap_{t>0} \Sigma_t$ , which is a subset of  $\Sigma_t$  for any  $t > 0$ . The subsequent result indicates that the instability of intermediate equilibria of the semiflow implies the unordering property of all intermediate fixed points of each time-t map.

<span id="page-17-0"></span>**Proposition 3.1.** *For any given*  $t > 0$ , *if the map*  $O_t$  *satisfies the bistability assumption*  $(AS<sup>7</sup>)$  *with*  $E = \Sigma$ *, then*  $Q_t$  *satisfies* (A5) *with*  $E = \Sigma_t$ *.* 

*Proof.* Let  $t_0 > 0$  be given. We first show that any two points  $u \in \Sigma \setminus \{0, \beta\}$  and  $v \in$  $\Sigma_{t_0} \setminus \{0, \beta\}$  are unordered. Assume for contradiction that u and v are ordered, say  $u < v$ . Then the eventual strong monotonicity implies that  $u \ll v$ . Since u is strongly unstable from above, there exist a unit vector  $e \in Int(\mathcal{X}^+)$  and  $\delta > 0$  such that  $Q_{t_0}[u+\delta e] \gg u+\delta e$ with  $u + \delta e \in [[u, v]]_{\mathcal{X}}$ . From [\[34,](#page-45-14) Theorem 1.2.1],  $(Q_{t_0})^n[u + \delta e]$  is eventually strongly increasing and converges to some  $\alpha \in \Sigma$ . Note that  $\alpha \in [[u, v]_{\mathcal{X}}]$  is strongly unstable from below. Hence, by the same arguments as in the proof of Proposition [2.1,](#page-5-1) we obtain a contradiction.

Next we show that  $\Sigma_{t_0} \setminus \Sigma$  is unordered. By the first step it suffices to prove that for any two ordered elements  $u < v$  in  $\Sigma_{t_0} \setminus \Sigma$ ,  $[u, v]$  $\chi \cap \Sigma \neq \emptyset$ . Indeed, by eventual strong monotonicity, we have  $u \ll v$ . Thus, we can choose a sequence  $\{u_n\}$  on the segment connecting u and v such that  $u \ll u_n \ll u_{n+1} \ll v$  for all  $n \ge 1$ . By [\[34,](#page-45-14) Theorem 1.3.7],  $\omega(u) \leq \omega(u_n) \leq \omega(u_{n+1}) \leq \omega(v)$  for all  $n \geq 1$ . Clearly,  $\omega(u) = \{Q_t u : t \in$ [0, t<sub>0</sub>]} and  $\omega(v) = \{Q_t v : t \in [0, t_0]\}$ , and hence  $u \leq \omega(u_n) \leq v$  for all  $n \geq 1$ . Note that  $\bigcup_{n\geq 1} \omega(u_n)$  is contained in the compact set  $\overline{Q_{t_0}[\mathcal{X}_{\beta}]}$ . In the compact metric space consisting of all nonempty compact subsets of  $\overline{Q_{t_0}[\mathcal{X}_{\beta}]}$  with Hausdorff distance  $d_H$ , the sequence { $\omega(u_n)$  :  $n \geq 1$ } has a convergent subsequence. Without loss of generality, we assume that for some nonempty compact set  $\varpi \subset \overline{Q_{t_0}[\mathcal{X}_{\beta}]}$ ,  $\lim_{n\to\infty} d_H(\omega(u_n), \varpi) = 0$ . Since each  $\omega(u_n)$  is invariant for the semiflow  ${Q_t}_{t\geq0}$ , so is the compact set  $\omega$ , that is,  $Q_t \varpi = \varpi$  for all  $t \ge 0$ . For any given  $x, y \in \omega$ , there exist sequences of points  $x_n, y_n \in \overline{\omega}(u_n)$  such that  $x_n \to x$  and  $y_n \to y$  as  $n \to \infty$ . Since  $\omega(u_n) \leq \omega(u_{n+1}),$ we have  $x_n \le y_{n+1}$  and  $y_n \le x_{n+1}$  for all  $n \ge 1$ . Letting  $n \to \infty$ , we then have  $x \le y$ and  $y \le x$ , and hence  $x = y$ . This implies that  $\varpi$  is a singleton, that is,  $\varpi = {\alpha}$ . By the invariance of  $\varpi$  for the semiflow,  $\alpha$  is an equilibrium. Since  $u < \omega(u_n) < v$  for all  $n \geq 1$ , it follows that  $\alpha \in [u, v]_X$ .

For a continuous-time semiflow  ${Q<sub>t</sub>}<sub>t>0</sub>$ , we need the following definition of traveling waves.

**Definition 3.4.**  $\psi(x+ct)$  with  $\psi \in C$  is said to be a *traveling wave with speed*  $c \in \mathbb{R}$  of the continuous-time semiflow { $Q_t$ <sub>t≥0</sub> if  $Q_t[\psi](x) = \psi(x + ct)$  for all  $x \in \mathbb{R}$  and  $t > 0$ . We say that  $\psi$  *connects* 0 *to*  $\beta$  if  $\psi(-\infty) = 0$  and  $\psi(\infty) = \beta$ .

<span id="page-18-0"></span>**Theorem 3.4.** Assume that for each  $t > 0$ , the map  $O_t$  satisfies assumptions (A1) and  $(A3)$ – $(A5)$  *with*  $E = \Sigma_t$ *, and the time-one map*  $Q_1$  *satisfies*  $(A6)$  *with*  $E = \Sigma$ *. Then there exists*  $c \in \mathbb{R}$  *such that*  $\{Q_t\}_{t>0}$  *admits a nondecreasing traveling wave with speed* c *and connecting* 0 *to*  $\beta$ .

*Proof.* Let  $e_0$ ,  $e_\beta$  and  $\delta$  be chosen as in [\(2.7\)](#page-7-0), [\(2.3\)](#page-6-2) and [\(3.1\)](#page-9-2), respectively. We proceed in three steps.

Firstly, we show that there exists  $s_k \downarrow 0$  such that each discrete semiflow  $\{Q_{s_k}^n\}_{n\geq 0}$ admits nondecreasing traveling waves  $\psi_{\pm,s_k}(x + c_{\pm,s_k}n)$  with  $c_{-,s_k} \geq c_{+,s_k}$  and

$$
0 < \psi_{-,s_k}(0) \leq \delta e_0 \quad \text{and} \quad \beta - \delta e_\beta \leq \psi_{+,s_k}(0) < \beta,
$$

but

$$
\psi_{-,s_k}(0) \notin [0, \delta e_0]]
$$
 $\chi$  and  $\psi_{+,s_k}(0) \notin [[\beta - \delta e_\beta, \beta] \chi$ .

Indeed, since for each  $s > 0$  the map  $Q_s$  satisfies (A1)–(A5), from the first two steps of the proof of Theorem [3.1](#page-11-0) we see that for the discrete semiflow  ${ (Q_s)^n}_{n \geq 0}$ , there exist two nondecreasing traveling waves  $\phi_{\pm,s}(x + c_{\pm,s}t)$  with the following properties:

- (1)  $\phi_{-,s}$  connects 0 to some  $\alpha_{-,s} \in E_s \setminus \{0\}$  and  $\phi_{+,s}$  connects some  $\alpha_{+,s} \in E_s \setminus \{\beta\}$ to  $\beta$ ;
- (2)  $\alpha_{-,s}$  and  $\alpha_{+,s}$  are ordered and  $c_{-,s} \geq c_{+,s}$ .

By a similar argument to that in [\[47,](#page-45-15) Theorem 1.3.7], both  $\alpha_{\pm,s}$  have subsequences  $\alpha_{\pm,s}$ which tend to two equilibria of the semiflow as  $s_k \to 0$ , say  $\alpha_-$  and  $\alpha_+$ . Since  $\alpha_{-,s}$  and  $\alpha_{+s}$  are ordered, it follows from Proposition [3.1](#page-17-0) that there are only three possibilities for the relation of  $\alpha_-\,$  and  $\alpha_+$ :

(i) 
$$
\beta = \alpha_- \ge \alpha_+
$$
; (ii)  $\alpha_- \ge \alpha_+ = 0$ ; (iii)  $\alpha_- = \alpha_+ \in E \setminus \{0, \beta\}.$ 

If  $\alpha_$  =  $\beta$ , then for sufficiently large k we can define

$$
a_{s_k} := \sup\{x \in \mathbb{R} : \phi_{-,s_k}(x) \in [0, \delta e_0]_{\mathcal{X}}\},
$$
  

$$
b_{s_k} := \inf\{x \in \mathbb{R} : \phi_{-,s_k}(x) \in [\beta - \delta e_\beta, \beta]_{\mathcal{X}}\}.
$$

Hence,  $\psi_{-,s}(x) := \psi_s(x + a_s)$  and  $\psi_{+,s}(x) := \psi_s(x + b_s)$  are the required traveling waves. If  $\alpha_+ = 0$ , then for sufficiently large k we can define

$$
a_{s_k} := \sup\{x \in \mathbb{R} : \phi_{+,s_k}(x) \in [0, \delta e_0] \chi\},
$$
  

$$
b_{s_k} := \inf\{x \in \mathbb{R} : \phi_{+,s_k}(x) \in [\beta - \delta e_\beta, \beta] \chi\}.
$$

Hence,  $\psi_{-,s_k}(x) := \phi_{+,s_k}(x + a_{s_k})$  and  $\psi_{+,s_k}(x) := \phi_{+,s_k}(x + b_{s_k})$  are the required traveling waves. If  $\alpha_$  =  $\alpha_$  ∈  $E \setminus \{0, \beta\}$ , then by Lemma [3.1](#page-9-1) we have  $\alpha_$  =  $\alpha_$  ∈  $E \setminus \{[0, \delta e_0] \times \bigcup [\beta - \delta e_\beta, \beta] \times \}$ . Consequently, for sufficiently large k we can define

$$
a_{s_k} := \sup\{x \in \mathbb{R} : \phi_{-,s_k}(x) \in [0, \delta e_0]_{\mathcal{X}}\},
$$
  

$$
b_{s_k} := \inf\{x \in \mathbb{R} : \phi_{-,s_k}(x) \in [\beta - \delta e_\beta, \beta]_{\mathcal{X}}\}.
$$

Hence,  $\psi_{-,s_k}(x) := \phi_{-,s_k}(x + a_{s_k})$  and  $\psi_{+,s_k}(x) := \phi_{-,s_k}(x + b_{s_k})$  are the required traveling waves.

Secondly, we show that there exists a subsequence, still denoted by  $s_k$ , such that  $\psi_{\pm,s_k} \to \psi_{\pm}$  in  $\mathcal{C}_{\beta}$  and  $(1/s_k)c_{\pm,s_k} \to c_{\pm} \in \mathbb{R}$ . Indeed, for each  $s_k > 0$ , there exists an integer  $m_k > 0$  such that  $m_k s_k > 2$ . Then

<span id="page-19-0"></span>
$$
\psi_{-,s_k} = T_{m_k c_{s_k}} \circ Q_{m_k s_k} [\psi_{-,s_k}] = Q_2 \circ Q_{m_k s_k - 2} \circ T_{m_k c_{s_k}} [\psi_{-,s_k}] \in Q_1 \circ Q_1 [\mathcal{C}_{\beta}]. \tag{3.6}
$$

Clearly, the compactness of  $Q_1$  implies that  $Q_1 \circ Q_1[\mathcal{C}_{\beta}]$  is precompact in  $\mathcal{C}_{\beta}$ . Thus, there exists a subsequence, still denoted by  $s_k$ , and nonincreasing functions  $\psi_-, \psi_+ \in C_\beta$  with  $0 < \psi_-(0) \leq \delta e_0$  and  $\beta - \delta e_\beta \leq \psi_+(0) < \beta$  such that  $\psi_{-,s_k} \to \psi_-$  and  $\psi_{+,s_k} \to \psi_+$ in  $C_{\beta}$ . Also we claim that  $\psi_{\pm,s_k}(\pm \infty)$  all exist. Indeed, from [\(3.6\)](#page-19-0) there exists  $\phi_{s_k} \in C_{\beta}$ such that  $Q_1 \circ Q_1[\phi_{s_k}] \to \psi_-\$ . Note that  $\{Q_1[\phi_{s_k}]\}_{k\geq 1}$  also has a convergent subsequence with limit  $\phi \in C_\beta$ . Hence, by the uniqueness of limit we have  $Q_1[\phi] = \psi_-\$ . Note that  $\psi_{-}(k) = Q_1[\phi](k) = Q_1[\phi(\cdot + k)](0)$  and  $\{Q_1[\phi(\cdot + k)]\}_{k \ge 1}$  has a convergent subsequence. Hence,  $\psi_-(\pm\infty)$  exist because  $\psi_-$  is nonincreasing. Similarly,  $\psi_+(\pm\infty)$  exist. Also,

$$
\psi_{-}(-\infty) \le \psi_{-}(0) \le \delta e_0 \quad \text{and} \quad \psi_{+}(\infty) \ge \psi_{+}(0) \ge \beta - \delta e_{\beta},\tag{3.7}
$$

but

<span id="page-20-2"></span>
$$
\psi_{-}(0) \notin [0, \delta e_0] | \chi
$$
 and  $\psi_{+}(0) \notin [[\beta - \delta e_{\beta}, \beta] \chi$ .

Consequently, by the monotonicity of  $\psi_{\pm}$ ,

$$
\psi_{-}(x) \notin [0, \delta e_0] |_{\mathcal{X}}, \forall x > 0, \quad \text{and} \quad \psi_{+}(x) \notin [[\beta - \delta e_{\beta}, \beta]_{\mathcal{X}}, \forall x < 0. \tag{3.8}
$$

Since  $\psi$  and  $\psi$  + are the limits of sequences of monotone functions with different translations, we can employ the same arguments as in the second step of the proof of Theo-rem [3.1](#page-11-0) to show that  $\psi_-(\infty)$  and  $\psi_+(-\infty)$  are ordered.

To prove that  $(1/s_k)c_{\pm,s_k}$  have convergent subsequences, we only need to prove that  $(1/s_k)c_{-,s_k}$  is bounded above and  $(1/s_k)c_{+,s_k}$  is bounded below because  $c_{-,s_k} \ge c_{+,s_k}$ . Assume for contradiction that some subsequence, still say  $(1/s_k)c_{-,s_k}$ , tends to  $\infty$ . Note that for each  $s > 0$  there exists  $n_s \in \mathbb{Z}^+$  such that the integer part of  $1/s$ , denoted by  $\langle 1/s \rangle$ , equals  $n_s$  and  $1/(n_s + 1) < s \leq 1/n_s$ . Hence,  $s \langle 1/s \rangle \rightarrow 1$  as  $s \rightarrow 0$ . Then

<span id="page-20-0"></span>
$$
\lim_{k \to \infty} \left\langle \frac{1}{s_k} \right\rangle c_{-,s_k} = \lim_{k \to \infty} \frac{1}{s_k} c_{-,s_k} \cdot s_k \left\langle \frac{1}{s_k} \right\rangle = \lim_{k \to \infty} \frac{1}{s_k} c_{-,s_k} = \infty.
$$

Thus, using the first observation in  $(3.8)$ , we have

$$
Q_1[\delta e_0] \ge Q_1[\psi_-(-\infty)] = Q_1[\psi_-(-\infty)](0) = \lim_{x \to -\infty} Q_1[\psi_-(+x)](0)
$$
  
\n
$$
= \lim_{x \to -\infty} Q_1[\psi_-](x) = \lim_{x \to -\infty} \lim_{k \to \infty} (Q_{s_k})^{(1/s_k)}[\psi_{-,s_k}](x)
$$
  
\n
$$
= \lim_{x \to -\infty} \lim_{k \to \infty} \psi_{-,s_k} \left(x + \left\langle \frac{1}{s_k} \right\rangle c_{-,s_k}\right) \ge \lim_{x \to -\infty} \lim_{y \to \infty} \lim_{k \to \infty} \psi_{-,s_k}(y)
$$
  
\n
$$
= \lim_{y \to \infty} \psi_-(y) = \psi_-(\infty) \notin [0, \delta e_0]]x, \tag{3.9}
$$

which contradicts  $Q_1[\delta e_0] \ll \delta e_0$ . Similarly, if  $(1/s_k)c_{+,s_k} \to -\infty$ , then the second observation in  $(3.8)$  implies that

<span id="page-20-1"></span>
$$
Q_1[\beta - \delta e_{\beta}] \leq Q_1[\psi_+(\infty)] = Q_1[\psi_+(\infty)](0) = \lim_{x \to \infty} Q_1[\psi_+(\cdot + x)](0)
$$
  
\n
$$
= \lim_{x \to \infty} Q_1[\psi_+](x) = \lim_{x \to \infty} \lim_{k \to \infty} (Q_{s_k})^{(1/s_k)}[\psi_{+,s_k}](x)
$$
  
\n
$$
= \lim_{x \to \infty} \lim_{k \to \infty} \psi_{+,s_k} \left(x + \left\langle \frac{1}{s_k} \right\rangle c_{+,s_k}\right) \leq \lim_{x \to \infty} \lim_{y \to -\infty} \lim_{k \to \infty} \psi_{+,s_k}(y)
$$
  
\n
$$
= \lim_{y \to -\infty} \psi_+(y) = \psi_+(-\infty) \notin [[\beta - \delta e_{\beta}, \beta]_X, \tag{3.10}
$$

which contradicts  $Q_1[\beta - \delta e_{\beta}] \gg \beta - \delta e_{\beta}$ . Consequently,  $(1/s_k)c_{\pm, s_k}$  are bounded.

Finally, we show that either  $\psi_-(x + c_+t)$  or  $\psi_+(x + c_+t)$  established in the second step is a traveling wave connecting 0 to  $\beta$ . Indeed, for any  $t > 0$ , there exist  $m_k \in \mathbb{Z}$  and  $r_k \in [0, s_k)$  such that  $t = m_k s_k - r_k$ . Clearly,  $r_k \to 0$  as  $k \to \infty$ . Then

$$
Q_t[\psi_{\pm}] = \lim_{k \to \infty} Q_{t+r_k}[\psi_{\pm,s_k}] = \lim_{k \to \infty} Q_{m_k s_k}[\psi_{\pm,s_k}] = \lim_{k \to \infty} \psi_{\pm,s_k}(\cdot + m_k c_{\pm,s_k})
$$
  
= 
$$
\lim_{k \to \infty} \psi_{\pm,s_k}(\cdot + (t+r_k)\frac{1}{s_k}c_{\pm,s_k}) = \psi_{\pm}(\cdot + c_{\pm}t),
$$

where the last equality follows from Proposition [7.2\(](#page-41-1)2). From the equality  $Q_t[\psi_+] =$  $\psi_{\pm}(\cdot + ct)$  for all  $t \ge 0$ , we see that  $\psi(\pm \infty)$  are equilibria. Recall that  $\psi_{-}(-\infty) \le$  $\delta e_0 \leq \psi(\infty)$  and  $\psi_+(\infty) \geq \beta - \delta e_\beta \geq \psi_+(-\infty)$ . Then  $\psi_-(-\infty) = 0$ ,  $\psi_+(\infty) = \beta$ , and there are only three possibilities for  $\psi_{-}(\infty)$  and  $\psi_{+}(-\infty)$ :

- (i)  $\beta = \psi_{-}(\infty) > \psi_{+}(-\infty);$
- (ii)  $\psi_{-}(\infty) > \psi_{+}(-\infty) = 0;$
- (iii)  $\psi_{-}(\infty) = \alpha = \psi_{+}(-\infty)$  for some  $\alpha \in \Sigma \setminus \{0, \beta\}.$

Since the time-one map  $Q_1$  satisfies (A6) with  $E = \Sigma$ , we can employ the same arguments as in the proof of Lemma [3.4](#page-10-2) to exclude (iii). Thus, either (i) or (ii) holds, completing the proof.  $\Box$ 

#### *3.5. Continuous-time semiflows in a discrete habitat*

In this case,  $\mathcal{T} = \mathbb{R}^+$  and  $\mathcal{H} = \mathbb{Z}$ . Let  $\beta \gg 0$  be an equilibrium of the semiflow  $\{Q_t\}_{t \geq 0}$ . We start with the definition of traveling waves for this case.

<span id="page-21-1"></span>**Definition 3.5.**  $\psi(i+ct)$  with  $\psi \in \mathcal{B}_{\beta}$  is said to be a *traveling wave with speed*  $c \in \mathbb{R}$  of the continuous-time semiflow  $\{Q_t\}_{t\geq 0}$  if  $Q_t[\psi](i) = \psi(i + ct)$  for all  $i \in \mathbb{Z}$  and  $t \geq 0$ . Clearly,  $\psi$  is continuous if  $c \neq 0$ .

For each  $t > 0$ , define  $\tilde{Q}_t : \mathcal{B}_\beta \to \mathcal{B}_\beta$  by  $\tilde{Q}_s[\phi](x) = Q_s[\phi(\cdot + x)](0)$ . Then it is easy to see the following result holds.

**Lemma 3.7.**  $\{\tilde{Q}_t\}_{t\geq 0}$  *has the following properties:* 

- (i)  $\tilde{Q}_0[\phi] = \phi$  *for all*  $\phi \in \mathcal{B}$ .
- (ii)  $\widetilde{Q}_t \circ \widetilde{Q}_s[\phi] = \widetilde{Q}_{t+s}[\phi]$  *for all t*,  $s \geq 0$  *and*  $\phi \in \mathcal{B}$ .
- (iii) *For fixed*  $x \in \mathbb{R}$ *, if*  $t_n \to t$  *and*  $\phi_n(i + x) \to \phi(i + x)$  *in*  $\mathcal X$  *for any*  $i \in \mathbb{Z}$ *, then*  $\tilde{Q}_{t_n}[\phi_n](x) \to \tilde{Q}_t[\phi](x)$  in X.

<span id="page-21-0"></span>We combine the ideas in the proofs of Theorems [3.2](#page-14-1) and [3.4](#page-18-0) to prove the following result for continuous-time semiflows in a discrete habitat.

**Theorem 3.5.** Let  $\mathcal{X} = C(M, \mathbb{R}^d)$ . Assume that for each  $t > 0$ , the map  $Q_t$  satisfies (A1) and (A3)–(A5) with  $E = \Sigma_t$ , and the time-one map  $Q_1$  satisfies (A6) with  $E = \Sigma$ . *Then there exists*  $c \in \mathbb{R}$  *such that*  $\{Q_t\}_{t>0}$  *admits a nondecreasing traveling wave with speed* c *and connecting* 0 *to* β*.*

*Proof.* Let  $\delta$ ,  $e_0$ ,  $e_\beta$  be chosen as in [\(3.1\)](#page-9-2), [\(2.7\)](#page-7-0) and [\(2.3\)](#page-6-2). We proceed in three steps.

Firstly, since for any  $s > 0$  the map  $Q_s$  satisfies (A1)–(A5), the proofs of Theorems [3.2](#page-14-1) and [3.4](#page-18-0) show that there exists  $s_k \downarrow 0$  such that  $\{ (Q_{s_k})^n \}_{n \geq 0}$  admits two nondecreasing traveling waves  $\tilde{\psi}_{\pm,s_k}(x + c_{\pm,s_k}n)$  with  $c_{-,s_k} \geq c_{+,s_k}$ , that is, there exists a countable subset  $\Theta_k$  such that

$$
\tilde{Q}_{s_k}[\tilde{\psi}_{\pm,s_k}](x) = \tilde{\psi}_{\pm,s_k}(x + c_{\pm,s_k}), \quad \forall x \in \mathbb{R} \setminus \Theta_k.
$$

Furthermore,  $\tilde{\psi}_{-,s_k}$  is left continuous and  $\tilde{\psi}_{+,s_k}$  is right continuous, with

$$
0 < \tilde{\psi}_{-,s_k}(0) \leq \delta e_0 \quad \text{and} \quad \beta - \delta e_\beta \leq \tilde{\psi}_{+,s_k}(0) < \beta,
$$

but

$$
\tilde{\psi}_{-,s_k}(0) \notin [0, \delta e_0]]_{\mathcal{X}}
$$
 and  $\tilde{\psi}_{+,s_k}(0) \notin [[\beta - \delta e_\beta, \beta]_{\mathcal{X}}$ .

Secondly, we show that for the above sequence  $s_k$ , there exist a countable set  $\tilde{\Gamma} \subset \mathbb{R}$ and a subsequence, still denoted by  $s_k$ , such that  $(1/s_k)c_{\pm,s_k} \to c_{\pm} \in \mathbb{R}$  and  $\tilde{\psi}_{\pm,s_k}(x)$ converges in X for all  $x \in \mathbb{R} \setminus \tilde{\Gamma}$ . Indeed, let  $\Theta = \bigcup_{k=1}^{\infty} \Theta_k$ . Hence,  $\Theta$  is countable and

$$
\tilde{Q}_{s_k}[\tilde{\psi}_{\pm,s_k}](x) = \tilde{\psi}_{\pm,s_k}(x + c_{\pm,s_k}), \quad \forall k \ge 1, \ x \in \mathbb{R} \setminus \Theta.
$$

From Proposition [7.4,](#page-41-4) there exists another countably dense set  $\Gamma \subset \mathbb{R}$  such that  $\Gamma \cap \Theta$  $= \emptyset$ . By the same arguments as in the proof of Theorem [3.2,](#page-14-1) the limits

$$
\tilde{\psi}_{-}(x) := \lim_{y \in \Gamma, y \uparrow x} \lim_{k \to \infty} \tilde{\psi}_{-,s_k}(y), \quad \tilde{\psi}_{+}(x) := \lim_{y \in \Gamma, y \downarrow x} \lim_{k \to \infty} \tilde{\psi}_{+,s_k}(y), \quad \forall x \in \mathbb{R},
$$

are well defined and all  $\tilde{\psi}_{\pm}(\pm\infty)$  exist. Furthermore,  $\tilde{\psi}_{-}(\infty)$  and  $\tilde{\psi}_{+}(-\infty)$  are ordered in  $X$ , and

$$
\tilde{\psi}_{-}(-\infty) \leq \psi_{-}(0) \leq \delta e_0
$$
 and  $\tilde{\psi}_{+}(\infty) \geq \tilde{\psi}_{+}(0) \geq \beta - \delta e_{\beta}$ ,

but

<span id="page-22-0"></span>
$$
\tilde{\psi}_{-}(0) \notin [0, \delta e_0] | \chi
$$
 and  $\tilde{\psi}_{+}(0) \notin [[\beta - \delta e_{\beta}, \beta] \chi$ .

Further,  $\tilde{\psi}_{\pm}(x^{\pm})$  exist for all  $x \in \mathbb{R} \setminus \Gamma$ . Hence, Theorem [7.1](#page-42-0) yields a countable subset  $\tilde{\Gamma}$ of  $\mathbb R$  such that

$$
\tilde{\psi}_{\pm,s_k}(x) \to \tilde{\psi}_{\pm}(x), \quad \forall x \in \mathbb{R} \setminus \tilde{\Gamma}.
$$
\n(3.11)

By similar arguments to the second step of the proof of Theorem [3.4,](#page-18-0)  $(1/s_k)c_{\pm,s_k}$  are bounded.

Finally, we prove that either  $\tilde{\psi}_-(x+c_-t)$  or  $\tilde{\psi}_+(x+c_+t)$  is a nondecreasing traveling wave connecting 0 to  $\beta$ . Indeed, from [\(3.11\)](#page-22-0) and Proposition [7.3,](#page-41-2) there exists a countable subset  $\Gamma_1$  of  $\mathbb R$  such that

$$
\tilde{\psi}_{\pm,s_k}(i+x) \to \tilde{\psi}_{\pm}(i+x), \quad \forall i \in \mathbb{Z}, \ x \in \mathbb{R} \setminus \Gamma_1.
$$

Hence, for any  $x \in \mathbb{R} \setminus \Gamma_1$  and  $t > 0$ ,

<span id="page-23-0"></span>
$$
\tilde{Q}_{t}[\tilde{\psi}_{-}](x) = Q_{t}[\tilde{\psi}_{-}(\cdot + x)](0) = \lim_{k \to \infty} Q_{t+r_{k}}[\tilde{\psi}_{-,s_{k}}(\cdot + x)](0)
$$
\n
$$
= \lim_{k \to \infty} Q_{m_{k}s_{k}}[\tilde{\psi}_{-,s_{k}}(\cdot + x)](0) = \lim_{k \to \infty} \tilde{Q}_{m_{k}s_{k}}[\tilde{\psi}_{-,s_{k}}](x)
$$
\n
$$
= \lim_{k \to \infty} (\tilde{Q}_{s_{k}})^{m_{k}}[\tilde{\psi}_{-,s_{k}}](x) = \lim_{k \to \infty} \tilde{\psi}_{-,s_{k}}(x + m_{k}c_{-,s_{k}})
$$
\n
$$
= \lim_{k \to \infty} \tilde{\psi}_{-,s_{k}}\left(x + (t + r_{k})\frac{1}{s_{k}}c_{-,s_{k}}\right).
$$
\n(3.12)

If  $c = 0$ , we can choose  $x_0$  such that

$$
Q_t[\tilde{\psi}_-(x_0+\cdot)](i)=\lim_{k\to\infty}\tilde{\psi}_{-,s_k}\bigg(x+(t+r_k)\frac{1}{s_k}c_{-,s_k}\bigg)=\tilde{\psi}_-(x_0+i),\quad\forall i\in\mathbb{Z}.
$$

If  $c_-\neq 0$ , there exists a countable subset  $\Gamma_2$  of R such that

$$
\tilde{Q}_t[\tilde{\psi}_-](x) = \lim_{k \to \infty} \tilde{\psi}_{-,s_k}\bigg(x + (t + r_k)\frac{1}{s_k}c_{-,s_k}\bigg) = \tilde{\psi}_-(x + c_-t), \quad \forall x \notin \Gamma_1, \ x + c_-t \in \Gamma_2.
$$

Without loss of generality, we assume that  $c_$  > 0. For any  $y \in \mathbb{R}$ , we can choose  $x_0 \in \mathbb{R}$ and  $t_0 \ge 0$  such that  $x_0 + c_0 = y$  and  $\tilde{\psi}_-(x)$  is continuous at  $x = x_0 + i$  for all  $i \in \mathbb{Z}$ . Now one can find  $x_{\pm,k} \in \mathbb{R} \setminus \Gamma_1$  and  $t_{\pm,k} \to t_0$  with  $y_{\pm,k} := x_{\pm} + c_- t_{\pm,k} \in \mathbb{R} \setminus \Gamma_2$  such that  $y_{-,k}$  ↑ y and  $y_{+,k}$  ↓ y. Note that

$$
\tilde{\psi}_{-}(y^{-}) := \lim_{k \to \infty} \tilde{\psi}_{-}(y_{-,k}) = \lim_{k \to \infty} Q_{t_{-,k}}[\tilde{\psi}_{-}(\cdot + x_{-,k})](0) = Q_{t_0}[\tilde{\psi}_{-}(\cdot + x_0)](0),
$$
  

$$
\tilde{\psi}_{-}(y^{+}) := \lim_{k \to \infty} \tilde{\psi}_{-}(y_{+,k}) = \lim_{k \to \infty} Q_{t_{+,k}}[\tilde{\psi}_{-}(\cdot + x_{+,k})](0) = Q_{t_0}[\tilde{\psi}_{-}(\cdot + x_0)](0).
$$

Thus,  $\tilde{\psi}_-(x)$  is continuous in  $x \in \mathbb{R}$ . Hence, again by Proposition [7.5](#page-41-3) and [\(3.12\)](#page-23-0), we have  $\tilde{Q}_t[\tilde{\psi}_-](x) = \tilde{\psi}_-(x+c_-t)$  for all  $x \in \mathbb{R}$  and  $t \geq 0$ . Therefore,  $\tilde{\psi}_-(x+c_-t)$  is a traveling wave connecting 0 to some  $\alpha_-\in \Sigma\setminus\{0\}$ . Similarly, we can construct the traveling wave  $\psi_+(x+c_+t)$  connecting some  $\alpha_+ \in \Sigma \setminus \{\beta\}$  to  $\beta$ . Moreover,  $\alpha_-$  and  $\alpha_+$  are ordered. The rest of the proof is essentially the same as for Theorem  $3.4$ .

## 4. Semiflows in a periodic habitat

A typical example of evolution systems in a periodic habitat is

<span id="page-23-1"></span>
$$
u_t = (d(x)u_x)_x + f(u), \quad t > 0, \ x \in \mathbb{R}, \tag{4.1}
$$

where  $d(x)$  is a positive periodic function of  $x \in \mathbb{R}$ . Under the assumption that f has exactly three ordered zeros  $0 < a < 1$  with  $f'(0) < 0, f'(a) > 0$  and  $f'(1) < 0$ , Xin [\[43\]](#page-45-7) employed perturbation methods to obtain the existence of a spatially periodic traveling wave  $V(x + ct, x)$  with  $V(-\infty, \cdot) = 0$  and  $V(\infty, \cdot) = 1$  provided that  $d(x)$ is sufficiently close to a positive constant in a certain sense (see also [\[42\]](#page-45-9)). For a general positive periodic function  $d(x)$ , the existence of such a traveling wave remains open. We will revisit this problem in Subsection 6.3. Below we extend the previous results in homogeneous habitats to periodic ones.

A map  $Q : \mathcal{E} \to \mathcal{E} \subset \mathcal{C}$  is said to be *spatially periodic* with a positive period  $r \in \mathcal{H}$  if  $Q \circ T_r = T_r \circ Q$ , where  $T_r$  is the *r*-translation operator. Similarly, a semiflow  $\{Q_t\}_{t \in \mathcal{T}}$  on  $\mathcal{E} \subset \mathcal{C}$  is said to be spatially periodic with a positive period  $r \in \mathcal{H}$  if  $Q_t \circ T_r = T_r \circ Q_t$ for all  $t \in \mathcal{T}$ , where  $\mathcal{H} = \mathbb{Z}$  or  $\mathbb{R}$  and  $\mathcal{T} = \mathbb{Z}^+$  or  $\mathbb{R}^+$ .

**Definition 4.1.** (i) An *r*-periodic function  $\beta(x)$  is said to be an *r*-periodic steady state of the map Q [semiflow  $\{Q_t\}_{t\in\mathcal{T}}$ ] if  $Q[\beta] = \beta [Q_t[\beta] = \beta$  for all  $t \in \mathcal{T}$ ].

(ii)  $V(x + ct, x)$  is said to be a *spatially r-periodic traveling wave with speed* c of the semiflow  $\{Q_t\}_{t\in\mathcal{T}}$  if  $Q_t[V(\cdot, \cdot)](x) = V(x + ct, x)$  and  $V(\cdot, x)$  is *r*-periodic in *x*. Moreover, we say that  $V(\xi, x)$  *connects* 0 *to*  $\beta(x)$  if  $\lim_{\xi \to -\infty} ||V(\xi, x)||_{\mathcal{X}} = 0$  and  $\lim_{\xi \to \infty} ||V(\xi, x) - \beta(x)||_{\mathcal{X}} = 0$  uniformly for  $x \in \mathcal{H}$ .

Motivated by [\[26,](#page-44-13) Section 5], we can regard a spatially periodic semiflow on  $\mathcal{E} \subset \mathcal{C}$  as a spatially homogeneous semiflow on another phase space. For any positive  $h \in \mathcal{H}$ , define  $[0, h]$  $\mathcal{H} := \{l \in \mathcal{H} : 0 \le l \le h\}$ . We use  $\mathcal{Y}$  to denote  $C([0, r]_{\mathcal{H}}, \mathcal{X})$  and  $\mathcal{S}$  to denote the set of all bounded functions from  $r\mathbb{Z}$  to  $\mathcal{Y}$ . Clearly,  $\mathcal{Y}$  can be regarded as a subspace of S. Let  $\mathcal{Y}^+ = C([0, r]_{\mathcal{H}}, \mathcal{X}^+)$  and  $\mathcal{S}^+$  be the set of all bounded functions from  $r\mathbb{Z}$ to  $\mathcal{Y}^+$ . We equip  $\mathcal Y$  with the norm  $||u||_{\mathcal{Y}} = \max{||u(x)||_{\mathcal{X}} : x \in [0, r]_{\mathcal{H}}}\}$ , and S with the compact-open topology. Thus,  $\mathcal Y$  is a Banach lattice with the norm  $\|\cdot\|_{\mathcal Y}$  and the cone  $\mathcal Y^+$ . Let

$$
\mathcal{K} := \{ f \in \mathcal{S} : f(ri)(r) = f(r(i+1))(0), \ \forall i \in \mathbb{Z} \}.
$$

It is easy to see that

$$
\mathcal{K} \cap \mathcal{Y} = \{ f \in \mathcal{S} : f(ri) \equiv f(ri) \text{ and } f(ri)(0) = f(ri)(r), \forall i, j \in \mathbb{Z} \}.
$$

For any  $\phi \in \mathcal{C}$ , define  $\tilde{\phi} \in \mathcal{S}$  by

<span id="page-24-0"></span>
$$
\tilde{\phi}(ri)(y) = \phi(ri + y), \quad \forall i \in \mathbb{Z}, \ y \in [0, r]_{\mathcal{H}}.
$$

Then we have the following observation.

**Lemma 4.1.** *For any*  $f \in \mathcal{K}$ , there exists a unique  $\phi_f \in \mathcal{C}$  such that  $\tilde{\phi}_f = f$ . Further, if  $f \in \mathcal{K} \cap \mathcal{Y}$ , then  $\phi_f$  is *r*-periodic.

*Proof.* For any  $x \in \mathcal{H}$ , we can find  $i \in \mathbb{Z}$  and  $y \in [0, r]$  with that  $x = ri + y$ . It is easy to see that this decomposition of x is unique when  $x \in \mathcal{H} \setminus r\mathbb{Z}$ , and is in two possible ways when  $x \in r\mathbb{Z}$ . More precisely, each  $x \in r\mathbb{Z}$  can be decomposed into either  $x = r(i + 1) + 0$  or  $x = ri + r$  for some  $i \in \mathbb{Z}$ . Note that  $f(r(i + 1))(0) = f(ri)(r)$ . Thus  $\phi_f(x) = \phi_f(r_i + y) := f(r_i)(y)$  is a well defined function in C. Clearly,  $\phi_f = f$ . If  $f(ri) = u$  for all  $i \in \mathbb{Z}$ , then  $\phi_f(ri) = u$  for all  $i \in \mathbb{Z}$ , which implies that  $\phi_f$  is  $r$ -periodic.  $\Box$ 

If we define  $F : \mathcal{C} \to \mathcal{K}$  by  $F(\phi) = \phi$ , then F is a homeomorphism between C and K. Let  $\beta(x)$  be a strongly positive r-periodic steady state of  $\{Q_t\}_{t\geq0}$ . With a little abuse of notation, we use  $C_\beta$  to denote the set { $\phi \in C : 0 \le \phi \le \beta$ }. Now we can define a semiflow  ${P_t}_{t \in \mathcal{T}}$  on  $\mathcal{K}_{\tilde{\beta}} := {f \in \mathcal{K} : 0 \le f \le \tilde{\beta}}$  by

<span id="page-25-1"></span>
$$
P_t[f] = F \circ Q_t[\phi_f], \quad \forall f \in \mathcal{K}_{\tilde{\beta}}, \ t \in \mathcal{T}.
$$
 (4.2)

Clearly,  $P_t \circ F = F \circ Q_t$  for all  $t \in \mathcal{T}$ , which implies that the semiflows  $\{Q_t\}_{t \in \mathcal{T}}$  and  ${P_t}_{t\in\mathcal{T}}$  are topologically conjugate. Moreover,  ${P_t}_{t\in\mathcal{T}}$  is spatially homogeneous and  $\tilde{\beta}$  is its equilibrium. Thus,  $\{Q_t\}_{t\in\mathcal{T}}$  on  $\mathcal{C}_{\beta}$  has a spatially *r*-periodic traveling wave if  ${P_t}_{t\in\mathcal{T}}$  on  $\mathcal{K}_{\tilde{\beta}}$  has a traveling wave. Before stating the main result, we first introduce the bistability assumption. Let  $\beta(x) \gg 0$  be an r-periodic steady state of  $\{Q_t\}_{t\in\mathcal{T}}$ . Assume that 0 is a trivial steady state. Define

$$
\Pi_{\beta} := \{ \phi \in \mathcal{C} : \phi(x) = \phi(x+r), \ 0 \le \phi(x) \le \beta(x), \ \forall x \in \mathcal{H} \}.
$$

As in Definition [2.1,](#page-4-1) we can define the strong stability of periodic steady states for a map Q in the space of periodic functions.

<span id="page-25-0"></span>**Definition 4.2.** A steady state  $\alpha \in \Pi_{\beta}$  is said to be *strongly stable from below* for the map  $Q: \Pi_{\beta} \to \Pi_{\beta}$  if there exist  $\delta_{\alpha}^+ > 0$  and a strongly positive element  $e_{\alpha}^+ \in \Pi_{\beta}$  such that

$$
Q[\alpha - \eta e_{\alpha}^{+}] \gg \alpha - \eta e_{\alpha}^{+}, \quad \forall \eta \in (0, \delta_{\alpha}^{+}].
$$
 (4.3)

The strong instability from below is defined by reversing the inequality [\(4.3\)](#page-25-0). Similarly, we can define strong stability (and instability) from above.

We need the following bistability assumption on the spatially  $r$ -periodic map  $Q$ .

 $(A5'')$  (*Bistability*) 0 and  $\beta \gg 0$  are strongly stable *r*-periodic steady states from above and below, respectively, for  $Q : \Pi_{\beta} \to \Pi_{\beta}$ , and the set of all intermediate r-periodic steady states is totally unordered in  $\Pi_{\beta}$ .

We note that a sufficient condition for the unordering property of all intermediate r-periodic steady states is:  $Q : \Pi_{\beta} \to \Pi_{\beta}$  is eventually strongly monotone and all intermediate fixed points are strongly unstable from both above and below.

<span id="page-25-2"></span>**Theorem 4.1.** Let  $X = C(M, \mathbb{R}^d)$ . Assume that for any  $t > 0$ , the map  $Q_t$  satisfies (A2)–(A4) and the bistability assumption (A5<sup>n</sup>). Further, assume that the map  $P_1 :=$  $FQ_1F^{-1}$  satisfies assumption (A6) with C and  $\beta$  replaced by K and  $\tilde{\beta}$ , respectively. Then *the spatially r-periodic semiflow*  ${Q_t}_{t\in\mathcal{T}}$  *has an r-periodic traveling wave*  $V(x, x + ct)$ *. Moreover,*  $V(x, \xi)$  *is nondecreasing in*  $\xi$  *and connects* 0 *to*  $\beta(x)$ *.* 

*Proof.* Fix  $t \ge 0$  and define  $P_t$  as in [\(4.2\)](#page-25-1). Then it is easy to see that  $P_t$  satisfies (A1)–(A5) with  $C_{\beta}$  replaced by  $K_{\tilde{\beta}}$ . From Theorems [3.2](#page-14-1) and [3.5,](#page-21-0)  $\{P_t\}_{t \in \mathcal{T}}$  admits a traveling wave  $U(x + ct)$  with U connecting 0 to  $\tilde{\beta}$ . By the definitions of traveling waves in a discrete habitat (see Definitions [3.2](#page-13-0) and [3.5\)](#page-21-1), we can find  $x_0 \in \mathbb{R}$  such that  $g := U(\cdot + x_0) \in \mathcal{K}_{\tilde{\beta}}$ and  $P_t[g](ri) = U(ri + ct + x_0)$  for all  $i \in \mathbb{Z}$ . By Lemma [4.1,](#page-24-0) we can find  $\psi, h_t \in \mathbb{C}$ such that  $\tilde{\psi} = g$  and  $\tilde{h}_t = U(\cdot + ct + x_0)$ , and hence  $P_t[\tilde{\psi}] = \tilde{h}_t$ . By the topological conjugacy of  $Q_t$  and  $P_t$ , we have  $Q_t[\psi] = h_t$ . Note that  $\tilde{\psi} = g = U(\psi + x_0) = \tilde{h}_0$ . Lemma [4.1](#page-24-0) yields  $\psi = h_0$ . If  $c = 0$ , then  $Q_t[\psi] = h_t \equiv h_0 = \psi$ , which implies that  $\psi$  is a traveling wave with speed zero. If  $c \neq 0$ , then we define  $V(\xi, x) := h_{(\xi-x)/c}(x)$ . Consequently,

$$
V(x+ct, x) = h_t(x) = Q_t[\psi](x) = Q_t[h_0](x) = Q_t[V(\cdot, \cdot)](x), \quad \forall x \in \mathcal{H}, t \ge 0.
$$

This completes the proof.  $\Box$ 

Remark 4.1. In applications of Theorem 4.1, to verify assumption (A6) for the map  $P_1$  it suffices to show that for any intermediate periodic steady state  $\alpha$ , the sum of the rightward spreading speed of  $Q_1 : [0, \alpha]_C \to [0, \alpha]_C$  and the leftward spreading speed of  $Q_1 : [\alpha, \beta]_C \to [\alpha, \beta]_C$  is positive since  $P_1$  and  $Q_1$  are topologically conjugate.

To finish this section, we note that the bistability structure can be obtained for equation  $(4.1)$  under appropriate conditions so that the existence result in [\[43,](#page-45-7) [42\]](#page-45-9) is improved (see the details in Subsection 6.3). Further, Theorem [4.1](#page-25-2) with  $\mathcal{H} = \mathbb{Z}$  and  $\mathcal{X} = \mathbb{R}$  can be used to rediscover the existence result of [\[14\]](#page-44-10) for a one-dimensional lattice equation under the bistability assumption.

#### 5. Semiflows with weak compactness

In assumption (A4) of Section 2, we assume that  $Q : C_{\beta} \to C_{\beta}$  is compact with respect to the compact-open topology. In this section, we establish the existence of bistable waves under some weaker compactness assumptions.

Fix  $\tau > 0$ . It is well known that the time-t solution map of time-delayed reactiondiffusion equations such as

<span id="page-26-0"></span>
$$
\frac{\partial u}{\partial t} = \frac{\partial^2 u}{\partial x^2} + f(u(t, x), u(t - \tau, x))
$$
\n(5.1)

is compact in the compact-open topology if and only if  $t > \tau$ , where the phase space C is chosen as  $C(\mathbb{R}, C([- \tau, 0], \mathbb{R}))$ . The first purpose of this section is to show that our results are still valid for this kind of evolution equations by introducing an alternative assumption (A4').

To state it, we need some notation for time-delayed evolution systems. Let  $\tau \in \mathcal{T}$ , let  $\mathcal F$  be a Banach lattice with the positive cone  $\mathcal F^+$  having nonempty interior,  $\beta \in \text{Int}(\mathcal F^+)$ , and  $\mathcal{X}_{\beta} = C([- \tau, 0], \mathcal{F}_{\beta})$ . Any  $\phi \in \mathcal{C}_{\beta}$  can be regarded as an element in  $C([- \tau, 0] \times$  $H, \mathcal{F}^+$ ). For any subset B of  $[-\tau, 0] \times H$ , we define  $\phi|_B$  as the restriction of  $\phi$  to B.

(A4<sup>'</sup>) (*Compactness*) There exists  $s \in (0, \tau]$  such that:

- (i)  $Q[\phi](\theta, x) = \phi(\theta + s, x)$  whenever  $\theta + s \leq 0$ .
- (ii) For any  $\epsilon \in (0, s)$ , the set  $Q[\mathcal{C}_{\beta}]|_{[-s+\epsilon, 0] \times \mathcal{H}}$  is precompact.
- (iii) For any subset  $\mathcal{J} \subset \mathcal{C}_{\beta}$  with  $\mathcal{J}(0, \cdot) \subset C(\mathcal{H}, \mathcal{Y}_{\beta})$  being precompact, the set  $Q[\mathcal{J}]|_{[-s,0]\times H}$  is precompact.

This assumption was motivated by  $[25,$  Assumption  $(AG')$ ]. Let us use equation  $(5.1)$  to explain (A4'). For any  $t > \tau$ , one can directly verify that the solution map  $Q_t$  satisfies (A4) by rewriting [\(5.1\)](#page-26-0) in integral form (see, e.g., [\[41\]](#page-45-16)); and for any  $t \in (0, \tau]$ , one can show that  $Q_t$  satisfies  $(A4')(i) \& (ii)$  by the same arguments. For  $(A4')(iii)$ , we provide a proof below.

Let  $T(0) = I$ , and for any  $t \in (0, \tau]$ , let  $T(t)$  be the time-t map of the heat equation  $u_t = \Delta u$ . Then [\(5.1\)](#page-26-0) can be written as

$$
u(t, x; \phi) = T(t)\phi(x) + \int_0^t T(t - s) f(u(s, u(s - \tau)))(x) ds,
$$

and hence  $Q_t[\phi](\theta, x) = u(t + \theta, x)$ . Note that for any  $\phi \in C_\beta$ ,  $T(t)\phi \rightarrow \phi$  in the compact-open topology as  $t \rightarrow 0$ . It then follows from the triangular inequality and the absolute continuity of integrals that for any compact subset  $\mathcal{H}_1 \subset \mathbb{R}$ , the set  $Q_t[\mathcal{J}]|_{[-t,0]\times \mathcal{H}_1}$  is equicontinuous, and hence  $Q[\mathcal{J}]|_{[-t,0]\times \mathbb{R}}$  is precompact in  $\mathcal{C}_{\beta}$ .

<span id="page-27-0"></span>**Lemma 5.1.** *Let*  $A_{\xi}$ ,  $\xi \geq 1$ , *be defined as in Section* 3 *and let*  $\beta \in Int(\mathcal{F}^+)$ *.* Assume that the map  $Q : C_{\beta} \rightarrow C_{\beta}$  satisfies (A4'). Then there exists an inte $g$ er  $m_0$  *such that*  $\bigcup_{\xi \in [1,1+\delta]} (Q \circ A_{\xi})^{m_0} [\mathcal{C}_{\beta}] \subset \mathcal{C}_{\beta}$  *is precompact when*  $\mathcal{H} = \mathbb{R}$ *, and*  $\bigcup_{\xi \in [1,2]} (\tilde{Q} \circ A_{\xi})^{m_0} [\mathcal{B}_{\beta}](x) \subset \mathcal{X}_{\beta}$  is precompact for any  $x \in \mathbb{R}$  when  $\mathcal{H} = \mathbb{Z}$ .

*Proof.* We only handle the case  $\mathcal{H} = \mathbb{R}$  since the proof for  $\mathcal{H} = \mathbb{Z}$  is essentially similar. Let s and  $\tau$  be as in (A4'). There exists  $m_0 \in \mathbb{N}$  such that  $s \in \left(\frac{1}{m_0+1}\tau, \frac{1}{m_0}\tau\right]$ . By assumption (A4')(i), for any  $\xi \ge 1$  and  $\phi_0 \in C_\beta$ ,

$$
\phi_1^{\xi}(\theta, x) := Q \circ A_{\xi}[\phi_0](\theta, x) = \begin{cases} \phi_0(\theta + s, \xi x), & \theta + s \le 0, \\ Q[\phi_0(\xi \cdot)](\theta, x), & \theta + s > 0 \end{cases}
$$

This implies that for any  $\xi \ge 1$  and  $\epsilon < s - \frac{1}{m_0 + 1} \tau$ ,

$$
\bigcup_{\xi\in[1,2]}Q\circ A_{\xi}[\mathcal{C}_{\beta}]\big|_{[-s+\epsilon,0]\times\mathbb{R}}\subset Q[\mathcal{C}_{\beta}]\big|_{[-s+\epsilon,0]\times\mathbb{R}}.
$$

Since  $Q[C_\beta]_{[-s+\epsilon,0]\times\mathbb{R}}$  is precompact, as assumed in (A4')(ii), it then follows that  $\bigcup_{\xi \in [1,2]} Q \circ A_{\xi}[\mathcal{C}_{\beta}](0,\cdot) \subset C(\mathbb{R},\mathcal{Y}_{\beta})$  is precompact. By  $(A4')$ (iii) and similar arguments, we have

$$
\phi_2^{\xi}(\theta, x) := Q \circ A_{\xi}[\phi_1^{\xi}](\theta, x) = \begin{cases} \phi_1^{\xi}(\theta + s, \xi x), & \theta + s \le 0, \\ Q[\phi_1^{\xi}(\xi \cdot)][(\theta, x), & \theta + s > 0, \end{cases}
$$

$$
= \begin{cases} \phi_0(\theta + 2s, \xi^2 x), & \theta + 2s \le 0, \\ Q[\phi_0(\xi \cdot)][(\theta + s, \xi x), & 0 < \theta + 2s \le s, \\ Q[\phi_1^{\xi}(\xi \cdot)][(\theta, x), & \theta + s > 0, \end{cases}
$$

This implies that  $\bigcup_{\xi \in [1,2]} (Q \circ A_{\xi})^2 [\mathcal{C}_{\beta}]_{[-2s+\epsilon,0]\times \mathbb{R}}$  is precompact. Consequently, the set  $\bigcup_{\xi \in [1,2]} (Q \circ A_{\xi})^2 [\mathcal{C}_{\beta}](0,\cdot) \subset C(\mathbb{R}, \mathcal{Y}_{\beta})$  is compact. By induction, we have

$$
\phi_{m_0+1}^{\xi}(\theta, x) := Q \circ A[\phi_{m_0}^{\xi}](\theta, x) = \begin{cases} \phi_{m_0}^{\xi}(\theta + s, \xi x), & \theta + s \le 0, \\ Q[\phi_{m_0}^{\xi}(\xi \cdot)](\theta, x), & \theta + s > 0, \end{cases}
$$
  
= ...

$$
= \begin{cases} Q[\phi_0(\xi \cdot)](\theta + (m_0 + 1)s, \xi^{m_0}x), & 0 < \theta + (m_0 + 1)s \le s, \\ Q[\phi_1^{\xi}(\xi \cdot)](\theta + m_0s, \xi^{m_0 - 1}x), & 0 < \theta + m_0s \le s, \\ \cdots \\ Q[\phi_{m_0 - 1}^{\xi}(\xi \cdot)](\theta + s, \xi x), & 0 < \theta + s \le s, \\ Q[\phi_{m_0}^{\xi}](\theta, x), & \theta + s > 0. \end{cases}
$$

This implies that  $\bigcup_{\xi \in [1,2]} (Q \circ A_{\xi})^{m_0} [\mathcal{C}_{\beta}]$  is precompact in  $\mathcal{C}_{\beta}$ .

<span id="page-28-1"></span>Theorem 5.1. *All results in Theorems* [3.1-](#page-11-0)[3.5](#page-21-0) *and* [4.1](#page-25-2) *are valid if we replace* (A4)  $with (A4').$ 

*Proof.* Following the proof of these theorems, we only need to modify the parts where we use the compactness assumption (A4). At these parts, we apply Lemma [5.1.](#page-27-0)  $\Box$ 

Note that the solution maps of the nonlocal dispersal equation

$$
u_t = J * u - u + f(u)
$$

satisfy neither (A4) nor (A4'). The second purpose of this section is to modify our theory so that it applies to such integro-differential systems in a homogeneous habitat.

Let M denote the set of all nondecreasing functions from  $\mathbb R$  to X, and  $\beta \in \mathcal{X}^+$ . We equip M with the compact-open topology. Assume that Q maps  $M_\beta$  to  $M_\beta$ . Let E denote the set of fixed points of Q restricted to  $\mathcal{X}_{\beta}$ . Suppose that 0 and  $\beta$  are in E. We impose the following assumptions on Q:

- (B1) (*Translation invariance*)  $T_y \circ Q[\phi] = Q \circ T_y[\phi]$  for all  $\phi \in \mathcal{M}_\beta$  and  $y \in \mathbb{R}$ .
- (B2) (*Continuity*)  $Q : \mathcal{M}_\beta \to \mathcal{M}_\beta$  is continuous in the sense that if  $\phi_n \to \phi$  in  $\mathcal{M}_\beta$ , then  $Q[\phi_n](x) \to Q[\phi](x)$  in  $\mathcal{X}_\beta$  for almost all  $x \in \mathbb{R}$ .
- (B3) (*Monotonicity*) Q is order preserving in the sense that  $Q[\phi] \geq Q[\psi]$  whenever  $\phi \geq \psi$  in  $\mathcal{M}_{\beta}$ .
- (B4) (*Weak compactness*) For any fixed  $x \in \mathbb{R}$ , the set  $Q[\mathcal{M}_{\beta}](x)$  is precompact in  $\mathcal{X}_{\beta}$ .
- (B5) (*Bistability*) The fixed points 0 and  $\beta$  are strongly stable from above and below, respectively, for the map  $Q : \mathcal{X}_{\beta} \to \mathcal{X}_{\beta}$ , and the set  $E \setminus \{0, \beta\} \subset \mathcal{X}_{\beta}$  is totally unordered.
- (B6) (*Counter-propagation*) For each  $\alpha \in E \setminus \{0, \beta\}$ ,  $c^*_{-}(\alpha, \beta) + c^*_{+}(0, \alpha) > 0$ .

Comparing assumptions  $(A1)$ – $(A6)$  and  $(B1)$ – $(B6)$ , one can find that the assumptions of translation invariance, monotonicity, bistability and counter-propagation are the same. The difference lies in the assumptions of continuity and compactness. Clearly, the compactness assumption (B4) is much weaker than (A4).

<span id="page-28-0"></span>**Theorem 5.2.** Let  $X = C(M, \mathbb{R}^d)$  and assume  $Q : \mathcal{M}_\beta \to \mathcal{M}_\beta$  satisfies (B1)–(B6). *Then there exist*  $c \in \mathbb{R}$  *and*  $\psi \in M_\beta$  *connecting* 0 *to*  $\beta$  *such that*  $Q[\psi](x) = \psi(x + c)$ *for all*  $x \in \mathbb{R}$ *.* 

*Proof.* We combine the proofs of Theorems [3.1](#page-11-0) and [3.2.](#page-14-1) More precisely, one can repeat the proof of Theorem [3.1](#page-11-0) except for the parts where the compactness assumption (A4) is used. For these parts, one needs to use the idea in Theorem  $3.2$ , where  $\ddot{Q}$  has the same compactness property as  $Q$ . In the rest of this section, we say  $\{Q_t\}_{t\geq 0}$  is a semiflow on  $\mathcal{M}_\beta$  provided that  $Q_0 = I$ ;  $Q_t \circ Q_s = Q_{t+s}$  for all  $t, s > 0$ ; and  $Q_{t_n}[\phi_n](x) \to Q_t[\phi](x)$  in  $\mathcal{X}_\beta$  for almost all  $x \in \mathbb{R}$ whenever  $t_n \to t$  and  $\phi_n \to \phi$  in  $\mathcal{M}_\beta$ .

**Theorem 5.3.** Let  $X = C(M, \mathbb{R}^d)$ . Assume that  $\{Q_t\}_{t \geq 0}$  is a semiflow on  $\mathcal{M}_{\beta}$ , and for *any*  $t > 0$ , the map  $Q_t$  satisfies (B1) and (B3)–(B6). Then there exist  $c \in \mathbb{R}$  and  $\psi \in \mathcal{M}_\beta$ *connecting* 0 *to β such that*  $Q_t[\psi](x) = \psi(x + ct)$  *for all*  $x \in \mathbb{R}$ *.* 

*Proof.* As in the proof of Theorem [5.2,](#page-28-0) we need to combine the proofs of Theorems [3.4](#page-18-0) and [3.5.](#page-21-0)  $\Box$ 

Similarly, we can define  $\omega$ -time periodic semiflows on  $\mathcal{M}_{\beta}$  and obtain the following result.

**Theorem 5.4.** Let  $X = C(M, \mathbb{R}^d)$ . Assume that  $\{Q_t\}_{t\geq 0}$  is an  $\omega$ -time periodic semiflow *on*  $M_\beta$ *. Let*  $\beta(t)$  *be a strongly positive periodic solution of*  $\{Q_t\}_{t>0}$  *restricted to*  $X_\beta$ *. Further, assume that the Poincaré map*  $Q_{\omega}$  *satisfies* (B1) *and* (B3)–(B6) *with*  $\beta = \beta(0)$ *. Then there exist*  $c \in \mathbb{R}$  *and*  $\phi(t, x)$  *with*  $\phi(t, -\infty) = 0$  *and*  $\phi(t, \infty) = \beta(t)$  *such that*  $Q_t[\psi](x) = \psi(t, x + ct)$  *for all*  $x \in \mathbb{R}$ *. Moreover,*  $\phi(t, \cdot) \in \mathcal{M}_\beta$  *and*  $\phi(t, \cdot)$  *is ω*-*periodic in*  $t \geq 0$ *.* 

## 6. Applications

In this section, we apply the abstract results to four kinds of monotone evolution systems: a time-periodic reaction-diffusion system, a parabolic system in a cylinder, a parabolic equation with variable diffusion, and a nonlocal and time-delayed reaction-diffusion equation.

#### *6.1. A time-periodic reaction-diffusion system*

Consider the time-periodic reaction-diffusion system

<span id="page-29-0"></span>
$$
\frac{\partial u}{\partial t} = A \Delta u + f(t, u), \quad x \in \mathbb{R}, \tag{6.1}
$$

where  $u = (u_1, ..., u_n)^T$ ,  $A = \text{diag}\{d_1, ..., d_n\}$  with each  $d_i > 0$  and  $f = (f_1, ..., f_n)^T$ is  $\omega$ -periodic in  $t \ge 0$  (i.e.,  $f(t, \cdot) = f(t + \omega, \cdot)$ ). The existence of periodic bistable traveling waves of  $(6.1)$  with  $n = 1$  was proved in [\[1\]](#page-43-2). Here we generalize this result to  $n > 1$ .

Let  $f \in C^1(\mathbb{R}_+ \times \mathbb{R}^n, \mathbb{R}^n)$ . In order to apply Theorem [3.3](#page-16-0) to system [\(6.1\)](#page-29-0), we choose  $C := C(\mathbb{R}, \mathbb{R}^n)$ ,  $\mathcal{X} := \mathbb{R}^n$ , and  $\mathcal{E} \subset C$  to be the set of all bounded functions from  $\mathbb{R}$ to  $\mathbb{R}^n$ . Using the solution maps  $\{T(t)\}_{t\geq 0}$  of the heat equation  $\frac{\partial u}{\partial t} = A \Delta u$ , we write [\(6.1\)](#page-29-0) in integral form:

$$
u(t; \phi) = T(t)\phi + \int_0^t T(t - s) f(s, u(s; \phi)) ds.
$$
 (6.2)

Define  $Q_t[\phi] := u(t; \phi)$  for all  $\phi \in \mathcal{E}$ . Let 0 and  $\beta \gg 0$  be two fixed points of the Poincaré map  $Q_{\omega}$  in X, and let E be the set of all spatially homogeneous fixed points of  $Q_{\omega}$  in  $\mathcal{X}_{\beta}$ . We impose the following assumptions:

- (C1) The Jacobian matrix  $Df(t, u)$  ( $D = D_u$ ) is cooperative and irreducible for all  $t \ge 0$ and  $u \ge 0$ . For any  $\omega$ -periodic solution  $p(t)$  of  $u' = f(t, u)$  with  $0 < p(0) < \beta$ , every off-diagonal entry of  $Df(t, p(t))$  is either strictly positive for all t or identically zero.
- (C2) The spatially homogeneous system  $u' = f(t, u)$  is of bistable type, that is, 0 and  $\beta$ are stable fixed points of  $Q_{\omega}$  in the sense that  $s\left(\frac{d}{du}Q_{\omega}[0]\right) < 0$  and  $s\left(\frac{d}{du}Q_{\omega}[\beta]\right) < 0$ , and any  $\alpha \in E \setminus \{0, \beta\}$  is unstable in the sense that  $s\left(\frac{d}{du}Q_{\omega}[\alpha]\right) > 0$ , where  $s(M)$  is the stability modulus of the matrix M defined by  $s(M) = \max\{Re \lambda :$  $\lambda$  is an eigenvalue}.

**Theorem 6.1.** Assume that (C1)–(C2) hold, and let  $\beta(t)$  be the periodic solution of  $u' =$  $f(t, u)$  *with*  $\beta(0) = \beta$ . Then there exists  $c \in \mathbb{R}$  *such that* [\(6.1\)](#page-29-0) *admits a time-periodic traveling wave*  $U(t, x + ct)$  *connecting* 0 *to*  $\beta(t)$ *.* 

*Proof.* It is easy to see that the discrete semiflow  $\{Q_{\omega}^n\}_{n\geq1}$  on  $\mathcal{C}_{\beta}$  satisfies (A1)–(A5) with  $Q = Q_{\omega}$ . Next we show that (A6) holds with  $Q = Q_{\omega}$ .

Note that for any  $\alpha \in E \setminus \{0, \beta\}, \{Q_{\omega}^n\}_{n \geq 1} : [\alpha, \beta]_{\mathcal{C}} \to [\alpha, \beta]_{\mathcal{C}}$  has monostable dynamics, where  $\alpha$  is unstable and  $\beta$  is stable. By the theory developed in [\[26\]](#page-44-13),  $Q_{\omega}$  admits leftward and rightward spreading speeds  $c^*(\alpha, \beta)$  and  $c^*_+(\alpha, \beta)$ . Since  $Q_\omega$  is reflectively invariant, we further have  $c^*_{-}(\alpha, \beta) = c^*_{+}(\alpha, \beta) =: c^*(\alpha, \beta)$ , which is called the spreading speed of this monostable subsystem. Note that  ${Q_{\omega}^n}_{n \ge 1} : [0, \alpha]_{\mathcal{C}} \to [0, \alpha]_{\mathcal{C}}$  also has monostable dynamics, where 0 is stable and  $\alpha$  is unstable. Similarly, this monostable subsystem also admits a spreading speed  $c^*(0, \alpha)$ .

Denote the periodic solution  $Q_t[\alpha]$  by  $\alpha(t)$ , and define

$$
\bar{f}(t, u) = f(t, \alpha(t) + u) - f(t, \alpha(t)).
$$
\n(6.3)

Clearly,  $\alpha(0) = \alpha$  and  $\bar{f}(t, 0) \equiv 0$ . Then the semiflow  $Q_t$  has the same dynamics on [ $\alpha$ ,  $\beta$ ]<sub>C</sub> as the semiflow  $\overline{Q}_t$  on  $C_{\beta-\alpha}$  associated with

$$
u_t = A \Delta u + \bar{f}(t, u). \tag{6.4}
$$

In particular, these two semiflows share the spreading speed  $c^*(\alpha, \beta)$ .

To give a lower bound on  $c^*(\alpha, \beta)$ , we employ a similar idea to that in the proof of [\[40,](#page-45-17) Lemma 4.1]. Choose  $\rho > 0$  such that the diagonal entries of  $D\bar{f}(t, 0) + \rho I$  are strictly positive. We define for each  $i$  the projection

$$
\{\Pi_i[u]\}_j = \begin{cases} u_j & \text{if } \{D\bar{f}(t,0) + \rho I\}_{ij} > 0, \\ 0 & \text{if } \{D\bar{f}(t,0) + \rho I\}_{ij} = 0, \end{cases}
$$
(6.5)

and

<span id="page-30-0"></span>
$$
\sigma(t) = \min \left\{ \frac{\partial \bar{f}_i}{\partial u_j}(t,0) + \rho I_{ij} : \frac{\partial \bar{f}_i}{\partial u_j}(t,0) + \rho I_{ij} > 0 \right\}.
$$

It is easy to see from assumption (C1) that  $\Pi_i$  is independent of  $t \in [0, \omega]$ , and  $\sigma :=$  $\inf{\lbrace \sigma(t) : t \in [0, \omega] \rbrace}$  is positive. Further,

$$
u \ge \Pi_i[u] \quad \text{and} \quad \bar{f}_i(t, u) \ge \bar{f}_i(t, \Pi_i[u]), \quad \forall t \ge 0, \ u \ge 0. \tag{6.6}
$$

By the differentiability of  $\bar{f} = (\bar{f}_1, \dots, \bar{f}_n)$ , for any  $\epsilon \in (0, 1)$ , there exists  $\eta = \eta(\epsilon) \gg 0$ in  $\mathbb{R}^n$  such that

$$
\bar{f}_i(t, \Pi_i[u]) \ge \nabla \bar{f}_i(t, 0) \cdot \Pi_i[u] - \sigma \epsilon \|\Pi_i[u]\|, \quad \forall t \in [0, \omega], \ u \in [0, \eta]. \tag{6.7}
$$

Note that

$$
\nabla \bar{f}_i(t,0) \cdot \Pi_i[u] = (D\bar{f}(t,0)u)_i, \quad \forall u \in \mathbb{R}^n,
$$
\n(6.8)

and

$$
\|\Pi_i[u]\| \le \sum_j (\Pi_i[u])_j \le \sum_j \frac{(D\bar{f}(t,0) + \rho I)_{ij}}{\sigma} (\Pi_i[u])_j
$$
  
= 
$$
\sum_j \frac{\nabla \bar{f}_i(t,0) \cdot \Pi_i[u] + \rho (\Pi_i[u])_i}{\sigma}
$$
  
= 
$$
\sigma^{-1}(D\bar{f}(t,0)u + \rho u)_i, \quad \forall u \ge 0.
$$
 (6.9)

Combining  $(6.6)$ – $(6.9)$ , we obtain

$$
\bar{f}(t, u) \ge D\bar{f}(t, 0)u - \epsilon [D\bar{f}(t, 0)u + \rho u], \quad \forall t \in [0, \omega], u \in [0, \eta]. \tag{6.10}
$$

For the above  $\eta = \eta(\epsilon)$ , there exists  $\gamma = \gamma(\epsilon) \gg 0$  in  $\mathbb{R}^n$  such that

<span id="page-31-0"></span>
$$
\bar{Q}_t[\phi] \le \bar{Q}_t[\gamma] \le \eta, \quad \forall \phi \in \mathcal{C}_{\gamma}, \ t \in [0, \omega]. \tag{6.11}
$$

Thus, the comparison theorem implies that

<span id="page-31-3"></span><span id="page-31-2"></span>
$$
\bar{Q}_t[\phi] \ge M_t^{\epsilon}[\phi], \quad \forall \phi \in \mathcal{C}_{\gamma}, \ t \in [0, \omega], \tag{6.12}
$$

where  $M_t^{\epsilon}$  is the solution map of

$$
u_t = A\Delta u + D\bar{f}(t,0)u - \epsilon(D\bar{f}(t,0) + \rho I)u.
$$
\n(6.13)

Let  $\rho_{\epsilon}(\mu)$  be the principal Floquet multiplier of the linear periodic cooperative and irreducible system

<span id="page-31-1"></span>
$$
\frac{dv}{dt} = [\mu^2 A + D\bar{f}(t,0)]v - \epsilon (D\bar{f}(t,0) + \rho I)v.
$$
 (6.14)

Since  $\lim_{\epsilon \to 0^+} \rho_{\epsilon}(0) = \rho_0(0) > 1$ , we can fix an  $\epsilon \in (0, 1)$  such that  $\rho_{\epsilon}(0) > 1$ . Let  $v(t, w)$  be the solution of [\(6.14\)](#page-31-1) satisfying  $v(0, w) = w \in \mathbb{R}^n$ . It is easy to see that  $u(t, x) = e^{-\mu x} v(t, w)$  is a solution of the linear periodic system [\(6.13\)](#page-31-2). Define  $\Phi_{\epsilon}(\mu) := \ln \rho_{\epsilon}(\mu)/\mu$ . From [\[25,](#page-44-14) Theorem 3.10] and [\(6.12\)](#page-31-3), we have

$$
c^*(\alpha, \beta) \ge \inf_{\mu > 0} \Phi_{\epsilon}(\mu). \tag{6.15}
$$

Now we prove that  $\Phi_{\epsilon}(\infty) = \infty$ . Let  $\lambda_{\epsilon}(\mu) = (1/\omega) \ln \rho_{\epsilon}(\mu)$ . By Floquet theory, there exists a positive  $\omega$ -periodic function  $\xi(t) := (\xi_1(t), \dots, \xi_n(t))^T$  such that  $v(t) :=$  $e^{\lambda_{\epsilon}(\mu)t} \xi(t)$  is a solution of [\(6.14\)](#page-31-1). In particular,

$$
\xi_1'(t) = (d_1\mu^2 - \lambda_{\epsilon}(\mu) - \epsilon \rho)\xi_1(t) + (1 - \epsilon) \sum_{i=1}^n \frac{\partial}{\partial u_i} \bar{f}_1(t, 0)\xi_i(t).
$$

Dividing by  $\xi_1(t)$  and integrating from 0 to  $\omega$  gives

$$
0 = (d_1\mu^2 - \lambda_{\epsilon}(\mu) - \epsilon \rho)\omega + (1 - \epsilon) \int_0^{\omega} \sum_{i=1}^n \frac{\partial}{\partial u_i} \bar{f}_1(t, 0) \xi_i(t) / \xi_1(t) dt, \quad \forall \mu > 0.
$$

Since the matrix  $D\bar{f}(t, 0)$  is cooperative and  $\xi(t)$  is positive, we obtain

$$
0 \ge (d_1\mu^2 - \lambda_{\epsilon}(\mu) - \epsilon \rho)\omega + (1 - \epsilon) \int_0^{\omega} \frac{\partial}{\partial u_1} \bar{f}_1(t, 0) dt.
$$

This implies that

$$
\Phi_{\epsilon}(\mu) = \frac{\omega \lambda_{\epsilon}(\mu)}{\mu} \ge d_1 \mu \omega + \frac{1}{\mu} \bigg[ -\epsilon \rho \omega + (1 - \epsilon) \int_0^{\omega} \frac{\partial}{\partial u_1} \bar{f}_1(t, 0) dt \bigg],
$$

and hence  $\Phi_{\epsilon}(\infty) = \infty$ . By [\[25,](#page-44-14) Lemma 3.8], we have  $\inf_{\mu>0} \Phi_{\epsilon}(\mu) > 0$ . Therefore,  $c^*(\alpha, \beta) > 0$ . Similarly, we can prove that  $c_*(0, \alpha) > 0$ . Thus, assumption (A6) with  $Q = Q_{\omega}$  holds. Consequently, Theorem [3.3](#page-16-0) completes the proof.

#### *6.2. A reaction-diffusion-advection system in a cylinder*

In this subsection, we consider the system

<span id="page-32-0"></span>
$$
\begin{cases} \frac{\partial u}{\partial t} = A \frac{\partial^2 u}{\partial x^2} + B \Delta_y u + E(y) \frac{\partial u}{\partial x} + f(u), & x \in \mathbb{R}, y \in \Omega \subset \mathbb{R}^{m-1}, t > 0, \\ \frac{\partial u}{\partial v} = 0 & \text{on } (0, \infty) \times \mathbb{R} \times \partial \Omega, \end{cases}
$$
(6.16)

where A, B are positive definite diagonal  $n \times n$  matrices, E is a diagonal matrix of smooth functions of y,  $\Omega$  is a bounded and convex open subset in  $\mathbb{R}^{m-1}$  with smooth boundary  $\partial \Omega$ ,  $\Delta_y = \sum_{i=1}^{m-1} \partial^2/\partial y_i^2$ , and v is the outer unit normal vector to  $\partial \Omega \times \mathbb{R}$ .

The existence of bistable traveling waves for  $(6.16)$  with  $n = 1$  was obtained in [\[11\]](#page-43-4). Here we extend this result to the case  $n \geq 2$ . Assume that  $f \in C^1(\mathbb{R}^n, \mathbb{R}^n)$  satisfies the following conditions:

- (D1) The Jacobian matrix  $Df(u)$  is cooperative and irreducible for all  $u \ge 0$ .
- (D2) f is of bistable type in the sense that it has exactly three ordered zeros:  $0 < a < \beta$ and  $s(Df(0)) < 0$ ,  $s(Df(a)) > 0$ ,  $s(Df(\beta)) < 0$ .

**Theorem 6.2.** Assume that (D1)–(D2) *hold. Then there exists*  $c \in \mathbb{R}$  *such that system* [\(6.16\)](#page-32-0) *admits a traveling wave connecting* 0 *to* β *with speed* c*.*

*Proof.* In order to employ Theorem [3.4,](#page-18-0) we choose  $\mathcal{X} := C(\bar{\Omega}, \mathbb{R}^n)$  and  $\mathcal{C} := C(\mathbb{R}, \mathcal{X})$ with the standard cones  $\chi^+$  and  $\mathcal{C}^+$ , respectively. Let  $G(t, x, y, w)$  be the Green function of the linear equation

<span id="page-33-0"></span>
$$
\begin{cases} \frac{\partial u}{\partial t} = A \frac{\partial^2 u}{\partial x^2} + B \Delta_y u + E(y) \frac{\partial u}{\partial x}, & x \in \mathbb{R}, y \in \Omega, t > 0, \\ \frac{\partial u}{\partial v} = 0 & \text{on } (0, \infty) \times \mathbb{R} \times \partial \Omega. \end{cases}
$$
(6.17)

Then the solution of [\(6.17\)](#page-33-0) with initial value  $u(0, \cdot) = \phi(\cdot) \in C$  can be expressed as

$$
u(t, x, y; \phi) = \int_{\mathbb{R}} \int_{\Omega} G(t, x - z, y, w) \phi(z, w) dw dz.
$$

Define  $T(t)\phi = u(t, \cdot; \phi)$  for  $\phi \in C_\beta$ . Using the variation of constants formula, we write [\(6.16\)](#page-32-0) subject to  $u(0, \cdot) = \phi(\cdot) \in C_\beta$  as an integral equation

$$
u(t, x, y; \phi) = T(t)[\phi](x, y) + \int_0^t T(t - s) f(u(s, x, y)) ds.
$$
 (6.18)

By linear operator theory, for any  $\phi \in C_\beta$ , system [\(6.16\)](#page-32-0) has a unique solution  $u(t; \phi)$  with  $u(0; \phi) = \phi$ , which exists globally on  $[0, \infty)$ . Define  $Q_t[\phi] := u(t, \phi)$ . Then  $\{Q_t\}_{t \geq 0}$  is a subhomogeneous semiflow on  $C_\beta$  (see [\[25,](#page-44-14) Section 5.3]). Also, assumption (D1) ensures that the semiflow  $\{Q_t\}_{t\geq 0}$  restricted to  $\mathcal{X}_\beta$  is strongly monotone (see [\[34\]](#page-45-14)). Further, it is easy to see that  $Q_t$ ,  $t \geq 0$ , satisfies assumptions (A1)–(A4). Since the domain  $\Omega$  is convex, it follows from the result in  $[24]$  that any nonconstant steady state of the x-independent system

<span id="page-33-1"></span>
$$
\begin{cases} \frac{\partial u}{\partial t} = B\Delta_y u + f(u), & y \in \Omega, \ t > 0, \\ \frac{\partial u}{\partial v} = 0 & \text{on } (0, \infty) \times \partial\Omega, \end{cases}
$$

is linearly unstable. This then implies that  $Q_t$  satisfies (A5'). Now it remains to show that  $(A6)$  holds for  $Q_1$ .

For each x-independent steady state  $\alpha = \alpha(y)$  in [[0,  $\beta$ ]] $\chi$ , system [\(6.16\)](#page-32-0) has monostable dynamics on  $[\alpha, \beta]_C$ . To better understand the dynamics of this subsystem, we make the transform  $u = w + \alpha$  and define  $g(w, y) := f(w + \alpha(y)) - f(\alpha(y))$ . Then the dynamics is equivalent to that of the following system on  $[0, \beta - \alpha]_C$ :

$$
\begin{cases} \frac{\partial w}{\partial t} = A \frac{\partial^2 w}{\partial x^2} + B \Delta_y w + E(y) \frac{\partial w}{\partial x} + g(w, y), & x \in \mathbb{R}, y \in \Omega, t > 0, \\ \frac{\partial w}{\partial y} = 0 & \text{on } (0, \infty) \times \mathbb{R} \times \partial \Omega. \end{cases}
$$
(6.19)

System [\(6.19\)](#page-33-1) has exactly two x-independent steady states  $S_1 := 0$  and  $S_2 := \beta - \alpha \gg 0$ . By the theory developed in [\[26\]](#page-44-13), [\(6.19\)](#page-33-1) has a leftward spreading speed  $\hat{c}^*$  in a strong sense. Let  $c^*_{-}(\alpha, \beta)$  be defined as in [\(2.6\)](#page-7-1) with  $Q = Q_1$ . Then  $c^*_{-}(\alpha, \beta) \ge \hat{c}^*_{-}$ .

To verify (A6) for  $Q_1$ , we first estimate the speed  $\hat{c}^*$ . Consider the linearized system of  $(6.19)$  at the equilibrium  $S_1$ :

<span id="page-33-2"></span>
$$
\begin{cases} \frac{\partial w}{\partial t} = A \frac{\partial^2 w}{\partial x^2} + B \Delta_y w + E(y) \frac{\partial w}{\partial x} + \frac{\partial g(0, y)}{\partial w} w, & x \in \mathbb{R}, y \in \Omega, t > 0, \\ \frac{\partial w}{\partial v} = 0 & \text{on } (0, \infty) \times \mathbb{R} \times \partial \Omega. \end{cases}
$$
(6.20)

Suppose  $w(t, x, y) := e^{\mu x} v(t, y)$  is a solution of [\(6.20\)](#page-33-2). Then  $v(t, y)$  satisfies the  $\mu$ -parameterized linear parabolic equation

<span id="page-34-0"></span>
$$
\begin{cases} \frac{\partial v}{\partial t} = B\Delta_y u + \left[\mu^2 A + \mu E(y) + \frac{\partial g(0, y)}{\partial w}\right] v, & y \in \Omega, \ t > 0, \\ \frac{\partial v}{\partial v} = 0 & \text{on } (0, \infty) \times \partial \Omega. \end{cases}
$$
(6.21)

Let  $\lambda^-(\mu)$  be the principal eigenvalue of [\(6.21\)](#page-34-0). Then by the results in [\[25,](#page-44-14) Section 3], it follows that  $\hat{c}^* \ge \inf_{\mu > 0} \lambda^+(\mu)/\mu$ , and  $\lambda^-(\mu)$  is convex in  $\mu$ . Moreover, it is easy to see that  $\lim_{\mu\to\infty} \frac{\lambda^-(\mu)}{\mu} = \infty$  and  $\lim_{\mu\to 0^+} \frac{\lambda^-(\mu)}{\mu} = \infty$ , and hence  $\lambda^-(\mu)/\mu$  attains its infimum at some  $\mu_1 \in (0, \infty)$ .

Similarly, system [\(6.16\)](#page-32-0) has monostable dynamics on [0,  $\alpha$ ]<sub>C</sub>. To better understand this dynamics of this subsystem, we make the transform  $w = \alpha - u$  and define  $h(w, y) :=$  $f(\alpha(y)) - f(\alpha(y) - w)$ . Then the dynamics is equivalent to that of the following system on  $[0, \alpha]_C$ :

$$
\begin{cases} \frac{\partial w}{\partial t} = A \frac{\partial^2 w}{\partial x^2} + B \Delta_y w + E(y) \frac{\partial w}{\partial x} + h(w, y), & x \in \mathbb{R}, y \in \Omega, t > 0, \\ \frac{\partial w}{\partial y} = 0 & \text{on } (0, \infty) \times \mathbb{R} \times \partial \Omega. \end{cases}
$$

By the same arguments, this system has a rightward spreading speed  $\hat{c}_+^*$  and  $c^*_{+}(0, \alpha) \geq \hat{c}^*_{+}$ . Also, by a similar procedure, we define  $\lambda^+(\mu)$  as the principal eigenvalue of the problem

<span id="page-34-1"></span>
$$
\begin{cases} \frac{\partial v}{\partial t} = B\Delta_y u + \left[\mu^2 A - \mu E(y) + \frac{\partial g(0, y)}{\partial w}\right] v, & y \in \Omega, \ t > 0, \\ \frac{\partial v}{\partial v} = 0 & \text{on } (0, \infty) \times \partial \Omega. \end{cases}
$$
(6.22)

Then  $\hat{c}_+^* \ge \inf_{\mu > 0} \lambda^+(\mu)/\mu$ , and  $\lambda^+(\mu)/\mu$  attains its infimum at some  $\mu_2 \in (0, \infty)$ . Clearly,  $\lambda^+(\mu) = \lambda^-(-\mu)$  because  $\frac{\partial g(0, y)}{\partial w} = \frac{\partial h(0, y)}{\partial w} = f'(\alpha(y)).$ 

From assumption (D2), we see that  $S_1$  [S<sub>2</sub>] is a linearly unstable [stable] steady state of the  $x$ -independent system

$$
\begin{cases} \frac{\partial w}{\partial t} = B\Delta_y w + g(w, y), & y \in \Omega, \ t > 0, \\ \frac{\partial w}{\partial y} = 0 & \text{on } (0, \infty) \times \partial\Omega. \end{cases}
$$
(6.23)

More precisely, letting  $\lambda_0$  be the principal eigenvalue of the linearized equation at  $S_1$ , we have  $\lambda_0 > 0$ . Obviously, equations [\(6.21\)](#page-34-0) and [\(6.22\)](#page-34-1) with  $\mu = 0$  are the same, and hence  $\lambda^+(0) = \lambda_0 = \lambda^-(0) > 0.$ 

With the information above, now we can show that  $Q_1$  satisfies (A6). Let  $\theta = \frac{\mu_2}{\mu_1 + \mu_2}$ and  $\lambda^+(\mu) = \lambda^-(-\mu)$ . Note that  $\theta\mu_1 + (1 - \theta)(-\mu_2) = 0$ . Hence,

$$
c_{-}^{*}(\alpha, \beta) + c_{+}^{*}(0, \alpha) \ge \frac{\lambda^{-}(\mu_{1})}{\mu_{1}} + \frac{\lambda^{+}(\mu_{2})}{\mu_{2}} = \frac{\mu_{1} + \mu_{2}}{\mu_{1}\mu_{2}} [\theta \lambda^{-}(\mu_{1}) + (1 - \theta)\lambda^{-}(-\mu_{2})]
$$
  

$$
\ge \frac{\mu_{1} + \mu_{2}}{\mu_{1}\mu_{2}} \lambda^{-}(\theta \mu_{1} + (1 - \theta)(-\mu_{2}))
$$
  

$$
= \frac{\mu_{1} + \mu_{2}}{\mu_{1}\mu_{2}} \lambda^{-}(0) = \frac{\mu_{1} + \mu_{2}}{\mu_{1}\mu_{2}} \lambda_{0} > 0.
$$
 (6.24)

Consequently, Theorem [3.4](#page-18-0) completes the proof.  $\Box$ 

## *6.3. A parabolic equation with periodic diffusion*

In this subsection, we study the existence of spatially periodic traveling waves of the parabolic equation

<span id="page-35-0"></span>
$$
u_t = (d(x)u_x)_x + f(u), \quad t > 0, \ x \in \mathbb{R}, \tag{6.25}
$$

where  $f(u) = u(1 - u)(u - a)$ ,  $a \in (0, 1)$ , and  $d(x)$  is a positive,  $C<sup>1</sup>$ , and r-periodic function on  $\mathbb R$  for some real  $r > 0$ .

For any  $\phi \in C(\mathbb{R}, [0, 1])$ , equation [\(6.25\)](#page-35-0) admits a unique solution  $u(t; \phi)$  with  $u(0;\phi) = \phi$ . Define  $Q_t$ :  $C(\mathbb{R}, [0, 1]) \rightarrow C(\mathbb{R}, [0, 1])$  by  $Q_t[\phi] = u(t; \phi)$ . Then  ${Q_t}_{t\geq0}$  is a continuous, compact and monotone semiflow on  $C(\mathbb{R}, [0, 1])$  equipped with the compact-open topology. Let  $C_{\text{per}}(\mathbb{R}, [0, 1])$  be the set of all continuous and r-periodic functions from R to [0, 1]. Then the semiflow  ${Q_t}_{t\geq0}$  restricted to  $C_{\text{per}}(\mathbb{R}, [0, 1])$  is strongly monotone. Choosing  $\mathcal{H} = \mathbb{R}$  and  $\mathcal{X} = \mathbb{R}$  in Theorem [4.1,](#page-25-2) one can easily verify that  $\{Q_t\}_{t>0}$  satisfies assumptions (A2)–(A4). If [\(6.25\)](#page-35-0) admits the bistability structure, then Proposition  $3.1$  implies  $(AS'')$  and a similar argument to that in the previous section shows that (A6) also holds. Thus, we focus on finding sufficient conditions on  $d(x)$  under which  $(6.25)$  admits the bistability structure.

Let  $\bar{u}$  be an *r*-periodic steady state of [\(6.25\)](#page-35-0). As in [\[10\]](#page-43-11), we define  $\lambda_1(\bar{u}, d)$  as the largest number such that there exists a function  $\phi > 0$  which satisfies

$$
\begin{cases}\n(d\phi_x)_x + f'(\bar{u})\phi = \lambda_1(\bar{u}, d)\phi, & x \in \mathbb{R}, \\
\phi \text{ is } r\text{-periodic and } \|\phi\|_{\infty} = 1.\n\end{cases}
$$
\n(6.26)

We call  $\lambda_1(\bar{u}, d)$  the *principal eigenvalue* of  $\bar{u}$ , and  $\phi$  the corresponding eigenfunction. We say  $\bar{u}$  is *linearly unstable* if  $\lambda_1(\bar{u}, d) > 0$ , and *linearly stable* if  $\lambda_1(\bar{u}, d) < 0$ . Define

$$
W := \{ \psi \in C^1(\mathbb{R}, \mathbb{R}) : \psi(x) = \psi(x+r), \ \forall x \in \mathbb{R} \}
$$

and equip it with the topology induced by the maximum norm. We say  $\psi \in W$  has *property* (P) if every possible nonconstant r-periodic steady state of [\(6.25\)](#page-35-0) with  $d = \psi$ is linearly unstable, that is,  $(6.25)$  with  $d = \psi$  admits no nonconstant r-periodic steady state  $\bar{u}$  such that  $\lambda_1(\bar{u}, \psi) \leq 0$ . Define

$$
Y := \{ \psi \in W : \psi(x) > 0 \text{ and } \psi \text{ has property (P)} \}.
$$

## <span id="page-35-1"></span>Lemma 6.1. *Any positive constant function is in* Y *.*

*Proof.* Let  $d(x) \equiv \overline{d}$  be given. If [\(6.25\)](#page-35-0) has no nonconstant r-periodic steady state, we are done. Let  $\bar{u}$  be a nonconstant r-periodic steady state of [\(6.25\)](#page-35-0). We need to prove  $\lambda_1(\bar{u}, \bar{d}) > 0$ . Assume for contradiction that  $\lambda_1(\bar{u}, \bar{d}) \leq 0$ . Let  $\phi$  be the positive eigenfunction associated with  $\lambda_1(\bar{u}, \bar{d})$ . Define  $M := \max_{0 \le x \le r} {\{\overline{u}_x | \neq \phi\}}$  and  $\psi(x, t) :=$  $e^{-\gamma t}(|\bar{u}_x|^2/\phi - M^2\phi)$ . It is easy to see that  $\psi(t, x) \le 0$  for all x and t. Let  $\xi := |\bar{u}_x|^2/\phi$ and  $\eta := M^2 \phi$ . Then

$$
\xi_x = (|\bar{u}_x|^2 \phi^{-1})_x = 2\bar{u}_x \bar{u}_{xx} \phi^{-1} - |\bar{u}_x|^2 \phi^{-2} \phi_x,
$$
  

$$
\xi_{xx} = 2\phi^{-3} [\bar{u}_{xx} \phi - \bar{u}_x \phi_x]^2 + \phi^{-3} [2\bar{u}_x \bar{u}_{xxx} \phi^2 - |\bar{u}_x|^2 \phi \phi_{xx}].
$$

Note that

$$
0 = [\bar{d}\bar{u}_{xx} + f(\bar{u})]_x = \bar{d}\bar{u}_{xxx} + f'(\bar{u})\bar{u}_x.
$$

It follows that

$$
e^{\gamma t}(\psi_t - \bar{d}\psi_{xx} + [\gamma - f'(\bar{u})]\psi)
$$
  
= -[\bar{d}\xi\_{xx} + f'(\bar{u})\xi] + \lambda\_1(\bar{u}, \bar{d})\eta  
= -2\bar{d}\phi^{-3}[\bar{u}\_{xx}\phi - \bar{u}\_x\phi\_x]^2 - \bar{d}\phi^{-3}[2\bar{u}\_x\bar{u}\_{xxx}\phi^2 - |\bar{u}\_x|^2\phi\phi\_{xx}] - f'(\bar{u})|\bar{u}\_x|^2\phi^{-1}  
+ \lambda\_1(\bar{u}, \bar{d})\eta  
 $\leq \bar{d}|\bar{u}_x|^2\phi^{-2}\phi_{xx} + f'(\bar{u})\phi^{-1}|\bar{u}_x|^2 - 2f'(\bar{u})|\bar{u}_x|^2\phi^{-1} - 2\bar{d}\phi^{-1}\bar{u}_x\bar{u}_{xxx} + \lambda_1(\bar{u}, \bar{d})\eta= \lambda_1(\bar{u}, \bar{d})\xi + \lambda_1(\bar{u}, \bar{d})\eta - 2\bar{u}_x\phi^{-1}[f'(\bar{u})\bar{u}_x + \bar{d}\bar{u}_{xxx}] = \lambda_1(\bar{u}, \bar{d})[\xi + \eta].$ 

Hence,  $\psi_t - \bar{d}\psi_{xx} + [\gamma - f'(\bar{u})]\psi \le 0$  because  $\lambda_1(\bar{u}, \bar{d}) \le 0$ .

Since  $\bar{u}$  is not a constant and  $\psi(t, x)$  is r-periodic in  $x \in \mathbb{R}$ , we can choose  $x_0$  such that  $\psi(x_0, t) = \psi(x_0 + r, t) = \min_{x \in \mathbb{R}} \psi(x, t) < 0$ , and hence  $\psi_x|_{x=x_0} = \psi_x|_{x=x_0+r} = 0$ . Thus,  $\psi(t, x)$  with  $x \in [x_0, x_0 + r]$  satisfies the equation

$$
\begin{cases} \psi_t - \bar{d}\psi_{xx} + [\gamma - f'(\bar{u})] \psi \le 0, & x \in (x_0, x_0 + r), \\ \psi_x|_{x = x_0} = \psi_x|_{x = x_0 + r} = 0, \end{cases}
$$
\n(6.27)

and  $\psi(t, x)$  attains its maximum 0 at  $(x^*, t)$  with  $x^* \in (x_0, x_0 + r)$ . By the strong maximum principle, we see that  $\psi(t, x) \equiv 0$ , which implies that  $\bar{u}_x / \phi$  is a constant. Since  $\bar{u}_x/\phi$  is r-periodic, it follows that  $\bar{u}_x \equiv 0$ , and hence  $\bar{u}$  is a constant, a contradiction.  $\Box$ 

**Remark [6.1](#page-35-1).** By the proof above, the conclusion of Lemma 6.1 is valid for any  $f \in C^1$ .

## <span id="page-36-1"></span>Lemma 6.2. Y *is nonempty and open in* W*.*

*Proof.* Clearly, Lemma [6.1](#page-35-1) implies that  $Y \neq \emptyset$ . Let  $d^* \in Y$ . We need to show that  $d^*$  is an interior point of Y. Assume for contradiction that there is a sequence of points  $d_n \in W \setminus Y$  such that  $d_n \to d^*$  in W as  $n \to \infty$ . Then [\(6.25\)](#page-35-0) with  $d = d_n$  admits a nonconstant *r*-periodic steady state  $u_n$  with the principal eigenvalue  $\lambda_1(u_n, d_n) \leq 0$ . Using the transformation  $v_n = d_n(u_n)_x$ , we see that  $(u_n, v_n)$  is a periodic solution of the ordinary differential system

<span id="page-36-0"></span>
$$
\begin{cases}\n(u_n)_x = v_n/d_n, \\
(v_n)_x = -f(u_n).\n\end{cases}
$$
\n(6.28)

By elementary phase plane arguments, it then follows that

$$
0 \le \inf_{x \in \mathbb{R}} u_n(x) \le a \le \sup_{x \in \mathbb{R}} u_n(x) \le 1, \quad \forall n \ge 1, x \in \mathbb{R}.
$$
 (6.29)

Thus, the sequence of functions  $((u_n)_x, (v_n)_x)$  is uniformly bounded and equicontinuous, and hence  $(u_n, v_n)$  has a uniformly convergent subsequence, still denoted by  $(u_n, v_n)$ .

Let  $(u^*, v^*)$  be the limiting function of  $(u_n, v_n)$ . Then  $u^*$  is an *r*-periodic steady state of  $(6.25)$  with  $d = d^*$ . It is easy to see from  $(6.29)$  that  $u^*$  is not the constant function 0 or 1.

Let  $\phi_n$  be the positive eigenfunction associated with  $\lambda_1(u_n, d_n)$ . Then

<span id="page-37-1"></span><span id="page-37-0"></span>
$$
(d_n(\phi_n)_x)_x + f'(u_n(x))\phi_n = \lambda_1(u_n, d_n)\phi_n.
$$
\n
$$
(6.30)
$$

Dividing [\(6.30\)](#page-37-0) by  $\phi_n$  and integrating from 0 to r, we obtain

$$
\int_0^r \frac{d_n [(\phi_n)_x]^2}{\phi_n^2} dx + \int_0^r f'(u_n(x)) dx = \lambda_1(u_n, d_n) r \le 0.
$$
 (6.31)

Since  $f \in C^1$  and  $f'(a) > 0$ , we see that  $u^*$  cannot be the constant a: otherwise, the uniform convergence of  $u_n$  to a implies that  $f'(u_n(x)) > 0$  for all  $x \in [0, r]$  and sufficiently large *n*, which contradicts [\(6.31\)](#page-37-1). Thus,  $u^*$  is a nonconstant *r*-periodic function. Since  $d^* \in Y$ , we have  $\lambda_1(u^*, d^*) > 0$ .

Note that  $u_n \to u^*$  in  $C(\mathbb{R}, \mathbb{R})$  and  $d_n \to d^*$  in W. By the variational characterization of the principal eigenvalue  $\lambda_1(u_n, d_n)$  (see [\[10,](#page-43-11) (5.2)]), it then follows that  $0 \geq \lambda_1(u_n, d_n) \to \lambda_1(u^*, d^*) > 0$ , a contradiction.

The following counter-example shows that the parabolic equation [\(6.25\)](#page-35-0) admits no bistability structure for the general periodic function  $d(x)$ .

**Lemma 6.3.** *Let either*  $f(u) = u(1 - u^2)$ *, or*  $f(u) = u(1 - u)(u - 1/2)$ *. Then there exists a positive function* d ∈ W *such that* [\(6.25\)](#page-35-0) *admits a pair of linearly stable, nonconstant, and* r*-periodic steady states.*

*Proof.* We only consider the case where  $f(u) = u(1 - u^2)$  since the other one can be obtained under appropriate scalings. Our proof is based on the main result in [\[21,](#page-44-18) Theorem 3]. Without loss of generality, we assume that  $r = 4$ . In what follows, we use some notation of [\[21\]](#page-44-18).

Let  $l \in (0, 1)$  be fixed and  $c^0$  be the step function on [-1, 1] defined by

$$
c^{0}(x) = \begin{cases} 1, & x \in [-1, -l] \cup (l, 1], \\ 0, & x \in (-l, l]. \end{cases}
$$
(6.32)

Define *D* := {(*x*, *y*) : *x* = ±*l*, *y* ∈ [0, 1]} ∪ graph of  $c^0$ . By [\[21,](#page-44-18) Theorem 3], for any positive even function  $c \in C^1([-1, 1], \mathbb{R}^+)$  which is sufficiently close to  $c^0$  (in the sense that the distance between  $D$  and the graph of  $c$  is small enough), the Neumann boundary problem

$$
\begin{cases} u_t = (cu_x)_x + u(1 - u^2), & x \in (-1, 1), \\ u_x(t, \pm 1) = 0, \end{cases}
$$
 (6.33)

admits an odd increasing steady state  $u_c$  which is linearly stable. That is, there exist  $\lambda_1$  < 0 and  $\phi > 0$  such that

$$
\begin{cases}\n(c\phi_x)_x + f'(u_c)\phi = \lambda_1 \phi, & x \in (-1, 1), \\
\phi_x(\pm 1) = 0.\n\end{cases}
$$
\n(6.34)

In particular, we can choose c such that  $c_x(-1) = c_x(1) = 0$ . Since c is even and f is odd, we see that  $v_c(x) := u_c(-x)$  is also a steady state, and  $\lambda_1$  is the corresponding eigenvalue with the positive eigenfunction  $\phi(-x)$ .

Now we can construct a linearly stable 4-periodic steady state of [\(6.25\)](#page-35-0). Define two 4-periodic functions by

$$
\tilde{d}(x) = \begin{cases} c(x), & x \in [-1, 1], \\ c(2 - x), & x \in (1, 3), \end{cases} \qquad w_1(x) = \begin{cases} u_c(x), & x \in [-1, 1], \\ u_c(2 - x), & x \in (1, 3). \end{cases}
$$

Then  $w_1(x)$  is a 4-periodic steady state of [\(6.25\)](#page-35-0) with  $d = d$ . Let the positive 4-periodic function  $\rho(x)$  be defined by

$$
\rho(x) = \begin{cases} \phi(x), & x \in [-1, 1], \\ \phi(2 - x), & x \in (1, 3). \end{cases}
$$

Then  $\lambda_1$  and  $\rho$  solve the eigenvalue problem

$$
\begin{cases} (\tilde{d}\rho_x)_x + f'(w_1)\rho = \lambda_1 \rho, & x \in \mathbb{R}, \\ \rho \text{ is } r\text{-periodic.} \end{cases}
$$

This implies that  $w_1$  is a linearly stable periodic steady state of [\(6.25\)](#page-35-0) with  $d = \tilde{d}$ . Similarly, so is  $w_2(x) := w_1(x + 2)$ .

As a consequence of Theorem [4.1,](#page-25-2) together with Lemmas [6.1](#page-35-1) and [6.2,](#page-36-1) we have the following result on the existence of bistable traveling waves for [\(6.25\)](#page-35-0).

<span id="page-38-0"></span>**Theorem 6.3.** Let  $\bar{d}$  be a given positive constant. Then there exists  $\delta_0 > 0$  such that *for any*  $d \in W$  *with*  $\|d - \overline{d}\|_W < \delta_0$ , [\(6.25\)](#page-35-0) *admits a spatially periodic traveling wave solution*  $u(t, x) := V(x + ct, x)$  *with some speed*  $c \in \mathbb{R}$  *and connecting* 0 *to* 1*. Moreover,*  $V(\xi, x)$  *is nondecreasing in*  $\xi$ *.* 

We remark that Theorem [6.3](#page-38-0) is a  $C^0$ -perturbation result in W, and hence it improves the existence result in [\[42,](#page-45-9) Theorem 3.1], where  $H<sup>s</sup>$ -perturbation is used for some  $s > 2$ .

## *6.4. A nonlocal and time-delayed reaction-diffusion equation*

Fix  $\tau > 0$ . Choose  $\mathcal{X} := C([- \tau, 0], \mathbb{R}), \mathcal{Y} := C(\mathbb{R}, \mathbb{R})$  and  $\mathcal{C} := C([- \tau, 0], \mathcal{Y})$ . We equip X with the maximum norm, and Y and C with norms similar to [\(2.1\)](#page-4-2). Define  $\mathcal{Y}_+ :=$  $C(\mathbb{R}, \mathbb{R}_+)$ . Let d be the metric in  $C(\mathcal{Y})$  induced by the norm. We are interested in bistable traveling waves of the nonlocal and time-delayed reaction-diffusion equation

<span id="page-38-1"></span>
$$
\begin{cases} \frac{\partial u(t,x)}{\partial t} = \frac{\partial^2 u(t,x)}{\partial x^2} + f(u_t)(x), \quad t > 0, \ x \in \mathbb{R}, \\ u_0 = \phi \in \mathcal{C}, \end{cases}
$$
(6.35)

where  $f : C \to Y$  is Lipschitz continuous and for each  $t \geq 0$ ,  $u_t \in C$  is defined by

$$
u_t(\theta, x) := u(t + \theta, x), \quad \forall \theta \in [-\tau, 0], \ x \in \mathbb{R}.
$$

If the functional f takes the form  $f(\phi)(x) = F(\phi(0, x), \phi(-\tau, x))$ , then [\(6.35\)](#page-38-1) becomes a local and time-delayed reaction-diffusion equation

<span id="page-39-0"></span>
$$
\frac{\partial u(t,x)}{\partial t} = \frac{\partial^2 u(t,x)}{\partial x^2} + F(u(t,x), u(t-\tau, x)).
$$
\n(6.36)

The bistable traveling waves of [\(6.36\)](#page-39-0) were studied in [\[31\]](#page-44-8). If  $f(\phi)(x) = -d\phi(0, x) +$  $\int_{\mathbb{R}} b(\phi(-\tau, y))k(x - y) dy$ , then [\(6.35\)](#page-38-1) becomes a nonlocal and time-delayed reactiondiffusion equation

<span id="page-39-1"></span>
$$
\frac{\partial u(t,x)}{\partial t} = \frac{\partial^2 u(t,x)}{\partial x^2} - du(t,x) + \int_{\mathbb{R}} b(u(t-\tau,y))k(x-y) \, dy. \tag{6.37}
$$

The existence, uniqueness and stability of bistable waves of [\(6.37\)](#page-39-1) were established in [\[29\]](#page-44-9).

Note that  $\mathbb R$  can be regarded as a subspace of  $\mathcal X$ , and the latter can also be regarded as a subspace of C. Define  $\bar{f}: \mathcal{X} \to \mathbb{R}$  by  $\bar{f}(\varphi) = f(\varphi)$  and  $\hat{f}: \mathbb{R} \to \mathbb{R}$  by  $\hat{f}(\xi) = f(\xi)$ . In order to obtain the existence of bistable waves for system  $(6.35)$ , we impose the following assumptions on the functional  $f$ :

(E1)  $0 < \alpha < \beta$  are three equilibria and there are no other equilibria between 0 and  $\beta$ . (E2) The functional  $f : C_\beta \to \mathcal{Y}$  is quasi-monotone in the sense that

$$
\lim_{h \to 0^+} \frac{1}{h} d\big([\phi(0) - \psi(0)] + h[f(\phi) - f(\psi)]; \mathcal{Y}_+\big) = 0 \quad \text{whenever } \phi \ge \psi \text{ in } \mathcal{C}_{\beta}.
$$

- (E3) The equilibria 0 and  $\beta$  are stable, and  $\alpha$  is unstable in the sense that  $\hat{f}'(0) < 0$ ,  $\hat{f}'(\alpha) > 0$  and  $\hat{f}'(\beta) < 0$ .
- (E4) For each  $\varphi \in \mathcal{X}_{\beta}$ , the derivative  $\bar{L}(\varphi) := D\bar{f}(\varphi)$  of  $\bar{f}$  can be represented as

$$
\bar{L}(\varphi)\chi = a(\varphi)\chi(0) + \int_{-\tau}^{0} \chi(\theta) d_{\theta}\eta(\varphi) =: a(\varphi)\chi(0) + L_{1}(\varphi)\chi,
$$

where  $\eta(\varphi)$  is a positive Borel measure on  $[-\tau, 0]$  and  $\eta(\varphi)([-\tau, -\tau + \epsilon]) > 0$  for all small  $\epsilon > 0$ .

(E5) For any small  $\epsilon > 0$ , there exists  $\delta \in (0, \beta)$  and a linear operator  $L_{\epsilon} : C_{\beta} \to Y$ such that  $L_{\epsilon} \phi \to Df(\alpha) \phi$  for all  $\phi \in C_{\beta}$  as  $\epsilon \to 0$  and

<span id="page-39-2"></span>
$$
f(\alpha + \phi) \ge L_{\epsilon}(\phi)
$$
 and  $f(\alpha - \phi) \le -L_{\epsilon}(\phi)$ ,  $\forall \phi \in C_{\delta}$ .

Using the solution maps  $\{T(t)\}_{t\geq 0}$  generated by the heat equation  $\frac{\partial u(t,x)}{\partial t} = \frac{\partial^2 u(t,x)}{\partial x^2}$  $\frac{u(t,x)}{\partial x^2}$ , we write system  $(6.35)$  in integral form

$$
\begin{cases} u(t, \cdot) = T(t)\phi(0, \cdot) + \int_0^t T(t - r) f(u_r(\cdot, \cdot)) dr, & t > 0, \\ u(\theta, \cdot) = \phi(\theta, \cdot), & \theta \in [-\tau, 0]. \end{cases}
$$
(6.38)

Note that traveling waves of system [\(6.38\)](#page-39-2) are those of system [\(6.35\)](#page-38-1). It then remains to show that  $(6.38)$  admits a bistable traveling wave.

Theorem 6.4. *Under assumptions* (E1)–(E5)*, system* [\(6.35\)](#page-38-1) *admits a nondecreasing traveling wave*  $\phi(x + ct)$  *with*  $\phi(-\infty) = 0$  *and*  $\phi(\infty) = \beta$ *.* 

*Proof.* From assumptions  $(E1)$ – $(E2)$ , we see that system  $(6.38)$  generates a monotone semiflow  $\{Q_t\}_{t>0}$  on  $\mathcal{C}_{\beta}$  with

$$
Q_t[\phi](\theta, x) = u_t(\theta, x; \phi), \quad \forall (\theta, x) \in [-\tau, 0] \times \mathbb{R},
$$

where  $u(t, x; \phi)$  is the unique solution of system [\(6.38\)](#page-39-2) satisfying  $u_0(\cdot, \cdot; \phi) = \phi \in C_\beta$ . By similar arguments to those in Section 5, the map  $Q_t$  satisfies (A4) if  $t > \tau$  and (A4') if  $t \in (0, \tau]$ .

Let  $\overline{Q}_t$  be the restriction of  $Q_t$  to  $\mathcal{X}_{\beta}$ . Denote the derivative  $D\overline{Q}_t[\hat{0}]$  of  $\overline{Q}_t$  by  $\overline{M}_{0,t}$ . Then  $\overline{M}_{0,t}$  is the solution map of the functional equation

<span id="page-40-1"></span>
$$
\frac{du}{dt} = \bar{L}(0)u_t = a(0)u(t) + L_1(0)u_t.
$$
\n(6.39)

By assumptions (E2) and (E4), system  $(6.39)$  admits a principal eigenvalue  $s_0$  with an associated eigenfunction  $v_0 := e^{s_0 \theta}$  (see [\[34,](#page-45-14) Theorem 5.5.1]). More precisely,  $\bar{M}_{0,t}[v_0] =$  $e^{s_0 t}v_0$ . Furthermore, [\[34,](#page-45-14) Corollary 5.5.2] implies that  $s_0 < 0$  since  $\hat{f}'(0) < 0$ . Therefore, there exists  $\delta_0(t) > 0$  such that

$$
\bar{Q}_t[\delta v_0] = \bar{Q}_t[0] + D\bar{Q}_t[0][\delta v_0] + o(\delta^2) = \delta \bar{M}_{0,t}[v_0] + o(\delta^2) \n= \delta e^{s_0 t} v_0 + o(\delta^2) = \delta v_0 + \delta [e^{s_0 t} - 1]v_0 + o(\delta^2) \ll \delta v_0, \quad \forall \delta \in (0, \delta_0(t)].
$$

Similarly, there exist  $\delta_{\alpha}(t)$ ,  $v_{\alpha}$  and  $\delta_{\beta}(t)$ ,  $v_{\beta}$  such that

$$
\overline{Q}_t[\beta - \delta v_\beta] \gg \beta - \delta v_\beta, \quad \forall \delta \in (0, \delta_\beta(t)],
$$

and

$$
\overline{Q}_t[\alpha + \delta v_\alpha] \gg \alpha + \delta v_\alpha, \quad \overline{Q}_t[\alpha - \delta v_\alpha] \ll \alpha - \delta v_\alpha, \quad \forall \delta \in (0, \delta_\alpha(t)].
$$

It remains to show that (A6) is also true. Indeed, from [\[25,](#page-44-14) Theorem 2.17] the solution semiflows  $\{Q_t\}_{t\geq 0}$  restricted to  $[0, \alpha]_{\mathcal{C}}$  and  $[\alpha, \beta]_{\mathcal{C}}$  admit spreading speeds  $c^*(0, \alpha)$  and  $c^*(\alpha, \beta)$ , respectively. Let  $M^t_{\epsilon}$  be the solution maps of the linear system

$$
\begin{cases} u(t,\cdot) = T(t)\phi(0,\cdot) + \int_0^t T(t-r)L_{\epsilon}(u_r(\cdot,\cdot)) dr, & t > 0, \\ u(\theta,\cdot) = \phi(\theta,\cdot), & \theta \in [-\tau,0]. \end{cases}
$$

Then assumption (E5) guarantees that  $Q_t[\phi] \geq M_t^{\epsilon}[\phi]$  when  $\phi \in C_\delta$ , where  $\delta = \delta(\epsilon)$ is defined in (E5). Therefore, we see from [\[25,](#page-44-14) Theorem 3.10] that  $c^*(0, \alpha) \geq \overline{c}$  and  $c^*(\alpha, \beta) \geq \bar{c}$ , where  $\bar{c}$  is a positive number determined by the linearized system of [\(6.35\)](#page-38-1) at  $u \equiv \alpha$ , and hence (A6) holds. Consequently, Theorem [5.1](#page-28-1) completes the proof.  $\Box$ 

<span id="page-40-0"></span>Remark 6.2. At this moment we are unable to present a general result on the uniqueness and global attractivity of bistable waves under the current abstract setting. However, one may use the convergence theorem for monotone semiflows (see [\[47,](#page-45-15) Theorem 2.2.4]) and similar arguments to the proofs of [\[47,](#page-45-15) Theorem 10.2.1] and [\[45,](#page-45-18) Theorem 3.1] to obtain the global attractivity (and hence uniqueness) of bistable waves for the four examples in this section.

# 7. Appendix

In this appendix, we present certain properties of Banach lattices and countable subsets of R, and some convergence results for sequences of monotone functions, including an abstract variant of Helly's theorem.

<span id="page-41-0"></span>Proposition 7.1. *Every Banach lattice* X *has the following properties:*

- (1) *For any*  $u, v \in \mathcal{X}$  *with*  $v \in \mathcal{X}^+$ *, if*  $-v \le u \le v$ *, then*  $||u||_{\mathcal{X}} \le ||v||_{\mathcal{X}}$ *.*
- (2) If  $u_k \to u$  and  $v_k \to v$  in X with  $u_k \ge v_k$ , then  $u \ge v$ .

<span id="page-41-1"></span>Proposition 7.2. *The space* C *has the following properties:*

- (1) Let  $\phi$  *be a monotone function in* C. If  $x_k \in \mathcal{H}$  *nondecreasingly tends to*  $x \in \mathcal{H} \cup \{\infty\}$ *and*  $\lim_{k\to\infty} \phi(x_k) = u \in \mathcal{X}$ , then  $\lim_{y \uparrow x} \phi(y) = u$ . A similar result holds if  $x_k$ *nonincreasingly tends to*  $x \in \mathcal{H} \cup \{-\infty\}$ *.*
- (2) Assume that  $h, h_k : \mathcal{H} \to \mathcal{H}$  are continuous and  $\phi_k \to \phi$  in C. If  $h_k(x) \to h(x)$ *uniformly for* x *in any bounded subset of*  $H$ *, then*  $\phi_k \circ h_k \to \phi \circ h$  *in* C.

<span id="page-41-2"></span>We omit the easy proofs of Propositions [7.1](#page-41-0) and [7.2.](#page-41-1)

**Proposition 7.3.** Assume that D is a countable subset of  $\mathbb{R}$ . Then for any  $c \in \mathbb{R}$ , there *exists another countable subset* A *of*  $\mathbb R$  *such that*  $(\mathbb R \setminus A) + cm \subset \mathbb R \setminus D$  *for all*  $m \in \mathbb Z^+$ *.* 

*Proof.* It suffices to show the set  $A := \{x \in \mathbb{R} : \text{there exists } m \text{ such that } x + cm \in D\}$  is countable. Indeed, we have  $A = \bigcup_{m=1}^{\infty} (D - cm)$ . This implies A is countable.

<span id="page-41-4"></span>**Proposition 7.4.** *For any countable subset*  $\Gamma_1$  *of*  $\mathbb{R}$ *, there exists another countable set*  $\Gamma_2$ *that is dense in*  $\mathbb R$  *and*  $\Gamma_1 \cap \Gamma_2 = \emptyset$ *.* 

*Proof.* Since  $\Gamma_1$  is countable and  $\bigcup_{\alpha \in \mathbb{R}} (\alpha + \mathbb{Q}) = \mathbb{R}$ , there must exist a sequence  $\alpha_n$ such that  $\Gamma_1 \subset \bigcup_{n=1}^{\infty} (\alpha_n + \mathbb{Q})$ . Note that  $\bigcup_{n=1}^{\infty} (\alpha_n + \mathbb{Q})$  is countable. Hence, there exists  $\alpha \in \mathbb{R}$  such that  $\alpha \notin \bigcup_{n=1}^{\infty} (\alpha_n + \mathbb{Q})$ . This means that  $\alpha - \alpha_n \notin \mathbb{Q}$  for all  $n \geq 1$ , and hence  $(\alpha + \mathbb{Q}) \cap (\alpha_n + \mathbb{Q}) = \emptyset$  for all  $n \geq 1$ . Define  $\Gamma_2 := \alpha + \mathbb{Q}$ . We then see that  $\Gamma_2$ is countable and dense in R, and  $\Gamma_1 \cap \Gamma_2 = \emptyset$ .

<span id="page-41-3"></span>**Proposition 7.5.** Assume that  $f, f_n : \mathbb{R} \to \mathcal{X}$  are nondecreasing and the set D is dense *in*  $\mathbb{R}$ *. If*  $s_n \to 0$ *,*  $f(s)$  *is continuous on* D *and*  $f_n(s) \to f(s)$  *for every*  $s \in D$ *, then*  $f_n(s + s_n) \to f(s)$  *for every*  $s \in D$ *.* 

*Proof.* Let  $s \in D$  be fixed. For any  $\delta > 0$ , since  $D - s$  is dense in R, we can choose  $\delta_+ \in (D-s) \cap (0, \delta)$  and  $\delta_- \in (D-s) \cap (-\delta, 0)$ . Clearly,  $s + \delta_+$ ,  $s + \delta_- \in D$ . Thus, there exists an integer  $N_{\delta}$  such that  $s + s_n \in (s + \delta_-, s + \delta_+)$  for all  $n \geq N_{\delta}$ . Since

$$
f_n(s+\delta_-)-f_n(s+\delta_+) \le f_n(s+s_n)-f_n(s) \le f_n(s+\delta_+) - f_n(s+\delta_-), \quad \forall n \ge N_\delta,
$$

we have

$$
|| f_n(s+s_n) - f_n(s) ||_{\mathcal{X}} \le || f_n(s+\delta_+) - f_n(s+\delta_-) ||_{\mathcal{X}}, \quad \forall n \ge N_\delta.
$$

It then follows that

$$
|| f_n(s + s_n) - f(s) ||_{\mathcal{X}} \le || f_n(s + s_n) - f_n(s) ||_{\mathcal{X}} + || f_n(s) - f(s) ||_{\mathcal{X}}
$$
  
\n
$$
\le || f_n(s + \delta_+) - f_n(s + \delta_-) ||_{\mathcal{X}} + || f_n(s) - f(s) ||_{\mathcal{X}}
$$
  
\n
$$
\le || f(s + \delta_+) - f_n(s + \delta_+) ||_{\mathcal{X}} + || f(s + \delta_+) - f(s + \delta_-) ||_{\mathcal{X}}
$$
  
\n
$$
+ || f_n(s + \delta_-) - f(s + \delta_-) ||_{\mathcal{X}} + || f_n(s) - f(s) ||_{\mathcal{X}}
$$

for all  $n > N_\delta$ . Now the pointwise convergence of  $f_n$  in D and the continuity of f on D complete the proof.  $\Box$ 

To end this section, we prove a convergence theorem for sequences of monotone functions from  $\mathbb R$  to the special Banach lattice  $C(M, \mathbb R^d)$  defined in Section 2, which is a variant of Helly's theorem [\[17,](#page-44-19) p. 165] for sequences of monotone functions from  $\mathbb R$  to  $\mathbb R$ .

<span id="page-42-0"></span>**Theorem 7.1.** Let D be a dense subset of  $\mathbb{R}$  and  $f_n, n \geq 1$ , be a sequence of nonde*creasing functions from*  $\mathbb R$  *to the Banach lattice*  $\mathcal X := C(M, \mathbb R^d)$ *. Assume that:* 

- (i) *For any*  $s \in D$ *,*  $f_n(s)$  *is convergent in*  $\mathcal{X}$ *.*
- (ii) *There exists a countable set*  $D_1 \subset \mathbb{R}$  *such that for any*  $s \in \mathbb{R} \setminus D_1$ *, the limits* lim<sub>m→∞</sub> lim<sub>n→∞</sub>  $f_n(s_{\pm,m})$  *exist in*  $\mathcal{X}$ *, where*  $s_{-,m}$   $\uparrow$  *s and*  $s_{+,m}$   $\downarrow$  *s with*  $s_{\pm,m}$   $\in$  *D.*

*Then*  $f_n(s)$  *is convergent in*  $X$  *for almost all*  $s \in \mathbb{R}$ *.* 

*Proof.* Due to assumption (ii), we can define  $f : \mathbb{R} \to \mathcal{X}$  by

$$
f(s) := \begin{cases} \lim_{x \uparrow s} \lim_{x \in D, n \to \infty} f_n(x), & s \in \mathbb{R} \setminus D_1, \\ \text{any value}, & s \in D_1. \end{cases}
$$
(7.1)

We first show that the set of discontinuity points of  $f$  is at most countable. Define

$$
A := \{ s \in \mathbb{R} \setminus D_1 : f(s^-), f(s^+) \text{ both exist} \}, \quad B := \{ s \in A : f(s^-) < f(s^+) \}.
$$

For any  $s \in B$ , there are  $x_1 \in M$  and  $1 \le i \le d$  such that  $(f(s^{-})(x_1))_i < (f(s^{+})(x_1))_i$ . Recall that M is compact, so there is a countable dense subset  $M_1$ . It then follows that there must be  $x_2 \in M_1$  such that  $(f(s^-)(x_2))_i < (f(s^+)(x_2))_i$ . Therefore,

$$
B = \bigcup_{i=1}^{m} \bigcup_{x \in M_1} \{ s \in A : (f(s^{-})(x))_i < (f(s^{+})(x))_i \}.
$$

Since for each fixed i and x,  $(f(s)(x))_i$  is a nondecreasing function from  $\mathbb{R} \setminus D_1$  to  $\mathbb{R}$ , the set  $\{s \in A : (f(s^-)(x))_i < (f(s^+)(x))_i\}$  is at most countable, and hence so is B.

Now we can prove the conclusion. Assume that  $s \in \mathbb{R}$  is a continuity point of f. For any  $\delta > 0$ , choose  $\delta_- \in D \cap (s - \delta, s)$  and  $\delta_+ \in D \cap (s, s + \delta)$ . Then we have

$$
f_n(\delta_-) - f_n(\delta_+) \le f_n(s) - f_n(\delta_-) \le f_n(\delta_+) - f_n(\delta_-), \quad \forall n \ge 1,
$$

which, together with Proposition  $7.1(2)$  $7.1(2)$ , implies that

<span id="page-42-1"></span>
$$
|| f_n(s) - f_n(\delta_-) ||_{\mathcal{X}} \le || f_n(\delta_+) - f_n(\delta_-) ||_{\mathcal{X}}, \quad \forall n \ge 1.
$$
 (7.2)

On the other hand, by [\(7.2\)](#page-42-1) and the triangular inequality we have

$$
|| f_n(s) - f(s) ||_{\mathcal{X}} \le || f_n(s) - f_n(\delta_-) ||_{\mathcal{X}} + || f_n(\delta_-) - f(s) ||_{\mathcal{X}}
$$
  
\n
$$
\le || f_n(\delta_+) - f_n(\delta_-) ||_{\mathcal{X}} + || f_n(\delta_-) - f(s) ||_{\mathcal{X}}
$$
  
\n
$$
\le || f_n(\delta_+) - f(\delta_+) ||_{\mathcal{X}} + || f(\delta_+) - f(\delta_-) ||_{\mathcal{X}}
$$
  
\n
$$
+ || f(\delta_-) - f_n(\delta_-) ||_{\mathcal{X}} + || f_n(\delta_-) - f(\delta_-) ||_{\mathcal{X}}
$$
  
\n
$$
+ || f(\delta_-) - f(s) ||_{\mathcal{X}}, \quad \forall n \ge 1.
$$

Now the pointwise convergence of  $f_n$  in D and the continuity of f at s complete the  $\Box$ 

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#### References

- <span id="page-43-2"></span>[1] Alikakos, N., Bates, P., Chen, X.: Periodic traveling waves and locating oscillating patterns in multidimensional domains. Trans. Amer. Math. Soc. 351, 2777–2805 (1999) [Zbl 0929.35067](http://www.zentralblatt-math.org/zmath/en/advanced/?q=an:0929.35067&format=complete) [MR 1467460](http://www.ams.org/mathscinet-getitem?mr=1467460)
- <span id="page-43-10"></span>[2] Aronson, D. G., Weinberger, H. F.: Multidimensional nonlinear diffusion arising in population dynamics. Adv. Math. 30, 33–76 (1978) [Zbl 0407.92014](http://www.zentralblatt-math.org/zmath/en/advanced/?q=an:0407.92014&format=complete) [MR 0511740](http://www.ams.org/mathscinet-getitem?mr=0511740)
- <span id="page-43-1"></span>[3] Bates, P., Chmaj, A.: A discrete convolution model for phase transitions. Arch. Ration. Mech. Anal. 150, 281–305 (1999) [Zbl 0956.74037](http://www.zentralblatt-math.org/zmath/en/advanced/?q=an:0956.74037&format=complete) [MR 1741258](http://www.ams.org/mathscinet-getitem?mr=1741258)
- <span id="page-43-0"></span>[4] Bates, P., Fife, P., Ren, X., Wang, X.: Traveling waves in a convolution model for phase transitions. Arch. Ration. Mech. Anal. 138, 105–136 (1997) [Zbl 0889.45012](http://www.zentralblatt-math.org/zmath/en/advanced/?q=an:0889.45012&format=complete) [MR 1463804](http://www.ams.org/mathscinet-getitem?mr=1463804)
- <span id="page-43-7"></span>[5] Berestycki, H., Hamel, F.: Front propagation in periodic excitable media. Comm. Pure Appl. Math. 55, 949–1032 (2002) [Zbl 1024.37054](http://www.zentralblatt-math.org/zmath/en/advanced/?q=an:1024.37054&format=complete) [MR 1900178](http://www.ams.org/mathscinet-getitem?mr=1900178)
- <span id="page-43-5"></span>[6] Berestycki, H., Hamel, F.: Non-existence of travelling front solutions of some bistable reaction-diffusion equations. Adv. Differential Equations 5, 723–746 (2000) [Zbl 0988.35081](http://www.zentralblatt-math.org/zmath/en/advanced/?q=an:0988.35081&format=complete) [MR 1750116](http://www.ams.org/mathscinet-getitem?mr=1750116)
- <span id="page-43-6"></span>[7] Berestycki, H., Hamel, F., Matano, H.: Bistable traveling waves around an obstacle. Comm. Pure Appl. Math. 62, 729–788 (2009) [Zbl 1172.35031](http://www.zentralblatt-math.org/zmath/en/advanced/?q=an:1172.35031&format=complete) [MR 2512611](http://www.ams.org/mathscinet-getitem?mr=2512611)
- <span id="page-43-8"></span>[8] Berestycki, H., Hamel, F., Nadirashvili, N.: The speed of propagation for KPP type problems. I. Periodic framework. J. Eur. Math. Soc. 7, 173–213 (2005) [Zbl 1142.35464](http://www.zentralblatt-math.org/zmath/en/advanced/?q=an:1142.35464&format=complete) [MR 2127993](http://www.ams.org/mathscinet-getitem?mr=2127993)
- <span id="page-43-9"></span>[9] Berestycki, H., Hamel, F., Nadirashvili, N.: The speed of propagation for KPP type problems. II. General domains. J. Amer. Math. Soc. 23, 1–34 (2010) [Zbl 1197.35073](http://www.zentralblatt-math.org/zmath/en/advanced/?q=an:1197.35073&format=complete) [MR 2552247](http://www.ams.org/mathscinet-getitem?mr=2552247)
- <span id="page-43-11"></span>[10] Berestycki, H., Hamel, F., Roques, L.: Analysis of the periodically fragmented environment model: I – Species persistence. J. Math. Biol. 51, 75–113 (2005) [Zbl 1066.92047](http://www.zentralblatt-math.org/zmath/en/advanced/?q=an:1066.92047&format=complete) [MR 2214420](http://www.ams.org/mathscinet-getitem?mr=2214420)
- <span id="page-43-4"></span>[11] Berestycki, H., Nirenberg, L.: Travelling fronts in cylinders. Ann. Inst. H. Poincare Anal. Non ´ Linéaire 9, 497-572 (1992) [Zbl 0799.35073](http://www.zentralblatt-math.org/zmath/en/advanced/?q=an:0799.35073&format=complete) [MR 1191008](http://www.ams.org/mathscinet-getitem?mr=1191008)
- <span id="page-43-3"></span>[12] Casten, R. G., Holland, C. J.: Instability results for reaction-diffusion equations with Neu-mann boundary conditions. J. Differential Equations 12, 266-273 (1978) [Zbl 0338.35055](http://www.zentralblatt-math.org/zmath/en/advanced/?q=an:0338.35055&format=complete) [MR 0480282](http://www.ams.org/mathscinet-getitem?mr=0480282)
- <span id="page-44-3"></span>[13] Chen, X.: Existence, uniqueness, and asymptotic stability of traveling waves in nonlocal evolution equations. Adv. Differential Equations 2, 125–160 (1997) [Zbl 1023.35513](http://www.zentralblatt-math.org/zmath/en/advanced/?q=an:1023.35513&format=complete) [MR 1424765](http://www.ams.org/mathscinet-getitem?mr=1424765)
- <span id="page-44-10"></span>[14] Chen, X., Guo, J. S., Wu, C. C.: Traveling waves in discrete periodic media for bistable dynamics. Arch. Ration. Mech. Anal. 189, 189–236 (2008) [Zbl 1152.37033](http://www.zentralblatt-math.org/zmath/en/advanced/?q=an:1152.37033&format=complete) [MR 2413095](http://www.ams.org/mathscinet-getitem?mr=2413095)
- <span id="page-44-5"></span>[15] Coutinho, R., Fernandez, B.: Fronts in extended systems of bistable maps coupled via convolutions. Nonlinearity 17, 23–47 (2004) [Zbl 1083.37059](http://www.zentralblatt-math.org/zmath/en/advanced/?q=an:1083.37059&format=complete) [MR 2023429](http://www.ams.org/mathscinet-getitem?mr=2023429)
- <span id="page-44-2"></span>[16] Coville, J.: Traveling fronts in asymmetric nonlocal reaction diffusion equation: The bistable and ignition cases. Preprint of CMM (2012)
- <span id="page-44-19"></span>[17] Doob, J. L.: Measure Theory, Springer, New York (1994) [Zbl 0791.28001](http://www.zentralblatt-math.org/zmath/en/advanced/?q=an:0791.28001&format=complete) [MR 1253752](http://www.ams.org/mathscinet-getitem?mr=1253752)
- <span id="page-44-6"></span>[18] Fang, J., Zhao, X.-Q.: Monotone wavefronts for partially degenerate reaction-diffusion systems. J. Dynam. Differential Equations 21, 663–680 (2009) [Zbl 1180.35275](http://www.zentralblatt-math.org/zmath/en/advanced/?q=an:1180.35275&format=complete) [MR 2575368](http://www.ams.org/mathscinet-getitem?mr=2575368)
- <span id="page-44-0"></span>[19] Fife, P. C., McLeod, J. B.: The approach of solutions of nonlinear diffusion equations to travelling front solutions. Arch. Ration. Mech. Anal. 65, 335–361 (1977) [Zbl 0361.35035](http://www.zentralblatt-math.org/zmath/en/advanced/?q=an:0361.35035&format=complete) [MR 0442480](http://www.ams.org/mathscinet-getitem?mr=0442480)
- <span id="page-44-1"></span>[20] Fife, P. C., McLeod, J. B.: A phase plane discussion of convergence to travelling fronts for nonlinear diffusion. Arch. Ration. Mech. Anal. 75, 281–314 (1980/81) [Zbl 0459.35044](http://www.zentralblatt-math.org/zmath/en/advanced/?q=an:0459.35044&format=complete) [MR 0607901](http://www.ams.org/mathscinet-getitem?mr=0607901)
- <span id="page-44-18"></span>[21] Fusco, G., Hale, J.: Stable equilibria in a scalar parabolic equation with variable diffusion. SIAM J. Math. Anal. 16, 1152–1164 (1985) [Zbl 0597.35040](http://www.zentralblatt-math.org/zmath/en/advanced/?q=an:0597.35040&format=complete) [MR 0807902](http://www.ams.org/mathscinet-getitem?mr=0807902)
- <span id="page-44-11"></span>[22] Guo, J. S., Hamel, F.: Front propagation for discrete periodic monostable equations. Math. Ann. 335, 489–525 (2006) [Zbl 1116.35063](http://www.zentralblatt-math.org/zmath/en/advanced/?q=an:1116.35063&format=complete) [MR 2221123](http://www.ams.org/mathscinet-getitem?mr=2221123)
- <span id="page-44-12"></span>[23] Guo, J. S., Wu, C.-C.: Uniqueness and stability of traveling waves for periodic monostable lattice dynamical system. J. Differential Equations 246, 3818–3833 (2009) [Zbl 1202.34024](http://www.zentralblatt-math.org/zmath/en/advanced/?q=an:1202.34024&format=complete) [MR 2514727](http://www.ams.org/mathscinet-getitem?mr=2514727)
- <span id="page-44-17"></span>[24] Kishimoto, K., Weinberger, H. F.: The spatial homogeneity of stable equilibria of some reaction-diffusion systems on convex domains. J. Differential Equations 58, 15–21 (1985) [Zbl 0599.35080](http://www.zentralblatt-math.org/zmath/en/advanced/?q=an:0599.35080&format=complete) [MR 0791838](http://www.ams.org/mathscinet-getitem?mr=0791838)
- <span id="page-44-14"></span>[25] Liang, X., Zhao, X.-Q.: Asymptotic speeds of spread and traveling waves for monotone semiflows with applications. Comm. Pure Appl. Math.  $60$ ,  $1-40$  ( $2007$ ) Zbl  $1106.76008$ [MR 2361306](http://www.ams.org/mathscinet-getitem?mr=2361306)
- <span id="page-44-13"></span>[26] Liang, X., Zhao, X.-Q.: Spreading speeds and traveling waves for abstract monostable evolution systems. J. Funct. Anal. 259, 857–903 (2010) [Zbl 1201.35068](http://www.zentralblatt-math.org/zmath/en/advanced/?q=an:1201.35068&format=complete) [MR 2270161](http://www.ams.org/mathscinet-getitem?mr=2270161)
- <span id="page-44-4"></span>[27] Lui, R.: Existence and stability of travelling wave solutions of a nonlinear integral operator. J. Math. Biol. 16, 199–220 (1983) [Zbl 03812247](http://www.zentralblatt-math.org/zmath/en/advanced/?q=an:03812247&format=complete) [MR 0693376](http://www.ams.org/mathscinet-getitem?mr=0693376)
- <span id="page-44-15"></span>[28] Lui, R.: Biological growth and spread modeled by systems of recursions. I. Mathematical theory. Math. Biosci. 93, 269–295 (1989) [Zbl 0706.92014](http://www.zentralblatt-math.org/zmath/en/advanced/?q=an:0706.92014&format=complete) [MR 0984281](http://www.ams.org/mathscinet-getitem?mr=0984281)
- <span id="page-44-9"></span>[29] Ma, S., Wu, J.: Existence, uniqueness and asymptotic stability of traveling wavefronts in a non-local delayed diffusion equation. J. Dynam. Differential Equations 19, 391–436 (2007) [Zbl 1157.35057](http://www.zentralblatt-math.org/zmath/en/advanced/?q=an:1157.35057&format=complete) [MR 2333414](http://www.ams.org/mathscinet-getitem?mr=2333414)
- <span id="page-44-7"></span>[30] Matano, H., Asymptotic behavior and stability of solutions of semilinear diffusion equations. Publ. RIMS Kyoto Univ. 15, 401–454 (1979) [Zbl 0445.35063](http://www.zentralblatt-math.org/zmath/en/advanced/?q=an:0445.35063&format=complete) [MR 0555661](http://www.ams.org/mathscinet-getitem?mr=0555661)
- <span id="page-44-8"></span>[31] Schaaf, K. W.: Asymptotic behavior and traveling wave solutions for parabolic functional differential equations. Trans. Amer. Math. Soc. 302, 587–615 (1987) [Zbl 0637.35082](http://www.zentralblatt-math.org/zmath/en/advanced/?q=an:0637.35082&format=complete) [MR 0891637](http://www.ams.org/mathscinet-getitem?mr=0891637)
- <span id="page-44-16"></span>[32] Schaefer, H. H.: Banach Lattices and Positive Operators, Springer, New York (1974) [Zbl 0296.47023](http://www.zentralblatt-math.org/zmath/en/advanced/?q=an:0296.47023&format=complete) [MR 0423039](http://www.ams.org/mathscinet-getitem?mr=0423039)
- <span id="page-45-10"></span><span id="page-45-0"></span>[33] Shen, W.: Traveling waves in time almost periodic structures governed by bistable nonlinearities, I–II. J. Differential Equations 159, 1–54, 55–101 (1999) [Zbl 0939.35016\(](http://www.zentralblatt-math.org/zmath/en/advanced/?q=an:0939.35016&format=complete)I) [Zbl 0939.35015\(](http://www.zentralblatt-math.org/zmath/en/advanced/?q=an:0939.35015&format=complete)II) [MR 1726918\(](http://www.ams.org/mathscinet-getitem?mr=1726918)I) [MR 1726919\(](http://www.ams.org/mathscinet-getitem?mr=1726919)II)
- <span id="page-45-14"></span>[34] Smith, H. L.: Monotone Dynamical Systems: An Introduction to the Theory of Competitive and Cooperative Systems. Math. Surveys Monogr. 41, Amer. Math. Soc., Providence, RI (1995) [Zbl 0821.34003](http://www.zentralblatt-math.org/zmath/en/advanced/?q=an:0821.34003&format=complete) [MR 1319817](http://www.ams.org/mathscinet-getitem?mr=1319817)
- <span id="page-45-5"></span>[35] Smith, H. L., Zhao, X.-Q.: Global asymptotic stability of traveling waves in delayed reaction-diffusion equations. SIAM J. Math. Anal. 31, 514–534 (2000) [Zbl 0941.35125](http://www.zentralblatt-math.org/zmath/en/advanced/?q=an:0941.35125&format=complete) [MR 1740724](http://www.ams.org/mathscinet-getitem?mr=1740724)
- <span id="page-45-4"></span>[36] Volpert, A. I., Volpert, V. A., Volpert, V. A.: Traveling Wave Solutions of Parabolic Systems. Transl. Math. Monogr. 140, Amer. Math. Soc. (1994) [Zbl 1001.35060](http://www.zentralblatt-math.org/zmath/en/advanced/?q=an:1001.35060&format=complete) [MR 1297766](http://www.ams.org/mathscinet-getitem?mr=1297766)
- <span id="page-45-6"></span>[37] Wang, Z., Li, W., Ruan, S.: Existence and stability of traveling wave fronts in reaction advection diffusion equations with nonlocal delay. J. Differential Equations 238, 153–200 (2007) [Zbl 05171430](http://www.zentralblatt-math.org/zmath/en/advanced/?q=an:05171430&format=complete) [MR 2334595](http://www.ams.org/mathscinet-getitem?mr=2334595)
- <span id="page-45-11"></span>[38] Weinberger, H. F.: Long-time behavior of a class of biological models. SIAM J. Math. Anal. 13, 353–396 (1982) [Zbl 0529.92010](http://www.zentralblatt-math.org/zmath/en/advanced/?q=an:0529.92010&format=complete) [MR 0653463](http://www.ams.org/mathscinet-getitem?mr=0653463)
- <span id="page-45-12"></span>[39] Weinberger, H. F.: On spreading speeds and traveling waves for growth and migration models in a periodic habitat. J. Math. Biol. 45, 511–548 (2002) [Zbl 1058.92036](http://www.zentralblatt-math.org/zmath/en/advanced/?q=an:1058.92036&format=complete) [MR 1943224](http://www.ams.org/mathscinet-getitem?mr=1943224)
- <span id="page-45-17"></span>[40] Weinberger, H. F., Lewis, M. A., Li, B.: Analysis of linear determinacy for spread in cooperative models. J. Math. Biol. 45, 183–218 (2002) [Zbl 1023.92040](http://www.zentralblatt-math.org/zmath/en/advanced/?q=an:1023.92040&format=complete) [MR 1930974](http://www.ams.org/mathscinet-getitem?mr=1930974)
- <span id="page-45-16"></span>[41] Wu, J.: Theory and Applications of Partial Functional Differential Equations. Springer, New York (1996) [Zbl 0870.35116](http://www.zentralblatt-math.org/zmath/en/advanced/?q=an:0870.35116&format=complete) [MR 1415838](http://www.ams.org/mathscinet-getitem?mr=1415838)
- <span id="page-45-9"></span>[42] Xin, J.: Front propagation in heterogeneous media. SIAM Rev. 42, 161–230 (2000) [Zbl 0951.35060](http://www.zentralblatt-math.org/zmath/en/advanced/?q=an:0951.35060&format=complete) [MR 1778352](http://www.ams.org/mathscinet-getitem?mr=1778352)
- <span id="page-45-7"></span>[43] Xin, X.: Existence and stability of traveling waves in periodic media governed by a bistable nonlinearity. J. Dynam. Differential Equations 3, 541–573 (1991) [Zbl 0769.35033](http://www.zentralblatt-math.org/zmath/en/advanced/?q=an:0769.35033&format=complete) [MR 1129560](http://www.ams.org/mathscinet-getitem?mr=1129560)
- <span id="page-45-8"></span>[44] Xin, X.: Existence and nonexistence of traveling waves and reaction-diffusion front propagation in periodic media. J. Statist. Phys. 73, 893–926 (1993) [Zbl 1102.35340](http://www.zentralblatt-math.org/zmath/en/advanced/?q=an:1102.35340&format=complete) [MR 1251222](http://www.ams.org/mathscinet-getitem?mr=1251222)
- <span id="page-45-18"></span>[45] Xu, D., Zhao, X.-Q.: Bistable waves in an epidemic model. J. Dynam. Differential Equations 16, 679–707 (2004) [Zbl 1073.34027](http://www.zentralblatt-math.org/zmath/en/advanced/?q=an:1073.34027&format=complete) [MR 2157846](http://www.ams.org/mathscinet-getitem?mr=2157846)
- <span id="page-45-1"></span>[46] Yagisita, H.: Existence of traveling wave solutions for a nonlocal bistable equation: an abstract approach. Publ. RIMS Kyoto Univ. 45, 955–979 (2009) [Zbl 1191.35094](http://www.zentralblatt-math.org/zmath/en/advanced/?q=an:1191.35094&format=complete) [MR 2597125](http://www.ams.org/mathscinet-getitem?mr=2597125)
- <span id="page-45-15"></span>[47] Zhao, X.-Q.: Dynamical Systems in Population Biology. Springer, New York (2003) [Zbl 1023.37047](http://www.zentralblatt-math.org/zmath/en/advanced/?q=an:1023.37047&format=complete) [MR 1980821](http://www.ams.org/mathscinet-getitem?mr=1980821)
- <span id="page-45-13"></span>[48] Zhao, X.-Q., Spatial dynamics of some evolution systems in biology. In: Recent Progress on Reaction-Diffusion Systems and Viscosity Solutions, Y. Du et al. (eds.), World Sci., 332–363 (2009) [Zbl 1182.35041](http://www.zentralblatt-math.org/zmath/en/advanced/?q=an:1182.35041&format=complete) [MR 2532932](http://www.ams.org/mathscinet-getitem?mr=2532932)
- <span id="page-45-2"></span>[49] Zinner, B.: Stability of traveling wavefronts for the discrete Nagumo equation. SIAM J. Math. Anal. 22, 1016–1020 (1991) [Zbl 0739.34060](http://www.zentralblatt-math.org/zmath/en/advanced/?q=an:0739.34060&format=complete) [MR 1112063](http://www.ams.org/mathscinet-getitem?mr=1112063)
- <span id="page-45-3"></span>[50] Zinner, B.: Existence of traveling wavefront solutions for the discrete Nagumo equation. J. Differential Equations 96, 1–27 (1992) [Zbl 0752.34007](http://www.zentralblatt-math.org/zmath/en/advanced/?q=an:0752.34007&format=complete) [MR 1153307](http://www.ams.org/mathscinet-getitem?mr=1153307)