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Boundary estimates for certain degenerate and singular parabolic equations

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Abstract. We study the boundary behavior of non-negative solutions to a class of degenerate/singular parabolic equations, whose prototype is the parabolic p -Laplace equation. Assuming that such solutions continuously vanish on some distinguished part of the lateral part S_T of a Lipschitz cylinder, we prove Carleson-type estimates, and deduce some consequences under additional assumptions on the equation or the domain. We then prove analogous estimates for non-negative solutions to a class of degenerate/singular parabolic equations of porous medium type.

Keywords. Degenerate and singular parabolic equations, Harnack estimates, boundary Harnack inequality, Carleson estimate

1. Introduction

In this paper we start the study of the boundary behavior of weak solutions to a class of degenerate/singular equations whose prototype is the parabolic p -Laplace equation

$$u_t - \operatorname{div}(|Du|^{p-2}Du) = 0, \quad (1.1)_0$$

where Dw denotes the gradient of w with respect to the space variables. More precisely, let E be an open set in \mathbb{R}^N and for $T > 0$ let E_T denote the cylindrical domain $E \times (0, T]$. Moreover let

$$S_T = \partial E \times [0, T], \quad \partial_P E_T = S_T \cup (E \times \{0\})$$

denote the lateral and the parabolic boundary respectively.

We shall consider quasi-linear parabolic partial differential equations of the form

$$u_t - \operatorname{div} \mathbf{A}(x, t, u, Du) = 0 \quad \text{weakly in } E_T \quad (1.1)$$

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where the function $\mathbf{A} : E_T \times \mathbb{R}^{N+1} \rightarrow \mathbb{R}^N$ is only assumed to be measurable and subject to the structure conditions

$$\begin{cases} \mathbf{A}(x, t, u, \xi) \cdot \xi \geq C_0|\xi|^p \\ |\mathbf{A}(x, t, u, \xi)| \leq C_1|\xi|^{p-1} \end{cases} \quad \text{a.e. } (x, t) \in E_T, \forall u \in \mathbb{R}, \forall \xi \in \mathbb{R}^N \quad (1.2)$$

where C_0 and C_1 are given positive constants, and $p > 1$. We refer to the parameters p, N, C_0, C_1 as our structural data, and we write $\gamma = \gamma(N, p, C_0, C_1)$ if γ can be quantitatively determined a priori only in terms of the above quantities. A function

$$u \in C([0, T]; L^2(E)) \cap L^p(0, T; W^{1,p}(E)) \quad (1.3)$$

is a *weak sub[super]solution* to (1.1)–(1.2) if for every subinterval $[t_1, t_2] \subset (0, T]$,

$$\int_E u \varphi dx \Big|_{t_1}^{t_2} + \int_{t_1}^{t_2} \int_E [-u\varphi_t + \mathbf{A}(x, t, u, Du) \cdot D\varphi] dx dt \leq [\geq] 0 \quad (1.4)$$

for all non-negative test functions

$$\varphi \in W^{1,2}(0, T; L^2(E)) \cap L^p(0, T; W_0^{1,p}(E)).$$

This guarantees that all the integrals in (1.4) are convergent.

Under the conditions (1.2), equation (1.1) is degenerate when $p > 2$ and singular when $1 < p < 2$, since the modulus of ellipticity $|Du|^{p-2}$ respectively tends to 0 or to $+\infty$ as $|Du| \rightarrow 0$. In the latter case, we further distinguish between singular *supercritical* range (when $2N/(N + 1) < p < 2$), and singular *critical and subcritical* range (when $1 < p \leq 2N/(N + 1)$). When $p = 2$, the equation is uniformly parabolic, and the theory is fairly complete.

For points in \mathbb{R}^N we use the notation (x', x_N) or $(x_1, \dots, x_{N-1}, x_N)$; $D'w$ stands for the gradient of w with respect to x' .

For $y \in \mathbb{R}^N$ and $\rho > 0$, $K_\rho(y)$ denotes the cube of edge 2ρ , centered at y with faces parallel to the coordinate planes. When y is the origin of \mathbb{R}^N we simply write K_ρ ; $K'_\rho(y')$ denotes the $(N - 1)$ -dimensional cube $\{x' : |x_i - y_i| < \rho, i = 1, \dots, N - 1\}$; we write for short $\{|x_i - y_i| < \rho\}$.

For $\theta > 0$ we also define

$$Q_\rho^-(\theta) = K_\rho \times (-\theta\rho^p, 0], \quad Q_\rho^+(\theta) = K_\rho \times (0, \theta\rho^p],$$

and for $(y, s) \in \mathbb{R}^N \times \mathbb{R}$,

$$(y, s) + Q_\rho^-(\theta) = K_\rho(y) \times (s - \theta\rho^p, s], \quad (y, s) + Q_\rho^+(\theta) = K_\rho(y) \times (s, s + \theta\rho^p].$$

Now fix $(x_0, t_0) \in E_T$ such that $u(x_0, t_0) > 0$ and construct the cylinders

$$(x_0, t_0) + Q_\rho^\pm(\theta) \quad \text{where} \quad \theta = \left(\frac{c}{u(x_0, t_0)} \right)^{p-2}, \quad (1.5)$$

and c is a given positive constant. These cylinders are “intrinsic” to the solution, since their height is determined by the value of u at (x_0, t_0) . Cylindrical domains of the form $K_\rho \times (0, \rho^p]$ reflect the natural, parabolic space-time dilations that leave the homogeneous, prototype equation (1.1)₀ invariant. The latter however is not homogeneous with respect to the solution u . The time dilation by a factor $u(x_0, t_0)^{2-p}$ is intended to restore the homogeneity. Most of the results we describe in this paper hold in such geometry.

Our reference domains are Lipschitz (respectively $C^{1,1}$, C^2) domains. We recall that a bounded domain $E \subset \mathbb{R}^N$ is said to be a *Lipschitz domain* if for each $y \in \partial E$ there exists a radius r_0 such that in an appropriate coordinate system,

$$\begin{aligned} E \cap K_{8r_0}(y) &= \{x = (x', x_N) \in \mathbb{R}^N : x_N > \Phi(x')\} \cap K_{8r_0}(y), \\ \partial E \cap K_{8r_0}(y) &= \{x = (x', x_N) \in \mathbb{R}^N : x_N = \Phi(x')\} \cap K_{8r_0}(y), \end{aligned}$$

where Φ is a Lipschitz function with $\|D'\Phi\|_{L^\infty} \leq L$. The quantities r_0 and L are independent of $y \in \partial E$. We say that L is the *Lipschitz constant* of E .

We define $C^{1,1}$ and C^2 domains analogously, requiring Φ to be of class $C^{1,1}$ or C^2 respectively. Correspondingly we set $M_{1,1} = \|\Phi\|_{1,1}$ and $M_2 = \|\Phi\|_2$, and we assume they are independent of $y \in \partial E$.

Finally, a *Lipschitz* (respectively $C^{1,1}$, C^2) *cylinder* is a cylindrical domain E_T whose cross section E is a Lipschitz (respectively $C^{1,1}$, C^2) domain.

We are interested in solutions u to (1.1)–(1.2) continuously *vanishing* on some distinguished part of the lateral part S_T of a cylinder. Our main goal is to show that near the boundary, u is controlled above in a non-tangential fashion. More precisely, this means that an inequality of the type

$$u \leq \gamma u(P_\rho) \tag{1.6}$$

holds in a box Ψ_ρ of size ρ , based on S_T , where P_ρ is a non-tangential point and γ depends only on the structural data. The first results of this kind are due to Carleson [11], for the Laplace equation in Lipschitz domains, and to Kemper [34] for the heat equation in domains, which are locally given by a function satisfying a mixed Lipschitz condition, with exponent 1 in the space variables and 1/2 in the time variable (also called parabolic Lipschitz domains). Since then, an inequality like (1.6) is known as a *Carleson estimate*.

There is another inequality naturally associated to (1.6), namely

$$u/v \approx u(P_r)/v(P_r). \tag{1.7}$$

Inequality (1.7) is known as the *boundary comparison principle* or the *boundary Harnack inequality*. For linear equations, it implies the Hölder continuity up to the boundary of the quotient u/v and that the vanishing speeds of u and v are the same.

Both (1.6) and (1.7) have been generalized to more general contexts and operators and they have become essential tools in analyzing the boundary behavior of non-negative solutions.

In the elliptic context we mention [33] for the Laplace operator in non-tangentially accessible domains, [8] and [3], [4], [23] for elliptic operators in divergence and non-divergence form, respectively, [40], [41] for the p -Laplace operator, and [12], [13] for the Kolmogorov operator.

We emphasize that for the Laplacian a Carleson estimate has been proved in [1] to be equivalent to the boundary Harnack principle. It would be quite interesting to explore this connection between the two inequalities also in the non-linear setting.

For parabolic operators, we cite [48], [25], [31], [26] for cylindrical domains, and [24] for parabolic Lipschitz domains.

A classical application of the two inequalities is to Fatou-type theorems, but even more remarkable is their role in the regularity theory of two-phase free boundary problems, as shown in the two seminal papers [6], [7], where a general strategy to attack the regularity of the free boundary governed by the Laplace operator has been set up.

This technique has been subsequently extended to stationary problems governed by variable coefficients linear and semilinear operators [10], [29], to fully nonlinear operators [27], [28], and to the p -Laplace operator [42], [43].

The free boundary regularity theory for two-phase parabolic problems is less developed. For Stefan type problems we mention [9], [30], [14] and the references therein. In particular [14] deals with a one-phase Stefan problem for the p -Laplacian when $p > 2$.

The present paper places itself exactly along this line of research. The Carleson estimate for our singular/degenerate equations is the first piece of information, very useful to analyze the regularity of (e.g.) Lipschitz or flat free boundaries. Thanks to recent developments in the field of Harnack inequalities for quasi-linear parabolic equations of p -Laplace type [17, 18, 20, 35], in Theorems 2.1 and 3.1 we extend estimate (1.6) to non-negative solutions to (1.1)–(1.2) in cylindrical Lipschitz domains. According to the theory developed in the above papers, a Carleson type estimate makes sense only for $p > 2N/(N + 1)$.

Indeed, in the critical and subcritical range, explicit counterexamples rule out the possibility of a Harnack inequality. Only so-called Harnack-type estimates are possible, where, however, the ratio of infimum to supremum in proper space-time cylinders depends on the solution itself (for more details, see [21, Chapter 6, §§11–15]).

The approach developed for linear elliptic equations in [8] and essentially at the same time adapted to linear parabolic equations in [48], to prove the Carleson estimate, is centered around two basic estimates for solutions: the Harnack inequality and the geometric decay of the oscillation of u up to the boundary. To see this, let us consider a non-negative solution in a cylinder, and assume further that the solution vanishes on a part of the lateral boundary, which we assume to be a part of the hyperplane $\{x_N = 0\}$. Fix a point \bar{P} at unit distance from the lateral boundary; by translation we may assume that $\bar{P} = (e_N, 0)$. To simplify even more, due to the homogeneity of the equation, without loss of generality we may assume that $u(\bar{P}) = 1$. A repeated application of the Harnack inequality in a dyadic fashion gives

$$u(P) \leq H^k u(\bar{P}) \quad \text{whenever} \quad \text{dist}(P, \partial E_T) \geq 2^{-k}, \quad (1.8)$$

for P in a boundary space-time box $\Psi^+ = E_T \cap K_1(0) \times (-2, -1)$. To continue, one defines a sequence of boundary space-time boxes Ψ_k^+ such that $\Psi_k^+ \subset \Psi_{k-1}^+ \subset \dots \subset \Psi^+$ for all $k \in \mathbb{N}$.

Now, suppose that (1.6) does not hold true, i.e. there is no constant $\gamma > 0$, which depends only on the data, such that $u(P) \leq \gamma$ for all $P \in \Psi_1^+$. Consequently, there

must exist $P_1 \in \Psi_1^+$ such that $u(P_1) > H^h$, where $h \in \mathbb{N}$. Inequality (1.8) implies that $\text{dist}(P_1, \partial E_T) < 2^{-h}$. If h is chosen large enough, by the geometric decay of the oscillation of u up to the boundary, one deduces the existence of $P_2 \in \Psi_2^+$ such that $u(P_2) > H^{h+1}$ and $\text{dist}(P_2, \partial E_T) < 2^{-(h+1)}$. Repeating this yields a sequence $\{P_j\}_{j=1}^\infty$ of points approaching the boundary, whereas the sequence $\{u(P_j)\}_{j=1}^\infty$ blows up: this contradicts the assumption that u vanishes continuously on the boundary, and we conclude

$$\sup_{\Psi_1^+} u \leq H^h,$$

which is just (1.6) in our setting.

Although the overall strategy in the non-linear setting follows the same kind of argument, its implementation presents a difficulty due to the lack of homogeneity of the equations, and there is also a striking difference between the singular and the degenerate case; this is already reflected in the intrinsic character of the interior Harnack inequality, and it is amplified when approaching the boundary through dyadically shrinking intrinsic cylinders. Concerning the Carleson estimate, its statement in the degenerate case can be considered as the intrinsic version of the analogous statement in the linear uniformly parabolic case. Things are different in the singular supercritical case, where, in general, one can only prove a somewhat weaker estimate (see Theorem 3.1), due to the possibility for a solution to extinguish in finite time. Indeed, the counterexamples in §3.2 show that one cannot do any better, unless some control of the interior oscillation of the solution is available (Corollary 3.1).

The difference between the two cases, degenerate and singular supercritical, becomes more evident when one considers the validity of a boundary Harnack principle, even in smooth cylinders. Solutions to the parabolic p -Laplace equations can vanish arbitrarily fast in the degenerate case $p > 2$ (see §2.1), so that no possibility exists to prove the boundary Harnack principle in its generality. On the other hand, in the singular case, the existence of suitable barriers provides linear behavior. Together with Carleson's estimate, this fact implies almost immediately a Hopf principle and the boundary Harnack inequality.

The last section of this note is devoted to extending all the previous results to non-negative solutions to a large class of degenerate/singular parabolic equations, whose prototype is the porous medium equation (see §4 for all the details).

In a forthcoming paper we plan to extend the boundary Harnack principle to Lipschitz cylinders.

Remark. After completing the paper, we learnt that Kuusi, Mingione and Nyström [39] independently proved a boundary Harnack inequality, which is similar to the one we give here in Theorem 3.8.

2. The degenerate case $p > 2$

2.1. Main results

To describe our main results we need to introduce some further notation. Let E_T be a Lipschitz cylinder and fix $(x_0, t_0) \in S_T$; in a neighborhood of such a point, the cross sec-

tion is represented by the graph $\{(x', x_N) : x_N = \Phi(x')\}$, where Φ is a Lipschitz function and $\|D'\Phi\|_\infty \leq L$. Without loss of generality, from here on we assume $\Phi(x'_0) = 0$ and $L \geq 1$.

For $\rho \in (0, r_0)$, let $x_\rho = (x'_0, 2L\rho)$, and let $P_\rho = P_\rho(x_0, t_0) = (x'_0, 2L\rho, t_0) \in E_T$ be such that $u(P_\rho) > 0$. Note that $\text{dist}(x_\rho, \partial E)$ is of order ρ . Set

$$\Psi_\rho^-(x_0, t_0) = E_T \cap \left\{ (x, t) : |x_i - x_{0,i}| < \frac{\rho}{4}, |x_N| < 2L\rho, t \in \left(t_0 - \frac{\alpha + \beta}{2}\theta\rho^p, t_0 - \beta\theta\rho^p \right) \right\}$$

where $\theta = [c/u(P_\rho)]^{p-2}$ with c given in Theorem 2.6 below, and $\alpha > \beta$ are positive parameters. We are now ready to state our main result in the degenerate case $p > 2$.

Theorem 2.1 (Carleson estimate, $p > 2$). *Let u be a non-negative weak solution to (1.1)–(1.2) in E_T . Assume that*

$$(t_0 - \theta(4\rho)^p, t_0 + \theta(4\rho)^p] \subset (0, T]$$

and that u vanishes continuously on

$$\partial E \cap \{|x_i - x_{0,i}| < 2\rho, |x_N| < 8L\rho\} \times (t_0 - \theta(4\rho)^p, t_0 + \theta(4\rho)^p).$$

Then there exist positive parameters $\alpha > \beta$ and a constant $\tilde{\gamma} > 0$, depending only on p, N, C_0, C_1, L , such that

$$u(x, t) \leq \tilde{\gamma}u(P_\rho) \quad \text{for every } (x, t) \in \Psi_\rho^-(x_0, t_0). \tag{2.1}$$

Remark 2.2. Without going too much into details here, let us point out that for the prototype equation (1.1)₀, estimate (2.1) could be extended from Lipschitz cylinders to a wider class of cylinders E_T , whose cross section E is a so-called *N.T.A. domain* (non-tangentially accessible domain). For more details, we refer the reader to [9, §12.3].

Weak solutions to (1.1) with zero Dirichlet boundary conditions on a Lipschitz domain are Hölder continuous up to the boundary (see, for example, [16, Chapter III, Theorem 1.2]). Combining this result with the previous Carleson estimate yields a quantitative estimate on the decay of u at the boundary, invariant by the intrinsic rescaling

$$x = x_0 + \rho y, \quad t = t_0 + \frac{\rho^p}{u(P_\rho)^{p-2}} \tau.$$

Corollary 2.3. *Under the same assumption of Theorem 2.1, we have*

$$0 \leq u(x, t) \leq \gamma \left(\frac{\text{dist}(x, \partial E)}{\rho} \right)^\mu u(P_\rho)$$

for every $(x, t) \in \Psi_{\rho/2}^-(x_0, t_0)$, where $\mu \in (0, 1)$ depends only on p, N, C_0, C_1, L .

If we restrict our attention to solutions to the model equation (1.1)₀ and to C^2 cylinders, the result of Corollary 2.3 can be strengthened.

Theorem 2.4 (Lipschitz decay). *Let E_T be a C^2 cylinder and u a non-negative weak solution to (1.1)₀ in E_T . Assume the other assumptions of Theorem 2.1 hold. Then there exist positive parameters $\alpha > \beta$ and a constant $\gamma > 0$, depending only on p , N and the C^2 -constant M_2 of E , such that*

$$0 \leq u(x, t) \leq \gamma \frac{\text{dist}(x, \partial E)}{\rho} u(P_\rho) \quad (2.2)$$

for every (x, t) in

$$E_T \cap \left\{ |x_i - x_{0,i}| < \frac{\rho}{4}, 0 < x_N < 2M_2\rho \right\} \times \left(t_0 - \frac{\alpha + 3\beta}{4} \theta \rho^p, t_0 - \beta \theta \rho^p \right).$$

Remark 2.5. Estimate (2.2) is not surprising, as it is well-known that under the assumptions of Theorem 2.4, solutions are Lipschitz continuous up to the boundary (see [16]). As a matter of fact, in [44, 45] Lieberman has proved $C^{1+\alpha}$ regularity up to the boundary for solutions of a proper p -Laplace type equation, with conormal and Dirichlet boundary conditions. Relying on the recent papers [5, 36, 37, 38], these results can be extended both to a wider class of degenerate equations with *differentiable* principal part, which have the same structure of the p -Laplacian, and to *less regular* $C^{1,\alpha}$ domains. We limited ourselves to (1.1)₀ and C^2 domains, mainly to avoid the introduction of further structural assumptions and technical details.

Notice that, in general, the bound below by zero in (2.2) cannot be improved. Indeed, when $p > 2$, two explicit solutions to the parabolic p -Laplacian in the half-space $\{x_N \geq 0\}$ that vanish at $x_N = 0$ are given by

$$u_1(x, t) = x_N, \quad u_2(x, t) = \frac{p-2}{p^{(p-1)/(p-2)}} (T-t)^{-1/(p-2)} x_N^{p/(p-2)}. \quad (2.3)$$

The power-like behavior, as exhibited in u_2 , is not the “worst” possible case. Indeed, let $E = \{-1 \leq x_i \leq 1, 0 \leq x_N \leq 1/4\}$, and consider the following Cauchy–Dirichlet problem in $E \times [0, T)$:

$$\begin{cases} u_t - \text{div}(|Du|^{p-2} Du) = 0, \\ u(x, 0) = C T^{-1/(p-2)} \exp(-1/x_N), \\ u(x', 0, t) = 0, \\ u(x', 1/4, t) = C (T-t)^{-1/(p-2)} e^{-4}, \\ u(x, t) = C (T-t)^{-1/(p-2)} \exp(-1/x_N), \quad x \in \partial E \cap \{0 < x_N < 1/4\}, \end{cases} \quad (2.4)$$

where

$$C = \frac{1}{2(p-1)(p-2)} \left(\frac{e(p-2)}{2p} \right)^{\frac{2p}{p-2}}.$$

It is easy to check that the function

$$u_3 = C(T - t)^{-1/(p-2)} \exp(-1/x_N), \quad x_N > 0, \tag{2.5}$$

is a supersolution to the above problem. Therefore, the *solution* to the same problem (which is obviously positive) lies below u_3 and approaches the zero boundary value at $x_N = 0$ at least with exponential speed.

Example (2.4)–(2.5) can be further generalized. Let $\gamma \in (0, 1)$, $E = \{x_N > 0\}$, $T = 2/\gamma - 1$. Then

$$u(x, t) = \left[\frac{p-2}{p-1} \gamma^{1/(p-1)} (t+1) \left(\gamma + \frac{x_N - 2}{t+1} \right)_+ \right]^{\frac{p-1}{p-2}} \tag{2.6}$$

is a solution to (1.1)₀ in E_T , and it vanishes not only on the boundary $\{x_N = 0\}$, but also in the set $\{0 < x_N < 2 - \gamma(t+1), 0 < t < T\}$, which has positive measure.

2.2. The interior Harnack inequality

As mentioned in the Introduction, our results are strongly based on the interior Harnack inequalities proved in [17, 18, 19, 35], which we recall here.

Theorem 2.6. *Let u be a non-negative weak solution to (1.1)–(1.2) in E_T for $p > 2$, and let $(x_0, t_0) \in E_T$ be such that $u(x_0, t_0) > 0$. There exist positive constants c and γ , depending only on p, N, C_0, C_1 , such that for all intrinsic cylinders $(x_0, t_0) + Q_{2\rho}^\pm(\theta)$ as in (1.5), contained in E_T ,*

$$\gamma^{-1} \sup_{K_\rho(x_0)} u(\cdot, t_0 - \theta\rho^p) \leq u(x_0, t_0) \leq \gamma \inf_{K_\rho(x_0)} u(\cdot, t_0 + \theta\rho^p). \tag{2.7}$$

The constants γ and c deteriorate as $p \rightarrow \infty$ in the sense that $\gamma(p), c(p) \rightarrow \infty$ as $p \rightarrow \infty$; however, they are stable as $p \rightarrow 2$.

Remark 2.7. In all the previously mentioned works, the requirement on the cylinder is that $(x_0, t_0) + Q_{4\rho}^\pm(\theta) \subset E_T$; by a proper adjustment of the parameters c and γ we can work under the more restrictive condition we are now assuming.

As already pointed out in the Introduction, in [48], the Carleson estimate is a consequence of the Harnack inequality of [47] and a geometric argument, based on the control of the oscillation. In particular, a key tool is the so-called Harnack chain, namely the control on the value of $u(x, t)$ by the value of $u(x_0, t_0)$ with $t < t_0$, thanks to the repeated application of the Harnack inequality.

In [18], an equivalent statement for solutions to (1.1)–(1.2) is given, but a careful examination of the proof shows that the result there actually holds only for solutions defined in $\mathbb{R}^N \times (0, T)$, and not in a smaller domain E_T . Although the correct form of the Harnack chain for solutions defined in E_T , when $E \subsetneq \mathbb{R}^N$, can be given, such a result is of no use in the proof of Carleson’s estimates, as there are two different, but equally important obstructions.

First of all, as (2.6) shows too, u can vanish and hence prevent any further application of the Harnack inequality. Indeed, let us consider the following two examples.

Let $\gamma \in (0, 1)$; the function

$$u(x, t) = \left[\frac{p-2}{p-1} \gamma^{1/(p-1)} (t+1) \left(\gamma + \frac{x_N - 2}{t+1} \right)_+ \right]^{\frac{p-1}{p-2}} + \left[\frac{p-2}{p-1} \gamma^{1/(p-1)} (t+1) \left(\gamma - \frac{x_N + 2}{t+1} \right)_+ \right]^{\frac{p-1}{p-2}}$$

is a solution to the parabolic p -Laplace equation in $\mathbb{R}^N \times (0, 2/\gamma - 1)$ and vanishes in the cone

$$\begin{cases} 0 < t < 2/\gamma - 1, \\ -(2 - \gamma(t+1)) < x_N < 2 - \gamma(t+1). \end{cases}$$

If we take (x, t) and (x_0, t_0) with $t < t_0$ on opposite sides of the cone, there is no way to build a Harnack chain that connects the two points.

Let $\gamma_p = (1/\lambda)^{1/(p-1)}(p-2)/p$, with $\lambda = N(p-2) + p$, consider the cylinder $\{x_N > 0\} \times (0, (2\gamma_p)^\lambda)$ and let $x_1 = (0, \dots, 0, 2)$, $x_2 = (0, \dots, 0, 6)$. The function

$$u(x, t) = t^{-N/\lambda} \left[1 - \gamma_p \left(\frac{|x - x_1|}{t^{1/\lambda}} \right)^{\frac{p}{p-1}} \right]^{\frac{p-1}{p-2}}_+ + t^{-N/\lambda} \left[1 - \gamma_p \left(\frac{|x - x_2|}{t^{1/\lambda}} \right)^{\frac{p}{p-1}} \right]^{\frac{p-1}{p-2}}_+$$

is a solution to the parabolic p -Laplace equation in the indicated cylinder and vanishes on its parabolic boundary. Notice that such a solution is the sum of two Barenblatt functions with poles respectively at x_1 and x_2 and masses $M_1 = M_2 = 1$: in the interval $0 < t < (2\gamma_p)^\lambda$ the support of u is given by two disjoint regions R_1 and R_2 , and only at time $T = (2\gamma_p)^\lambda$ does the support of u finally become a simply connected set. Once more, taking (x, t) and (x_0, t_0) respectively in R_1 and R_2 , there is no way to connect them with a Harnack chain. As a matter of fact, before the two supports touch, each Barenblatt function does not *feel* in any way the presence of the other one. In particular, we can change the mass of the two Barenblatt functions; this will modify the time T at which the two supports touch, but up to T , there is no way one Barenblatt component can detect the change performed on the other one.

On the other hand, one could think that if we have a solution vanishing on a flat piece of the boundary and strictly positive everywhere in the interior, then one could build a Harnack chain extending arbitrarily close to the boundary. However, this is not the case, as clearly shown by the following example.

Let us consider a domain $E \subset \mathbb{R}^N$ with a part of its boundary coinciding with the hyperplane $\{x_N = 0\}$, and let $\Gamma = \partial E \cap \{x_N = 0\}$. Let $\bar{T} > 0$ be given and consider a non-negative solution u to

$$\begin{cases} u_t - \operatorname{div}(|Du|^{p-2} Du) = 0 & \text{in } E_{\bar{T}}, \\ u > 0 & \text{in } E_{\bar{T}}, \\ u = 0 & \text{on } \Gamma \times (0, \bar{T}]. \end{cases}$$

Let u be such that its value is bounded above by the distance to the flat boundary piece raised to some given power $a > 0$, i.e.

$$u(x, t) \leq \gamma \operatorname{dist}(x, \Gamma)^a, \quad a > 0, (x, t) \in E_{\bar{T}}, \tag{2.8}$$

where $\gamma > 0$ is a suitable parameter. Indeed, this is the case of solutions given in (2.3) and (2.5) with $\bar{T} = T/2$ for example, and therefore such a situation does take place.

Let $(x_0, t_0) = (x'_0, x_{0,N}, t_0) \in E_T$ be such that $\operatorname{dist}(x_0, \Gamma) = 1$. The goal is to form a Harnack chain of dyadic non-tangential cylinders approaching the boundary, while the chain stays inside $E_{\bar{T}}$; we want to control the size of the time interval we need to span in order to complete the chain. Let

$$\begin{aligned} u_0 &= u(x_0, t_0), \\ r_k &= 2^{-k}, \\ x_k &= (\hat{x}'_0, 2^{-k}), \\ t_k &= t_0 - c^{p-2} \sum_{i=0}^{k-1} u_i^{2-p} r_i^p, \\ u_k &= u(x_k, t_k) \approx (2^{-k})^a \end{aligned}$$

for $k = 1, 2, \dots$. Assuming that at each step one can use Harnack’s inequality, we get an estimate on the size of t_k from above,

$$t_k \leq t_0 - c^{p-2} \sum_{i=0}^{k-1} (2^{-ai})^{2-p} 2^{-ip} \leq t_0 - c^{p-2} \sum_{i=0}^{k-1} 2^{ai(p-2)-ip},$$

which diverges to $-\infty$ as $k \rightarrow \infty$ and $x_k \rightarrow \Gamma$ if $a \geq p/(p - 2)$. Considering the solution u_2 from (2.3), we see that the above dyadic Harnack chain would diverge for such a solution as $a = p/(p - 2)$.

Notice that this is a counterexample to the use of the Harnack chain in the *proof* of the Carleson estimate, but not a counterexample to the Carleson estimate itself.

Remark 2.8. The infinite length of the time interval needed to reach the boundary is just one face (i.e. consequence) of the finite speed of propagation when $p > 2$. Points (x, t) that lie *inside* a proper p -paraboloid centered at (x_0, t_0) can be reached, starting from (x_0, t_0) ; if u_0 is very small, and therefore the p -paraboloid is very narrow, with small values of r one ends up with very large values of t . On the other hand, points (x, t) that lie *outside* the same p -paraboloid centered at (x_0, t_0) cannot be reached.

Remark 2.9. We conjecture that $p/(p - 2)$ is a sort of threshold exponent: when $a < p/(p - 2)$ in (2.8), the regularizing effects of the diffusion kick in, and eventually the solution becomes linear, allowing for more precise bounds from below in (2.2); when a is larger, the time evolution part wins, and we have much less regularity. We probably need further information on the behavior of u on the lower base of the cylinder in order to make all the previous heuristics more rigorous. This will be the object of future investigation.

Eventually, as stated in Theorem 2.1, the Carleson estimate does hold true. However, as we could not make use of any form of Harnack chain in its proof, we resort to a contradiction argument, detailed in §2.4.2.

2.3. Hölder continuity and oscillation control

In this subsection we consider $p > 1$, since the statements are the same in both cases $p > 2$ and $1 < p < 2$. It is well known that locally bounded weak solutions to (1.1)–(1.2) are locally Hölder continuous. For the full statement and proof, see [16, Chapters III and IV], or [32] (when $p > 2$) for a simpler approach, based on the same ideas used in [17]. The local Hölder continuity is a consequence of the following lemma. Note the slight difference between the cases $p > 2$ and $1 < p < 2$.

In the following let u be a weak solution to (1.1)–(1.2) in E_T for $1 < p < \infty$. Fix a point in E_T , which, up to translation, we may take to be the origin of \mathbb{R}^{N+1} . For $\rho > 0$ consider the cylinder

$$Q_* = \begin{cases} K_\rho \times (-\rho^2, 0] & \text{if } p > 2, \\ K_\rho \times (-\rho^p, 0] & \text{if } 1 < p < 2, \end{cases}$$

with vertex at $(0, 0)$, and set

$$\mu_0^+ = \sup_{Q_*} u, \quad \mu_0^- = \inf_{Q_*} u, \quad \omega_0 = \operatorname{osc}_{Q_*} u = \mu_0^+ - \mu_0^-.$$

Depending on ω_0 , construct the cylinder

$$Q_0 = K_\rho \times (-\theta_0 \rho^p, 0], \quad \text{where } \theta_0 = (c/\omega_0)^{p-2}.$$

If $p > 2$, c being the constant in (1.5), we assume that

$$\omega_0 > c\rho. \tag{2.9}$$

If $1 < p < 2$, we assume that $\omega_0 \leq 1$. The previous two conditions ensure that $Q_0 \subset Q_*$, and the following lemma holds.

Lemma 2.10 (Hölder continuity, [17]). *There exist constants $\epsilon, \delta \in (0, 1)$ and $c \geq 1$, depending only on p, N, C_0, C_1 , such that, setting*

$$\omega_n = \delta\omega_{n-1}, \quad \theta_n = (c/\omega_n)^{p-2}, \quad \rho_n = \epsilon\rho_{n-1} \quad \text{and} \quad Q_n = Q_{\rho_n}^-(\theta_n),$$

for all non-negative integers n we have $Q_{n+1} \subset Q_n$ and

$$\operatorname{osc}_{Q_n} u \leq \omega_n. \tag{2.10}$$

Remark 2.11. Given u and $p > 2$, we can directly assume (2.9), since otherwise there is nothing to prove; moreover, in this case, we have $c = 1$. If $1 < p < 2$, then c is a suitable constant greater than 1. Also, note that (2.10) yields

$$\operatorname{osc}_{Q_n} u \leq \delta^n \operatorname{osc}_{Q_*} u.$$

Therefore, the lemma builds a sequence of intrinsic cylinders, where the oscillation is controlled by a proper power of an absolute constant; the starting cylinder is the only one that does not have an intrinsic size.

In the proof of Theorems 2.1–2.4 we need the following two lemmas. The former is the well-known reflection principle, whose proof is standard. The interested reader can refer to [46, Lemma 2.7] or [5, Lemma 2.8]. The latter is an alternative De Giorgi-type lemma with initial data, taken from [21], to which we refer for the proof.

Lemma 2.12. *Let*

$$\begin{aligned} Q_p &= \{x : |x_i| < 1, 0 < x_N < 2, t_1 < t < t_2\}, \\ Q_n &= \{x : |x_i| < 1, -2 < x_N < 0, t_1 < t < t_2\}, \end{aligned}$$

and u be a non-negative weak solution to (1.1)–(1.2) in Q_p with $u = 0$ on $\partial Q_p \cap \{x : x_N = 0\}$. Let

$$Q' = \{x : |x_i| < 1, |x_N| < 2, t_1 < t < t_2\},$$

and define \tilde{u} , $\tilde{\mathbf{A}}$ as

$$\begin{aligned} \tilde{u}(x', x_N, t) &= \begin{cases} u(x', x_N, t) & \text{if } x_N \geq 0, \\ -u(x', -x_N, t) & \text{if } x_N < 0, \end{cases} \\ \tilde{A}_i(x', x_N, t) &= \begin{cases} A_i(x', x_N, t) & \text{if } x_N \geq 0, \\ -A_i(x', -x_N, t) & \text{if } x_N < 0, \end{cases} \quad i = 1, \dots, N-1, \\ \tilde{A}_N(x', x_N, t) &= \begin{cases} A_N(x', x_N, t) & \text{if } x_N \geq 0, \\ A_N(x', -x_N, t) & \text{if } x_N < 0, \end{cases} \end{aligned}$$

where $\tilde{\mathbf{A}}(x', x_N, t) = \tilde{\mathbf{A}}(x', x_N, t, \tilde{u}(x', x_N, t), D\tilde{u}(x', x_N, t))$. Then \tilde{u} is a weak solution in Q' to (1.1)–(1.2) with \mathbf{A} replaced by $\tilde{\mathbf{A}}$.

Lemma 2.13. *Let u be a non-negative weak supersolution to (1.1)–(1.2) in E_T , and let $a \in (0, 1)$. Let $Q^+ = K_{2\rho}(y) \times (s, s + \theta(2\rho)^p]$, and ξ be a positive number such that*

$$u(x, s) \geq \xi \quad \text{for a.e. } x \in K_{2\rho}(y)$$

and

$$\frac{|[u < \xi] \cap Q^+|}{|Q^+|} \leq \delta \frac{\xi^{2-p}}{\theta} \quad (2.11)$$

for a constant $\delta \in (0, 1)$ depending only on the data and a , and independent of ξ , ρ , and θ . Then

$$u \geq a\xi \quad \text{a.e. in } K_\rho(y) \times (s, s + \theta(2\rho)^p].$$

Notice that (2.11) is automatically satisfied for $\theta = \delta/\xi^{p-2}$.

2.4. Proofs of Theorems 2.1–2.4

2.4.1. *Flattening the boundary.* If we introduce the new variables

$$y_i = x_i, \quad i = 1, \dots, N - 1, \quad y_N = x_N - \Phi(x'),$$

the portion $\partial E \cap \{|x_i - x_{0,i}| < 2\rho, |x_N| < 8L\rho\}$ coincides with a portion of the hyperplane $y_N = 0$. Let $\tilde{K}_{2\rho}(x_0) = \{|x_i - x_{0,i}| < 2\rho, |x_N| < 4L\rho\}$. We orient y_N so that $E \cap \tilde{K}_{2\rho}(x_0) \subset \{y_N > 0\}$. It is easy to see that, in the new variables, (1.1) becomes

$$u_t - \operatorname{div}_y \tilde{\mathbf{A}}(y, t, u, D_y u) = 0,$$

and $\tilde{\mathbf{A}}(y, t, u, D_y u)$ satisfies the same kind of structural conditions as given in (1.2). We refer for more details to [16, Chapter X, §2].

Denoting again by x the transformed variables y , under the previously described change of variables, u is still a solution to an equation of type (1.1)–(1.2), and satisfies a homogeneous Dirichlet condition on a *flat* boundary. This is a notion which might not be standard in the literature, and we introduce it mainly to simplify the notation in the following. Consider the sets $\tilde{K}_{2\rho}(x_0)$ and

$$K_{2\rho}^*(x_0) = \{|x_i - x_{0,i}| < 2\rho, 0 < x_N < 4L\rho\}. \tag{2.12}$$

Definition 2.14. We say that the boundary of E is *flat with respect to x_N* if the portion $\partial E \cap \tilde{K}_{2\rho}(x_0)$ coincides with the portion of the hyperplane $\{x_N = 0\} \cap \tilde{K}_{2\rho}(x_0)$, and $K_{2\rho}^*(x_0) \subset E$. We orient x_N so that $E \cap \tilde{K}_{2\rho}(x_0) \equiv K_{2\rho}^*(x_0) \subset \{x_N > 0\}$.

Therefore, proving Theorem 2.1 reduces to the proof of the following lemma.

Lemma 2.15. *Let u be a non-negative weak solution to (1.1)–(1.2) in E_T for $p > 2$. Take $(x_0, t_0) \in S_T$, $\rho \in (0, r_0)$, let $P_\rho = (x'_0, 2L\rho, t_0)$, and assume that $u(P_\rho) > 0$, ∂E is flat with respect to x_N , and $(t_0 - \theta(4\rho)^p, t_0 + \theta(4\rho)^p) \subset (0, T]$, where $\theta = [c/u(P_\rho)]^{p-2}$ with c given by Theorem 2.6. Suppose that u vanishes continuously on $(\partial E \cap K_{2\rho}(x_0)) \times (t_0 - \theta(4\rho)^p, t_0 + \theta(4\rho)^p]$. Then there exist positive parameters $\alpha > \beta$ and a constant $\tilde{\gamma} > 0$, depending only on p, N, C_0, C_1 , such that*

$$u(x, t) \leq \tilde{\gamma} u(P_\rho) \tag{2.13}$$

for all $(x, t) \in \{|x_i - x_{0,i}| < \rho/4, 0 < x_N < 2L\rho\} \times (t_0 - \frac{\alpha+\beta}{2}\theta\rho^p, t_0 - \beta\theta\rho^p]$.

2.4.2. *Intrinsic rescaling and Harnack-based upper bounds.* The change of variable

$$x \rightarrow \frac{x - x_0}{2L\rho}, \quad t \rightarrow u(P_\rho)^{p-2} \frac{t - t_0}{\rho^p}$$

maps $Q_\rho^*(4^p\theta) = K_{2\rho}^*(x_0) \times (t_0 - \theta(4\rho)^p, t_0 + \theta(4\rho)^p]$ into

$$\hat{Q} = \{|y_i| < 1/L, 0 < y_N < 2\} \times (-4^p c^{p-2}, 4^p c^{p-2}],$$

x_ρ into $y_0 = (0, \dots, 0, 1)$, $\tilde{K}_\rho(x_0)$ into $\tilde{K}_1 = \{|y_i| < 1/(2L), |y_N| < 1\}$, $K_{2\rho}^*(x_0)$ into $K_2^*(y_0) = \{|y_i| < 1/L, 0 < y_N < 2\}$ and the portion $S_T \cap Q_\rho^*(4^p\theta)$ of the lateral boundary into

$$\Xi = \{(y', 0) : |y_i| < 1/L\} \times (-4^p c^{p-2}, 4^p c^{p-2}].$$

Denote again by (x, t) the transformed variables and let $y_0 = (0, \dots, 0, 1)$. Then the rescaled function

$$v_\rho(x, t) = \frac{1}{u(P_\rho)} u\left(2L\rho x + x_0, t_0 + \frac{t\rho^p}{u(P_\rho)^{p-2}}\right)$$

is a non-negative weak solution to

$$\partial_t v_\rho - \operatorname{div} \mathbf{A}_\rho(x, t, v, Dv) = 0$$

in \hat{Q} , where $v_\rho(y_0, 0) = 1$, and it is easy to see that \mathbf{A}_ρ satisfies structure conditions analogous to (1.2). Both here, and later on when dealing with a similar change of variable in the singular case, we drop the subscript ρ in v_ρ and \mathbf{A}_ρ for the sake of simplicity.

To avoid further technical complications, without loss of generality, we assume $L = 1$. The proof reduces to showing that there exists a constant $\tilde{\gamma}$, depending only on p, N, C_0, C_1 , such that

$$v(x, t) \leq \tilde{\gamma} \tag{2.14}$$

for all $(x, t) \in \{|x_i| < 1/2, 0 < x_N < 1\} \times (-\frac{\alpha+\beta}{2}c^{p-2}, -\beta c^{p-2}]$. In the following we denote again the rescaled function v by u .

Set

$$K = \{|x_i| < 1/2, 0 < x_N < 1\}, \quad Q = K \times [-\alpha c^{p-2}, -\beta c^{p-2}].$$

Suppose there exists $P_1 \in Q$ such that

$$u(P_1) \geq \gamma^{\lfloor k_0/\log_2(3/2) \rfloor + m},$$

where $\lfloor a \rfloor$ stands for the integer part of the real number a , $k_0 \in \mathbb{N}$ is sufficiently large, m will be fixed later on, and γ is the constant that appears in the Harnack inequality (2.7). We claim that

$$0 < x_{1,N} < (1/2)^{k_0}.$$

Indeed, if not, then $x_{1,N} \geq (1/2)^{k_0}$, and by repeated application of the Harnack inequality we will show that this yields a contradiction. This procedure will also determine the values of α and β .

With respect to space variables, the worst possible case for P_1 is when $x_{1,N} = (1/2)^{k_0}$, and $x_{1,1} = \dots = x_{1,N-1} = \pm 1/2$. For simplicity let us assume $x_{1,1} = \dots = x_{1,N-1} = 1/2$.

If we want to repeatedly apply the Harnack inequality, in this way getting closer and closer to $y_0 = (0, \dots, 0, 1)$, we need to evaluate the largest possible size of ρ in (2.7) at each step. The situation will then be the following:

$$\begin{aligned} x_{1,N} &= \left(\frac{1}{2}\right)^{k_0}, & \text{dist}(x_1, \partial E) &= \left(\frac{1}{2}\right)^{k_0}, & \rho_1 &= \frac{1}{2} \left(\frac{1}{2}\right)^{k_0}, \\ x_{2,N} &= \frac{3}{2} \left(\frac{1}{2}\right)^{k_0}, & \text{dist}(x_2, \partial E) &= \frac{3}{2} \left(\frac{1}{2}\right)^{k_0}, & \rho_2 &= \frac{1}{2} \frac{3}{2} \left(\frac{1}{2}\right)^{k_0}, \\ & \vdots & & & & \\ x_{j,N} &= \left(\frac{1}{2}\right)^{k_0} \left(\frac{3}{2}\right)^j, & \text{dist}(x_j, \partial E) &= \left(\frac{1}{2}\right)^{k_0} \left(\frac{3}{2}\right)^j, & \rho_j &= \frac{1}{2} \left(\frac{1}{2}\right)^{k_0} \left(\frac{3}{2}\right)^j. \end{aligned}$$

We need to determine the value of j at which we stop. We obviously need

$$\left(\frac{1}{2}\right)^{k_0} \left(\frac{3}{2}\right)^j = 1, \quad j \approx \frac{k_0}{\log_2(3/2)}.$$

Besides getting to y_0 , we need to have a full cube about it, where u is bounded below, and we also need to take into account the other coordinates, and not just x_N . Therefore, as j has to belong to \mathbb{N} , we eventually let

$$j = \tilde{k}_0 + 3,$$

where we have set $\tilde{k}_0 = \lfloor k_0 / \log_2(3/2) \rfloor$. Correspondingly, by repeated application of the Harnack inequality, we conclude that

$$\forall x \in K_{1/2}(y_0) \quad u(x, t_f) \geq \gamma^{\tilde{k}_0+m-j} = \gamma^{m-3}, \quad (2.15)$$

and now the main point becomes the evaluation of the interval where t_f can range. When dealing with the time variable, it is easy to see that we have two extreme situations.

The first extreme case is when $t_1 \approx -\beta c^{p-2}$ and

$$u(P_1) = \gamma^{\tilde{k}_0+m}, \quad u(P_2) = \gamma^{\tilde{k}_0+m-1}, \quad \dots, \quad u(P_j) = \gamma^{\tilde{k}_0+m-j},$$

that is, when the Harnack inequality gives the exact growth of u .

The second extreme case is when $t_1 \approx -\alpha c^{p-2}$ and

$$u(P_1) \gg \gamma^{\tilde{k}_0+m}, \quad u(P_2) \gg \gamma^{\tilde{k}_0+m}, \quad \dots, \quad u(P_j) \gg \gamma^{\tilde{k}_0+m},$$

that is, the function u is very large and its actual decrease (if any) cannot be evaluated.

In the latter situation, we can directly assume that $t_f = -\alpha c^{p-2}$. Let us evaluate what happens in the former case. By the repeated application of the Harnack inequality, we have

$$\begin{aligned}
 t_f &= -\beta c^{p-2} + \left(\frac{c}{\gamma^{\tilde{k}_0+m}}\right)^{p-2} \left[\frac{1}{2} \left(\frac{1}{2}\right)^{k_0}\right]^p \\
 &\quad + \dots + \left(\frac{c}{\gamma^{\tilde{k}_0+m-j}}\right)^{p-2} \left[\frac{1}{2} \left(\frac{1}{2}\right)^{k_0} \left(\frac{3}{2}\right)^j\right]^p \\
 &= -\beta c^{p-2} + \left(\frac{c}{\gamma^{\tilde{k}_0+m}}\right)^{p-2} \left(\frac{1}{2}\right)^{(k_0+1)p} \left(1 + \gamma^{p-2} \left(\frac{3}{2}\right)^p\right. \\
 &\quad \left. + \dots + \gamma^{(p-2)j} \left(\frac{3}{2}\right)^{jp}\right) \\
 &= -\beta c^{p-2} + \left(\frac{c}{\gamma^{\tilde{k}_0+m}}\right)^{p-2} \left(\frac{1}{2}\right)^{(k_0+1)p} \sum_{i=0}^j \left[\gamma^{p-2} \left(\frac{3}{2}\right)^p\right]^i \\
 &= -\beta c^{p-2} + \left(\frac{c}{\gamma^{\tilde{k}_0+m}}\right)^{p-2} \left(\frac{1}{2}\right)^{(k_0+1)p} \frac{(3/2)^{p(j+1)} \gamma^{(p-2)(j+1)} - 1}{(3/2)^p \gamma^{p-2} - 1},
 \end{aligned}$$

and we deduce that

$$\begin{aligned}
 t_f &< -\beta c^{p-2} + 2 \left(\frac{c}{\gamma^{\tilde{k}_0+m}}\right)^{p-2} \left(\frac{1}{2}\right)^{(k_0+1)p} \frac{(3/2)^{p(j+1)} \gamma^{(p-2)(j+1)}}{(3/2)^p \gamma^{p-2}} \\
 &< -\beta c^{p-2} + 2 \left(\frac{c}{\gamma^{\tilde{k}_0+m}}\right)^{p-2} \left(\frac{1}{2}\right)^{(k_0+1)p} \left(\frac{3}{2}\right)^{jp} \gamma^{j(p-2)} \\
 &< -\beta c^{p-2} + \left(\frac{c}{\gamma^{\tilde{k}_0+m-j}}\right)^{p-2} \left(\frac{9}{4}\right)^p = -\beta c^{p-2} + \left(\frac{c}{\gamma^{m-3}}\right)^{p-2} \left(\frac{9}{4}\right)^p,
 \end{aligned}$$

where we have taken into account the value of j . Therefore, we conclude that

$$t_f \in \left[-\alpha c^{p-2}, \left(-\beta + \frac{(9/4)^p}{\gamma^{(m-3)(p-2)}} \right) c^{p-2} \right].$$

Correspondingly, for such a t_f , we have not only (2.15), but also

$$\forall x \in K_{1/4}(y_0) \quad u(x, t_f) \geq \gamma^{m-3-l},$$

where $l \in \mathbb{N}$ is to be found. The next calculations will determine l , and consequently m, α, β , in order to have a contradiction and prove the claim.

We apply Lemma 2.13, setting $\xi = \gamma^{m-3-l}, a = \gamma^{-1}$; the only role played by l is to provide a *smaller* initial lower bound on u , and therefore generate a *longer* cylinder Q^+ , where the information propagates, and such that $(y_0, 0) \in Q^+$. Since $u(y_0, 0) = 1$, we have a contradiction if we end up with $u(y_0, 0) > 1$. Hence, we need to have

- (i) $-\alpha < -\beta$: we do not want the lower and upper bases of the cylinder Q to coincide;
- (ii) $-\beta + (9/4)^p / \gamma^{(m-3)(p-2)} \leq 0$: we need to be below the reference point $(y_0, 0)$ with respect to the time variable;

- (iii) $-\alpha c^{p-2} + (\delta/\gamma^{m-3-l})^{p-2}(1/4)^p \geq 0$: the cylinder Q^+ should include the reference point $(y_0, 0)$;
 (iv) $\gamma^{m-4-l} > 1$: this yields $u(y_0, 0) > 1$.

First of all, choose l such that $\gamma^{(p-2)l} = 10^p c^{p-2} / \delta^{p-2}$, and then $m = l + 5$; then condition (iv) is satisfied. Notice that both l and m depend only on the data, but not on k_0 . Finally let

$$\beta = \frac{(9/4)^p}{\gamma^{(m-3)(p-2)}}, \quad \alpha = \frac{(10/4)^p}{\gamma^{(m-3)(p-2)}}.$$

In this way, conditions (i)–(iv) are all satisfied, and we have obtained the desired contradiction. Therefore, we conclude that if $u(P_1) > \gamma^{\tilde{k}_0+m}$, then $x_{1,N} < 1/2^{k_0}$. This obviously implies that

$$x_{1,N} < 1/2^{k_0} \quad \text{whenever} \quad u(P_1) > \gamma^{m(k_0+1)}. \quad (2.16)$$

From here on, we set $\gamma^m = H$.

Remark 2.16. Instead of using Lemma 2.13, the propagation of the bound below, which generates the contradiction, can be proved by a further application of the Harnack inequality.

2.4.3. *End of the proof of Theorem 2.1.* Let

$$Q_* = \left\{ |x_i| < \frac{1}{4}, 0 < x_N < 1, -\frac{\alpha + \beta}{2} c^{p-2} < t \leq -\beta c^{p-2} \right\},$$

$$Q^* = \left\{ |x_i| < \frac{1}{4}, -1 < x_N < 1, -\frac{\alpha + \beta}{2} c^{p-2} < t \leq -\beta c^{p-2} \right\}.$$

Starting from Q_* , Q^* is built by reflection; after extending u to Q^* as in Lemma 2.12, u is still a (signed) solution to (1.1)–(1.2).

Now let $P_1 = (x_1, t_1) = (x_{1,1}, \dots, x_{1,N}, t_1) \in Q_*$ be such that

$$u(P_1) \geq H^{k_0+1}; \quad (2.17)$$

by (2.16), we must have $0 < x_{1,N} < 1/2^{k_0}$. Set

$$Q(P_1) = \{|x - x_1| < 2^{-k_0} \epsilon^{-s}, t_1 - (2^{-k_0} \epsilon^{-s})^p < t < t_1\},$$

where ϵ is the quantity claimed by Lemma 2.10 and $s \in \mathbb{N}$ is to be fixed. Let

$$\omega_0^{(1)} = \operatorname{osc}_{Q(P_1)} u.$$

Without loss of generality, we may assume $\epsilon < 1/2$. Thanks to (2.17) and the construction of u by odd reflection, we have

$$\omega_0^{(1)} \geq 2H^{k_0+1}. \quad (2.18)$$

Moreover, if k_0 is large enough, we have $Q(P_1) \subset Q^*$. Set $\sigma_0^{(1)} = 2^{-k_0}\epsilon^{-s}$ and consider

$$Q_0^{(1)} = K_{\sigma_0^{(1)}}(x_1) \times (t_1 - \theta_0^{(1)}(\sigma_0^{(1)})^p, t_1), \quad \text{where } \theta_0^{(1)} = (c/\omega_0^{(1)})^{p-2}.$$

It is apparent that $Q_0^{(1)} \subset Q(P_1) \subset Q^*$. Notice that we do not need to assume (2.9) here, since, by construction, the cylinders are all correctly nested in one another. By Lemma 2.10, we can build a sequence

$$\omega_n^{(1)} = \delta\omega_{n-1}^{(1)}, \quad \theta_n^{(1)} = (c/\omega_n^{(1)})^{p-2}, \quad \sigma_n^{(1)} = \epsilon\sigma_{n-1}^{(1)}, \quad Q_n^{(1)} = Q_{\sigma_n^{(1)}}(\theta_n^{(1)}),$$

for all non-negative integers n . Such a sequence satisfies

$$Q_{n+1}^{(1)} \subset Q_n^{(1)}, \quad \text{osc}_{Q_n^{(1)}} u \leq \omega_n^{(1)}.$$

By iteration

$$\text{osc}_{Q_n^{(1)}} u \leq \delta^n \omega_0^{(1)} = \delta^n \text{osc}_{Q(P_1)} u, \quad \text{so} \quad \text{osc}_{Q(P_1)} u \geq \frac{1}{\delta^n} \text{osc}_{Q_n^{(1)}} u.$$

If we now choose $n = s$, and s such that $\delta^{-s} > H^{10}$, by the choice of $\sigma_0^{(1)}$ we obtain

$$\omega_0^{(1)} \geq 2H^{k_0+11},$$

and this obviously improves the previous lower bound given by (2.18). As u has been built by odd reflection, we conclude there must exist $P(x_2, t_2) = (x_{2,1}, \dots, x_{2,N}, t_2) \in Q(P_1)$ such that

$$u(P_2) \geq H^{k_0+11}.$$

As before, by (2.16), $0 < x_{2,N} < 1/2^{k_0+10}$, and also $t_1 - (2^{-k_0}\epsilon^{-s})^p < t_2 < t_1$. Set

$$Q(P_2) = \{|x - x_2| < 2^{-k_0-10}\epsilon^{-s}, t_2 - (2^{-k_0-10}\epsilon^{-s})^p < t < t_2\}.$$

Once more, provided k_0 is large enough, we can assume that $Q(P_2) \subset Q^*$. Arguing as before, we conclude there exists $P_3(x_3, t_3) \in Q(P_2)$ such that

$$u(P_3) \geq H^{k_0+21}.$$

Consequently, $0 < x_{3,N} < 2^{-k_0-20}$. By induction, we get $\{P_q(x_q, t_q)\}$ such that

$$u(P_q) \geq H^{k_0+1+10(q-1)} \tag{2.19}$$

and

$$0 < x_{q,N} < 2^{-k_0-10(q-1)}.$$

Notice that k_0 depends on α, β, c , and s , and therefore, due the definition of these quantities, on the data p, N, C_0, C_1 . Now choose k_0 so large that

$$t_1 - \sum_{q=1}^{\infty} (2^{-k_0-10(q-1)}\epsilon^{-s})^p > -\alpha c^{p-2},$$

and for all $i = 1, \dots, N - 1$,

$$x_{1,i} - \sum_{q=1}^{\infty} (2^{-k_0-10(q-1)} \epsilon^{-s}) > -\frac{1}{2}, \quad x_{1,i} + \sum_{q=1}^{\infty} (2^{-k_0-10(q-1)} \epsilon^{-s}) < \frac{1}{2}.$$

It is the need to satisfy these requirements that forces $|x_i| < 1/4$ in the definition of Q_* and Q^* . Indeed, the value of k_0 determines the value of $\tilde{\gamma}$ in (2.14), and therefore we must be able to choose k_0 independently of the solution or any other geometrical conditions. If $|x_i|$ were larger than $1/4$, then k_0 would also depend on the distance in space of P_1 to the boundary of the cube of edge $1/2$.

Once the previous conditions are satisfied, the sequence $\{P_q\}$ is contained in a fixed cylinder of Q_* . Together with (2.19), this leads to a contradiction, since the sequence approaches the boundary, and the corresponding values of u grow arbitrarily large, whereas u is assumed to vanish continuously at the boundary. \square

2.4.4. Proof of Corollary 2.3. Let

$$\mathcal{K} = \overline{\Psi_{\rho/2}^-(x_0, t_0)} \cap (\partial E \times (0, T)), \quad M = \sup_{\Psi_{\rho/2}^-(x_0, t_0)} u.$$

By [16, Chapter III, Theorem 1.2], for all $(x, t) \in \Psi_{\rho/2}^-(x_0, t_0)$ we have

$$u(x, t) \leq \gamma M \left(\frac{\inf_{(y,s) \in \mathcal{K}} (|x - y| + M^{(p-2)/p} |t - s|^{1/p})}{\rho} \right)^\mu.$$

Since we are dealing with a Lipschitz cylinder, in the infimum above we can take $t \equiv s$, which reduces the above to

$$u(x, t) \leq \gamma M \left(\frac{\text{dist}(x, \partial E)}{\rho} \right)^\mu.$$

By (2.1) and possibly a further application of the Harnack inequality, we conclude the proof. \square

2.4.5. Proof of Theorem 2.4. Since the boundary is of class C^2 , its portion $\partial E \cap K_{2\rho}(x_0)$ can now be represented by $x_N = \Phi(x')$, where Φ is a function of class C^2 satisfying

$$\Phi(x'_0) = 0, \quad D\Phi(x'_0) = 0, \quad \|D'\Phi\|_\infty \leq M_2, \quad \|(D')^2\Phi\|_\infty \leq M_2,$$

and $M_2 \in (0, 1)$ is a suitable parameter, provided ρ is small enough.

If we now introduce the new variables

$$y_i = x_i, \quad i = 1, \dots, N - 1, \quad y_N = x_N - \Phi(x'),$$

the portion $\partial E \cap K_{2\rho}(x_0)$ of the boundary coincides with the portion of the hyperplane $y_N = 0$ within $K_{2\rho}(y_0)$. We orient y_N so that $E \cap K_{2\rho}(y_0) \subset \{y_N > 0\}$. Denoting again

by x the transformed variable y , we proceed as in the proof of Theorem 2.1. Consequently, (1.1)₀ is rewritten as

$$u_t - \operatorname{div} \mathbf{A}(x, Du) = 0, \quad (2.20)$$

where, just as in the proof of Theorem 2.1,

$$\begin{cases} \mathbf{A}(x, \boldsymbol{\eta}) \cdot \boldsymbol{\eta} \geq C_0 |\boldsymbol{\eta}|^p \\ |\mathbf{A}(x, \boldsymbol{\eta})| \leq C_1 |\boldsymbol{\eta}|^{p-1} \end{cases} \quad \text{a.e. } (x, t) \in E_T, \quad (2.21)$$

C_0, C_1 are positive constants depending only on N, p, M_2 , and

$$\boldsymbol{\eta}' = (\eta_1, \dots, \eta_{N-1}), \quad |\boldsymbol{\eta}|^2 = \sum_{k=1}^N |\eta_k|^2.$$

Our solution to (2.20) vanishes continuously on $K_{2\rho}(x_0) \cap \{x_N = 0\}$. Moreover, by lengthy but quite straightforward calculations,

$$\frac{\partial A_i}{\partial \eta_k} \xi_i \xi_k \geq C_2 |\boldsymbol{\eta}|^{p-2} |\xi|^2, \quad (2.22)$$

$$\left| \frac{\partial A_i}{\partial \eta_k} \right| \leq C_3 |\boldsymbol{\eta}|^{p-2}, \quad \left| \frac{\partial A_i}{\partial x_k} \right| \leq C_4 |\boldsymbol{\eta}|^{p-1}, \quad (2.23)$$

where C_2, C_3, C_4 depend only on N, p , and M_2 . Finally, due to its definition, \mathbf{A} satisfies the homogeneity condition

$$\mathbf{A}(x, \boldsymbol{\eta}) = |\boldsymbol{\eta}|^{p-1} \mathbf{A}(x, \boldsymbol{\eta}/|\boldsymbol{\eta}|). \quad (2.24)$$

Therefore, proving Theorem 2.4 reduces to proving

Lemma 2.17. *Let u be a non-negative weak solution to*

$$u_t - \operatorname{div} \mathbf{A}(x, Du) = 0 \quad \text{in } E_T$$

for $p > 2$, where \mathbf{A} satisfies the structure conditions (2.21)–(2.24). Take $(x_0, t_0) \in S_T$, $\rho \in (0, r_0)$, let $P_\rho = (x'_0, 2M_2\rho, t_0)$, and assume that $u(P_\rho) > 0$, ∂E is flat with respect to x_N , and $(t_0 - \theta(4\rho)^p, t_0 + \theta(4\rho)^p) \subset (0, T]$, where $\theta = [c/u(P_\rho)]^{p-2}$. Suppose that u vanishes continuously on $(\partial E \cap K_{2\rho}(x_0)) \times (t_0 - \theta(4\rho)^p, t_0 + \theta(4\rho)^p)$. Then there exists a constant $\gamma > 0$, depending only on $N, p, C_i, i = 0, \dots, 4$, such that

$$0 \leq u(x', x_N, t) \leq \gamma(x_N/\rho)u(P_\rho)$$

for all $(x, t) \in \{|x_i - x_{0,i}| < \rho/2, 0 < x_N < 2M_2\rho\} \times (t_0 - \frac{\alpha+3\beta}{4}\theta\rho^p, t_0 - \beta\theta\rho^p]$.

Proof. We use the same argument of [22, Theorem 4.1]. By suitable rescaling and translation, assume $x_0 = 0$, $\rho = 1$, and let $\mathcal{M} = \sup_{\Psi_1^-} u$. By (2.13), $\mathcal{M} \leq \tilde{\gamma}u(P_1)$. Let $y = (0, \dots, 0, -1)$, $t_1 = t_0 - \beta\theta$, and consider the function

$$\eta_k(x, t) = \exp[-k(|x - y| - 1)] \exp[u(P_1)^{p-2}(t - t_1)],$$

and the set

$$\mathcal{N}_k = \{(x, t) : x_N > 0, 1 < |x - y| < 1 + 1/k, t_1 - (1 + \delta)\beta\theta < t < t_1\},$$

for some small enough, positive parameter δ . We assume k is so large that $\mathcal{N}_k \subset \Psi_1^-$. If we choose $C = \tilde{\gamma} \max\{(1 - e^{-1})^{-1}; (1 - e^{-(1+\delta)\beta c^{p-2}})^{-1}\}$, where $\tilde{\gamma}$ is the constant of (2.1), it is easy to verify that $\Theta_k(x, t) = Cu(P_1)(1 - \eta_k(x, t))$ satisfies $u \leq \Theta_k$ on the parabolic boundary of \mathcal{N}_k . Provided we choose k as the largest positive root of

$$C_0(p - 1)k^p - bk^{p-1} - C^{2-p} = 0,$$

where b is a positive quantity that depends only on M_2 , and C_0 is the constant of (2.21), relying on (2.21)–(2.24), it is a matter of straightforward calculations to verify that Θ_k is a supersolution to (2.20) in \mathcal{N}_k .

By the comparison principle, $u \leq \Theta_k$ in \mathcal{N}_k . In particular, for all $0 < x_N < 1/k$,

$$\begin{aligned} u(0, \dots, 0, x_N, t_1) &\leq \Theta_k(0, \dots, 0, x_N, t_1) = Cu(P_1)(1 - e^{-kx_N}) \\ &\leq \tilde{\gamma}x_N u(P_1). \end{aligned}$$

On the other hand, if $x_N \geq 1/k$, then

$$u(0, \dots, 0, x_N, t_1) \leq \mathcal{M} \leq \tilde{\gamma}kx_N u(P_1).$$

The same argument can be repeated using any $t_1 \in (t_0 - \frac{\alpha+3\beta}{4}\theta, t_0 - \beta\theta]$, and switching back to the original coordinates, we conclude. \square

Remark 2.18. Instead of relying on the comparison principle, Theorem 2.4 could be proved using the L^∞ -estimates for the gradient Du , as shown, for example, in [16, Chapter VIII] (see also [37, 38]). However, similar estimates are not known for solutions to equation (4.1)₀, and one does not expect them to hold true for solutions to (4.1)–(4.2); therefore, we opted for an approach that works in both cases.

3. The singular supercritical case $2N/(N + 1) < p < 2$

3.1. A weak Carleson estimate

We consider our first result in the singular supercritical case $2N/(N + 1) < p < 2$. Let $E_T, u, (x_0, t_0), \rho, x_\rho, P_\rho$ be as in Theorem 2.1 and set

$$I(t_0, \rho, h) = (t_0 - h\rho^p, t_0 + h\rho^p).$$

Moreover, let u be a weak solution to (1.1)–(1.2) such that

$$0 < u \leq M \quad \text{in } E_T, \quad (3.1)$$

and assume that

$$I(t_0, 9\rho, M^{2-p}) \subset (0, T]. \quad (3.2)$$

Then we define

$$\begin{aligned} \tilde{\Psi}_\rho(x_0, t_0) &= E_T \cap \{(x, t) : |x_i - x_{0,i}| < 2\rho, |x_N| < 4L\rho, t \in I(t_0, 9\rho, \eta_\rho^{2-p})\}, \\ \bar{\Psi}_\rho(x_0, t_0) &= E_T \cap \{(x, t) : |x_i - x_{0,i}| < \rho/4, |x_N| < 2L\rho, t \in I(t_0, \rho, \eta_\rho^{2-p})\}, \end{aligned}$$

where η_ρ is the first root of the equation

$$\max_{\tilde{\Psi}_\rho(x_0, t_0)} u = \eta_\rho. \quad (3.3)$$

Notice that both the functions $y_1(\eta_\rho) = \max_{\tilde{\Psi}_\rho(x_0, t_0)} u$ and $y_2(\eta_\rho) = \eta_\rho$ are increasing. Moreover

$$\begin{cases} y_1(0) \geq u(P_\rho) > 0, \\ y_2(0) = 0, \end{cases} \quad \text{and} \quad \begin{cases} y_1(M) \leq M, \\ y_2(M) = M. \end{cases}$$

Therefore, it is immediate that at least one root of (3.3) actually exists. Moreover, by (3.2), $\tilde{\Psi}_\rho(x_0, t_0) \subset E_T$.

As already mentioned in the Introduction, we can only provide a weak form of the Carleson estimate, expressed by the following theorem.

Theorem 3.1 (Carleson-type estimate, weak form, $2N/(N+1) < p < 2$). *Let u be a weak solution to (1.1)–(1.2) that satisfies (3.1). Assume that (3.2) holds true and u vanishes continuously on*

$$\partial E \cap \{|x_i - x_{0,i}| < 2\rho, |x_N| < 8L\rho\} \times I(t_0, 9\rho, M^{2-p}).$$

Then there exist constants $\gamma > 0$ and $\alpha \in (0, 1)$, depending only on p, N, C_0, C_1, L , such that

$$u(x, t) \leq \gamma \left(\frac{\text{dist}(x, \partial E)}{\rho} \right)^\alpha \sup_{\tau \in I(t_0, \rho, 2\eta_\rho^{2-p})} u(x_\rho, \tau)$$

for every $(x, t) \in \bar{\Psi}_\rho(x_0, t_0)$.

If we let

$$\Psi_{\rho, M}(x_0, t_0) = E_T \cap \{(x, t) : |x_i - x_{0,i}| < \rho/4, |x_N| < 2L\rho, t \in I(t_0, \rho, M^{2-p})\},$$

we have a second statement.

Corollary 3.2. *Under the same assumptions of Theorem 3.1, we have*

$$u(x, t) \leq \gamma \left(\frac{\text{dist}(x, \partial E)}{\rho} \right)^\alpha \sup_{\tau \in I(t_0, \rho, 2M^{2-p})} u(x_\rho, \tau)$$

for every $(x, t) \in \Psi_{\rho, M}(x_0, t_0)$.

The quantity η_ρ is known only qualitatively through (3.3), whereas M is a datum. Therefore, Corollary 3.2 can be viewed as a quantitative version of a purely qualitative statement. On the other hand, since η_ρ could be attained in P_ρ , Theorem 3.1 gives the sharpest possible statement, and is genuinely intrinsic.

Moreover, compared to Theorem 2.1 and Corollary 2.3, Theorem 3.1 combines two distinct statements in a single one (mainly for simplicity), and presents two fundamental differences: when $p > 2$, the value of u at a point above controls the values of u below, whereas when $2N/(N + 1) < p < 2$, the **maximum** of u over a proper time interval centered at t_0 controls the values of u **both above and below** the time level t_0 . These are consequences of the different statements of the Harnack inequality in the two cases. In fact, the following theorem is proved in [20] (see also [21] for a thorough presentation).

For fixed $(x_0, t_0) \in E_T$ and $\rho > 0$, we set $\mathcal{M} = \sup_{K_\rho(x_0)} u(x, t_0)$, and require that

$$K_{8\rho}(x_0) \times I(t_0, 8\rho, \mathcal{M}^{2-p}) \subset E_T. \tag{3.4}$$

Theorem 3.3 (Harnack inequality). *Let u be a non-negative weak solution to (1.1)–(1.2) in E_T for $p \in (2N/(N + 1), 2)$. Then there exist constants $\bar{\epsilon} \in (0, 1)$ and $\bar{\gamma} > 1$, depending only on p, N, C_0, C_1 , such that for all intrinsic cylinders $(x_0, t_0) + Q_{8\rho}^\pm(\theta)$ for which (3.4) holds,*

$$\bar{\gamma}^{-1} \sup_{K_\rho(x_0)} u(\cdot, \sigma) \leq u(x_0, t_0) \leq \bar{\gamma} \inf_{K_\rho(x_0)} u(\cdot, \tau) \tag{3.5}$$

for any pair of time levels σ, τ in the range

$$t_0 - \bar{\epsilon} u(x_0, t_0)^{2-p} \rho^p \leq \sigma, \tau \leq t_0 + \bar{\epsilon} u(x_0, t_0)^{2-p} \rho^p. \tag{3.6}$$

The constants $\bar{\epsilon}$ and $\bar{\gamma}^{-1}$ tend to zero both as $p \rightarrow 2$ and as $p \rightarrow 2N/(N + 1)$.

Remark 3.4. Compared to the degenerate case, we now have $c = 1$ for the size of the intrinsic cylinders. The upper bound \mathcal{M} has only the qualitative role to ensure that $(x_0, t_0) + Q_{8\rho}^\pm(\mathcal{M})$ are contained within the domain of definition of u .

3.2. A counterexample

Can we improve the result of Theorem 3.1, namely can we replace the supremum of u on $I(t_0, \rho, 2\eta_\rho^{2-p})$ with the pointwise value $u(P_\rho)$? This would certainly be possible if there existed a constant γ , depending only on the data p, N, C_0, C_1 , such that

$$\forall t \in I(t_0, \rho, 2\eta_\rho^{2-p}) \quad u(x_\rho, t) \leq \gamma u(P_\rho).$$

From a geometrical point of view, this amounts to building a Harnack chain connecting (x_ρ, t) and P_ρ , for all $t \in I(t_0, \rho, 2\eta_\rho^{2-p})$. In general, without further assumptions on u , this is not possible, as the following counterexample shows.

Let u be the unique non-negative solution to

$$\begin{cases} u \in C(\mathbb{R}_+; L^2(E)) \cap L^p(\mathbb{R}_+; W_0^{1,p}(E)), \\ u_t - \operatorname{div}(|Du|^{p-2}Du) = 0 \quad \text{in } E_T, \\ u(\cdot, 0) = u_0 \in C(\bar{E}), \end{cases}$$

with $u_0 > 0$ in E , and $u_0|_{\partial E} = 0$. By [16, Chapter VII, Proposition 2.1], there exists a finite time T_* , depending only on N, p, u_0 , such that $u(\cdot, t) \equiv 0$ for all $t \geq T_*$. By the results of [16, Chapter IV], $u \in C(\bar{E} \times (0, T_*))$. Suppose now that at time $t = T_* + 1$, we modify the boundary value and for any $t > T_* + 1$ we let $u(\cdot, t)|_{\partial E} = g(\cdot, t)$, where g is continuous and strictly positive. It is immediate to verify that u becomes strictly positive for any $t > T_* + 1$. Therefore, the positivity set for u is disconnected, $u(x, t) \equiv 0$ for all $(x, t) \in \bar{E} \times (T_*, T_* + 1)$, and if (x_ρ, t) and P_ρ lie on opposite sides of the vanishing layer for u , by the intrinsic nature of Theorem 3.3, there is no way to connect them with a Harnack chain.

This counterexample allows u to vanish identically for t in a proper interval, but by suitably modifying the boundary values, it is clear that we can have u strictly positive, and as close to zero as we want. Therefore, the impossibility of connecting two arbitrary points by a Harnack chain does not depend on the vanishing of u , but it is a general property of solutions to (1.1)–(1.2) whenever $E \not\cong \mathbb{R}^N$. Moreover, by properly adjusting the boundary value, one can even create an arbitrary number of oscillations for u between positivity and null regions.

We have considered solutions to the p -Laplace equation just for simplicity; everything continues to hold if we consider the same boundary value problem for (1.1)–(1.2).

Notice that if we deal with weak solutions to (1.1)–(1.2) in $\mathbb{R}^N \times (0, T]$, then we do not have boundary values any more, the situation previously discussed cannot occur, and therefore any two points (x, t) and (x_0, t_0) can always be connected by a Harnack chain provided both $u(x, t)$ and $u(x_0, t_0)$ are strictly positive, and $0 < t - t_0 < (\epsilon/8^p)t_0$, as discussed in [21, Chapter 7, Proposition 4.1]. The sub potential lower bound discussed there is then a property of weak solutions given in the whole $\mathbb{R}^N \times (0, T)$.

The Harnack inequality given in Theorem 3.3 is time-insensitive, and its constants are not stable as $p \rightarrow 2$. A different statement, analogous to the one given in Theorem 2.6, could be given, and in such a case the constants would be stable (see [21, Chapter 6] for a thorough discussion of the two possible forms). However, the eventual result is the same, and independently of the kind of Harnack inequality one considers, two points (x, t) and (x_0, t_0) of positivity for u cannot be connected by a Harnack chain.

Notice that we have a sort of dual situation: when $1 < p < 2$ the support of u can be disconnected in time; when $p > 2$, as we discussed in §2.2, the support can be disconnected in space.

Strictly speaking, the previous counterexample only shows that we cannot replace the line with a point, but per se it does not rule out that a strong form of Carleson's estimate

holds true all the same. However, if one tries to adapt to the singular supercritical context the standard proof based on the Harnack inequality and the boundary Hölder continuity (as we did, for example, in the degenerate context), then one quickly realizes that, in order to have the cylinders of Lemma 2.10 inside the reference cylinder, one needs to know in advance the oscillation of u ; this suggests that only a control in terms of the supremum taken in a proper set can be feasible.

3.3. A strong Carleson estimate

Compared to the statement of Theorem 3.1, a stronger form is indeed possible, provided we allow the parameter γ to depend not only on the data, but also on the oscillation of u . Let $E_T, u, (x_0, t_0), \rho, P_\rho$ be as in Theorem 2.1, and for $k = 0, 1, \dots$ set

$$\begin{aligned} \rho_k &= (7/8)^k \rho, & \sigma_k &= \rho_k / \gamma^{k(2-p)/p}, \\ x_{\rho_k} &= (x'_0, 2L\rho_k), & P_{\rho_k} &= (x'_0, 2L\rho_k, t_0), \\ \Psi_{\rho_k, M}(x_0, t_0) &= E_T \cap \{(x, t) : |x_i - x_{0,i}| < \rho_k/4, |x_N| < 2L\rho_k, t \in I(t_0, \sigma_k, M^{2-p})\}, \\ m_0 &= \inf_{\tau \in I(t_0, \rho, 2M^{2-p})} u(x_\rho, \tau), & M_0 &= \sup_{\tau \in I(t_0, \rho, 2M^{2-p})} u(x_\rho, \tau). \end{aligned}$$

Corollary 3.5 (Carleson-type estimate, strong form, $2N/(N + 1) < p < 2$). *Let u be a weak solution to (1.1)–(1.2) such that $0 < u \leq M$ in E_T . Assume that $I(t_0, 9\rho, M^{2-p}) \subset (0, T]$ and that u vanishes continuously on*

$$\partial E \cap \{|x_i - x_{0,i}| < 2\rho, |x_N| < 8L\rho\} \times I(t_0, 9\rho, M^{2-p}).$$

Then there exists a constant γ , depending only on $p, N, C_0, C_1, L, M/m_0$, such that

$$u(x, t) \leq \gamma u(P_{\rho_k}) \tag{3.7}$$

for every $(x, t) \in \Psi_{\rho_k, M}(x_0, t_0)$ and all $k = 0, 1, \dots$

Remark 3.6. The strong form of the Carleson-type estimate is derived from Corollary 3.2. An analogous statement can be derived from Theorem 3.1.

Remark 3.7. Estimate (3.7) has the same structure as the backward Harnack inequality for caloric functions that vanish just on a disk at the boundary (see [9, Theorem 13.7, p. 234]). This is not surprising, because (3.7) is indeed a backward Harnack inequality, due to the specific nature of the Harnack inequality for the singular case. However, it is worth mentioning that things are not completely equivalent; indeed, the constants we have in the time-insensitive Harnack inequality (3.5)–(3.6) are not stable (and cannot be stabilized), and therefore the result for caloric functions cannot be recovered from the singular case by simply letting $p \rightarrow 2$ (as is the case for many other results).

Another striking difference compared to the degenerate case appears when we consider $C^{1,1}$ cylinders and (mainly for simplicity) the prototype equation (1.1)₀. In this case, indeed, weak solutions vanishing on the lateral part enjoy linear behavior at the boundary with implications expressed in the following result. Note that the role of L in the definition of $\Psi_{\rho, M}$ is now played by $M_{1,1}$.

Theorem 3.8. *Let $2N/(N + 1) < p < 2$. Assume E_T is a $C^{1,1}$ cylinder, and (x_0, t_0) , ρ, P_ρ are as in Theorem 2.1. Let u, v be two weak solutions to (1.1)₀ in E_T satisfying the hypotheses of Theorem 3.1, that is, $0 < u, v \leq M$ in E_T . Then there exist positive constants $\bar{s}, \gamma, \beta, 0 < \beta \leq 1$, depending only on $N, p, M_{1,1}$, and $\rho_0, c_0 > 0$, depending also on the oscillation of u , such that the following properties hold.*

(a) Hopf principle:

$$|Du| \geq c_0 \quad \text{in } \Psi_{\rho_0, M}(x_0, t_0). \tag{3.8}$$

(b) Boundary Harnack inequality:

$$\gamma^{-1} \frac{\inf_{\tau \in I(t_0, \rho, 2M^{2-p})} u(x_\rho, \tau)}{\sup_{\tau \in I(t_0, \rho, 2M^{2-p})} v(x_\rho, \tau)} \leq \frac{u(x, t)}{v(x, t)} \leq \gamma \frac{\sup_{\tau \in I(t_0, \rho, 2M^{2-p})} u(x_\rho, \tau)}{\inf_{\tau \in I(t_0, \rho, 2M^{2-p})} v(x_\rho, \tau)} \tag{3.9}$$

for all $(x, t) \in \{x \in K_{\bar{s}\rho/4}(x_0) \cap E : \text{dist}(x, \partial E) < \bar{s}\rho/8\} \times I(t_0, \rho, \frac{1}{2}M^{2-p})$ with $\rho < \rho_0$.

(c) The quotient u/v is Hölder continuous with exponent β in $\Psi_{\rho_0/2, M}(x_0, t_0)$.

Remark 3.9. Since

$$\frac{\sup_{\tau \in I(t_0, \rho, 2M^{2-p})} u(x_\rho, \tau)}{\inf_{\tau \in I(t_0, \rho, 2M^{2-p})} v(x_\rho, \tau)} \leq \frac{M_{0,u}u(P_\rho)}{m_{0,u}} \frac{M_{0,v}}{m_{0,v}v(P_\rho)},$$

$$\frac{\inf_{\tau \in I(t_0, \rho, 2M^{2-p})} u(x_\rho, \tau)}{\sup_{\tau \in I(t_0, \rho, 2M^{2-p})} v(x_\rho, \tau)} \leq \frac{m_{0,u}u(P_\rho)}{M_{0,u}} \frac{m_{0,v}}{M_{0,v}v(P_\rho)},$$

the boundary Harnack inequality (3.9) can be rewritten as

$$\tilde{\gamma}^{-1} \frac{u(P_\rho)}{v(P_\rho)} \leq \frac{u(x, t)}{v(x, t)} \leq \tilde{\gamma} \frac{u(P_\rho)}{v(P_\rho)}$$

where now $\tilde{\gamma}$ depends not only on $N, p, M_{1,1}$, but also on $M_{0,u}/m_{0,u}$ and $M_{0,v}/m_{0,v}$.

Remark 3.10. Note that (a) implies that near a part of the lateral boundary where a non-negative solution vanishes, the parabolic p -Laplace operator is uniformly elliptic. Since we do not have an estimate at the boundary of the type $|Du(x, t)| \geq u(x, t)/\text{dist}(x, \partial E)$, (a) and (c) hold only in a small neighborhood of S_T , whose size depends on the solution, as both c_0 and the oscillation of the gradient Du depend on the oscillation of u ; this is precisely the meaning of ρ_0 . Moreover, as will be clear from the proof, we give only a qualitative dependence of β on the various quantities.

The proof relies on suitable estimates from above and below, which were originally proved in [22, §4] for solutions to singular porous medium equations in C^2 domains by building explicit barriers; they were later extended in [49] to solutions to doubly non-linear singular equations in $C^{1,\alpha}$ domains. Unfortunately, it recently turned out that there is a flaw in the argument, and, as pointed out in [39] in the context of p -Laplacian type equations, $C^{1,1}$ domains is the most general assumption one can have in order to build

barriers; that this is the threshold below which one cannot go had already been shown in the elliptic context in [2].

We recast these estimates in the lemma below, in a form tailored to our purposes. Indeed, the Hopf principle and a weak version of the boundary Harnack inequality follow easily from these estimates. Our improvement lies in the use of the Carleson estimates, which allow a more precise bound for $u(x, t)/v(x, t)$ in terms of $u(P_\rho)/v(P_\rho)$. The restriction to $2N/(N + 1) < p < 2$ comes into play only in this last step.

Thus, let ∂E be of class $C^{1,1}$ and u be a non-negative weak solution to (1.1)₀ in E_T for $1 < p < 2$. Assume that $u \leq M$ in E_T . For $x \in \mathbb{R}^N$, set $d(x) = \text{dist}(x, \partial E)$, and for $s > 0$, let

$$E^s = \{x \in E : s/2 \leq d(x) \leq 2s\}.$$

Lemma 3.11. *Let $\tau \in (0, T)$ and fix $x_0 \in \partial E$. Assume that u vanishes on*

$$\partial E \cap K_{2\rho}(x_0) \times (\tau, T).$$

For every $\nu > 0$, there exist positive constants γ_1, γ_2 , and $0 < \bar{s} < 1/2$, depending only on $N, p, \nu, M_{1,1}$, such that for all $\tau + \nu M^{2-p} \rho^p < t < T$, and for all $x \in E \cap K_{2\bar{s}\rho}(x_0)$ with $d(x) < \bar{s}\rho$,

$$\gamma_2 \frac{d(x)}{\rho} \inf_{K_{2\rho}(x_0) \cap E^{\bar{s}\rho} \times (\tau, T)} u \leq u(x, t) \leq \gamma_1 \frac{d(x)}{\rho} \sup_{E \cap K_{2\rho}(x_0) \times (\tau, T)} u.$$

From the above lemma, Theorem 3.8 follows rather easily.

3.3.1. Proof of Theorem 3.8. In order to simplify the proof, we assume $M_{1,1} \approx 1$. Otherwise, the anisotropies of the domains we have to take into account would make the reading particularly burdensome.

(a) The Hopf principle (3.8) is an easy consequence of the linear growth estimate, the interior Harnack inequality, and the regularity up to the boundary of Du [16, Chapters IX and X].

(b) To prove the boundary Harnack inequality, let $(x_0, t_0) \in S_T, \rho \in (0, r_0), x_\rho = (x'_0, 2M_{1,1}\rho), P_\rho = (x_\rho, t_0)$. In Lemma 3.11, let

$$\tau = t_0 - M^{2-p} \rho^p, \quad T = t_0 + M^{2-p} \rho^p, \quad \nu = 1/2,$$

and define

$$V_{1/2,\rho}(x_0, t_0) = \{x \in K_{\bar{s}\rho/4}(x_0) \cap E : d(x) < \bar{s}\rho/8\} \times I(t_0, \rho, \frac{1}{2}M^{2-p}),$$

where \bar{s} is the quantity claimed by Lemma 3.11 when $\nu = 1/2$. Since both u and v satisfy the assumptions of Lemma 3.11, for every $(x, t) \in V_{1/2,\rho}(x_0, t_0)$ we can write

$$\begin{aligned} \gamma_2 d(x) \mu_u(\bar{s}\rho) \leq \rho u(x, t) \leq \gamma_1 d(x) M_u(\rho), \\ \gamma_2 d(x) \mu_v(\bar{s}\rho) \leq \rho v(x, t) \leq \gamma_1 d(x) M_v(\rho), \end{aligned} \tag{3.10}$$

where

$$M_u(\rho) = \sup_{(E \cap K_{\rho/4}(x_0)) \times I(t_0, \rho, M^{2-p})} u, \quad \mu_u(\bar{s}\rho) = \inf_{(E^{\bar{s}\rho/8} \cap K_{\rho/4}(x_0)) \times I(t_0, \rho, M^{2-p})} u,$$

and analogously for v ; (3.10) yields

$$\frac{\gamma_2}{\gamma_1} \frac{\mu_u(\bar{s}\rho)}{M_v(\rho)} \leq \frac{u(x, t)}{v(x, t)} \leq \frac{\gamma_1}{\gamma_2} \frac{M_u(\rho)}{\mu_v(\bar{s}\rho)}, \quad \forall (x, t) \in V_{1/2, \rho}(x_0, t_0). \quad (3.11)$$

Notice that (3.11) holds for every $1 < p < 2$. Restricting p to the range $(2N/(N + 1), 2)$ allows us to apply Corollary 3.2 and Theorem 3.3.

Concerning $M_u(\rho)$, by Corollary 3.2 we have

$$M_u(\rho) \leq \gamma \sup_{\tau \in I(t_0, \rho, 2M^{2-p})} u(x_\rho, \tau).$$

On the other hand, $\mu_u(\bar{s}\rho)$ is attained at some point $(x_*, t_*) \in K_{\rho/4}(x_0) \cap E^{\bar{s}\rho/8} \times I(t_0, \rho, M)$. By the elliptic Harnack inequality of Theorem 3.3,

$$\gamma_3 u(x_\rho, t_*) \leq u(x_*, t_*) \leq \gamma_4 u(x_\rho, t_*),$$

where γ_3, γ_4 depend on $N, p, M_{1,1}$ (as a matter of fact, they depend on \bar{s} too, but once v is fixed, \bar{s} depends only on these quantities), and therefore

$$\mu_u(\bar{s}\rho) \geq \gamma_5 \inf_{\tau \in I(t_0, \rho, 2M^{2-p})} u(x_\rho, \tau),$$

where, once more, γ_5 depends only on $N, p, M_{1,1}$. Combining the previous estimates for u and the analogous ones for v yields

$$\gamma^{-1} \frac{\inf_{\tau \in I(t_0, \rho, 2M^{2-p})} u(x_\rho, \tau)}{\sup_{\tau \in I(t_0, \rho, 2M^{2-p})} v(x_\rho, \tau)} \leq \frac{u(x, t)}{v(x, t)} \leq \gamma \frac{\sup_{\tau \in I(t_0, \rho, 2M^{2-p})} u(x_\rho, \tau)}{\inf_{\tau \in I(t_0, \rho, 2M^{2-p})} v(x_\rho, \tau)}$$

for all $(x, t) \in V_{1/2, \rho}(x_0, t_0)$.

(c) We prove the Hölder continuity of the quotient u/v up to the boundary. First notice that, given $(x, t) \in S_T \cap \Psi_{\rho_0/2, M}(x_0, t_0)$ and denoting by $z_x = x + r\nu_x$ a point along the normal to ∂E at x , we can write

$$\frac{u(z_x, t)}{v(z_x, t)} = \frac{u(z_x, t) - u(x, t)}{v(z_x, t) - v(x, t)} = \frac{\frac{\partial u}{\partial \nu_x}(c_x, t)}{\frac{\partial v}{\partial \nu_x}(c_x, t)}$$

with a suitable c_x . Since u vanishes on the $C^{1,1}$ boundary, the tangential component of the gradient vanishes, and $\frac{\partial u}{\partial \nu_x} \equiv |Du|$. Letting $r \rightarrow 0$, by the Hopf principle and the C^α regularity of Du and Dv (see [16, Chapters IX and X]), we infer that u/v has a Hölder continuous trace on $S_T \cap \Psi_{\rho_0/2, M}(x_0, t_0)$. Moreover, if $(y, s) \in S_T \cap \Psi_{\rho_0/2, M}(x_0, t_0)$ and $(x, t) \in \Psi_{\rho_0/2, M}(x_0, t_0)$, we have, once more by Hopf's principle and the C^α regularity of Du and Dv ,

$$\left| \frac{u(x, t)}{v(x, t)} - \frac{u(y, s)}{v(y, s)} \right| \leq Cd((x, t), (y, s))^\alpha, \quad (3.12)$$

where d denotes the parabolic distance. When both (x, t) and (y, s) belong to $\Psi_{\rho_0/2, M}(x_0, t_0)$, strictly speaking we should distinguish three cases:

- $d((x, t), (y, s)) \simeq d((x, t), S_T) \simeq d((y, s), S_T)$; in that case, the interior Hölder continuity suffices to conclude;
- $d((x, t), (y, s)) \gg d((x, t), S_T), d((x, t), (y, s)) \gg d((y, s), S_T)$; in that case, on the left-hand side of (3.12) we can add and subtract u/v evaluated at the boundary;
- $d((x, t), (y, s)) \simeq d((x, t), S_T), d((x, t), (y, s)) \gg d((y, s), S_T)$; in that case, the desired result is a straightforward consequence of the triangle inequality.

In all three instances, the proof is quite standard and basically relies on (3.12) and on interior Hölder continuity of u/v . Notice that here we can use the classical parabolic distance, as the p -Laplacian is now uniformly elliptic, thanks to the Hopf principle. \square

3.4. Proof of Theorem 3.1

Although the overall strategy of the proof of Theorem 3.1 is the same as of the proof of Theorem 2.1, the singular case requires an adapted renormalization argument to control the vertical size of suitable dyadic cylinders. We mainly concentrate on the differences between the two proofs. The Hölder continuity up to the boundary employed in Corollary 2.3 also holds in the supercritical singular case, and the change of variables introduced in §2.4.1 works for any $p > 1$. Therefore, proving Theorem 3.1 reduces to proving the following lemma.

Lemma 3.12. *Let u be a weak solution to (1.1)–(1.2) such that $0 < u \leq M$ in E_T . Assume that ∂E is flat with respect to x_N and that*

$$I(t_0, 9\rho, M^{2-p}) \subset (0, T]. \quad (3.13)$$

Suppose that u vanishes on

$$\partial E \cap \{|x_i - x_{0,i}| < 2\rho, |x_N| < 8L\rho\} \times I(t_0, 9\rho, M^{2-p}).$$

Then there exists a constant $\gamma > 0$, depending only on p, N, C_0, C_1 , such that

$$u(x, t) \leq \gamma \sup_{\tau \in I(t_0, \rho, 2\eta_\rho^{2-p})} u(x_\rho, \tau)$$

for all $(x, t) \in \{|x_i - x_{0,i}| < \rho/4, 0 < x_N < 2L\rho\} \times I(t_0, \rho, \eta_\rho^{2-p})$.

Proof. Since the previous flattening of the boundary does not affect the value of M , we still have $0 < u \leq M$ in E_T for some $M > 0$. By (3.13) we can define

$$\Psi_\rho^* = K_{2\rho}^*(x_0) \times I(t_0, 9\rho, \eta_\rho^{2-p}) \subset E_T,$$

where $K_{2\rho}^*(x_0)$ has been defined in (2.12), and η_ρ is the first root of the equation

$$\max_{\Psi_\rho^*} u = \eta_\rho.$$

We have already shown that at least one solution of this equation exists. Moreover, we have $\Psi_\rho^* \subset E_T$. The change of variable

$$x \rightarrow \frac{x - x_0}{2L\rho}, \quad t \rightarrow \frac{1}{\eta_\rho^{2-p}} \frac{t - t_0}{\rho^p}$$

maps $K_{2\rho}^*(x_0) \times I(t_0, 9\rho, \eta_\rho^{2-p})$ into $\tilde{Q} = \{|y_i| < 1/L, 0 < y_N < 2\} \times (-9^p, 9^p]$, x_ρ into $y_0 = (0, \dots, 0, 1)$, $\tilde{K}_\rho(x_0)$ into $\tilde{K}_1 = \{|y_i| < 1/(2L), |y_N| < 1\}$, $K_{2\rho}^*(x_0)$ into $K_2^*(y_0) = \{|y_i| < 1/L, 0 < y_N < 2\}$, and the portion $S_T \cap \overline{\Psi_\rho^*}$ of the lateral boundary into

$$\Xi = \{(y', 0) : |y_i| < 1/L\} \times (-9^p, 9^p].$$

After denoting again by (x, t) the transformed variables, and letting $y_0 = (0, \dots, 0, 1)$ in the remainder of the proof, the rescaled function

$$v(x, t) = \frac{1}{\eta_\rho} u(2L\rho x + x_0, t_0 + \eta_\rho^{2-p} \rho^p t)$$

is a non-negative weak solution to

$$\partial_t v - \operatorname{div} \tilde{\mathbf{A}}(x, t, v, Dv) = 0,$$

in \tilde{Q} , where it is easy to see that $\tilde{\mathbf{A}}$ satisfies structure conditions analogous to (1.2), and

$$\forall(x, t) \in \tilde{Q} \quad 0 \leq v \leq 1.$$

As in the proof of Theorem 2.1, in order to simplify the notation and without loss of generality, from here on we assume $L = 1$.

Since the structure conditions have changed, we will denote by ϵ_* and γ_* the corresponding constants of the Harnack inequality, claimed by Theorem 3.3. We will repeatedly apply (3.5)–(3.6); due to all our assumptions, the only condition we need to take into account each time is that $K_{8R}(x_*) \subseteq E$, where $x_* \in \{|x_i| < 1/2, 0 < x_N < 1\}$, and R depends on the context.

The following argument closely resembles the proof given in the elliptic context, the main difference being the need to control the size of the time interval. On the other hand, we cannot simply repeat the argument used in the degenerate case, as it heavily relies on the fact that $p > 2$.

Consider the set $F_0 = \{(x', x_N) : |x_i| < 1/2, x_N = 1\}$, the points $P_{h'}$, whose coordinates are given by

$$\begin{aligned} x_i &= h_i/8, \quad h_i = -3, -2, \dots, 3, \quad i = 1, \dots, N - 1, \\ x_N &= 1, \end{aligned}$$

and the $(N - 1)$ -dimensional cubes $K_{1/4}(P_{h'}) \cap \{x_N = 1\}$. Notice that

- the $(N - 1)$ -dimensional cubes give an equal-size decomposition of F_0 ;
- due to their size and their distance from the boundary, $K_2(P_{h'}) \subseteq E$ and therefore we can apply (3.5)–(3.6).

Consequently, for any

$$t_* \in [-\epsilon_* v^{2-p}(y_0, 0), \epsilon_* v^{2-p}(y_0, 0)], \quad (3.14)$$

by (3.5) we have

$$v(x, t_*) \leq \gamma_* v(P_{h'}, t_*)$$

for all $x \in K_{1/4}(P_{h'}) \cap \{x_N = 1\}$. On the other hand, it is easy to see that

$$v(P_{h'}, t_*) \leq \gamma_* v(y_0, t_*).$$

Therefore, for any $x \in F_0$,

$$v(x, t_*) \leq \gamma_*^4 v(y_0, t_*).$$

Consider the slab

$$S_0 = \{(x', x_N) : |x_i| < 1/2, 7/8 < x_N < 1\}.$$

As noticed above, $K_2(\bar{x}) \subseteq E$ for any $\bar{x} \in F_0$; therefore we can apply (3.5) at time level t_* , and conclude that

$$\forall \bar{x} \in F_0, \forall x \in K_{1/4}(\bar{x}) \quad v(x, t_*) \leq \gamma_* v(\bar{x}, t_*).$$

Consequently, for all $x \in S_0$,

$$v(x, t_*) \leq \gamma_*^5 v(y_0, t_*). \quad (3.15)$$

Estimate (3.15) holds in particular for any $x \in F_1$, where

$$F_1 = \{(x', x_N) : |x_i| < 1/2, x_N = 7/8\}.$$

We can then iterate and conclude that

$$\begin{aligned} \forall x \in S_k &= \{(x', x_N) : |x_i| < 1/2, (7/8)^{k+1} < x_N < (7/8)^k\} \\ v(x, t_*) &\leq \gamma_*^{k+5} v(y_0, t_*). \end{aligned} \quad (3.16)$$

On the other hand, by (3.5)–(3.6), for any t_* as in (3.14),

$$v(y_0, t_*) \leq \gamma_* v(y_0, 0). \quad (3.17)$$

Combining (3.4)–(3.17) finally yields

$$v(x, t) \leq \gamma_*^{k+6} v(y_0, 0)$$

for all $(x, t) \in S_k \times [-\epsilon_* v^{2-p}(y_0, 0), \epsilon_* v^{2-p}(y_0, 0)]$.

For any $\tau \in [-2, 2]$ (and not just for $\tau = 0$) we can repeat the same argument and obtain

$$v(x, t) \leq \gamma_*^{k+6} v(y_0, \tau)$$

for all $(x, t) \in S_k \times [\tau - \epsilon_* v^{2-p}(y_0, \tau), \tau + \epsilon_* v^{2-p}(y_0, \tau)]$ and any $\tau \in [-2, 2]$, provided that $v(y_0, \tau) > 0$. Therefore, setting $\mathcal{M}_2 = \sup_{\tau \in [-2, 2]} v(y_0, \tau)$ yields

$$v(x, t) \leq \gamma_*^{k+6} \mathcal{M}_2 \tag{3.18}$$

for all $(x, t) \in S_k \times [-2, 2]$. This plays the role of (2.16) in the singular framework. Now let

$$\begin{aligned} Q_* &= \{|x_i| < 1/4, 0 < x_N < 1, t \in [-1, 1]\}, \\ Q^* &= \{|x_i| < 1/4, -1 < x_N < 1, t \in [-1, 1]\}, \\ \tilde{Q}_* &= \{|x_i| < 1/2, 0 < x_N < 1, t \in [-2, 2]\}, \\ \tilde{Q}^* &= \{|x_i| < 1/2, -1 < x_N < 1, t \in [-2, 2]\}. \end{aligned}$$

Notice that we need to assume $|x_i| < 1/4$ in the definition of Q_* and Q^* for two closely connected reasons:

- the first point P_1 we are going to choose must lie in a suitably small cylinder, so that the sequence $\{P_m\}$ is all contained in \tilde{Q}_* ;
- as in the degenerate case, we need to choose k_0 that depends only on the data.

We extend u from \tilde{Q}_* to \tilde{Q}^* by odd reflection; by Lemma 2.12, u is still a (signed) solution to (1.1)–(1.2).

Suppose there exists $P_1 = (x_1, t_1) = (x_{1,1}, \dots, x_{1,N}, t_1) \in Q_*$ such that

$$u(P_1) \geq \gamma_*^{k_0+6} \mathcal{M}_2; \tag{3.19}$$

by (3.18), we must have $0 < x_{1,N} < (7/8)^{k_0}$, $|x_i| < 1/4$, $t_1 \in [-1, 1]$. Consider the cylinder

$$Q(P_1) = \{|x - x_1| < 2(7/8)^{k_0+1} \epsilon^{-l}, t_1 - (2(7/8)^{k_0+1} \epsilon^{-l})^p < t \leq t_1\},$$

where ϵ is as in Lemma 2.10, and $l \in \mathbb{N}$ is to be determined. Without loss of generality, we may assume $\epsilon < 7/8$. Now let

$$\omega_0^{(1)} = \operatorname{osc}_{Q(P_1)} v.$$

We do not know the precise value of $\omega_0^{(1)}$, but thanks to (3.19), the construction of u by odd reflection, and the normalization of v , we surely have

$$2\mathcal{M}_2 \gamma_*^{k_0+6} \leq \omega_0^{(1)} \leq 2. \tag{3.20}$$

Provided k_0 is large enough, we have $Q(P_1) \subset \tilde{Q}^*$. Set $\sigma_0^{(1)} = 2(7/8)^{k_0+1} \epsilon^{-l}$ and consider

$$Q_0^{(1)} = K_{\sigma_0^{(1)}}(x_1) \times (t_1 - \theta_0^{(1)} (\sigma_0^{(1)})^p, t_1), \quad \text{where } \theta_0^{(1)} = (\omega_0^{(1)} / A)^{2-p},$$

and A is the quantity denoted by c in Lemma 2.10, which we assume to be larger than 2 without loss of generality. It is apparent that $Q_0^{(1)} \subset Q(P_1) \subset \tilde{Q}^*$. By Lemma 2.10, we can build a sequence

$$\omega_n^{(1)} = \delta \omega_{n-1}^{(1)}, \quad \theta_n^{(1)} = (\omega_n^{(1)}/A)^{2-p}, \quad \sigma_n^{(1)} = \epsilon \sigma_{n-1}^{(1)}, \quad Q_n^{(1)} = Q_{\sigma_n^{(1)}}(\theta_n^{(1)}),$$

for all non-negative integers n . Such a sequence satisfies

$$Q_{n+1}^{(1)} \subset Q_n^{(1)}, \quad \operatorname{osc}_{Q_n^{(1)}} v \leq \omega_n^{(1)}.$$

By iteration

$$\operatorname{osc}_{Q_0^{(1)}} v \leq \delta^n \omega_0^{(1)} = \delta^n \operatorname{osc}_{Q(P_1)} v \Rightarrow \operatorname{osc}_{Q(P_1)} v \geq \frac{1}{\delta^n} \operatorname{osc}_{Q_0^{(1)}} v.$$

If we now choose $n = l$, and l such that $\delta^{-l} > \gamma_*^5$, by the choice of $\sigma_0^{(1)}$ we conclude that

$$\omega_0^{(1)} \geq 2\gamma_*^{k_0+11} \mathcal{M}_2,$$

and this obviously improves the previous lower bound given by (3.20). As v has been built by odd reflection, we conclude there must exist $P(x_2, t_2) = (x_{2,1}, \dots, x_{2,N}, t_2) \in Q(P_1)$ such that

$$v(P_2) \geq \gamma_*^{k_0+11} \mathcal{M}_2.$$

As before, by (3.18), we have $0 < x_{2,N} < (7/8)^{k_0+10}$, $|x_i| < 1/4$, and also $t_1 - (2(7/8)^{k_0+1} \epsilon^{-l})^p < t_2 < t_1$. Set

$$Q(P_2) = \{|x - x_2| < 2(7/8)^{k_0+10} \epsilon^{-l}, t_2 - (2(7/8)^{k_0+10} \epsilon^{-l})^p < t < t_2\}.$$

Once more, provided k_0 is large enough, we can assume that $Q(P_2) \subset \tilde{Q}^*$. From now on, we proceed as in § 2.4.3. By induction, we obtain a sequence $\{P_m = (x_m, t_m)\}$ such that

$$v(P_m) \geq \gamma_*^{k_0+6+5(m-1)} \mathcal{M}_2 \quad (3.21)$$

and

$$0 < x_{m,N} < (7/8)^{k_0+1+5(m-1)}.$$

Provided we choose k_0 large enough, the sequence $\{P_m\}$ is all contained in the fixed cylinder \tilde{Q}_* ; since $0 < v < 1$ and $\mathcal{M}_2 \in (0, 1]$ is a fixed quantity, (3.21) eventually leads to a contradiction. Therefore, there exists $\tilde{\gamma}$, which depends only on the data, such that

$$\forall (x, t) \in Q_* \quad v(x, t) \leq \tilde{\gamma} \mathcal{M}_2. \quad (3.22)$$

We now switch back to the original variables. We conclude that there exists a constant $\tilde{\gamma}$, depending only on the data, such that

$$\begin{aligned} \forall (x, t) \in \{|x_i - x_{0,i}| < \rho/4, 0 < x_N < 2L\rho\} \times I(t_0, \rho, \eta_\rho^{2-p}) \\ u(x, t) \leq \tilde{\gamma} \sup_{\tau \in I(t_0, \rho, 2\eta_\rho^{2-p})} u(x_\rho, \tau). \end{aligned} \quad \square$$

3.4.1. *Proof of Corollary 3.2.* A close inspection of the previous proof shows that all the arguments continue to hold true if we replace the qualitative parameter η_ρ directly with M . \square

3.5. *Proof of Corollary 3.5*

Taking into account the notation of §3.4, it is enough to prove the following result.

Lemma 3.13. *Let u be a weak solution to (1.1)–(1.2) such that $0 < u \leq M$ in E_T . Assume that ∂E is flat with respect to x_N and that $I(t_0, 9\rho, M^{2-p}) \subset (0, T]$. Suppose that u vanishes on*

$$\partial E \cap \{|x_i - x_{0,i}| < 2\rho, |x_N| < 8L\rho\} \times I(t_0, 9\rho, M^{2-p}).$$

Then there exists a constant $\gamma > 0$, depending on $p, N, C_0, C_1, M/m_0$, such that

$$u(x, t) \leq \gamma u(P_{\rho_k})$$

for all $(x, t) \in \Psi_{\rho_k, M}(x_0, t_0)$ and $k = 0, 1, \dots$

Proof. We go back to the proof of Lemma 3.12, at (3.22) with η_ρ replaced by M . Let $(x_0, t_0) \in S_T$ and assume $\text{dist}(x, \partial E) = \rho$ (we now take $L = 1/2$, without loss of generality, in order to simplify the notation), and $I(t_0, 9\rho, M^{2-p}) \subset (0, T]$. We let

$$\begin{aligned} I_k &= (t_0, \sigma_k, 2M^{2-p}), & M_k &= \sup_{\tau \in I_k} u(x_{\rho_k}, \tau), \\ m_k &= \inf_{\tau \in I_k} u(x_k, \tau), & u_k &= u(x_{\rho_k}, t_0). \end{aligned}$$

By the weak form of the Carleson estimate given in Corollary 3.2, for all (x, t) in $\Psi_{\rho, M}(x_0, t_0)$ we have

$$u(x, t) \leq \gamma \sup_{\tau \in I_0} u(x_\rho, \tau),$$

which implies

$$u(x, t) \leq \gamma \frac{M_0}{m_0} u(P_\rho) \leq \gamma \frac{M}{m_0} u(P_\rho).$$

Analogously, working in a smaller box, for all $(x, t) \in \Psi_{\rho_k, M}(x_0, t_0)$,

$$u(x, t) \leq \gamma \sup_{\tau \in I_k} u(x_{\rho_k}, \tau) \leq \gamma \frac{M_k}{m_k} u(P_{\rho_k}).$$

To prove the lemma, we will show by induction that

$$M_k/m_k \leq \gamma^{2(N_0+1)} \quad \text{and} \quad N_k \leq N_0 \tag{3.23}$$

where

$$N_0 = 2 \left(\frac{M}{m_0} \right)^{2-p} \frac{8^p}{\bar{\epsilon}}, \quad N_k = 2 \left(\frac{M}{\gamma^k m_k} \right)^{2-p} \frac{8^p}{\bar{\epsilon}} \quad \text{for } k = 1, 2, \dots$$

Let us now consider $k = 1$. In order to cover the segment $\{x_\rho\} \times [t_0, t_0 + 2M^{2-p}\rho^p]$ (and the same can be said for $\{x_\rho\} \times [t_0 - 2M^{2-p}\rho^p, t_0]$), we need at most N_0 steps, where N_0 is given by

$$t_0 + \bar{\epsilon} N_0 m_0^{2-p} (\rho/8)^p = t_0 + 2M^{2-p}\rho^p,$$

which yields the above formula for N_0 . Without loss of generality, we can assume that $N_0 \in \mathbb{N}$, possibly by a slight modification of $\bar{\epsilon}$.

Now consider x_{ρ_1} ; by the elliptic Harnack inequality, for all $t \in I_1$ we have

$$\begin{aligned} u(x_{\rho_1}, t) \geq \gamma^{-1} u(x_\rho, t) &\Rightarrow m_1 \geq \gamma^{-1} m_0 \Rightarrow \frac{1}{\gamma m_1} \leq \frac{1}{m_0}, \\ u(P_{\rho_1}) \geq \gamma^{-1} u(P_\rho) &\Rightarrow u(P_{\rho_1}) \geq \gamma^{-1} m_0. \end{aligned}$$

Taking once more into account the Harnack inequality, we obtain

$$\begin{aligned} M_1 &\leq \sup_{t \in I_0} u(x_{\rho_1}, t) \leq \gamma^{N_0+1} u(P_\rho), \\ m_1 &\geq \inf_{t \in I_0} u(x_{\rho_1}, t) \geq \gamma^{-(N_0+1)} u(P_\rho), \end{aligned}$$

which yields

$$M_1/m_1 \leq \gamma^{2(N_0+1)}.$$

Moreover, to cover $\{x_{\rho_1}\} \times [t_0, t_0 + 2M^{2-p}\sigma_1^p]$, we need to take at most N_1 steps, where N_1 is given by

$$t_0 + \bar{\epsilon} N_1 m_1^{2-p} \left(\frac{7}{8^2}\rho\right)^p = t_0 + 2M^{2-p} \left(\frac{7}{8\gamma^{(2-p)/p}}\rho\right)^p,$$

which yields

$$N_1 = 2 \left(\frac{M}{\gamma m_1}\right)^{2-p} \frac{8^p}{\bar{\epsilon}} \leq 2 \left(\frac{M}{m_0}\right)^{2-p} \frac{8^p}{\bar{\epsilon}},$$

and therefore $N_1 \leq N_0$.

Let us now assume that $M_k/m_k \leq \gamma^{2(N_0+1)}$ and $N_k \leq N_0$ for some $k = 1, 2, \dots$. Now consider $x_{\rho_{k+1}} \in E$; by the elliptic Harnack inequality, for all $t \in I_{k+1}$ we have

$$\begin{aligned} u(x_{\rho_{k+1}}, t) \geq \gamma^{-1} u(x_{\rho_k}, t) &\Rightarrow m_{k+1} \geq \gamma^{-1} m_k, \\ u(P_{\rho_{k+1}}) \geq \gamma^{-1} u(P_{\rho_k}) &\Rightarrow u(P_{\rho_{k+1}}) \geq \gamma^{-1} m_k. \end{aligned}$$

Taking once more into account the Harnack inequality, we obtain

$$\begin{aligned} M_{k+1} &\leq \sup_{t \in I_k} u(x_{\rho_{k+1}}, t) \leq \gamma^{N_k+1} u(P_{\rho_k}), \\ m_{k+1} &\geq \inf_{t \in I_k} u(x_{\rho_{k+1}}, t) \geq \gamma^{-(N_k+1)} u(P_{\rho_k}), \end{aligned}$$

which yields

$$M_{k+1}/m_{k+1} \leq \gamma^{2(N_k+1)} \leq \gamma^{2(N_0+1)},$$

and the ratio has not grown.

Moreover, to cover $\{x_{\rho_{k+1}}\} \times [t_0, t_0 + 2M^{2-p}\sigma_{k+1}^p]$, we need to take at most N_{k+1} steps, where N_{k+1} is given by

$$t_0 + \bar{\epsilon} N_{k+1} m_{k+1}^{2-p} \left(\frac{\rho_k}{8} \rho\right)^p = t_0 + 2M^{2-p}(\sigma_k)^p,$$

which yields

$$N_{k+1} = 2 \left(\frac{M}{\gamma^{k+1} m_{k+1}} \right)^{2-p} \frac{8^p}{\bar{\epsilon}} \leq N_k.$$

By the induction principle we now have (3.23). Hence

$$M_k/m_k \leq \gamma^{2(N_0+1)} \quad \forall k = 1, 2, \dots,$$

where N_0 depends only on $p, N, C_0, C_1, M/m_0$. \square

4. The porous medium equation

Consider quasi-linear parabolic partial differential equations of the form

$$u_t - \operatorname{div} \mathbf{A}(x, t, u, Du) = 0 \quad \text{weakly in } E_T, \quad (4.1)$$

where the function $\mathbf{A} : E_T \times \mathbb{R}^{N+1} \rightarrow \mathbb{R}^N$ is only assumed to be measurable and subject to the structure conditions

$$\begin{cases} \mathbf{A}(x, t, u, \xi) \cdot \xi \geq m C_0 |u|^{m-1} |\xi|^2 \\ |\mathbf{A}(x, t, u, \xi)| \leq m C_1 |u|^{m-1} |\xi| \end{cases} \quad \text{a.e. } (x, t) \in E_T, \forall u \in \mathbb{R}, \forall \xi \in \mathbb{R}^N, \quad (4.2)$$

where C_0 and C_1 are given positive constants, and $m > 0$. The prototype of this class of parabolic equations is

$$u_t - \Delta(|u|^{m-1}u) = 0 \quad \text{weakly in } E_T. \quad (4.1)_0$$

For simplicity, we limit ourselves to the definition of non-negative solutions: a non-negative function

$$u \in C([0, T]; L^2(E)) \text{ with } \begin{cases} u^{(m+1)/2} \in L^2(0, T; W^{1,2}(E)) & \text{if } m > 1, \\ u^m \in L^2(0, T; W^{1,2}(E)) & \text{if } 0 < m < 1 \end{cases}$$

is a weak sub[super]solution to (4.1)–(4.2) if for every subinterval $[t_1, t_2] \subset (0, T)$,

$$\int_E u \varphi dx \Big|_{t_1}^{t_2} + \int_{t_1}^{t_2} \int_E [-u \varphi_t + \mathbf{A}(x, t, u, Du) \cdot D\varphi] dx dt \leq [\geq] 0 \quad (4.3)$$

for all non-negative test functions

$$\varphi \in W^{1,2}(0, T; L^2(E)) \cap L^2(0, T; W_0^{1,2}(E)).$$

This guarantees that all the integrals in (4.3) are convergent.

In (4.3), the symbol Du has to be understood in the sense of the following definition:

$$Du = \begin{cases} \frac{2}{m+1} \mathbf{1}_{\{u>0\}} u^{(1-m)/2} Du^{(m+1)/2} & \text{if } m > 1, \\ \frac{1}{m} \mathbf{1}_{\{u>0\}} u^{1-m} Du^m & \text{if } 0 < m < 1. \end{cases}$$

The parameters m, N, C_0, C_1 are the data; the partial differential equation (4.1)–(4.2) is degenerate when $m > 1$, and singular when $0 < m < 1$, since its modulus of ellipticity $|u|^{m-1}$ tends to 0 or to $+\infty$, respectively, as $|u| \rightarrow 0$. In the latter case, we further distinguish between the singular supercritical range (when $(N-2)_+/N < m < 1$), and singular critical and subcritical range (when $0 < m \leq (N-2)_+/N$).

We are interested in solutions to (4.1)–(4.2) *continuously* vanishing on some distinguished part of the lateral boundary S_T of a space-time cylinder, and we aim at extending to such solutions the results proved in the previous sections for solutions to (1.1)–(1.2). We will not give the full proofs, as most of the arguments can be reproduced almost verbatim. We will limit ourselves to discuss the changes that need to be done.

When dealing with this kind of problems, the cylinders to be considered are the ones already defined in §1, provided the height scales as $\theta\rho^2$ with $\theta = (c/u(x_0, t_0))^{m-1}$ and c a suitable positive parameter.

Intrinsic Harnack inequalities, of the same kind as considered in Theorems 2.6 and 3.3, can be proved respectively for $m > 1$ and $(N-2)_+/N < m < 1$. For the exact statements we refer to [17, 20, 21]. For a general introduction to the porous medium equation, see [15, 50].

4.1. The degenerate case $m > 1$

For a Lipschitz cylinder E_T with Lipschitz constant L , let $(x_0, t_0), \rho, P_\rho$ be as in §2.1, assume $u(P_\rho) > 0$, and set

$$\begin{aligned} \Psi_\rho^-(x_0, t_0) \\ = E_T \cap \{(x, t) : |x_i - x_{0,i}| < \rho/4, |x_N| < 2L\rho, t \in (t_0 - \frac{\alpha+\beta}{2}\theta\rho^2, t_0 - \beta\theta\rho^2)\}, \end{aligned}$$

where $\theta = [c/u(P_\rho)]^{m-1}$ with c given by the Harnack inequality for $m > 1$ (see [21, Chapter 5]), and $\alpha > \beta$ are positive parameters.

Theorem 4.1 (Carleson estimate, $m > 1$). *Let u be a non-negative weak solution to (4.1)–(4.2) in E_T . Assume that*

$$(t_0 - \theta(4\rho)^2, t_0 + \theta(4\rho)^2) \subset (0, T]$$

and that u vanishes continuously on

$$\partial E \cap \{|x_i - x_{0,i}| < 2\rho, |x_N| < 8L\rho\} \times (t_0 - \theta(4\rho)^2, t_0 + \theta(4\rho)^2).$$

Then there exist positive parameters $\alpha > \beta$ and a constant $\gamma > 0$, depending only on m, N, C_0, C_1, L , such that

$$u(x, t) \leq \gamma u(P_\rho) \quad \text{for every } (x, t) \in \Psi_\rho^-(x_0, t_0). \quad (4.4)$$

Proof. As in the proof of Theorem 2.1, the flattening of the boundary leaves the structure of the equation unchanged; relabeling the variables as before, we end up with

$$\begin{cases} \mathbf{A}(x, t, u, \xi) \cdot \xi \geq m C_0^* |u|^{m-1} |\xi|^2 \\ |\mathbf{A}(x, t, u, \xi)| \leq m C_1^* |u|^{m-1} |\xi| \end{cases} \quad \text{a.e. } (x, t) \in E_T, \forall u \in \mathbb{R}, \forall \xi \in \mathbb{R}^N,$$

where C_0^* and C_1^* are positive constants that depend only on C_0, C_1 , and the Lipschitz constant L . The remaining part of the proof proceeds as in Theorem 2.1. \square

Remark 4.2. As for the p -Laplacian, let us point out that also for the prototype equation (4.1)₀, estimate (4.4) could be extended from Lipschitz cylinders to a wider class of cylinders E_T whose cross section E is an N.T.A. domain.

Weak solutions to (4.1) with zero Dirichlet boundary conditions on a Lipschitz domain are Hölder continuous up to the boundary (see, for example, [16, Chapter III, Appendix]). Combining this result with the previous Carleson estimate yields a quantitative estimate on the decay of u at the boundary, invariant by the intrinsic rescaling

$$x = x_0 + \rho y, \quad t = t_0 + \frac{\rho^2}{u(P_\rho)^{m-1}} \tau.$$

Corollary 4.3. *Under the same assumptions of Theorem 4.1 we have*

$$u(x, t) \leq \gamma \left(\frac{\text{dist}(x, \partial E)}{\rho} \right)^\mu u(P_\rho)$$

for every $(x, t) \in \Psi_{\rho/2}^-(x_0, t_0)$, where $\gamma > 0$ and $\mu \in (0, 1)$ depend only on m, N, C_0, C_1, L .

If we restrict our attention to solutions to the model equation (4.1)₀ and to C^2 cylinders, the result of Corollary 4.3 can be strengthened.

Theorem 4.4 (A sharper decay). *Let E_T be a C^2 cylinder, and u a non-negative weak solution to (4.1)₀ in E_T . Assume the other assumptions of Theorem 4.1 hold. Then there exist positive parameters $\alpha > \beta$ and a constant $\gamma > 0$, depending only on m, N , and the C^2 -constant M_2 of E , such that for every (x, t) in*

$$E_T \cap \{|x_i - x_{0,i}| < \rho/4, 0 < x_N < 2M_2\rho\} \times \left(t_0 - \frac{\alpha + 3\beta}{4} \theta \rho^2, t_0 - \beta \theta \rho^2 \right]$$

we have

$$0 \leq u(x, t) \leq \gamma \left(\frac{\text{dist}(x, \partial E)}{\rho} \right)^{1/m} u(P_\rho), \tag{4.5}$$

Proof. In this context the barrier is

$$\Theta_k(x, t) = C u(P_\rho) (1 - \eta_k(x, t))^{1/m},$$

where

$$\eta_k(x, t) = \exp(-k(|x - y| - 1)) \exp(u(P_\rho)^{m-1}(t - t_1)),$$

and the constant C is chosen so that $u \leq \Theta_k$ on the parabolic boundary of \mathcal{N}_k . The remainder of the proof is as in Theorem 2.4; see also [22, Theorem 4.1]. \square

Notice that, in general, the bound below by zero in (4.5) cannot be improved. Indeed, when $m > 1$, two explicit solutions to (4.1)₀ in the half-space $\{x_N \geq 0\}$ that vanish at $x_N = 0$ are given by

$$u_1(x, t) = x_N^{1/m}, \quad u_2(x, t) = \left(\frac{m-1}{m+1}\right)^{1/(m-1)} (T-t)^{-1/(m-1)} x_N^{2/(m-1)}.$$

As in (2.6), we can even have solutions that vanish on a set of positive measure. For $\gamma \in (0, 1)$, consider, for example,

$$u(x, t) = \left[\frac{m-1}{m} \gamma(t+1) \left(\gamma + \frac{x_N - 2}{t+1} \right)_{+} \right]^{1/(m-1)},$$

which solves (4.1)₀ in $\{x_N > 0\} \times (0, 2/\gamma - 1]$.

Finally, all the remarks in §2.2 about the obstruction to the definition of a useful Harnack chain hold for the porous medium equation as well, without any significant change.

4.2. The singular supercritical case $(N - 2)_+/N < m < 1$

Let $E_T, u, (x_0, t_0), \rho, P_\rho$ be as in Theorem 4.1, and set

$$\tilde{I}(t_0, \rho, h) = (t_0 - h\rho^2, t_0 + h\rho^2).$$

Moreover, let u be a weak solution to (4.1)–(4.2) such that

$$0 < u \leq M \quad \text{in } E_T, \tag{4.6}$$

and assume that

$$\tilde{I}(t_0, 9\rho, M^{1-m}) \subset (0, T]. \tag{4.7}$$

Then we define

$$\begin{aligned} \tilde{\Psi}_\rho(x_0, t_0) &= E_T \cap \{(x, t) : |x_i - x_{0,i}| < 2\rho, |x_N| < 4L\rho, t \in \tilde{I}(t_0, 9\rho, \eta_\rho^{1-m})\}, \\ \tilde{\Psi}_\rho(x_0, t_0) &= E_T \cap \{(x, t) : |x_i - x_{0,i}| < \rho/4, |x_N| < 2L\rho, t \in \tilde{I}(t_0, \rho, \eta_\rho^{1-m})\}, \end{aligned}$$

where η_ρ is the first root of the equation

$$\max_{\tilde{\Psi}_\rho(x_0, t_0)} u = \eta_\rho. \tag{4.8}$$

As in §3, it is immediate that at least one root of (4.8) exists. Moreover, by (4.7), $\tilde{\Psi}_\rho(x_0, t_0) \subset E_T$. Finally, for $k = 0, 1, \dots$ let

$$\begin{aligned} \rho_k &= (7/8)^k \rho, \quad \sigma_k = \rho_k / \gamma^{k(1-m)/2}, \\ x_{\rho_k} &= (x'_0, 2L\rho_k), \quad P_{\rho_k} = (x'_0, 2L\rho_k, t_0), \\ \Psi_{\rho, M}(x_0, t_0) &= E_T \cap \{(x, t) : |x_i - x_{0,i}| < \rho/4, |x_N| < 2L\rho, t \in \tilde{I}(t_0, \rho, M^{1-m})\}, \\ \Psi_{\rho_k, M}(x_0, t_0) &= E_T \cap \{(x, t) : |x_i - x_{0,i}| < \rho_k/4, |x_N| < 2L\rho_k, t \in \tilde{I}(t_0, \sigma_k, M^{1-m})\}, \\ m_0 &= \inf_{\tau \in \tilde{I}(t_0, \rho, 2M^{1-m})} u(x_\rho, \tau), \quad M_0 = \sup_{\tau \in \tilde{I}(t_0, \rho, 2M^{1-m})} u(x_\rho, \tau). \end{aligned}$$

In the following statement, we give both the weak and the strong form of the Carleson estimates for solutions to a singular, supercritical porous medium equation.

Theorem 4.5 (Carleson estimate, $(N - 2)_+/N < m < 1$). *Let u be a weak solution to (4.1)–(4.2) that satisfies (4.6). Assume that (4.7) holds true and u vanishes continuously on*

$$\partial E \cap \{|x_i - x_{0,i}| < 2\rho, |x_N| < 8L\rho\} \times \tilde{I}(t_0, 9\rho, M^{1-m}).$$

Then there exist constants $\gamma > 0$ and $\alpha \in (0, 1)$, depending only on m, N, C_0, C_1, L , such that

$$u(x, t) \leq \gamma \left(\frac{\text{dist}(x, \partial E)}{\rho} \right)^\alpha \sup_{\tau \in I(t_0, \rho, 2M^{1-m})} u(x_\rho, \tau)$$

for every $(x, t) \in \tilde{\Psi}_\rho(x_0, t_0)$. Moreover, for every $(x, t) \in \Psi_{\rho, M}(x_0, t_0)$,

$$u(x, t) \leq \gamma \left(\frac{\text{dist}(x, \partial E)}{\rho} \right)^\alpha \sup_{\tau \in I(t_0, \rho, 2M^{1-m})} u(x_\rho, \tau).$$

Finally, there exists a constant $\hat{\gamma}$, depending on $m, N, C_0, C_1, L, M/m_0$, such that

$$u(x, t) \leq \hat{\gamma} u(P_{\rho_k})$$

for every $(x, t) \in \Psi_{\rho_k}(x_0, t_0)$ and all $k = 0, 1, \dots$.

Another difference from the degenerate case appears when we consider $C^{1,1}$ cylinders and (mainly for simplicity) the prototype equation (4.1)₀. In this case, indeed, weak solutions vanishing on the lateral part enjoy a proper power-like behavior at the boundary. As pointed out in §3.3, the following lemma was originally proved in [22, §4] for solutions in C^2 domains; it can be extended to solutions in $C^{1,1}$ domains, working as in [39, Lemma 3.1] for the parabolic p -Laplacian. The notation is as in Lemma 3.11.

Lemma 4.6. *Let $\tau \in (0, T)$ and fix $x_0 \in \partial E$. Assume that u vanishes on $\partial E \cap K_{2\rho}(x_0) \times (\tau, T)$. For every $\nu > 0$, there exist positive constants γ_1, γ_2 and $0 < \bar{s} < 1/2$, depending only on $N, m, \nu, M_{1,1}$, such that for all $\tau + \nu M^{1-m} \rho^2 < t < T$, and all $x \in E \cap K_{2\bar{s}\rho}(x_0)$ with $d(x) < \bar{s}\rho$,*

$$\gamma_2 \frac{d(x)}{\rho} \inf_{(E^{\bar{s}\rho} \cap K_{2\rho}(x_0)) \times (\tau, T)} u \leq u(x, t)^m \leq \gamma_1 \frac{d(x)}{\rho} \sup_{(E \cap K_{2\rho}(x_0)) \times (\tau, T)} u.$$

The implications of Lemma 4.6 are stated in the following result.

Theorem 4.7. *Let $(N - 2)_+/N < m < 1$. Assume E_T is a $C^{1,1}$ cylinder, and $(x_0, t_0), \rho, P_\rho$ are as in Theorem 4.1. Let u, v be weak solutions to (4.1)₀ in E_T satisfying the hypotheses of Theorem 4.5, that is, $0 < u, v \leq M$ in E_T . Then there exist positive constants $\bar{s}, \gamma, \beta, 0 < \beta \leq 1$, depending only on $N, m, M_{1,1}$, and $\rho_0, c_0 > 0$, depending also on the oscillation of u , such that the following properties hold.*

(a) Hopf principle:

$$|Du^m| \geq c_0 \quad \text{in } \Psi_{\rho_0, M}(x_0, t_0).$$

(b) Boundary Harnack inequality:

$$\gamma^{-1} \frac{\inf_{\tau \in \tilde{I}(t_0, \rho, 2M^{1-m})} u(x_\rho, \tau)}{\sup_{\tau \in \tilde{I}(t_0, \rho, 2M^{1-m})} v(x_\rho, \tau)} \leq \frac{u(x, t)}{v(x, t)} \leq \gamma \frac{\sup_{\tau \in \tilde{I}(t_0, \rho, 2M^{1-m})} u(x_\rho, \tau)}{\inf_{\tau \in \tilde{I}(t_0, \rho, 2M^{1-m})} v(x_\rho, \tau)} \quad (4.9)$$

for all $(x, t) \in \{x \in K_{\bar{s}\rho/4} \cap E : d(x) < \bar{s}\rho/8\} \times \tilde{I}(t_0, \rho, \frac{1}{2}M^{1-m})$ with $\rho < \rho_0$.

Remark 4.8. Proceeding as in Remark 3.9, the boundary Harnack inequality (4.9) can be rewritten as

$$\gamma_*^{-1} \frac{u(P_\rho)}{v(P_\rho)} \leq \frac{u(x, t)}{v(x, t)} \leq \gamma_* \frac{u(P_\rho)}{v(P_\rho)},$$

where γ_* depends not only on $N, p, M_{1,1}$, but also on $M_{0,u}/m_{0,u}$ and $M_{0,v}/m_{0,v}$.

Remark 4.9. As in Theorem 3.8, the Hopf principle holds true only in a small neighborhood of the boundary.

Remark 4.10. Since we do not have Hölder regularity estimates for the gradient of u , we cannot proceed as in Theorem 3.8 to prove the Hölder continuity of the ratio.

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