

The exponent in the orthogonality catastrophe for Fermi gases

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Dedicated to Hajo Leschke on the occasion of his 70th birthday.

Abstract. We quantify the asymptotic vanishing of the ground-state overlap of two non-interacting Fermi gases in d -dimensional Euclidean space in the thermodynamic limit. Given two one-particle Schrödinger operators in finite-volume which differ by a compactly supported bounded potential, we prove a power-law upper bound on the ground-state overlap of the corresponding non-interacting N -Fermion systems. We interpret the decay exponent γ in terms of scattering theory and find $\gamma = \pi^{-2} \|\arcsin |T_E/2|\|_{\text{HS}}^2$, where T_E is the transition matrix at the Fermi energy E . This exponent reduces to the one predicted by Anderson [Phys. Rev. **164**, 352–359 (1967)] for the exact asymptotics in the special case of a repulsive point-like perturbation.

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1. Introduction

We consider two quantum systems, each consisting of N non-interacting Fermions in a box of side length L in d -dimensional Euclidean space \mathbb{R}^d , with $d \in \mathbb{N}$. The single-particle Hamiltonians of the two systems differ by a local perturbation potential V . As a signature of inequivalent representations of the canonical commutation relations, the overlap $\langle \Phi_L^N, \Psi_L^N \rangle$ of the N -Fermion ground states Φ_L^N and Ψ_L^N must vanish in the thermodynamic limit $L \rightarrow \infty$, $N \rightarrow \infty$, $N/L^d \rightarrow \text{const.} > 0$, see [8, Chapter IV] and [13, Chapter II.1.1]. A quantitative version of this behaviour in terms of a power law

$$|\langle \Phi_L^N, \Psi_L^N \rangle|^2 \sim L^{-\gamma} \quad (1)$$

was predicted by P. W. Anderson in 1967. In [1] he presented a brief computation for the case of a point-like perturbation V in $d = 3$ dimensions and arrived at the upper bound

$$|\langle \Phi_L^N, \Psi_L^N \rangle|^2 \leq L^{-\gamma_1} \quad (2)$$

with

$$\gamma_1 = \pi^{-2}(\sin \delta)^2. \quad (3)$$

Here, δ is the (single-particle) scattering phase shift caused by the point interaction at the Fermi energy. Nowadays, this behaviour is often referred to as Anderson's orthogonality catastrophe in the physics literature. A mathematical proof for a generalisation of (2) and (3) was given recently in [11]. Allowing for a bounded, compactly supported, non-negative perturbation V in \mathbb{R}^d , it is shown there that (2) holds with

$$\gamma_1 = \frac{1}{\pi^2} \|T_E/2\|_{\text{HS}}^2, \quad (4)$$

where T_E denotes the transition matrix of scattering theory and $\|\cdot\|_{\text{HS}}$ the Hilbert-Schmidt norm for operators on the Hilbert space of the energy shell corresponding to the Fermi energy E . In the special case considered in [1], (4) reduces to (3). The principal strategy of the argument in [11] is to rewrite the overlap determinant as $|\langle \Phi_L^N, \Psi_L^N \rangle|^2 = \det A = \exp(\text{tr} \ln A)$ and to expand the logarithm in a series of non-negative terms

$$|\langle \Phi_L^N, \Psi_L^N \rangle|^2 = \exp \left\{ - \sum_{n \in \mathbb{N}} \frac{1}{n} \text{tr}((I - A)^n) \right\}, \quad (5)$$

see Lemma 3.1 below. A similar idea was used by M. Kac [16] in his proof of the Szegő limit theorem for Toeplitz determinants which is, in a way, an analogue to (1).

By dropping all but the first term $\text{tr}(I - A)$ of the series, which is called *Anderson integral* in the physics literature, one arrives at an upper bound. The main work of [11] consists in deriving a lower bound of the form $\text{tr}(I - A) \geq \gamma_1 \ln L$ for the Anderson integral with γ_1 given by (4). There are only few other mathematically rigorous works on Anderson’s orthogonality catastrophe [20, 9, 18, 10]. It is shown in [20] that (4) in fact provides the exact coefficient in the asymptotics $\text{tr}(I - A) \sim \gamma_1 \ln L$ of the Anderson integral in the thermodynamic limit for one-dimensional systems. We refer to [20, 11] and references therein for a brief description of the relevance of the orthogonality catastrophe in physics and for a discussion of the theoretical approaches in the physics literature.

In a second paper [2] in 1967, P. W. Anderson notes as an aside that the true asymptotics (1) of the overlap involves an exponent γ for which “... the main difference from the previous result [i.e. (3)] is to replace $(\sin \delta)^2$ by δ^2 .” After some controversies about the correctness of interchanging limits [26, 14], Anderson’s result (1) was confirmed in the case of a point interaction V with the decay exponent

$$\gamma = \pi^{-2} \delta^2$$

by theoretical-physics methods [14]. A mathematical proof was given recently in [10]. For reasons of comparison, we remark that the particle number N in [14] refers to the number of s -orbital states below the Fermi energy and thus $N \sim L$. Related results in the context of the Kondo problem in the physics literature can be found in [23, 34].

The purpose of the present paper is a mathematical contribution towards the exact asymptotics (1). We will prove in Theorem 2.2 that, in the presence of a rather general background potential V_0 , a bounded, compactly supported, non-negative perturbation potential V in \mathbb{R}^d causes the power-law decay

$$|\langle \Phi_L^N, \Psi_L^N \rangle|^2 \leq L^{-\gamma + o(L^0)}$$

of the overlap for almost every Fermi energy $E \in \mathbb{R}$ along subsequences $L \rightarrow \infty$. The decay exponent is given by

$$\gamma = \frac{1}{\pi^2} \|\arcsin|T_E/2|\|_{\text{HS}}^2. \tag{6}$$

We refer to Theorem 2.2 for the precise statement. In proving (6), we obtain a result on the trace of a product of spectral projections of two Schrödinger operators which may be interesting by itself, see Theorem 3.4.

Clearly, when comparing (6) to (4), we infer $\gamma_1 \leq \gamma$, and the two exponents are related in the spirit of Anderson's rule quoted above. In view of [10] and of the physicists' results, we conjecture that the exponent γ governs the true asymptotics (1) of the overlap whenever the modulus of the (appropriately defined) scattering phases does not exceed $\pi/2$.

The proof of Theorem 2.2 relies on the representation (5) of the overlap. We determine the dominant behaviour of each term in the n -sum in (5), because each term contributes to the asymptotics. In order to treat the terms with $n > 1$ we have to deal with additional issues. One is the non-positivity of certain trace expressions, another one is to compute the multi-dimensional integral

$$\int_{(0,\infty)^{2n}} du_1 \dots du_{2n} \frac{e^{-(u_1 + \dots + u_{2n})}}{(u_1 + u_2) \dots (u_{2n-1} + u_{2n})}, \quad (7)$$

which contributes to the asymptotics of the n th term in (5). Subsequently, the values of these integrals show up in the Taylor expansion of the function $x \mapsto (\arcsin x)^2$. We compute the integral (7) in Section 4.5 by identifying it with the first diagonal matrix element of the $(2n - 1)$ th power of the Hilbert matrix.

Since V causes scattering, the exponent γ is typically expected to be strictly positive. In the appendix, we prove this in the case without a background potential.

After we completed this paper, Frank and Pushnitski [7] established results on the asymptotics for traces of regularised projections of infinite-volume operators. Their work is partly a generalisation of our analysis in Sections 4.3 to 4.5. In particular, their consequent use of Hankel operators is conceptually valuable and leads to a simplification of proofs. From this point of view it is also less surprising that (a unitary equivalent operator to) the Hilbert matrix appears in our Section 4.5 when we compute the multi-dimensional integral (7).

2. Setup and main result

Let $d \in \mathbb{N}$, $\Lambda_1 \subseteq \mathbb{R}^d$ be open and bounded with $0 \in \Lambda_1$ and for $L > 1$, define $\Lambda_L := L \cdot \Lambda_1$.

Let the negative Laplacian $-\Delta_L$ be supplied with Dirichlet boundary conditions on Λ_L . We define two multiplication operators V_0 and V acting on $L^2(\Lambda_L)$, corresponding to real-valued functions on \mathbb{R}^d with the properties

$$\begin{aligned} \max\{V_0, 0\} \in K_{\text{loc}}(\mathbb{R}^d), \quad \max\{-V_0, 0\} \in K(\mathbb{R}^d), \\ V \in L^\infty(\mathbb{R}^d), \quad V \geq 0, \quad \text{supp } V \subseteq \Lambda_1 \text{ compact.} \end{aligned} \quad (V)$$

Here, we have written $K(\mathbb{R}^d)$ and $K_{\text{loc}}(\mathbb{R}^d)$ for the Kato class and the local Kato class, respectively, see [29]. The *finite-volume one-particle Schrödinger operators* $H_L := -\Delta_L + V_0$ and $H'_L := H_L + V$ are self-adjoint and densely defined in the Hilbert space $L^2(\Lambda_L)$. The *infinite-volume operators* $H := -\Delta + V_0$ and $H' := H + V$ are self-adjoint and densely defined in the Hilbert space $L^2(\mathbb{R}^d)$. Birman's theorem, see [4, Theorem 2] or [25, Theorem XI.10], is applicable by virtue of [29, Theorem B.9.1] and guarantees the existence and completeness of the wave operators for the pair H, H' . In particular, their absolutely continuous spectra are the same, i.e.

$$\sigma_{\text{ac}}(H) = \sigma_{\text{ac}}(H').$$

The assumptions (V) on V_0 and V , together with [5, Theorem 6.1], imply that the semigroup operators e^{-tH_L} and $e^{-tH'_L}$ generated by the finite-volume one-particle operators H_L and H'_L are trace class for every $t > 0$, and, a fortiori, compact. In particular, H_L and H'_L are bounded from below and have purely discrete spectra. We write $\lambda_1^L \leq \lambda_2^L \leq \dots$ and $\mu_1^L \leq \mu_2^L \leq \dots$ for their non-decreasing sequences of eigenvalues, counting multiplicities, and $(\varphi_j^L)_{j \in \mathbb{N}}$ and $(\psi_k^L)_{k \in \mathbb{N}}$ for the corresponding sequences of normalised eigenfunctions with an arbitrary choice of basis vectors in any eigenspace of dimension greater than one.

Given $N \in \mathbb{N}$, the *induced (non-interacting) finite-volume N -particle Schrödinger operators* \hat{H}_L and \hat{H}'_L act on the totally antisymmetric subspace $\bigwedge_{j=1}^N L^2(\Lambda_L)$ of the N -fold tensor product space and are given by

$$\hat{H}_L^{(l)} := \sum_{j=1}^N I^{\otimes j-1} \otimes H_L^{(l)} \otimes I^{\otimes N-j}.$$

The corresponding ground states are given by the totally antisymmetrised products

$$\Phi_L^N := \frac{1}{\sqrt{N!}} \varphi_1^L \wedge \dots \wedge \varphi_N^L, \quad \Psi_L^N := \frac{1}{\sqrt{N!}} \psi_1^L \wedge \dots \wedge \psi_N^L.$$

In order to avoid ambiguities from possibly degenerate eigenspaces and to realise a given *Fermi energy* $E \in \mathbb{R}$ in the thermodynamic limit, we choose the *number of particles* as

$$N_L(E) := \#\{j \in \mathbb{N} : \lambda_j^L \leq E\} \in \mathbb{N}_0, \tag{8}$$

which is the eigenvalue counting function of H_L at E .

The quantity of interest is the *ground-state overlap*

$$\mathcal{S}_L(E) := \langle \Phi_L^{N_L(E)}, \Psi_L^{N_L(E)} \rangle_{N_L(E)} = \det(\langle \varphi_j^L, \psi_k^L \rangle)_{j,k=1,\dots,N_L(E)}, \tag{9}$$

in particular its asymptotic behaviour as $L \rightarrow \infty$. In (9), $\langle \cdot, \cdot \rangle_N$ stands for the scalar product on the N -fermion space $\bigwedge_{j=1}^N L^2(\Lambda_L)$, and $\langle \cdot, \cdot \rangle$ for the one on the single-particle space $L^2(\Lambda_L)$. If $N_L(E) = 0$, we set $\mathcal{S}_L(E) := 1$.

Remark 2.1. The particular choice (8) of $N_L(E)$ as an eigenvalue counting function turns out to be technically useful when conducting the thermodynamic limit, see Lemma 3.3 below. The *particle density* $\rho(E)$ of the two non-interacting fermion systems in the thermodynamic limit coincides with the integrated density of states

$$\rho(E) = \lim_{L \rightarrow \infty} \frac{N_L(E)}{L^d |\Lambda_1|} \tag{10}$$

of the single-particle Schrödinger operator H (which is the same as the integrated density of states of H'), provided the limit exists. Here, $|\Lambda_1|$ denotes the Lebesgue measure of $\Lambda_1 \subseteq \mathbb{R}^d$. Situations where the limit (10) is known to exist include periodic V_0 , or V_0 vanishing at infinity. If the limit (10) does not exist, then this is due to the occurrence of more than one accumulation point, because the assumptions on V_0 in (V), together with [29, Theorem C.7.3], imply $\limsup_{L \rightarrow \infty} N_L(E)/L^d < \infty$ for every $E \in \mathbb{R}$. We will study the asymptotic behaviour of the overlap $\mathcal{S}_L(E)$ as $L \rightarrow \infty$ regardless of the existence of the limit (10).

The main result of this paper is an upper bound on the ground-state overlap $\mathcal{S}_L(E)$ for large L . Throughout we use the convention $\ln 0 := -\infty$. The terms null set and almost-every (a.e.) refer to Lebesgue measure if not specified otherwise.

Theorem 2.2 (orthogonality catastrophe). *Assume conditions (V). Let $(L_m)_{m \in \mathbb{N}}$ be a sequence in $(0, \infty)$ with $L_m \rightarrow \infty$. Then there exist a subsequence $(L_{m_k})_{k \in \mathbb{N}}$, a null set $\mathcal{N} \subseteq \mathbb{R}$ of exceptional Fermi energies and a function $\gamma: \mathbb{R} \setminus \mathcal{N} \rightarrow [0, \infty)$ such that for every $E \in \mathbb{R} \setminus \mathcal{N}$ the ground-state overlap (9) obeys*

$$|\mathcal{S}_{L_{m_k}}(E)| \leq \exp\left(-\frac{1}{2}\gamma(E) \ln L_{m_k} + o(\ln L_{m_k})\right) = L_{m_k}^{-\gamma(E)/2+o(1)} \tag{11}$$

as $k \rightarrow \infty$. Equivalently,

$$\limsup_{k \rightarrow \infty} \frac{\ln |\mathcal{S}_{L_{m_k}}(E)|}{\ln L_{m_k}} \leq -\frac{\gamma(E)}{2}. \tag{12}$$

The decay exponent γ is given by

$$\gamma(E) := \frac{1}{\pi^2} \|\arcsin|T_E/2|\|_{\text{HS}}^2. \tag{13}$$

Here, $T_E := S_E - I_E$ is the transition matrix, S_E is the scattering matrix for the pair (H, H') and energy E , and $\|\cdot\|_{\text{HS}}$ denotes the Hilbert–Schmidt norm on the fibre Hilbert space \mathcal{H}_E , on which T_E and S_E are defined.

Remarks 2.3. (i) We refer to Subsection 4.6 for a more precise definition of the scattering-theoretic quantities T_E and S_E .

(ii) In proving Theorem 2.2, we obtain a result on the asymptotics of the trace $\text{tr}\{(1_{(-\infty, E]}(H_L)1_{(E, \infty)}(H'_L))^n\}$ as $L \rightarrow \infty$, which may be interesting by itself; see Theorem 3.4.

(iii) The reason for passing to a subsequence $(L_{m_k})_{k \in \mathbb{N}}$ in Theorem 2.2 originates from Lemma 3.3 below. What stands behind it is the lack of known a.e.-bounds on the finite-volume spectral shift function for the pair of operators H_L, H'_L , which hold uniformly in the limit $L \rightarrow \infty$. This unfortunate fact has been noticed many times in the literature, see e.g. [15], and the pathological behaviour of the spectral shift function found in [17] illustrates that this is a delicate issue. However, in certain special situations such a.e.-bounds are known, and our result can be strengthened. More precisely, we have

Theorem 2.2'. Assume the situation of Theorem 2.2 with $d = 1$, or replace the perturbation potential V in Theorem 2.2 by a finite-rank operator $V = \sum_{\nu=1}^n \langle \phi_\nu, \cdot \rangle \phi_\nu$ with compactly supported $\phi_\nu \in L^2(\mathbb{R}^d)$ for $\nu = 1, \dots, n$, or consider the lattice problem on \mathbb{Z}^d corresponding to the situation in Theorem 2.2. Then the ground-state overlap (9) obeys

$$|\mathcal{S}_L(E)| \leq \exp\left(-\frac{1}{2}\gamma(E) \ln L + o(\ln L)\right) = L^{-\gamma(E)/2+o(1)}$$

for a.e. $E \in \mathbb{R}$ as $L \rightarrow \infty$. Equivalently,

$$\limsup_{L \rightarrow \infty} \frac{\ln|\mathcal{S}_L(E)|}{\ln L} \leq -\frac{\gamma(E)}{2}$$

for a.e. $E \in \mathbb{R}$.

Remarks 2.4. (i) In [11], similar statements to Theorem 2.2 and Theorem 2.2' were proved, in particular, the bound

$$\limsup_{L \rightarrow \infty} \frac{\ln|\mathcal{S}_L(E)|}{\ln L} \leq -\frac{\gamma_1(E)}{2}, \tag{14}$$

with the exponent

$$\gamma_1(E) = \frac{1}{\pi^2} \|T_E/2\|_{\text{HS}}^2. \tag{15}$$

Note that $\gamma_1(E)$, which is called $\gamma(E)$ in [11], is strictly smaller than $\gamma(E)$ whenever both are non-zero. The bigger exponent $\gamma(E)$ is due to treating all terms in a series expansion of $\ln|\mathcal{S}_L(E)|$ (see equation (17) below) instead of only the *Anderson integral*, which is the first term of the series and gives rise to $\gamma_1(E)$.

(ii) Another mathematical work dealing with AOC is [20]. That paper proves the *exact* asymptotics of the Anderson integral in the special case $d = 1$ and $V_0 = 0$. In particular, this yields a bound on the overlap as in (14) with the same non-optimal $\gamma_1(E)$ given by (15). The paper also provides a lower bound on $\mathcal{S}_L(E)$ with a smaller decay exponent [20, Corollary 5.6].

3. Series expansion of the overlap

In order to expand the ground-state overlap as a series, we introduce the orthogonal projections

$$P_L^N := \sum_{j=1}^N \langle \varphi_j^L, \cdot \rangle \varphi_j^L \quad \text{and} \quad \Pi_L^N := \sum_{k=1}^N \langle \psi_k^L, \cdot \rangle \psi_k^L \tag{16}$$

for $N \in \mathbb{N}_0$, i.e. the projections on the eigenspaces of the first N eigenvalues. Using those, we can prove the following lemma.

Lemma 3.1. *Let $L > 1$, $E \in \mathbb{R}$ and assume that $\mathcal{S}_L(E) \neq 0$. Then*

$$|\mathcal{S}_L(E)|^2 = \exp\left(-\sum_{n=1}^{\infty} \frac{1}{n} \text{tr}\{(P_L^{N_L(E)}(I - \Pi_L^{N_L(E)}))^n\}\right), \tag{17}$$

where we take the trace of operators on the Hilbert space $L^2(\Lambda_L)$.

Proof. For brevity, set $N := N_L(E)$. If $N = 0$, the assertion is true by definition. Otherwise, define the $N \times N$ -matrix $M := (\langle \varphi_j^L, \psi_k^L \rangle)_{j,k=1,\dots,N}$. Then $\mathcal{S}_L(E) = \det M$ and $|\mathcal{S}_L(E)|^2 = \det(MM^*)$. For $1 \leq j, \ell \leq N$, the (j, ℓ) -th entry of MM^* is

$$(MM^*)_{j,\ell} = \sum_{k=1}^N \langle \varphi_j^L, \psi_k^L \rangle \langle \psi_k^L, \varphi_\ell^L \rangle = \langle \varphi_j^L, \Pi_L^N \varphi_\ell^L \rangle = \langle \varphi_j^L, P_L^N \Pi_L^N P_L^N \varphi_\ell^L \rangle.$$

Since $\mathcal{S}_L(E) \neq 0$ by assumption and therefore $MM^* > 0$, we have $0 \leq P_L^N (I - \Pi_L^N) P_L^N < 1$. Moreover, being of finite rank, $P_L^N (I - \Pi_L^N) P_L^N$ is a trace class operator. Thus, we compute

$$\begin{aligned} |\mathcal{S}_L(E)|^2 &= \det(I - P_L^N (I - \Pi_L^N) P_L^N) \\ &= \exp(\text{tr}\{\ln(I - P_L^N (I - \Pi_L^N) P_L^N)\}) \\ &= \exp\left(-\text{tr}\left\{\sum_{n=1}^{\infty} \frac{1}{n} (P_L^N (I - \Pi_L^N) P_L^N)^n\right\}\right) \\ &= \exp\left(-\sum_{n=1}^{\infty} \frac{1}{n} \text{tr}\{(P_L^N (I - \Pi_L^N))^n\}\right), \end{aligned}$$

where we used the expansion $\ln(1 - x) = -\sum_{n=1}^{\infty} x^n/n$ for the logarithm, which converges absolutely for $|x| < 1$. □

Remark 3.2. Lemma 3.1 will be the starting point of our estimates for $|\mathcal{S}_L(E)|$. Equation (17) can be written as

$$-\ln|\mathcal{S}_L(E)| = \frac{1}{2} \sum_{n=1}^{\infty} \frac{1}{n} \text{tr}\{(P_L^{N_L(E)} (I - \Pi_L^{N_L(E)}))^n\}. \tag{18}$$

The trace expressions in (18) are non-negative, so any truncation of the series yields a lower bound on $-\ln|\mathcal{S}_L(E)|$, and therefore an upper bound on the overlap. Keeping only the term for $n = 1$, one recovers the so-called *Anderson integral*, which was estimated in [11].

In the sequel, we will find an upper bound on $|\mathcal{S}_L(E)|$ by bounding each individual term of (18) from below.

We begin by recasting the orthogonal projections (16) as functions of H_L and H'_L in the sense of the spectral calculus. The projections in (16) are not necessarily functions of H_L and H'_L , since the N th eigenvalues might be of multiplicity higher than one. The choice of $N_L(E)$ in (8), together with a convergence result of the spectral shift function, allows us to rewrite them, at the cost of passing to a subsequence of lengths.

Lemma 3.3. For $n \in \mathbb{N}$, $L > 1$ and $E \in \mathbb{R}$, define

$$\mathcal{F}_L^n(E) := \text{tr}\{(1_{(-\infty, E]}(H_L)1_{(E, \infty)}(H'_L))^n\} \tag{19}$$

and

$$\mathcal{J}_L^n(E) := \text{tr}\{(P_L^{N_L(E)}(I - \Pi_L^{N_L(E)}))^n\}.$$

Then

(i) Assume (V) and let $(L_m)_{m \in \mathbb{N}} \subset (0, \infty)$ be a sequence of increasing lengths with $L_m \uparrow \infty$. Then there exists a subsequence $(L_{m_k})_{k \in \mathbb{N}}$ such that for a.e. Fermi energy $E \in \mathbb{R}$

$$|\mathcal{F}_{L_{m_k}}^n(E) - \mathcal{J}_{L_{m_k}}^n(E)| = o(\ln L_{m_k}) \tag{20}$$

as $k \rightarrow \infty$.

(ii) Assume the situation of Theorem 2.2'. Then

$$\sup_{L > 1} \sup_{E \in \mathbb{R}} |\mathcal{F}_L^n(E) - \mathcal{J}_L^n(E)| < \infty.$$

Proof. For fixed $L > 1$ and $E \in \mathbb{R}$, the definition of $N_L(E)$ in (8) implies

$$\lambda_{N_L(E)}^L \leq E < \lambda_{N_L(E)+1}^L \leq \mu_{N_L(E)+1}^L$$

if we set $\lambda_0^L := -\infty$. This allows us to write

$$P_L^{N_L(E)} = 1_{(-\infty, E]}(H_L) \tag{21}$$

and

$$\begin{aligned} I - \Pi_L^{N_L(E)} &= 1_{(E, \infty)}(H'_L) - \sum_{k=1}^{N_L(E)} 1_{(E, \infty)}(\mu_k^L) \langle \psi_k^L, \cdot \rangle \psi_k^L \\ &=: 1_{(E, \infty)}(H'_L) - Q. \end{aligned}$$

The operator Q is an orthogonal projection with trace

$$\begin{aligned} \text{tr } Q &= \#\{k \in \{1, \dots, N_L(E)\} : \mu_k^L > E\} \\ &= N_L(E) - \#\{k \in \mathbb{N} : \mu_k^L \leq E\} \\ &=: \xi_L(E) \end{aligned}$$

equal to the finite-volume spectral-shift function at the Fermi energy.

Using $A^n - B^n = \sum_{k=1}^n B^{k-1}(A - B)A^{n-k}$ for bounded operators A and B , we write the difference of operator powers on the left-hand side of (20) as

$$\begin{aligned} &(P_L^{N_L(E)}1_{(E, \infty)}(H'_L))^n - (P_L^{N_L(E)}(I - \Pi_L^{N_L(E)}))^n \\ &= \sum_{k=1}^n (P_L^{N_L(E)}(I - \Pi_L^{N_L(E)}))^{k-1} P_L^{N_L(E)} Q (P_L^{N_L(E)}1_{(E, \infty)}(H'_L))^{n-k}, \end{aligned} \tag{22}$$

where we also use (21). We estimate the traces of the operators on the right-hand side of (22) by bounding the operator norms of all projections, except for Q , by 1. We then arrive at $n\xi_L(E)$ as an upper bound for (22). The claim follows by exploiting the weak convergence of ξ_L as $L \rightarrow \infty$ [15, Theorem 1.4] in the situation of (i), or using the uniform boundedness of ξ_L in the situation of (ii). We refer to [11, Lemma 3.9] for a detailed argument. \square

Having established (20), we will prove a diverging lower bound for $\text{tr}\{(1_{(-\infty, E]}(H_L)1_{(E, \infty)}(H'_L))^n\}$ as $L \rightarrow \infty$. There will be no restriction to particular sequences of lengths from now on. The following theorem is the main ingredient of the proof.

Theorem 3.4. *Assume the situation of Theorem 2.2 or Theorem 2.2'. Then there exists a null set $\mathcal{N} \subset \mathbb{R}$ of exceptional Fermi energies such that*

$$\text{tr}\{(1_{(-\infty, E]}(H_L)1_{(E, \infty)}(H'_L))^n\} \geq nJ_{2n} \text{tr}(|T_E/(2\pi)|^{2n}) \ln L + o(\ln L) \quad (23)$$

for every $E \in \mathbb{R} \setminus \mathcal{N}$ and every $n \in \mathbb{N}$ as $L \rightarrow \infty$. The error term $o(\ln L)$ depends on n and E , and we introduced the constant

$$J_{2n} := \pi^{2(n-1)} 2^{2n-1} \frac{[(n-1)!]^2}{(2n)!}. \quad (24)$$

Remarks 3.5. (i) In the next section, we will spell out explicitly the proof of Theorem 3.4 for the situation of Theorem 2.2 only. It follows from Corollary 4.25, Theorem 4.26 and Theorem 4.32. The proof is fully analogous (and even simpler) in the remaining situations of Theorem 2.2', where V is a finite-rank operator.

(ii) The constant J_{2n} will emerge as the value of a $2n$ -dimensional integral which we calculate using the spectral representation of the Hilbert matrix, see Subsection 4.5 below.

Given Theorem 3.4, we are now in a position to prove Theorem 2.2.

Proof of Theorem 2.2. Let $M \in \mathbb{N}$. Let \mathcal{N} be the null set from Theorem 3.4. Let $E \in \mathbb{R} \setminus \mathcal{N}$. We start from Lemma 3.1 and Lemma 3.3, which imply

$$\begin{aligned} -\ln|\mathcal{S}_{L_{m_k}}(E)| &\geq \frac{1}{2} \sum_{n=1}^M \frac{1}{n} \text{tr}\{(P_{L_{m_k}}^{N_{L_{m_k}}(E)}(I - \Pi_{L_{m_k}}^{N_{L_{m_k}}(E)}))^n\} \\ &= \frac{1}{2} \sum_{n=1}^M \frac{1}{n} \text{tr}\{(1_{(-\infty, E]}(H_{L_{m_k}})1_{(E, \infty)}(H'_{L_{m_k}}))^n\} + o(\ln L_{m_k}) \end{aligned}$$

for a subsequence $(L_{m_k})_{k \in \mathbb{N}}$, as $k \rightarrow \infty$, with an M -dependent error term $o(\ln L_{m_k})$. By Theorem 3.4, this gives

$$-\ln|\mathcal{S}_{L_{m_k}}(E)| \geq \frac{1}{2} \operatorname{tr} \left\{ \sum_{n=1}^M J_{2n} |T_E / (2\pi)|^{2n} \right\} \ln L_{m_k} + o(\ln L_{m_k}) \quad (25)$$

as $k \rightarrow \infty$, with an M -dependent error term $o(\ln L_{m_k})$. The constants J_{2n} show up in the series expansion [12, eq. 1.645 2]

$$\sum_{n=1}^{\infty} J_{2n} x^{2n} = \pi^{-2} (\arcsin(\pi x))^2 \quad \text{for } |x| \leq \frac{1}{\pi}.$$

Therefore, monotone convergence and the functional calculus yield

$$\lim_{M \rightarrow \infty} \operatorname{tr} \left\{ \sum_{n=1}^M J_{2n} |T_E / (2\pi)|^{2n} \right\} = \pi^{-2} \|\arcsin|T_E / 2|\|_{\text{HS}}^2.$$

Since (25) is valid for every $M \in \mathbb{N}$, we infer

$$\limsup_{k \rightarrow \infty} \frac{\ln|\mathcal{S}_{L_{m_k}}(E)|}{\ln L_{m_k}} \leq -\frac{1}{2} \pi^{-2} \|\arcsin|T_E / 2|\|_{\text{HS}}^2 = -\frac{\gamma(E)}{2},$$

which proves (12). For (11), note that by the definition of the limit superior for every $\varepsilon > 0$ there is $k_0 \in \mathbb{N}$ such that

$$\frac{\ln|\mathcal{S}_{L_{m_k}}(E)|}{\ln L_{m_k}} \leq -\frac{\gamma(E)}{2} + \varepsilon$$

for all $k \geq k_0$, which implies the claim. □

It remains to prove Theorem 3.4.

4. Proof of Theorem 3.4

4.1. An integral representation for $\operatorname{tr}\{(f(H_L)g(H'_L))^n\}$. Throughout this subsection, $n \in \mathbb{N}$, $L > 1$ and $E \in \mathbb{R}$ are all fixed. Using the eigenvalue equations of H_L and H'_L , we rewrite trace expressions like (19).

Lemma 4.1. *Let $f, g: \mathbb{R} \rightarrow [0, 1]$ be measurable functions with compact supports $\operatorname{supp} f \subseteq (-\infty, E]$ and $\operatorname{supp} g \subseteq (E, \infty)$. Then*

$$\operatorname{tr}\{(f(H_L)g(H'_L))^n\} = \sum_{\alpha, \beta \in \mathbb{N}^n} \prod_{j=1}^n \left(f(\lambda_{\alpha_j}^L) g(\mu_{\beta_j}^L) \frac{\langle \varphi_{\alpha_j}^L, V \psi_{\beta_j}^L \rangle \langle \psi_{\beta_j}^L, V \varphi_{\alpha_{j+1}}^L \rangle}{(\mu_{\beta_j}^L - \lambda_{\alpha_j}^L)(\mu_{\beta_j}^L - \lambda_{\alpha_{j+1}}^L)} \right), \quad (26)$$

for multi-indices $\alpha = (\alpha_1, \dots, \alpha_n) \in \mathbb{N}^n$ with the convention $\alpha_{n+1} := \alpha_1$.

Proof. We begin noting that

$$f(H_L) = \sum_{j \in \mathbb{N}} f(\lambda_j^L) \langle \varphi_j^L, \cdot \rangle \varphi_j^L, \quad g(H'_L) = \sum_{k \in \mathbb{N}} g(\mu_k^L) \langle \psi_k^L, \cdot \rangle \psi_k^L. \quad (27)$$

To ease notation, we employ the bra-ket notation in the next formula, writing $\langle \varphi, \cdot \rangle \varphi =: |\varphi\rangle\langle\varphi|$ for $\varphi \in L^2(\Lambda_L)$. Then (27) implies

$$(f(H_L)g(H'_L))^n = \sum_{\alpha, \beta \in \mathbb{N}^n} \left(\prod_{j=1}^n f(\lambda_{\alpha_j}^L) g(\mu_{\beta_j}^L) \right) \prod_{j=1}^n |\varphi_{\alpha_j}^L\rangle\langle\varphi_{\alpha_j}^L, \psi_{\beta_j}^L\rangle\langle\psi_{\beta_j}^L| \quad (28)$$

and

$$\begin{aligned} & \text{tr}\{(f(H_L)g(H'_L))^n\} \\ &= \sum_{\alpha, \beta \in \mathbb{N}^n} \left(\prod_{j=1}^n f(\lambda_{\alpha_j}^L) g(\mu_{\beta_j}^L) \right) \prod_{j=1}^n \langle \varphi_{\alpha_j}^L, \psi_{\beta_j}^L \rangle \langle \psi_{\beta_j}^L, \varphi_{\alpha_{j+1}}^L \rangle, \end{aligned} \quad (29)$$

where we used the convention $\alpha_{n+1} := \alpha_1$ for $\alpha \in \mathbb{N}^n$. Now, we note that the eigenvalue equations imply

$$\lambda_j^L \langle \varphi_j^L, \psi_k^L \rangle = \langle H_L \varphi_j^L, \psi_k^L \rangle = \mu_k^L \langle \varphi_j^L, \psi_k^L \rangle - \langle \varphi_j^L, V \psi_k^L \rangle$$

for $j, k \in \mathbb{N}$, and therefore

$$\langle \varphi_j^L, \psi_k^L \rangle = \frac{\langle \varphi_j^L, V \psi_k^L \rangle}{\mu_k^L - \lambda_j^L} \quad (30)$$

whenever $\lambda_j^L \neq \mu_k^L$. Since f and g have disjoint supports, (30) and (29) yield the claim. \square

Remark 4.2. In analogy to [11], one might be tempted to define a spectral correlation “measure” by

$$\begin{aligned} & \mu_L^{2n}(A_1 \times \dots \times A_n \times B_1 \times \dots \times B_n) \\ & := \text{tr}\{(1_{A_1}(H_L)V1_{B_1}(H'_L)V \dots 1_{A_n}(H_L)V1_{B_n}(H'_L)V)\} \end{aligned} \quad (31)$$

for $n \in \mathbb{N}$, $L > 1$ and bounded $A_1, \dots, A_n, B_1, \dots, B_n \in \text{Borel}(\mathbb{R})$, which was done for the case $n = 1$ in [11]. Lemma 4.1 would then read

$$\text{tr}\{(f(H_L)g(H'_L))^n\} = \int_{\mathbb{R}^n \times \mathbb{R}^n} d\mu_L^{2n}(x, y) \prod_{j=1}^n \frac{f(x_j)g(y_j)}{(y_j - x_j)(y_j - x_{j+1})}.$$

However, (31) is not necessarily non-negative for $n \geq 2$, and therefore we cannot mimick the proof of [11].

Next, we rewrite the right-hand side of (26) using a variation of an integral formula that goes back to Feynman and Schwinger.

Lemma 4.3 (Feynman–Schwinger parametrization). *Let $x_1, \dots, x_n \in (0, \infty)$. Then*

$$\frac{1}{x_1 \cdots x_n} = \int_0^\infty dt t^{n-1} \int_{(0, \infty)^n} du |u|_1 e^{-|u|_1} e^{-tu \cdot x}, \quad (32)$$

where $u \cdot x = \sum_{j=1}^n u_j x_j$ denotes the Euclidean scalar product and $|u|_1 := \sum_{j=1}^n |u_j|$ the 1-norm on \mathbb{R}^n .

Proof. For any measurable function $f: (0, \infty)^n \rightarrow (0, \infty)$ the coarea formula implies

$$\int_{(0, \infty)^n} du f(u) = \int_0^\infty dt \int_{\mathcal{M}} \frac{dS(\xi)}{\sqrt{n}} t^{n-1} f(t\xi), \quad (33)$$

where dS stands for integration with respect to the surface measure on $\mathcal{M} := \{\xi \in (0, \infty)^n : |\xi|_1 = 1\}$. Let $r > 0$. Starting from $x_j^{-1} = \int_0^\infty du_j e^{-u_j x_j}$, we compute using (33)

$$\begin{aligned} \frac{1}{x_1 \cdots x_n} &= \int_{(0, \infty)^n} du e^{-u \cdot x} \\ &= \int_0^\infty dt \int_{\mathcal{M}} \frac{dS(\xi)}{\sqrt{n}} t^{n-1} e^{-t\xi \cdot x} \\ &= \int_0^\infty dt \int_{\mathcal{M}} \frac{dS(\xi)}{\sqrt{n}} t^{n-1} r^n e^{-rt\xi \cdot x}, \end{aligned}$$

which is r -independent. Given any measurable function $g: (0, \infty) \rightarrow (0, \infty)$ with $\int_0^\infty dr \frac{g(r)}{r} = 1$, we therefore get

$$\begin{aligned} \frac{1}{x_1 \cdots x_n} &= \int_0^\infty dr g(r) \int_0^\infty dt \int_{\mathcal{M}} \frac{dS(\xi)}{\sqrt{n}} t^{n-1} r^{n-1} e^{-rt\xi \cdot x} \\ &= \int_0^\infty dt t^{n-1} \int_{(0, \infty)^n} du g(|u|_1) e^{-tu \cdot x}, \end{aligned}$$

where we used the Fubini–Tonelli theorem and (33) with $f(u) = g(|u|_1) e^{-tu \cdot x}$. Choosing $g(r) := r e^{-r}$ finishes the proof. \square

We use (32) to rewrite the right-hand side of (26).

Lemma 4.4. *Let $f, g: \mathbb{R} \rightarrow [0, 1]$ be measurable functions with compact supports $\text{supp } f \subseteq (-\infty, E]$ and $\text{supp } g \subseteq (E, \infty)$. Then,*

$$\begin{aligned} & \text{tr}\{(f(H_L)g(H'_L))^n\} \\ &= \int_0^\infty dt t^{2n-1} \int_{(0,\infty)^n \times (0,\infty)^n} d(u, v) (|u|_1 + |v|_1) e^{-|u|_1 - |v|_1} \\ & \quad \text{tr} \left\{ \prod_{j=1}^n \sqrt{V} f(H_L) e^{(u_j + v_{j-1})t(H_L - E)} V g(H'_L) e^{-(u_j + v_j)t(H'_L - E)} \sqrt{V} \right\} \end{aligned} \tag{34}$$

with the convention $v_0 := v_n$ for $v \in \mathbb{R}^n$.

Proof. Let $x \in (-\infty, 0]^n$, $y \in (0, \infty)^n$ and define $x_{n+1} := x_1$. Then, by (32),

$$\begin{aligned} & \frac{1}{\prod_{j=1}^n (y_j - x_j)(y_j - x_{j+1})} \\ &= \int_0^\infty dt t^{2n-1} \int_{(0,\infty)^n \times (0,\infty)^n} d(u, v) (|u|_1 + |v|_1) e^{-|u|_1 - |v|_1} \\ & \quad \exp \left(-t \sum_{j=1}^n (u_j (y_j - x_j) + v_j (y_j - x_{j+1})) \right) \end{aligned}$$

and

$$\sum_{j=1}^n (u_j (y_j - x_j) + v_j (y_j - x_{j+1})) = \sum_{j=1}^n ((u_j + v_j) y_j - (u_j + v_{j-1}) x_j)$$

for $u, v \in (0, \infty)^n$. Now, let $\alpha, \beta \in \mathbb{N}^n$. Setting $x_j = \lambda_{\alpha_j}^L - E$ and $y_j = \mu_{\beta_j}^L - E$, we can write the denominator in (26) as

$$\begin{aligned} & \frac{1}{\prod_{j=1}^n (\mu_{\beta_j}^L - \lambda_{\alpha_j}^L)(\mu_{\beta_j}^L - \lambda_{\alpha_{j+1}}^L)} \\ &= \int_0^\infty dt t^{2n-1} \int_{(0,\infty)^n \times (0,\infty)^n} d(u, v) (|u|_1 + |v|_1) e^{-|u|_1 - |v|_1} \\ & \quad \prod_{j=1}^n e^{-(u_j + v_j)t(\mu_{\beta_j}^L - E)} e^{(u_j + v_{j-1})t(\lambda_{\alpha_j}^L - E)}. \end{aligned} \tag{35}$$

The sums over α and β in (26) contain only finitely many terms, due to the compact supports of f and g . Therefore these sums can be interchanged with the integrals from (35). This results in

$$\begin{aligned} & \text{tr}\{(f(H_L)g(H'_L))^n\} \\ &= \int_0^\infty dt t^{2n-1} \int_{(0,\infty)^n \times (0,\infty)^n} d(u, v) (|u|_1 + |v|_1) e^{-|u|_1 - |v|_1} \\ & \quad \sum_{\alpha, \beta \in \mathbb{N}^n} \prod_{j=1}^n (f(\lambda_{\alpha_j}^L) e^{(u_j + v_{j-1})t(\lambda_{\alpha_j}^L - E)} g(\mu_{\beta_j}^L) e^{-(u_j + v_j)t(\mu_{\beta_j}^L - E)} \\ & \quad \langle \varphi_{\alpha_j}^L, V\psi_{\beta_j}^L \rangle \langle \psi_{\beta_j}^L, V\varphi_{\alpha_{j+1}}^L \rangle), \end{aligned}$$

from which the assertion follows. □

4.2. Smoothing and infinite-volume operators. Throughout this subsection, $a \in (0, 1)$ and $n \in \mathbb{N}$ are fixed. We also fix a cut-off energy $E_0 \geq 1$ and a Fermi energy $E \in [-E_0 + 1, E_0 - 1]$.

The goal is to apply Lemma 4.4 using suitable functions f and g and to rewrite the right-hand side of (34) as a trace involving the infinite-volume operators H and H' . Switching from finite-volume to infinite-volume operators constitutes the core of the argument. The technical tool to implement this switch to infinite-volume objects is the Helffer–Sjöstrand formula, which supplies the proof of Lemma 4.8 below. Since it is applicable to sufficiently smooth functions only, we define appropriately smoothed versions of indicator functions.

Definition 4.5. Given a length $L > 1$, we say that $\chi_L^\pm \in C_c^\infty(\mathbb{R})$ are *smooth cut-off functions at energy E* , if they obey

$$\begin{aligned} 1_{[E+2L^{-a}, E_0]} &\leq \chi_L^+ \leq 1_{(E+L^{-a}, E_0+1)}, \\ 1_{[-E_0, E-2L^{-a}]} &\leq \chi_L^- \leq 1_{(-E_0-1, E-L^{-a})}, \end{aligned}$$

and if there exist L -independent constants $c_k > 0$ for $k \in \mathbb{N}_0$, such that

$$\chi_L^\pm(E \pm L^{-a} \pm x) \leq c_0 L^a x$$

for all $x \in [0, L^{-a})$ and

$$\left| \frac{d^k}{dx^k} \chi_L^\pm(E \pm L^{-a} \pm x) \right| \leq \begin{cases} c_k L^{ak} & \text{if } 0 \leq x < L^{-a}, \\ c_k & \text{otherwise,} \end{cases}$$

for every $k \in \mathbb{N}$ and $x \in \mathbb{R}$. We choose the smooth decay of χ_L^+ in $[E_0, E_0 + 1]$ independently of L , and analogously for χ_L^- . Clearly such functions exist. Figure 1 illustrates the behaviour of χ_L^\pm .

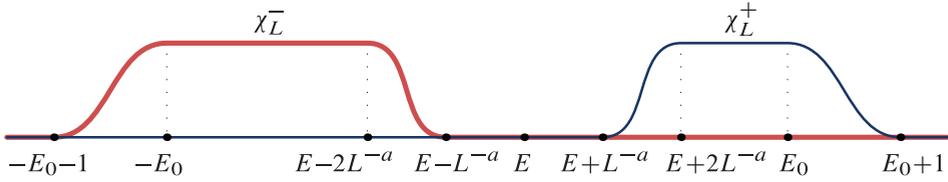


Figure 1. Sketch of the smooth cut-off functions χ_L^\pm .

We are interested in a lower bound for the left-hand side of (23) which is proportional to $\ln L$ up to subdominant corrections.

Lemma 4.6. *Let $L > 1$. Then*

$$\begin{aligned} & \text{tr}\{(1_{(-\infty, E]}(H_L)1_{(E, \infty)}(H'_L))^n\} \\ & \geq \int_{(0, \infty)^n \times (0, \infty)^n} d(u, v) (|u|_1 + |v|_1) e^{-|u|_1 - |v|_1} \int_0^\infty dt t^{2n-1} \\ & \quad \text{tr} \left\{ \prod_{j=1}^n \sqrt{V} \chi_L^-(H_L) e^{(u_j + v_{j-1})t(H_L - E)} V \chi_L^+(H'_L) e^{-(u_j + v_j)t(H'_L - E)} \sqrt{V} \right\}. \end{aligned} \tag{36}$$

Proof. The inequalities

$$1_{(-\infty, E]} \geq \chi_L^- \quad \text{and} \quad 1_{(E, \infty)} \geq \chi_L^+, \tag{37}$$

together with the cyclicity of the trace, imply

$$\text{tr}\{(1_{(-\infty, E]}(H_L)1_{(E, \infty)}(H'_L))^n\} \geq \text{tr}\{(\chi_L^-(H_L)\chi_L^+(H'_L))^n\}.$$

Together with Lemma 4.4, this yields the claim. □

Remark 4.7. In the sequel, we determine the exact asymptotics of the right-hand side of (36). Thus it is only the smoothing introduced in Lemma 4.6 which prevents us from determining the exact asymptotics in Theorem 3.4.

The following technical lemma constitutes the core of the arguments in the present subsection.

Lemma 4.8. *For $L > 1$, $t \geq 0$ and $x \in \mathbb{R}$ we define*

$$f_L^t(x) := \chi_L^-(x)e^{t(x-E)} \quad \text{and} \quad g_L^t(x) := \chi_L^+(x)e^{-t(x-E)},$$

and h_L^t will stand for either f_L^t or g_L^t . Let $M \in \mathbb{N} \setminus \{1\}$. Then there are constants $c > 0$, $L_0 > 1$ and a polynomial Q_M of degree $M + 1$ with non-negative coefficients, such that for every $t \geq 0$, every $L \geq L_0$ and every $\varepsilon \in (0, 1 - a)$ the estimate

$$\begin{aligned} & \|\sqrt{V}(h_L^t(H_L^{(t)}) - h_L^t(H^{(t)}))\sqrt{V}\| \\ & \leq Q_M(t/L^a)(L^{a-M(1-a-\varepsilon)} + L^{d+a(M+1)}e^{-cL^\varepsilon}) \end{aligned}$$

holds.

Proof. The proof essentially follows the ideas of a part of the proof in [11, Lemma 3.14]. The idea is to apply the Helffer–Sjöstrand formula to estimate the difference of resolvents, cf. formulae (4.13) and (4.16) in [11]. For a detailed exposition, see [19]. □

Before we prove the main assertion of this subsection, we need a spectral-gap estimate.

Lemma 4.9. *There is a constant $C > 0$ such that for every $L > 1$ and every $t \geq 0$ we have*

$$\text{tr}\{\sqrt{V}h_L^t(H_{(L)}^{(t)})\sqrt{V}\} \leq Ce^{-tL^{-a}},$$

where $h_L^t \in C_c^\infty(\mathbb{R})$ is as in Lemma 4.8.

Proof. For the first assertion, note that there is a bounded interval $I \subseteq \mathbb{R}$ such that $h_L^t \leq 1_I e^{-tL^{-a}}$ for all $t \geq 0$ and $L > 1$. Thus,

$$\begin{aligned} \text{tr}\{\sqrt{V}h_L^t(H_{(L)}^{(t)})\sqrt{V}\} & \leq e^{-tL^{-a}} \text{tr}\{\sqrt{V}1_I(H_{(L)}^{(t)})\sqrt{V}\} \\ & \leq e^{-tL^{-a}} e^{\sup I} \text{tr}\{\sqrt{V}e^{-H_{(L)}^{(t)}}\sqrt{V}\} \\ & \leq e^{-tL^{-a}} e^{\sup I} \text{tr}\{\sqrt{V}e^{-H^{(t)}}\sqrt{V}\} \end{aligned}$$

for $t \geq 0$ and $L > 1$. The last inequality and the finiteness of $\text{tr}\{\sqrt{V}e^{-H^{(t)}}\sqrt{V}\}$ follow from [5, Theorem 6.1]. □

The next lemma accomplishes the transition from finite-volume to infinite-volume operators.

Lemma 4.10. *For $L > 1$ and $t \geq 0$, let $f_L^t, g_L^t \in C_c^\infty(\mathbb{R})$ as in Lemma 4.8. Let $u, v \in (0, \infty)^n$. Then*

$$\int_0^\infty dt \, t^{2n-1} \operatorname{tr} \left\{ \left| \prod_{j=1}^n \sqrt{V} f_L^{(u_j+v_{j-1})t} (H_L) V g_L^{(u_j+v_j)t} (H'_L) \sqrt{V} \right. \right. \\ \left. \left. - \prod_{j=1}^n \sqrt{V} f_L^{(u_j+v_{j-1})t} (H) V g_L^{(u_j+v_j)t} (H') \sqrt{V} \right| \right\} \quad (38)$$

$$\leq \frac{1}{(|u|_1 + |v|_1)^{2n}} o(1)$$

as $L \rightarrow \infty$, where the $o(1)$ -term does not depend on u or v . We also used the convention $v_0 := v_n$.

Proof. To shorten formulas, we introduce a vector $\alpha \in (0, \infty)^{2n}$ via

$$\alpha_{2j-1} := u_j + v_{j-1} \quad \text{and} \quad \alpha_{2j} := u_j + v_j \quad (39)$$

for $1 \leq j \leq n$ and operators

$$A_k^{(L)} := \begin{cases} \sqrt{V} f_L^{\alpha_k t} (H_{(L)}) \sqrt{V} & \text{for } k \text{ odd,} \\ \sqrt{V} g_L^{\alpha_k t} (H'_{(L)}) \sqrt{V} & \text{for } k \text{ even,} \end{cases}$$

for $1 \leq k \leq 2n$. The difference of operator products in (38) is then

$$\prod_{j=1}^{2n} A_j^L - \prod_{j=1}^{2n} A_j = \sum_{k=1}^{2n} A_1 \dots A_{k-1} (A_k^L - A_k) A_{k+1}^L \dots A_{2n}^L. \quad (40)$$

The trace norm of this difference can be estimated using Lemma 4.9: There is a constant $C > 0$ such that

$$\operatorname{tr} \left\{ \left| \prod_{j=1}^{2n} A_j^L - \prod_{j=1}^{2n} A_j \right| \right\} \leq \sum_{k=1}^{2n} \|A_k^L - A_k\| \left(\prod_{j=1}^{k-1} \operatorname{tr}\{|A_j|\} \right) \left(\prod_{j=k+1}^{2n} \operatorname{tr}\{|A_j^L|\} \right) \\ \leq C^{2n-1} \sum_{k=1}^{2n} \|A_k^L - A_k\| e^{-(|\alpha|_1 - \alpha_k)tL^{-a}},$$

where $|\alpha|_1 = \alpha_1 + \dots + \alpha_{2n}$ denotes the 1-norm of $\alpha \in (0, \infty)^{2n}$. We estimate the k th term in this sum. Let $\varepsilon \in (0, 1 - a)$ and $M \in \mathbb{N}$. For L sufficiently large,

Lemma 4.8 implies

$$\begin{aligned} & \|A_k^L - A_k\| e^{-(|\alpha|_1 - \alpha_k)tL^{-a}} \\ & \leq Q_M(\alpha_k t/L^a)(L^{a-M(1-a-\varepsilon)} + L^{d+a(M+1)}e^{-cL^\varepsilon})e^{-(|\alpha|_1 - \alpha_k)tL^{-a}}, \end{aligned} \tag{41}$$

where $Q_M(x) = \sum_{\ell=0}^{M+1} q_\ell x^\ell$ is the polynomial in Lemma 4.8 with non-negative coefficients q_ℓ . Integrating (41) yields

$$\begin{aligned} & \int_0^\infty dt t^{2n-1} \|A_k^L - A_k\| e^{-(|\alpha|_1 - \alpha_k)tL^{-a}} \\ & \leq (L^{a-M(1-a-\varepsilon)} + L^{d+a(M+1)}e^{-cL^\varepsilon}) \\ & \quad \sum_{\ell=0}^{M+1} q_\ell \int_0^\infty dt t^{2n-1} e^{-(|\alpha|_1 - \alpha_k)tL^{-a}} \frac{\alpha_k^\ell t^\ell}{L^{a\ell}} \\ & = (L^{a-M(1-a-\varepsilon)} + L^{d+a(M+1)}e^{-cL^\varepsilon}) \\ & \quad \sum_{\ell=0}^{M+1} \frac{q_\ell \Gamma(2n + \ell) \alpha_k^\ell L^{a(2n+\ell)}}{L^{a\ell} (|\alpha|_1 - \alpha_k)^{2n+\ell}}, \end{aligned} \tag{42}$$

where Γ denotes Euler’s Gamma Function. The definition of $\alpha \in (0, \infty)^{2n}$ in (39) yields $|\alpha|_1 = 2(|u|_1 + |v|_1)$, and thus $|\alpha|_1 - \alpha_k \geq |u|_1 + |v|_1 \geq \alpha_k$. This makes the right-hand side of (42) smaller than

$$C_M \frac{L^{2na}}{(|u|_1 + |v|_1)^{2n}} (L^{a-M(1-a-\varepsilon)} + L^{d+a(M+1)}e^{-cL^\varepsilon})$$

with some constant $C_M > 0$ depending on Q_M and n . For given $\varepsilon < 1 - a$, we can choose M large enough for the L -terms to vanish as $L \rightarrow \infty$. □

Using Lemma 4.10, we can rewrite the right-hand side of (36).

Corollary 4.11. *The estimate*

$$\begin{aligned} & \text{tr}\{(1_{(-\infty, E]}(H_L)1_{(E, \infty)}(H'_L))^n\} \\ & \geq \int_{(0, \infty)^n \times (0, \infty)^n} d(u, v) (|u|_1 + |v|_1) e^{-|u|_1 - |v|_1} \int_0^\infty dt t^{2n-1} \\ & \quad \text{tr} \left\{ \prod_{j=1}^n \sqrt{V} \chi_L^-(H) e^{(u_j + v_{j-1})t(H-E)} V \chi_L^+(H') e^{-(u_j + v_j)t(H'-E)} \sqrt{V} \right\} + o(1) \end{aligned} \tag{43}$$

holds as $L \rightarrow \infty$.

Proof. The claim follows from Lemma 4.6 and Lemma 4.10, which imply that the integral

$$\int_{(0,\infty)^n \times (0,\infty)^n} d(u, v) (|u|_1 + |v|_1) e^{-|u|_1 - |v|_1} \int_0^\infty dt t^{2n-1} \left(\text{tr} \left\{ \prod_{j=1}^n \sqrt{V} f_L^{(u_j + v_{j-1})t} (H_L) V g_L^{(u_j + v_j)t} (H'_L) \sqrt{V} \right\} - \text{tr} \left\{ \prod_{j=1}^n \sqrt{V} f_L^{(u_j + v_{j-1})t} (H) V g_L^{(u_j + v_j)t} (H') \sqrt{V} \right\} \right)$$

vanishes in the limit $L \rightarrow \infty$, because

$$\int_{(0,\infty)^n \times (0,\infty)^n} d(u, v) \frac{e^{-|u|_1 - |v|_1}}{(|u|_1 + |v|_1)^{2n-1}} = \frac{1}{(2n-1)!} < \infty,$$

as can be seen from the coarea formula. □

Remark 4.12. Comparing the smooth cut-off functions χ_L^\pm with the ones in [11, Def. 3.13], the difference is that the cut-off functions there have E as the boundary of their support, while the ones here have distance L^{-a} between E and their support. To compensate for this, the t -integral has been cut off at $t = L^{-a}$ in [11, Lemma 3.11], which yields a lower bound for $n = 1$. For $n \geq 2$, it is not immediately clear if the integrand in (43) is positive, so cutting off the integration might not result in a lower bound; this is the reason for choosing the cut-off functions differently from those in [11].

4.3. Infinite-volume trace expressions. Throughout this subsection, we fix $a \in (0, 1)$, $n \in \mathbb{N}$ and a cut-off energy $E_0 \geq 1$.

In Corollary 4.11, we gave a lower bound on the n th term of (18) in which only infinite-volume operators occur. In order to control the errors in that step, it was necessary to introduce smoothed versions of indicator functions in (37). In the present subsection, our aim is to replace these smoothed functions with discontinuous ones, which will allow us to determine the asymptotics of the resulting expression.

We introduce measures $\mu^1, \nu^1: \text{Borel}(\mathbb{R}) \rightarrow [0, \infty]$ defined by

$$\mu^1(A) := \text{tr}\{\sqrt{V} 1_A(H) \sqrt{V}\}, \quad \nu^1(B) := \text{tr}\{\sqrt{V} 1_B(H') \sqrt{V}\} \tag{44}$$

for $A, B \in \text{Borel}(\mathbb{R})$. The expressions in (44) are finite for bounded Borel sets as a consequence of [29, Theorem B.9.2].

The absolutely continuous parts of the measures μ^1 and ν^1 will turn out to be important. To define their densities in an applicable manner, we use a limiting absorption principle due to Birman and Èntina.

Proposition 4.13 ([4, Lemma 4.3]). *There exists a null set $\mathcal{N}_0 \subset \mathbb{R}$ such that the limits*

$$A(E) := \lim_{\varepsilon \downarrow 0} \frac{1}{2\varepsilon} \sqrt{V} 1_{(E-\varepsilon, E+\varepsilon)}(H) \sqrt{V}, \quad (45)$$

$$B(E) := \lim_{\varepsilon \downarrow 0} \frac{1}{2\varepsilon} \sqrt{V} 1_{(E-\varepsilon, E+\varepsilon)}(H') \sqrt{V}$$

exist in trace class for all $E \in \mathbb{R} \setminus \mathcal{N}_0$ and define non-negative trace class operators $A(E)$ and $B(E)$.

In the next lemma we identify the densities of the absolutely continuous parts of μ^1 and ν^1 . The proof of this lemma follows directly from the definitions.

Lemma 4.14. *The functions $E \mapsto \operatorname{tr} A(E)$, respectively $E \mapsto \operatorname{tr} B(E)$, are locally integrable Lebesgue densities of the absolutely continuous parts of μ^1 , respectively ν^1 .*

We will need an auxiliary statement for the main result of this subsection.

Lemma 4.15. *Let μ be a locally finite Borel measure on \mathbb{R} . Let $c_0 > 0$ and $0 < \varepsilon < \delta < c_0$. Then for a.e. $x_0 \in \mathbb{R}$ there is a constant C , depending on x_0 , c_0 and μ , such that for all $t > 0$*

$$\int_{[x_0, x_0+\delta]} d\mu(x) e^{-t(x-x_0)} \leq C \frac{1 - e^{-t\delta}}{t}$$

and

$$\int_{[x_0+\varepsilon, x_0+\delta]} d\mu(x) e^{-t(x-x_0)} \leq C e^{-t\varepsilon/2} \frac{1 - e^{-t\delta/2}}{t/2} \leq C \frac{e^{-t\varepsilon/2}}{t/2}.$$

The exceptional set of values of x_0 for which the assertion does not hold depends neither on c_0 , ε nor δ .

Proof. The constant

$$C := \sup_{\eta \in (0, c_0)} \frac{1}{\eta} \mu([x_0, x_0 + \eta])$$

is finite for a.e. $x_0 \in \mathbb{R}$. We compute using Tonelli's theorem

$$\begin{aligned} & \int_{[x_0, x_0 + \delta]} d\mu(x) e^{-t(x-x_0)} \\ &= \int_{[x_0, x_0 + \delta]} d\mu(x) \left(e^{-t\delta} + t \int_x^{x_0 + \delta} d\xi e^{-t(\xi-x_0)} \right) \\ &= \delta e^{-t\delta} \frac{1}{\delta} \mu([x_0, x_0 + \delta]) + t \int_{x_0}^{x_0 + \delta} d\xi \int_{[x_0, \xi]} d\mu(x) e^{-t(\xi-x_0)} \\ &\leq C\delta e^{-t\delta} + t \int_{x_0}^{x_0 + \delta} d\xi e^{-t(\xi-x_0)} \frac{\xi - x_0}{\xi - x_0} \mu([x_0, \xi]) \\ &\leq C\delta e^{-t\delta} + Ct \int_0^\delta d\xi \xi e^{-t\xi} \\ &= C \frac{1 - e^{-t\delta}}{t}. \end{aligned}$$

The second assertion follows from the first one and the bound $e^{-t(x-x_0)} \leq e^{-t\varepsilon/2} e^{-t(x-x_0)/2}$ for $\varepsilon \leq x - x_0 \leq \delta$. □

Definition 4.16. (i) For $k \in \mathbb{N}$, we define

$$I_k := \int_{(0, \infty)^k} du \frac{|u|_1 e^{-|u|_1}}{\prod_{j=1}^k (u_j + u_{j+1})},$$

where $u_{k+1} := u_1$ for $u \in \mathbb{R}^k$.

(ii) We define discontinuous L -independent functions $\chi^\pm: \mathbb{R} \rightarrow [0, 1]$ by

$$\chi^- := \max\{\chi_L^-, 1_{[-E_0, E)}\} \quad \text{and} \quad \chi^+ := \max\{\chi_L^+, 1_{(E, E_0]}\}.$$

Remarks 4.17. (i) The integral I_k will be discussed further in Subsection 4.5; in particular, I_k is finite for every $k \in \mathbb{N}$.

(ii) The functions χ_L^\pm converge pointwise to χ^\pm as $L \rightarrow \infty$. They are obtained from replacing the smooth L -dependent part by a discontinuous step at E . Figure 2 illustrates the behaviour of χ^\pm .

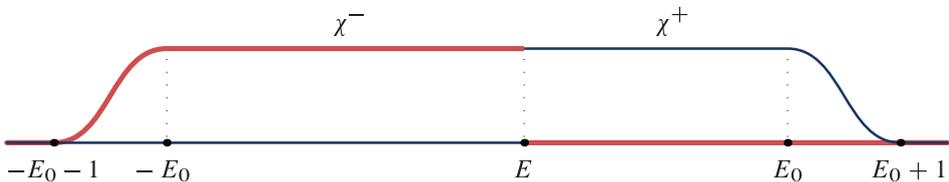


Figure 2. Sketch of the discontinuous cut-off functions χ^\pm .

The following lemma is the main result of the current section.

Lemma 4.18. *There is a null set $\mathcal{N} \subseteq \mathbb{R}$ which does not depend on a, n and E_0 , such that for every $E \in [-E_0, E_0] \setminus \mathcal{N}$,*

$$\int_0^\infty dt t^{2n-1} \int_{(0,\infty)^n \times (0,\infty)^n} d(u, v) (|u|_1 + |v|_1) e^{-|u|_1 - |v|_1} \left[\operatorname{tr} \left\{ \prod_{j=1}^n \sqrt{V} \chi_L^-(H) e^{(u_j + v_{j-1})t(H-E)} \sqrt{V} \right. \right. \\ \left. \left. \sqrt{V} \chi_L^+(H') e^{-(u_j + v_j)t(H'-E)} \sqrt{V} \right\} \right. \\ \left. - e^{-tL^{-a}} \operatorname{tr} \left\{ \prod_{j=1}^n \sqrt{V} \chi^-(H) e^{(u_j + v_{j-1})t(H-E)} \sqrt{V} \right. \right. \\ \left. \left. \sqrt{V} \chi^+(H') e^{-(u_j + v_j)t(H'-E)} \sqrt{V} \right\} \right] = O(1) \tag{46}$$

as $L \rightarrow \infty$, where the $O(1)$ -term depends on a, n, E and E_0 .

Proof. First, notice that if f_j, g_j are bounded measurable functions of compact support for $1 \leq j \leq n$, then

$$\operatorname{tr} \left\{ \prod_{j=1}^n |\sqrt{V} f_j(H) V g_j(H') \sqrt{V}| \right\} \leq \prod_{j=1}^n \operatorname{tr} \left\{ \sqrt{V} f_j(H) \sqrt{V} \right\} \operatorname{tr} \left\{ \sqrt{V} g_j(H') \sqrt{V} \right\} \\ = \int_{\mathbb{R}^n} d\mu^n(x) \int_{\mathbb{R}^n} dv^n(y) \prod_{j=1}^n f_j(x_j) g_j(y_j), \tag{47}$$

where we wrote μ^n and ν^n for the n -fold product measure of μ^1 and ν^1 , respectively.

For brevity, let $\delta := L^{-a}$. We introduce a vector $\alpha \in (0, \infty)^{2n}$ via

$$\alpha_{2j-1} := u_j + v_{j-1} \quad \text{and} \quad \alpha_{2j} := u_j + v_j$$

for $1 \leq j \leq n$ and operators

$$A_k^{(L)} := \begin{cases} \sqrt{V} \chi_{(L)}^-(H) e^{\alpha_k t(H-E)} \sqrt{V} & \text{for } k \text{ odd,} \\ \sqrt{V} \chi_{(L)}^+(H') e^{-\alpha_k t(H'-E)} \sqrt{V} & \text{for } k \text{ even} \end{cases}$$

for $1 \leq k \leq 2n$. The difference of operator products in (46) then equals

$$\prod_{j=1}^{2n} A_j^L - e^{-t\delta} \prod_{j=1}^{2n} A_j = e^{-t\delta} \left(\prod_{j=1}^{2n} A_j^L - \prod_{j=1}^{2n} A_j \right) + (1 - e^{-t\delta}) \prod_{j=1}^{2n} A_j^L, \quad (48)$$

where as in (40),

$$\prod_{j=1}^{2n} A_j^L - \prod_{j=1}^{2n} A_j = \sum_{k=1}^{2n} A_1 \dots A_{k-1} (A_k^L - A_k) A_{k+1}^L \dots A_{2n}^L. \quad (49)$$

We will treat the two terms on the right-hand side of (48) individually. For the first term, we estimate the k th term in (49). We will carry out the argument in the case where k is even. The argument is similar for odd k . Since $0 \leq \chi^+ - \chi_L^+ \leq 1_{[E, E+2\delta]}$, $\chi_{(L)}^- \leq 1_{[-E_0-1, E]}$ and $\chi_{(L)}^+ \leq 1_{[E, E_0+1]}$, (47) implies

$$\begin{aligned} & \text{tr}\{|A_1 \dots A_{k-1} (A_k^L - A_k) A_{k+1}^L \dots A_{2n}^L|\} \\ & \leq \int_{[-E_0-1, E]^n} d\mu^n(x) \int_{[E, E_0+1]^n} d\nu^n(y) 1_{[E, E+2\delta]}(y_k) \\ & \quad \exp\left(-t \sum_{j=1}^n ((u_j + v_j)(y_j - E) - (u_j + v_{j-1})(x_j - E))\right) \\ & \leq C \frac{1 - e^{-2(u_k+v_k)t\delta}}{t} \frac{1}{t^{2n-1} \prod_{j=1}^n (u_j + v_j)(u_j + v_{j-1})}, \end{aligned}$$

where C is some finite constant and the last inequality follows for a.e. $E \in [-E_0, E_0]$ from applying Lemma 4.15 to every integral and the estimate $1 - e^{-tx} \leq 1$ to all but the k th term. Using the bound $1 - e^{-2(u_k+v_k)t\delta} \leq 2(u_k + v_k)t\delta$, we conclude

$$e^{-t\delta} \text{tr} \left\{ \left| \prod_{j=1}^{2n} A_j^L - \prod_{j=1}^{2n} A_j \right| \right\} \leq 4C \frac{\delta e^{-t\delta}}{t^{2n-1} \prod_{j=1}^n (u_j + v_j)(u_j + v_{j-1})},$$

and therefore

$$\begin{aligned} & \int_0^\infty dt t^{2n-1} \int_{(0, \infty)^n \times (0, \infty)^n} d(u, v) (|u|_1 + |v|_1) e^{-|u|_1 - |v|_1} \\ & \quad e^{-t\delta} \text{tr} \left\{ \left| \prod_{j=1}^{2n} A_j^L - \prod_{j=1}^{2n} A_j \right| \right\} \\ & \leq 4C \int_0^\infty dt \delta e^{-t\delta} \int_{(0, \infty)^n \times (0, \infty)^n} d(u, v) \frac{(|u|_1 + |v|_1)^2 e^{-|u|_1 - |v|_1}}{\prod_{j=1}^n (u_j + v_j)(u_j + v_{j-1})}. \end{aligned}$$

Here, the t -integral yields 1 for every $\delta > 0$, and the (u, v) -integral is finite since

$$\begin{aligned} & \int_{(0,\infty)^n \times (0,\infty)^n} d(u, v) \frac{(|u|_1 + |v|_1)^2 e^{-|u|_1 - |v|_1}}{\prod_{j=1}^n (u_j + v_j)(u_j + v_{j-1})} \\ & \leq \int_{(0,\infty)^n \times (0,\infty)^n} d(u, v) \frac{(|u|_1 + |v|_1) e^{-|u|_1/2 - |v|_1/2}}{\prod_{j=1}^n (u_j + v_j)(u_j + v_{j-1})} \\ & = 2I_{2n} < \infty, \end{aligned}$$

with I_{2n} as in Definition 4.16. This shows that the integral of the trace norm of the first term on the right-hand side of (48) yields an error that remains finite as $L \rightarrow \infty$.

The trace norm of the second term on the right-hand side of (48) is

$$\begin{aligned} & (1 - e^{-t\delta}) \operatorname{tr} \left\{ \left| \prod_{j=1}^{2n} A_j^L \right| \right\} \\ & \leq (1 - e^{-t\delta}) \int_{[-E_0-1, E-\delta]^n} d\mu^n(x) \int_{[E+\delta, E_0+1]^n} dv^n(y) \\ & \quad \exp \left(-t \sum_{j=1}^n ((u_j + v_j)(y_j - E) - (u_j + v_{j-1})(x_j - E)) \right) \\ & \leq C(1 - e^{-t\delta}) \prod_{j=1}^n \frac{e^{-(u_j + v_{j-1})t\delta/2} e^{-(u_j + v_j)t\delta/2}}{(u_j + v_{j-1})(u_j + v_j)(t/2)^2} \\ & = 2^{2n} C \frac{(1 - e^{-t\delta}) e^{-(|u|_1 + |v|_1)t\delta} t^{-2n}}{\prod_{j=1}^n (u_j + v_j)(u_j + v_{j-1})}, \end{aligned}$$

where the first inequality is a consequence of $\chi_{(L)}^- \leq 1_{[-E_0-1, E]}$ and $\chi_{(L)}^+ \leq 1_{[E, E_0+1]}$, and the second inequality follows for a.e. $E \in [-E_0, E_0]$ from Lemma 4.15. Now, we perform the t - and (u, v) -integration

$$\begin{aligned} & \int_{(0,\infty)^n \times (0,\infty)^n} d(u, v) (|u|_1 + |v|_1) e^{-|u|_1 - |v|_1} \int_0^\infty dt t^{2n-1} \\ & \quad \frac{(1 - e^{-t\delta}) e^{-(|u|_1 + |v|_1)t\delta} t^{-2n}}{\prod_{j=1}^n (u_j + v_j)(u_j + v_{j-1})} \\ & = \int_{(0,\infty)^n \times (0,\infty)^n} d(u, v) \int_0^\infty dt \frac{(1 - e^{-t}) e^{-(1+t)(|u|_1 + |v|_1)} (|u|_1 + |v|_1)}{t \prod_{j=1}^n (u_j + v_j)(u_j + v_{j-1})} \\ & = I_{2n} \int_0^\infty dt \frac{1 - e^{-t}}{t^2 + t} < \infty, \end{aligned}$$

where we performed the successive changes of variables $t\delta \rightsquigarrow t$ and $(1 + t)(u, v) \rightsquigarrow (u, v)$. □

Corollary 4.19. *For a.e. $E \in [-E_0, E_0]$, we have*

$$\begin{aligned} & \text{tr}\{(1_{(-\infty, E]}(H_L)1_{(E, \infty)}(H'_L))^n\} \\ & \geq \int_{(0, \infty)^n \times (0, \infty)^n} d(u, v) (|u|_1 + |v|_1) e^{-|u|_1 - |v|_1} \int_0^\infty dt t^{2n-1} e^{-tL^{-a}} \\ & \quad \text{tr}\left\{\prod_{j=1}^n \sqrt{V} \chi^-(H) e^{(u_j + v_{j-1})t(H-E)} V \chi^+(H') e^{-(u_j + v_j)t(H'-E)} \sqrt{V}\right\} \\ & + O(1) \end{aligned} \tag{50}$$

as $L \rightarrow \infty$. The null set of exceptional energies does not depend on a, n and E_0 .

Proof. We combine Corollary 4.11 and Lemma 4.18. □

4.4. The logarithmic divergence. Throughout this subsection, we fix $a \in (0, 1)$, $n \in \mathbb{N}$ and $E_0 \geq 1$.

The goal is to determine the asymptotics of the right-hand side of (50).

Lemma 4.20. *For a.e. $E \in [-E_0, E_0]$,*

$$\begin{aligned} A_t(E) & := \sqrt{V} t e^{t(H-E)} \chi^-(H) \sqrt{V} \longrightarrow A(E), \\ B_t(E) & := \sqrt{V} t e^{-t(H'-E)} \chi^+(H') \sqrt{V} \longrightarrow B(E) \end{aligned} \tag{51}$$

as $t \rightarrow \infty$, where the convergences are in trace class. Moreover,

$$\begin{aligned} \sup_{t \geq 0} \|A_t(E)\| & \leq \sup_{t \geq 0} \text{tr} A_t(E) < \infty, \\ \sup_{t \geq 0} \|B_t(E)\| & \leq \sup_{t \geq 0} \text{tr} B_t(E) < \infty. \end{aligned} \tag{52}$$

Proof. We follow [11, Lemma 3.16] and treat the operator $B_t(E)$; the assertions for $A_t(E)$ can be proved using analogous arguments. Recall that $B_t(E)$ is non-negative. To prove (51), we show (1) convergence of the trace norms and (2) weak convergence of the operators. Together, this implies convergence in trace class via [30, Addendum H].

For the trace norms, we compute

$$\begin{aligned} \text{tr} B_t(E) & = \text{tr}\{\sqrt{V} t e^{-t(H'-E)} \chi^+(H') \sqrt{V}\} \\ & = \int_{[E, E_0]} dv^1(y) t e^{-t(y-E)} + \int_{[E_0, E_0+1]} dv^1(y) \chi^+(y) t e^{-t(y-E)}, \end{aligned} \tag{53}$$

where the second term converges to zero as $t \rightarrow \infty$ for $E < E_0$. The first term can be written as

$$\int_{[E, E_0]} dv^1(y) t e^{-t(y-E)} = (v_{E_0}^1 * \varrho_t)(E), \tag{54}$$

where we introduced the restricted (finite) measure $v_{E_0}^1(M) := v^1(M \cap [-E_0, E_0])$ for $M \in \text{Borel}(\mathbb{R})$ and the approximation of the identity $x \mapsto \varrho_t(x) := t e^{tx} 1_{(-\infty, 0)}(x)$. As $t \rightarrow \infty$, the convolution in (54) converges for a.e. $E \in [-E_0, E_0]$ to $\frac{dv_{E_0}^1}{dE} = \text{tr } B(E)$, see e.g. [21, Subsection 2.4.1]. Thus, the trace norm of $B_t(E)$ converges to that of $B(E)$ as $t \rightarrow \infty$. This, together with the continuity of $[0, \infty) \ni t \mapsto \text{tr } B_t(E)$, which can be seen from (53), implies (52).

For the weak convergence, take some dense countable set $\mathcal{D} \subseteq L^2(\mathbb{R}^d)$. Then by a similar delta-argument as above,

$$\lim_{t \rightarrow \infty} \langle \varphi, B_t(E) \psi \rangle = \langle \varphi, B(E) \psi \rangle$$

for all $\varphi, \psi \in \mathcal{D}$ and all $E \in [-E_0, E_0]$ outside a null set depending on \mathcal{D} . Together with (52), this proves weak convergence to $B(E)$ for a.e. $E \in [-E_0, E_0]$, see [31, Theorem 4.26]. □

The following quantity will enter the asymptotics we set out to prove.

Definition 4.21. For $E \in \mathbb{R} \setminus \mathcal{N}_0$, let

$$\eta_{2n}(E) := \text{tr} \{ (A(E)B(E))^n \} \tag{55}$$

and extend it trivially to a function $\eta_{2n}: \mathbb{R} \rightarrow [0, \infty)$. The non-negativity of (55) can be seen from the cyclicity of the trace.

The next corollary will show that the trace expression on the right-hand side of (50), times an appropriate power of t , converges to $\eta_{2n}(E)$ in the $t \rightarrow \infty$ limit.

Corollary 4.22. Let $\alpha_1, \dots, \alpha_n, \beta_1, \dots, \beta_n > 0$. Then for a.e. $E \in [-E_0, E_0]$

$$t^{2n} \text{tr} \left\{ \prod_{j=1}^n \sqrt{V} \chi^-(H) \alpha_j e^{\alpha_j t(H-E)} V \chi^+(H') \beta_j e^{-\beta_j t(H'-E)} \sqrt{V} \right\} \longrightarrow \eta_{2n}(E)$$

as $t \rightarrow \infty$.

Proof. Using the notation of Lemma 4.20, we have to show

$$\left| \operatorname{tr} \left\{ \prod_{j=1}^n A_{\alpha_j t}(E) B_{\beta_j t}(E) \right\} - \operatorname{tr} \{ (A(E)B(E))^n \} \right| \longrightarrow 0 \tag{56}$$

as $t \rightarrow \infty$. By Lemma 4.20, $\operatorname{tr} |A_{\alpha_j t}(E) - A(E)| \rightarrow 0$ and $\operatorname{tr} |B_{\beta_j t}(E) - B(E)| \rightarrow 0$ as $t \rightarrow \infty$, while $\sup_{t \geq 0} \|A_t(E)\|$ and $\sup_{t \geq 0} \|B_t(E)\|$ are finite. Writing the difference of operator products in (56) as in (40), this proves the corollary. \square

Lemma 4.23. *Let $f \in L^1_{\text{loc}}(\mathbb{R})$ and suppose $\lim_{t \rightarrow \infty} f(t)$ exists. Then*

$$\lim_{t \rightarrow \infty} f(t) = - \lim_{s \downarrow 0} \frac{1}{\ln s} \int_1^\infty dt t^{-1} e^{-st} f(t).$$

Proof. Take a compact interval $[s_0, c] \subseteq (0, \infty)$. Then

$$\frac{d}{ds} \int_1^\infty dt t^{-1} e^{-st} f(t) = - \int_1^\infty dt e^{-st} f(t) \tag{57}$$

for $s \in [s_0, c]$, because $|\frac{d}{ds} t^{-1} e^{-st} f(t)| \leq e^{-s_0 t} |f(t)|$, which is integrable on $[1, \infty)$. Therefore (57) holds for all $s > 0$. If $\lim_{s \downarrow 0} \int_1^\infty dt t^{-1} e^{-st} f(t)$ exists, then $\lim_{t \rightarrow \infty} f(t) = 0$ and the assertion holds. Otherwise,

$$\begin{aligned} - \lim_{s \downarrow 0} \frac{1}{\ln s} \int_1^\infty dt t^{-1} e^{-st} f(t) &= \lim_{s \downarrow 0} \frac{1}{1/s} \int_1^\infty dt e^{-st} f(t) \\ &= \lim_{s \downarrow 0} s \int_0^\infty dt e^{-st} f(t) \\ &= \lim_{t \rightarrow \infty} f(t), \end{aligned}$$

where the last equality is the statement of the classical final-value theorem, see [6, Theorem 34.3]. \square

We are now ready to compute the asymptotics of the right-hand side of (50).

Theorem 4.24. *For a.e. $E \in [-E_0, E_0]$,*

$$\begin{aligned} &\lim_{L \rightarrow \infty} \frac{1}{a \ln L} \int_0^\infty dt t^{2n-1} e^{-tL^{-a}} \int_{(0, \infty)^n \times (0, \infty)^n} d(u, v) (|u|_1 + |v|_1) e^{-|u|_1 - |v|_1} \\ &\operatorname{tr} \left\{ \prod_{j=1}^n \sqrt{V} \chi^-(H) e^{(u_j + v_{j-1})t(H-E)} V \chi^+(H') e^{-(u_j + v_j)t(H'-E)} \sqrt{V} \right\} \\ &= I_{2n} \eta_{2n}(E). \end{aligned} \tag{58}$$

Proof. Let $u, v \in (0, \infty)^n$ and define

$$Z(u, v) := \prod_{j=1}^n (u_j + v_{j-1})(u_j + v_j).$$

Using the notation of Lemma 4.20, we see that

$$\begin{aligned} Z(u, v)t^{2n} \operatorname{tr} & \left\{ \prod_{j=1}^n \sqrt{V} \chi^-(H) e^{(u_j+v_{j-1})t(H-E)} \sqrt{V} \right. \\ & \left. \sqrt{V} \chi^+(H') e^{-(u_j+v_j)t(H'-E)} \sqrt{V} \right\} \\ & = \operatorname{tr} \left\{ \prod_{j=1}^n A_{(u_j+v_{j-1})t}(E) B_{(u_j+v_j)t}(E) \right\}, \end{aligned}$$

where

$$\left| \operatorname{tr} \left\{ \prod_{j=1}^n A_{(u_j+v_{j-1})t}(E) B_{(u_j+v_j)t}(E) \right\} \right| \leq \left(\sup_{t \geq 0} \operatorname{tr} \{A_t(E)\} \sup_{t \geq 0} \operatorname{tr} \{B_t(E)\} \right)^n \tag{59}$$

$< \infty$.

By Corollary 4.22,

$$\lim_{t \rightarrow \infty} \operatorname{tr} \left\{ \prod_{j=1}^n A_{(u_j+v_{j-1})t}(E) B_{(u_j+v_j)t}(E) \right\} = \eta_{2n}(E)$$

for all $u, v \in (0, \infty)^n$. By Remark 4.17 (i),

$$I_{2n} = \int_{(0, \infty)^n \times (0, \infty)^n} d(u, v) \frac{(|u|_1 + |v|_1) e^{-|u|_1 - |v|_1}}{Z(u, v)} < \infty. \tag{60}$$

Equations (59) and (60) supply the assumptions of the dominated convergence theorem. It yields the convergence

$$\lim_{t \rightarrow \infty} f(t) = I_{2n} \eta_{2n}(E)$$

for

$$\begin{aligned} f(t) := & \int_{(0, \infty)^n \times (0, \infty)^n} d(u, v) (|u|_1 + |v|_1) e^{-|u|_1 - |v|_1} t^{2n} \\ & \operatorname{tr} \left\{ \prod_{j=1}^n \sqrt{V} \chi^-(H) e^{(u_j+v_{j-1})t(H-E)} \sqrt{V} \right. \\ & \left. \sqrt{V} \chi^+(H') e^{-(u_j+v_j)t(H'-E)} \sqrt{V} \right\} \end{aligned}$$

$$= \int_{(0,\infty)^n \times (0,\infty)^n} d(u, v) \frac{(|u|_1 + |v|_1)e^{-|u|_1 - |v|_1}}{Z(u, v)} \operatorname{tr} \left\{ \prod_{j=1}^n A_{(u_j + v_{j-1})t}(E) B_{(u_j + v_j)t}(E) \right\},$$

where $t > 0$. The assertion (58) follows from

$$- \lim_{L \rightarrow \infty} \frac{1}{\ln(L^{-a})} \int_0^\infty dt t^{-1} e^{-tL^{-a}} f(t) = \lim_{t \rightarrow \infty} f(t),$$

which is a consequence of Lemma 4.23 and of

$$\sup_{L > 1} \int_0^1 dt t^{-1} e^{-tL^{-a}} f(t) < \infty. \quad \square$$

Corollary 4.25. *For a.e. $E \in \mathbb{R}$, the estimate*

$$\operatorname{tr}\{(1_{(-\infty, E]}(H_L)1_{(E, \infty)}(H'_L))^n\} \geq \ln L I_{2n} \eta_{2n}(E) + o(\ln L)$$

holds, where the $o(\ln L)$ -term depends on n and E .

Proof. We deduce from Theorem 4.24, Corollary 4.19 and from the arbitrariness of E_0 that

$$\liminf_{L \rightarrow \infty} \frac{\operatorname{tr}\{(1_{(-\infty, E]}(H_L)1_{(E, \infty)}(H'_L))^n\}}{\ln L} \geq a I_{2n} \eta_{2n}(E) \quad (61)$$

for arbitrary $a \in (0, 1)$ and a.e. $E \in \mathbb{R}$. Thus (61) holds for $a = 1$ and a.e. $E \in \mathbb{R}$. By definition of the limit inferior, this implies the claim. \square

4.5. A multi-dimensional integral related to the Hilbert matrix. In this subsection, we compute the coefficient of $\eta_{2n}(E)$ in the asymptotics in Corollary 4.25, i.e., we compute the integral

$$I_n = \int_{(0,\infty)^n} du \frac{|u|_1 e^{-|u|_1}}{\prod_{j=1}^n (u_j + u_{j+1})}$$

in Definition 4.16 (i). Here, we use the convention $u_{n+1} = u_1$ for $u \in \mathbb{R}^n$.

We prove

Theorem 4.26. *Let $n \in \mathbb{N}_{\geq 2}$. Then*

$$I_n = (2\pi)^{n-2} \frac{\left(\Gamma\left(\frac{n}{2}\right)\right)^2}{\Gamma(n)}.$$

This implies

$$I_{2n} = (2\pi)^{2n-2} \frac{[(n-1)!]^2}{(2n-1)!} = n J_{2n}$$

for $n \in \mathbb{N}$, where J_{2n} was defined in (24).

We begin with an elementary lemma.

Lemma 4.27. *Let $n \in \mathbb{N}_{\geq 2}$. Then*

$$I_n = \frac{n}{2} \int_{(0,\infty)^n} du \frac{e^{-|u|_1}}{\prod_{j=1}^{n-1} (u_j + u_{j+1})}. \quad (62)$$

Proof. Using the symmetry of I_n in the components of u , we compute

$$\begin{aligned} I_n &= \frac{1}{2} \int_{(0,\infty)^n} du e^{-|u|_1} \frac{2|u|_1}{\prod_{j=1}^n (u_j + u_{j+1})} \\ &= \frac{1}{2} \sum_{k=1}^n \int_{(0,\infty)^n} du e^{-|u|_1} \frac{u_k + u_{k+1}}{\prod_{j=1}^n (u_j + u_{j+1})} \\ &= \frac{n}{2} \int_{(0,\infty)^n} du \frac{e^{-|u|_1}}{\prod_{j=1}^{n-1} (u_j + u_{j+1})}. \quad \square \end{aligned}$$

In the sequel, we will work with the Rosenblum–Rovnyak integral operator $T: L^2((0, \infty)) \rightarrow L^2((0, \infty))$, see [27] and [28], defined by

$$(Tf)(x) := \int_0^\infty dy \frac{e^{-(x+y)/2}}{x+y} f(y) \quad (63)$$

for $f \in L^2((0, \infty))$ and $x \in (0, \infty)$. This operator can be explicitly diagonalised. Following [33, Section 4.2], we define the Kontorovich–Lebedev transform, i.e. the unitary operator $U: L^2((0, \infty)) \rightarrow L^2((0, \infty))$,

$$(Uf)(k) := \pi^{-1} \sqrt{k \sinh(2\pi k)} |\Gamma(1/2 - ik)| \int_0^\infty dx x^{-1} W_{0,ik}(x) f(x) \quad (64)$$

for $f \in L^2((0, \infty))$ and $k \in (0, \infty)$, where $W_{0,ik}$ denotes the Whittaker function, see [22, Section 13.14] or [12, Section 9.22–9.23]. Then, the spectral representation due to Rosenblum reads

$$(UTf)(k) = \frac{\pi}{\cosh(k\pi)}(Uf)(k) \tag{65}$$

for $f \in L^2((0, \infty))$ and $k \in (0, \infty)$, see [33, Proposition 4.1].

Proof of Theorem 4.26. Let $n \in \mathbb{N}_{\geq 2}$. From (62) and (63), we see that

$$\frac{2}{n}I_n = \langle \phi_0, T^{n-1}\phi_0 \rangle_{L^2((0, \infty))} \tag{66}$$

with $\phi_0(x) := e^{-x/2}$. From (65) and (66), we obtain

$$\frac{2}{n}I_n = \langle U\phi_0, UT^{n-1}\phi_0 \rangle_{L^2((0, \infty))} = \int_0^\infty dk |(U\phi_0)(k)|^2 \left(\frac{\pi}{\cosh(k\pi)}\right)^{n-1}. \tag{67}$$

In order to compute $U\phi_0$, we employ the classical formula

$$|\Gamma(1/2 - ik)|^2 = \frac{\pi}{\cosh(k\pi)} \tag{68}$$

for $k \in \mathbb{R}$, which is a consequence of the reflection formula for the Gamma function, and

$$\int_0^\infty dx x^{-1} W_{0,ik}(x)e^{-x/2} = \frac{\pi}{\cosh(k\pi)} \quad (k > 0), \tag{69}$$

which follows from the special case $z = 1/2$ and $\nu = \kappa = 0$ in [22, eq. 13.23.4]. From (64), (68) and (69), we deduce

$$|(U\phi_0)(k)|^2 = 2\pi k \frac{\sinh(k\pi)}{(\cosh(k\pi))^2}$$

for $k > 0$. Inserting this into (67) yields

$$\frac{2}{n}I_n = 2\pi^{n-2} \int_0^\infty dk k \frac{\sinh k}{(\cosh k)^{n+1}} = \frac{2\pi^{n-2}}{n} \int_0^\infty dk \frac{1}{(\cosh k)^n},$$

where we applied the substitution $k \rightsquigarrow k/\pi$ and integrated by parts. This integral can be evaluated using the substitution $x = (\cosh k)^{-2}$:

$$\frac{2}{n}I_n = \frac{\pi^{n-2}}{n} \int_0^1 dx x^{n/2-1}(1-x)^{-1/2} = \frac{\pi^{n-2}}{n} B(n/2, 1/2),$$

where B denotes Euler’s Beta function. The claim follows from [12, eq. 8.384.4 and eq. 8.384.1]. □

Remark 4.28. The Rosenblum–Rovnyak operator is the special case $T = \mathcal{H}_0$ in [27, eq. (2.3)] and is unitarily equivalent to the Hilbert matrix

$$H: \ell^2(\mathbb{N}_0) \longrightarrow \ell^2(\mathbb{N}_0),$$

given by

$$(Hc)_j = \sum_{k=0}^{\infty} \frac{c_k}{j+k+1}$$

for $j \in \mathbb{N}_0$ and $c \in \ell^2(\mathbb{N}_0)$. In analogy to (66), the representation

$$I_n = \frac{n}{2} \langle e^{(0)}, H^{n-1} e^{(0)} \rangle_{\ell^2(\mathbb{N}_0)}$$

holds with $e^{(0)} := (1, 0, \dots) \in \ell^2(\mathbb{N}_0)$.

4.6. Relations to scattering theory. In order to complete the proof of Theorem 3.4 we need to relate the coefficient $\eta_{2n}(E)$ in Definition 4.21 to the transition matrix from scattering theory. We begin with a definition.

Definition 4.29. Let $\mathcal{H}_{ac}(H)$ be the absolutely continuous subspace of the self-adjoint operator H . Then $\mathcal{H}_{ac}(H)$ can be decomposed into a direct integral

$$\int_{\sigma_{ac}(H)}^{\oplus} dE \mathcal{H}_E$$

where \mathcal{H}_E is a Hilbert space for every $E \in \sigma_{ac}(H)$. The operator H acts on \mathcal{H}_E by multiplication with the identity, see [32, §1.5]. This means that a vector $f \in \mathcal{H}_{ac}(H)$ corresponds to a vector-valued function $E \mapsto f_E \in \mathcal{H}_E$, and Hf corresponds to $E \mapsto E f_E$.

The transition matrix T_E acts as a bounded operator on \mathcal{H}_E . Moreover, we have the following representation.

Lemma 4.30. *The limit*

$$\Phi_{\pm}(E) := \lim_{\varepsilon \downarrow 0} (I + \sqrt{V}(E \pm i\varepsilon - H')^{-1} \sqrt{V})$$

exists in the sense of convergence in operator norm for a.e. $E \in \mathbb{R}$. Moreover, there exists an operator $U(E): \mathcal{H}_{ac}(H) \rightarrow \mathcal{H}_E$ such that $U(E)^* U(E)$ is the identity on $\text{ran } \sqrt{A(E)}$ and the transition matrix $T_E: \mathcal{H}_E \rightarrow \mathcal{H}_E$ satisfies

$$T_E = U(E) \tilde{T}(E) U(E)^*,$$

where

$$\tilde{T}(E) := -2\pi i \sqrt{A(E)} \Phi_+(E) \sqrt{A(E)}.$$

Proof. This is a result in abstract scattering theory, see e.g. [4, §7], [32, §5.5], or [3, p. 394]. A detailed proof is given in [19]. \square

Corollary 4.31. *The identity*

$$\tilde{T}(E)^* \tilde{T}(E) = (2\pi)^2 \sqrt{A(E)} B(E) \sqrt{A(E)} \tag{70}$$

holds for a.e. $E \in \mathbb{R}$. In particular,

$$\|T_E\|_{2n}^{2n} = (2\pi)^{2n} \operatorname{tr}\{(A(E)B(E))^n\} = (2\pi)^{2n} \eta_{2n}(E) \tag{71}$$

for every $n \in \mathbb{N}$, where $\|T_E\|_{2n} := \sqrt[2n]{\operatorname{tr}|T_E|^{2n}}$ is the $2n$ -Schatten norm of T_E .

Proof. The operators $A(E)$ and $B(E)$ can be expressed as the operator limits

$$-\pi A(E) = \lim_{\varepsilon \downarrow 0} \operatorname{Im}(\sqrt{V}(E + i\varepsilon - H)^{-1} \sqrt{V})$$

$$-\pi B(E) = \lim_{\varepsilon \downarrow 0} \operatorname{Im}(\sqrt{V}(E + i\varepsilon - H')^{-1} \sqrt{V})$$

which exist for a.e. $E \in \mathbb{R}$, see [4, Lemma 4.5]. From this and the second resolvent identity $(z - H')^{-1} - (z - H)^{-1} = (z - H')^{-1} V (z - H)^{-1}$ for $z \in \mathbb{C} \setminus \mathbb{R}$, the statement

$$\begin{aligned} & \Phi_+(E)^* A(E) \Phi_+(E) \\ &= -\frac{1}{\pi} \lim_{\varepsilon \downarrow 0} \{ (I + \sqrt{V}(E - i\varepsilon - H')^{-1} \sqrt{V}) (\operatorname{Im}(\sqrt{V}(E + i\varepsilon - H)^{-1} \sqrt{V})) \\ & \quad (I + \sqrt{V}(E + i\varepsilon - H')^{-1} \sqrt{V}) \} \\ &= \frac{1}{2\pi i} \lim_{\varepsilon \downarrow 0} \{ \sqrt{V} (I + (E - i\varepsilon - H')^{-1} V) \\ & \quad ((E - i\varepsilon - H)^{-1} - (E + i\varepsilon - H)^{-1}) \\ & \quad (I + V(E + i\varepsilon - H')^{-1}) \sqrt{V} \} \\ &= \frac{1}{2\pi i} \lim_{\varepsilon \downarrow 0} \{ \sqrt{V} ((E - i\varepsilon - H')^{-1} - (E + i\varepsilon - H')^{-1}) \sqrt{V} \} \\ &= B(E) \end{aligned}$$

follows and yields (70). The unitary equivalence on $\operatorname{ran} \sqrt{A(E)}$ in Lemma 4.30 then implies (71). \square

Corollary 4.31 yields the following theorem.

Theorem 4.32. *Let $n \in \mathbb{N}$. For a.e. $E \in \mathbb{R}$*

$$\eta_{2n}(E) = \text{tr}(|T_E/(2\pi)|^{2n}),$$

where $T_E : \mathcal{H}_E \rightarrow \mathcal{H}_E$ is the transition matrix for the energy E .

A. Positivity of the exponent

Here we consider the special case $V_0 = 0$ and show that the decay exponent $\gamma(E)$ in (13) is strictly positive for a.e. $E > 0$. Throughout this appendix, we assume that $V \neq 0$ satisfies (V).

Theorem A.1. *Let $V_0 = 0$. Let $E > 0$. Then the operator $A(E)$ from (45) has the integral kernel*

$$A(E; x, y) = \frac{E^{d/2-1}}{2(2\pi)^d} \sqrt{V(x)} \sqrt{V(y)} \int_{\mathbb{S}^{d-1}} dS(\xi) e^{i\sqrt{E}\xi \cdot (x-y)} \quad (72)$$

for a.e. $x, y \in \mathbb{R}^d$. Here, dS stands for integration with respect to the surface measure on the unit sphere $\mathbb{S}^{d-1} \subseteq \mathbb{R}^d$.

Proof. Let $\varepsilon > 0$ and $f \in L^2(\mathbb{R}^d)$. Then, using the Fourier transform and spherical coordinates, we compute for a.e. $x \in \mathbb{R}^d$

$$\begin{aligned} & (\sqrt{V} 1_{(E-\varepsilon, E+\varepsilon)}(-\Delta) \sqrt{V} f)(x) \\ &= \frac{\sqrt{V(x)}}{(2\pi)^d} \int_{\mathbb{R}^d} dk \int_{\mathbb{R}^d} dy 1_{(E-\varepsilon, E+\varepsilon)}(|k|^2) e^{ik \cdot (x-y)} \sqrt{V(y)} f(y) \\ &= \frac{\sqrt{V(x)}}{2(2\pi)^d} \int_{\mathbb{R}^d} dy \sqrt{V(y)} f(y) \int_{E-\varepsilon}^{E+\varepsilon} dr r^{d/2-1} \int_{\mathbb{S}^{d-1}} dS(\xi) e^{i\sqrt{r}\xi \cdot (x-y)}, \end{aligned} \quad (73)$$

and therefore

$$\begin{aligned} & \lim_{\varepsilon \downarrow 0} \frac{1}{2\varepsilon} (\sqrt{V} 1_{(E-\varepsilon, E+\varepsilon)}(-\Delta) \sqrt{V} f)(x) \\ &= \frac{\sqrt{V(x)}}{2(2\pi)^d} \int_{\mathbb{R}^d} dy \sqrt{V(y)} f(y) E^{d/2-1} \int_{\mathbb{S}^{d-1}} dS(\xi) e^{i\sqrt{E}\xi \cdot (x-y)}, \end{aligned}$$

because the integrand in (73) is continuous in r . This implies (72). \square

Corollary A.2. *Let $d \geq 2$. Let $V_0 = 0$ and $V \neq 0$. Then for any $E > 0$ the operator $A(E)$ from (45) has infinite rank.*

Proof. We first show that the set of functions

$$\{\mathbb{R}^d \ni x \mapsto \sqrt{V(x)}e^{i\xi \cdot x} : \xi \in \mathbb{R}^d\} \tag{74}$$

is linearly independent. For this, notice that $\{\mathbb{C} \ni z \mapsto e^{isz} : s \in \mathbb{R}\}$ is linearly independent, since for $z = -ix$, these functions have different asymptotic behaviour for $x \rightarrow \infty$. Given a finite non-empty set $M \subseteq \mathbb{R}$ and $c_s \neq 0$ for $s \in M$, the analytic function $\mathbb{C} \ni z \mapsto \sum_{s \in M} c_s e^{isz}$ is therefore not identically zero, and thus $\mathbb{R} \ni x \mapsto \sum_{s \in M} c_s e^{isx}$ is zero only on a discrete subset of \mathbb{R} .

Given another finite non-empty set $M \subseteq \mathbb{R}^d$ and $c_\xi \neq 0$ for $\xi \in M$, define $F: \mathbb{R}^d \rightarrow \mathbb{C}$ via $F(x) := \sum_{\xi \in M} c_\xi e^{i\xi \cdot x}$. We show that $F^{-1}(\{0\}) \subseteq \mathbb{R}^d$ is a null set. Since F is continuous, this preimage is measurable with measure

$$\int_{\mathbb{R}^d} dx \, 1_{F^{-1}(\{0\})}(x) = \int_{\mathbb{S}^{d-1}} dS(\eta) \int_0^\infty dr \, r^{d-1} 1_{\{0\}}(F(r\eta)) = 0,$$

where the r -integral is zero since for $\eta \in \mathbb{S}^{d-1}$ fixed the function $r \mapsto F(r\eta) = \sum_{\xi \in M} c_\xi e^{ir\xi \cdot \eta}$ is zero only on a discrete subset of \mathbb{R} , as shown above. To show that the set (74) is linearly independent, it suffices to show that

$$\{x \in \mathbb{R}^d : \sqrt{V(x)}F(x) \neq 0\} = \{x \in \mathbb{R}^d : V(x) \neq 0\} \cap \{x \in \mathbb{R}^d : F(x) \neq 0\}$$

has positive measure. This is the case, since the first set in the intersection has positive measure and the second set is the complement of the null set $F^{-1}(\{0\})$.

Now, let $f \in \ker A(E)$. Then

$$0 = \langle f, A(E)f \rangle = \frac{E^{d/2-1}}{2(2\pi)^d} \int_{\mathbb{S}^{d-1}} dS(\xi) \left| \int_{\mathbb{R}^d} dx \sqrt{V(x)}e^{i\sqrt{E}\xi \cdot x} f(x) \right|^2,$$

and therefore

$$\int_{\mathbb{R}^d} dx \sqrt{V(x)}e^{i\sqrt{E}\xi \cdot x} f(x) = 0 \tag{75}$$

for a.e. $\xi \in \mathbb{S}^{d-1}$. Since the left-hand side of (75) is continuous in ξ , (75) holds in fact for all $\xi \in \mathbb{S}^{d-1}$. Since $f \in \ker A(E)$ was arbitrary, we conclude that

$$\{\mathbb{R}^d \ni x \mapsto \sqrt{V(x)}e^{i\sqrt{E}\xi \cdot x} : \xi \in \mathbb{S}^{d-1}\} \subseteq (\ker A(E))^\perp.$$

Since \mathbb{S}^{d-1} is an infinite set for $d \geq 2$, the set of functions on the left-hand side is infinite and linearly independent, and thus $\dim(\ker A(E))^\perp = \infty$. Since the coimage $(\ker A(E))^\perp$ of the linear map $A(E)$ is isomorphic to $\text{ran } A(E)$ (the restriction $A(E)|_{(\ker A(E))^\perp} : (\ker A(E))^\perp \rightarrow \text{ran } A(E)$ being bijective), this shows $\dim \text{ran } A(E) = \infty$. □

Remark A.3. We expect Corollary A.2 to generalise to the situation of non-zero background potentials V_0 with suitable decay by using generalised eigenfunctions due to Ikebe and Povzner (see [29, §C5] and references therein) in place of $e^{i\sqrt{E}\xi \cdot x}$.

The infinite rank of $A(E)$ implies positivity of $\gamma(E)$.

Theorem A.4. *Let $d \geq 2$ and $V_0 = 0$. Then the transition matrix T_E corresponding to the pair $H = -\Delta$ and $H' = -\Delta + V$ has infinite rank for a.e. $E > 0$. In particular, T_E is non-zero and therefore*

$$\gamma(E) = \pi^{-2} \|\arcsin|T_E/2|\|_{\text{HS}}^2 > 0$$

for a.e. $E > 0$.

Proof. By Lemma 4.30, it suffices to show that $\tilde{T}(E) = -2\pi i \sqrt{A(E)}\Phi_+(E) \sqrt{A(E)}$ has infinite rank, where $\Phi_{\pm}(E) = \lim_{\varepsilon \downarrow 0} (I + \sqrt{V}(E \pm i\varepsilon - H')^{-1} \sqrt{V})$. We show that its imaginary part $\text{Im } \tilde{T}(E) = \frac{1}{2i}(\tilde{T}(E) - \tilde{T}(E)^*)$ has infinite rank. For brevity, set $R := \lim_{\varepsilon \downarrow 0} \sqrt{V}(E + i\varepsilon - H')^{-1} \sqrt{V}$. Recall that by the limiting absorption principle, this limit exists in operator norm for a.e. $E > 0$; in particular, R is compact. We fix such an $E > 0$ from now on. Then

$$\begin{aligned} \text{Im } \tilde{T}(E) &= \frac{1}{2i}(-2\pi i \sqrt{A(E)}\Phi_+(E) \sqrt{A(E)} - 2\pi i \sqrt{A(E)}\Phi_-(E) \sqrt{A(E)}) \\ &= -2\pi \sqrt{A(E)}(I + \text{Re } R) \sqrt{A(E)}. \end{aligned}$$

Since $\text{Re } R$ is compact, we can write it as $\text{Re } R = R_1 + R_2$ where $\|R_1\| < 1/2$ and R_2 has finite rank. Thus

$$-\frac{1}{2\pi} \text{Im } \tilde{T}(E) = \sqrt{A(E)}(I + R_1) \sqrt{A(E)} + \tilde{A}$$

where \tilde{A} is a finite rank operator. Now, since $I + R_1 \geq I - \frac{1}{2}I = \frac{1}{2}I$, we get

$$\sqrt{A(E)}(I + R_1) \sqrt{A(E)} \geq \frac{1}{2}A(E).$$

By Corollary A.2, this operator has infinite rank. □

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References

- [1] P. W. Anderson, Infrared catastrophe in Fermi gases with local scattering potentials. *Phys. Rev. Lett.* **18** (1967), 1049–1051.
- [2] P. W. Anderson, Ground state of a magnetic impurity in a metal, *Phys. Rev.* **164** (1967), 352–359.
- [3] H. Baumgärtel and M. Wollenberg, *Mathematical scattering theory*. Mathematische Lehrbücher und Monographien, II. Abteilung: Mathematische Monographien, 59. Akademie-Verlag, Berlin, 1983. Licensed ed., Operator Theory: Advances and Applications, 9. Birkhäuser-Verlag, Basel, 1983. [MR 0781534](#) [MR 0699113](#) (licensed ed.) [Zbl 0536.47007](#) [Zbl 0536.47008](#) (licensed ed.)
- [4] M. Š. Birman and S. B. Èntina, The stationary method in the abstract theory of scattering. *Izv. Akad. Nauk SSSR Ser. Mat.* **31** (1967), 401–430. In Russian. English translation, *Math. USSR Izv.* **1** (1967), 391–420. [MR 0209895](#) [Zbl 0173.16501](#)
- [5] K. Broderix, D. Hundertmark, and H. Leschke, Continuity properties of Schrödinger semigroups with magnetic fields. *Rev. Math. Phys.* **12** (2000), 181–225. [MR 1756112](#) [Zbl 0961.81006](#)
- [6] G. Doetsch, *Introduction to the theory and application of the Laplace transformation*. Translated from the second German edition by W. Nader. Springer-Verlag, Berlin etc., 1974. [MR 0344810](#) [Zbl 0278.44001](#)
- [7] R. L. Frank and A. Pushnitski, The spectral density of a product of spectral projections. *J. Funct. Anal.* **268** (2015), 3867–3894. [MR 3341967](#) [Zbl 1312.81071](#)
- [8] K. O. Friedrichs, *Mathematical aspects of the quantum theory of fields*. Reprinted from *Comm. Pure Appl. Math.* Interscience Publishers, New York and London, 1953. [MR 0056468](#) [Zbl 0053.32602](#)
- [9] M. Gebert, Finite-size energy of non-interacting Fermi gases. *Math. Phys. Anal. Geom.* **18** (2015), article 27. [MR 3412472](#) [Zbl 1332.81056](#)
- [10] M. Gebert, The asymptotics of an eigenfunction-correlation determinant for Dirac- δ perturbations. *J. Math. Phys.* **56** (2015), article id. 072110. [MR 3405952](#) [Zbl 1327.82079](#)
- [11] M. Gebert, H. Küttler and P. Müller, Anderson’s orthogonality catastrophe. *Comm. Math. Phys.* **329** (2014), 979–998. [MR 3212877](#) [Zbl 1294.82013](#)
- [12] I. S. Gradshteyn and I. M. Ryzhik, *Table of integrals, series, and products*. Translated from the Russian. Translation edited and with a preface by D. Zwillinger and V. Moll. Eighth edition. Revised from the seventh edition. Elsevier/Academic Press, Amsterdam, 2015. [MR 3307944](#) [Zbl 1300.65001](#)
- [13] R. Haag, *Local quantum physics*. Fields, particles, algebras. Second revised and enlarged edition. Springer-Verlag, Berlin, 1996. [MR 1405610](#) [Zbl 0857.46057](#)
- [14] D. R. Hamann, Orthogonality catastrophe in metals. *Phys. Rev. Lett.* **26** (1971), 1030–1032.

- [15] P. D. Hislop and P. Müller, The spectral shift function for compactly supported perturbations of Schrödinger operators on large bounded domains. *Proc. Amer. Math. Soc.* **138** (2010), 2141–2150. [MR 2596053](#) [Zbl 1190.81127](#)
- [16] M. Kac, Toeplitz matrices, translation kernels and a related problem in probability theory. *Duke Math. J.* **21** (1954), 501–509. [MR 0062867](#) [Zbl 0056.10201](#)
- [17] W. Kirsch, Small perturbations and the eigenvalues of the Laplacian on large bounded domains. *Proc. Amer. Math. Soc.* **101** (1987), 509–512. [MR 0908658](#) [Zbl 0643.35074](#)
- [18] H. K. Knörr, P. Otte, and W. Spitzer, Anderson’s orthogonality catastrophe in one dimension induced by a magnetic field. *J. Phys. A* **48** (2015), article id. 325202. [MR 3376020](#) [Zbl 1332.82084](#)
- [19] H. Küttler, *Anderson’s orthogonality catastrophe*. Ph.D. thesis. Ludwig-Maximilians-Universität München, München, 2014.
- [20] H. Küttler, P. Otte, and W. Spitzer, Anderson’s orthogonality catastrophe for one-dimensional systems. *Ann. Henri Poincaré* **15** (2014), 1655–1696. [MR 3245883](#) [Zbl 1304.81164](#)
- [21] C. Muscalu and W. Schlag, *Classical and multilinear harmonic analysis*. Vol. I. Cambridge Studies in Advanced Mathematics, 137. Cambridge University Press, Cambridge, 2013. [MR 3052498](#) [Zbl 1281.42002](#)
- [22] NIST Digital Library of Mathematical Functions. Release 1.0.6 of 2013-05-06. Online companion to [24]. <http://dlmf.nist.gov/>
- [23] P. Nozières and C. T. de Dominicis, Singularities in the x-ray absorption and emission of metals. iii. one-body theory exact solution. *Phys. Rev.* **178** (1969), 1097–1107.
- [24] F. W. J. Olver, D. W. Lozier, R. F. Boisvert, and C. W. Clark (eds.), *NIST Handbook of mathematical functions*. U.S. Department of Commerce, National Institute of Standards and Technology, Washington, D.C., and Cambridge University Press, Cambridge, 2010. Print companion to [22]. [MR 2723248](#) [Zbl 1198.00002](#)
- [25] M. Reed and B. Simon, *Methods of modern mathematical physics*. III. Scattering theory. Academic Press, New York and London, 1979. [MR 0529429](#) [Zbl 0405.47007](#)
- [26] N. Rivier and E. Simanek, Exact calculation of the orthogonality catastrophe in metals. *Phys. Rev. Lett.* **26** (1971), 435–438.
- [27] M. Rosenblum, On the Hilbert matrix. I. *Proc. Amer. Math. Soc.* **9** (1958), 137–140. [MR 0094626](#) [Zbl 0080.10504](#)
- [28] J. Rovnyak, The Hilbert matrix as a singular integral operator. *Acta Sci. Math. (Szeged)* **31** (1970), 347–350. [MR 0291899](#) [Zbl 0203.14301](#)
- [29] B. Simon, Schrödinger semigroups. *Bull. Amer. Math. Soc. (N.S.)* **7** (1982), 447–526. Erratum, *ibid* **11** (1984), 426. [MR 0670130](#) [MR 0752806](#) (erratum) [Zbl 0524.35002](#)
- [30] B. Simon, *Trace ideals and their applications*. Second edition. Mathematical Surveys and Monographs, 120. American Mathematical Society, Providence, R.I., 2005. [MR 2154153](#) [Zbl 0423.47001](#)

- [31] J. Weidmann, *Linear operators in Hilbert spaces*. Translated from the German by J. Szücs. Graduate Texts in Mathematics, 68. Springer-Verlag, Berlin etc., 1980. [MR 0566954](#) [Zbl 0434.47001](#)
- [32] D. R. Yafaev, *Mathematical scattering theory*. General theory. Translated from the Russian by J. R. Schulenberger. Translations of Mathematical Monographs, 105. American Mathematical Society, Providence, R.I., 1992. [MR 1180965](#) [Zbl 0761.47001](#)
- [33] D. R. Yafaev, A commutator method for the diagonalization of Hankel operators. *Funktional. Anal. i Prilozhen.* **44** (2010), 65–79. In Russian. English translation, *Funct. Anal. Appl.* **44** (2010), 295–306. [MR 2768565](#) [Zbl 1271.47019](#)
- [34] G. Yuval and P. W. Anderson, Exact results for the Kondo problem: one-body theory and extension to finite temperature. *Phys. Rev. B* **1** (1970), 1522–1528.

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