Mixed data in inverse spectral problems for the Schrödinger operators

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Abstract. We consider the Schrödinger operator on a finite interval with an L^1 -potential. We prove that the potential can be uniquely recovered from one spectrum and subsets of another spectrum and point masses of the spectral measure (or norming constants) corresponding to the first spectrum. We also solve this Borg–Marchenko-type problem under some conditions on two spectra, when missing part of the second spectrum and known point masses of the spectral measure have different index sets.

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1. Introduction

In this paper, we consider the Schrödinger (Sturm-Liouville) equation

$$Lu = -u'' + qu = zu$$

on the interval $(0, \pi)$ with the boundary conditions

 $u(0) \cos \alpha - u'(0) \sin \alpha = 0,$ $u(\pi) \cos \beta + u'(\pi) \sin \beta = 0,$

and a real-valued potential $q \in L^1(0, \pi)$. The spectrum $\sigma_{\alpha,\beta}$ of the Schrödinger operator *L* corresponding to these boundary conditions defines a discrete subset of the real line, bounded from below, diverging to $+\infty$.

Direct spectral problems aim to get spectral information from the potential. In inverse spectral problems, the goal is to recover the potential from spectral information, such as the spectrum, the norming constants, the spectral measure or Weyl–Titchmarsh *m*-function. These notions are discussed in Section 2.

The first inverse spectral result on Schrödinger operators is given by Ambarzumian [1]. He considered continuous potential with Neumann boundary conditions at both endpoints ($\alpha = \beta = \pi/2$) and showed that $q \equiv 0$ if the spectrum consists of squares of integers.

Later Borg [12] proved that an L^1 -potential is uniquely recovered from two spectra corresponding to various pairs of boundary conditions and sharing the same boundary conditions at π ($\beta_1 = \beta_2$), one of which should be Dirichlet boundary condition at 0 ($\alpha_1 = 0$). Levinson [34] extended Borg's result by removing the restriction of Dirichlet boundary condition at 0.

Furthermore, Marchenko [38] observed that the spectral measure (or Weyl–Titchmarsh *m*-function) uniquely recovers an L^1 -potential.

Another classical result is due to Hochstadt and Lieberman [27], which says that if the first half of an L^1 -potential is known, one spectrum recovers the whole.

Statements of these classical results are given in Section 3.1.

Gesztesy, Simon and del Rio [13] generalized Levinson's theorem to three spectra, by showing two thirds of the union of three spectra is sufficient spectral data to recover an L^1 -potential.

Later on, Gesztesy and Simon [20] observed that extra smoothness conditions on the potential change required spectral data to recover the potential. They proved that the knowledge of the eigenvalues can be replaced by information on the derivatives of the potential. In addition, they [20] also generalized the Hochstadt– Lieberman theorem in the sense that more than the first half of an L^1 -potential and a sufficiently large subset of a spectrum recover the potential.

Afterwards, Amour, Raoux, and Faupin [3, 4] proved similar results using extra information on the smoothness of the potential.

In a remarkable result, Horváth [28] characterized unique recovery of a potential in terms of completeness of an exponential system depending on given eigenvalues and known part of the potential. This observation opened a new path [5, 28, 31, 37] by connecting inverse spectral problems and completeness of exponential systems. Moreover, Horváth and Sáfár [31] proved similar results in terms of a cosine system. The cosine system depends on subsets of eigenvalues and norming constants and their spectral data consists of these two subsets.

Recently, Makarov and Poltoratski [37] gave a version of Horváth's theorem [28] in terms of exterior Beurling–Malliavin density by combining Horváth's result and the Beurling–Malliavin theorem. In the same paper, they obtained another characterization result, which is an uncertainty version of Borg's theorem. As their spectral data, they considered a set of intervals known to include two spectra and characterized the inverse spectral problem in terms of a convergence criterion on this set of intervals.

All of these results mentioned above are discussed in Section 3.2.

Classical theorems of Borg, Levinson, Marchenko, Hochstadt, and Lieberman led to various other inverse spectral results on Schrödinger operators (see [2, 19, 23, 24, 25, 26, 29, 39, 41, 42, 44, 45, 46, 47, 48, 49, 50, 51, 52] and references therein). These problems can be divided into two groups. In Borg–Marchenko-type spectral problems, one tries to recover the potential from spectral data. However, Hochstadt–Lieberman-type (or mixed) spectral problems recover the potential using a mixture of partial information on the potential and spectral data.

In the present paper, our interest is on regular Schrödinger operators with summable potentials on a finite interval. However, many problems with locally summable potentials [14, 15, 16, 17, 21, 30, 32] or on various settings such as half-line [17, 19, 20, 22, 42, 44], real-line [17, 19, 21, 22, 44] or graphs [8, 9, 10, 11, 53] are solved.

Borg's, Levinson's and Hochstadt and Lieberman's theorems suggest that one spectrum gives exactly one half of the full spectral information required to recover the potential. Recalling the fact that the spectral measure is a discrete measure supported on a spectrum, the same can be said for the set of point masses of the spectral measure. As follows from Marchenko's theorem, the set of point masses of the spectral measure (or the set of norming constants) gives exactly one half of the full spectral information required to recover the potential.

These observations allow us to formulate the following question:

Inverse problem. Do one spectrum and partial information on another spectrum and the set of point masses of the spectral measure corresponding to the first spectrum recover the potential?

This Borg–Marchenko-type problem can be seen as a combination of Levinson's and Marchenko's results. In the present paper, we answer this question positively. First, we give a proof with the most common boundary conditions, Dirichlet (u = 0) and Neumann (u' = 0). Theorem 4.2 solves this inverse spectral problem when given part of the point masses of the spectral measure corresponding to the Dirichlet–Dirichlet spectrum matches with the missing part of the Neumann–Dirichlet spectrum, i.e. they share same index sets. In Theorem 4.6 and Theorem 4.8, we consider the non-matching index sets case with some restrictions on two spectra.

In order to deal with general boundary conditions we introduce a more general m-function in Section 4.3. With this m-function, we extend Theorem 4.2 in Theorem 4.11 to general boundary conditions. In Theorem 4.13 and Theorem 4.14 we consider the non-matching index sets case.

The paper is organized as follows.

- In Section 2.1 we discuss spectra of Schrödinger operators and their asymptotics for various boundary conditions.
- In Section 2.2 we define Weyl–Titchmarsh *m*-function and spectral measure for Schrödinger operators.
- In Section 3.1 we recall statements of the classical results of Ambarzumian, Borg, Levinson, Marchenko, Hochstadt, and Lieberman.
- In Section 3.2 we discuss some recent results in the finite interval setting with summable potential.
- In Section 4.1 we give a representation of Weyl–Titchmarsh *m*-function as an infinite product and prove the inverse spectral problem mentioned above with Dirichlet–Dirichlet, Neumann–Dirichlet boundary conditions.
- In Section 4.2 we consider the same problem in the non-matching index sets case.
- In Section 4.3 we introduce a more general *m*-function and solve the inverse spectral problem corresponding to this *m*-function with general boundary conditions in both the matching and non-matching index sets cases.
- In Appendix A we list all definitions and theorems from complex function theory used in this paper.

2. Preliminaries

2.1. One-dimensional Schrödinger operator on a finite interval. As it was defined in the introduction, we consider the Schrödinger equation

$$Lu = -u'' + qu = zu \tag{2.1}$$

on the interval $(0, \pi)$ associated with the boundary conditions

$$u(0)\cos\alpha - u'(0)\sin\alpha = 0,$$
 (2.2)

$$u(\pi)\cos\beta + u'(\pi)\sin\beta = 0, \qquad (2.3)$$

where $\alpha, \beta \in [0, \pi)$ and the potential $q \in L^1(0, \pi)$ is real-valued.

The spectrum $\sigma_{\alpha,\beta}$ of the Schrödinger operator

$$L: u \longmapsto -u'' + qu,$$

with $a \in L^1$ and boundary conditions (2.2), (2.3) is a discrete real sequence, bounded from below. Adding a positive constant to the potential q, shifts the spectrum by the same constant. This allows us to assume wlog $\sigma_{\alpha,\beta} \subset \mathbb{R}_+$. Throughout the paper we assume $\mathbb{N} = \{1, 2, 3, ...\}$. Asymptotic behavior of the spectrum $\sigma_{\alpha,\beta} = \{a_n\}_{n \in \mathbb{N}}$, depending on the signs of α and β , is given below.

If $\alpha \neq 0, \beta \neq 0$, then

$$a_n = (n-1)^2 + \frac{2}{\pi} [\cot(\beta) + \cot(\alpha)] + \frac{1}{\pi} \int_0^{\pi} q(x) dx + \alpha_n$$
(2.4)

where $\alpha_n = o(1)$ as $n \to +\infty$.

If $\alpha = 0, \beta = 0$, then

$$a_n = n^2 + \frac{1}{\pi} \int_0^{\pi} q(x) dx + \alpha_n$$
 (2.5)

where $\alpha_n = o(1)$ as $n \to +\infty$.

If $\alpha \neq 0, \beta = 0$, then

$$a_n = \left(n - \frac{1}{2}\right)^2 + \frac{2}{\pi}\cot(\alpha) + \frac{1}{\pi}\int_0^{n} q(x)dx + \alpha_n$$
(2.6)

where $\alpha_n = o(1)$ as $n \to +\infty$.

If $\alpha = 0, \beta \neq 0$, then

$$a_n = \left(n - \frac{1}{2}\right)^2 + \frac{2}{\pi}\cot(\beta) + \frac{1}{\pi}\int_0^{\pi} q(x)dx + \alpha_n$$
(2.7)

where $\alpha_n = o(1)$ as $n \to +\infty$.

In the case $q \in L^2(0, \pi)$, the same asymptotics are valid with $\{\alpha_n\}_{n \in \mathbb{N}} \in l^2$.

One can find these results in the classical texts on Schrödinger operators, for instance [35] or [36].

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2.2. Weyl–Titchmarsh *m*-function and the spectral measure. Let us choose the boundary condition (2.2) and introduce two solutions $s_z(t)$ and $c_z(t)$ of (2.1) satisfying the initial conditions

$$s_z(0) = \sin(\alpha), \quad s'_z(0) = \cos(\alpha),$$

$$c_z(0) = \cos(\alpha), \quad c'_z(0) = -\sin(\alpha).$$

Definition 2.1. The norming constant τ_{α} , for the eigenvalue a_n is defined as

$$\tau_{\alpha}(a_n) := \int_0^{\pi} |s_{a_n}(t)|^2 dt.$$

Note that $s_z(t)$ and $c_z(t)$ are linearly independent solutions and their Wronskian satisfies $W(c_z, s_z) = 1$. This allows us to represent $u_z(t)$, a solution of (2.1) with boundary conditions $u_z(\pi) = \sin \beta$, $u'_z(\pi) = -\cos \beta$, as

$$u_z(t) = c_z(t) + m_{\alpha,\beta}(z)s_z(t),$$

where

$$m_{\alpha,\beta}(z) = -\frac{W(c_z, u_z)}{W(s_z, u_z)}.$$

This is how we derive the *m*-function.

Definition 2.2. Weyl–Titchmarsh *m*-function with the boundary conditions (2.2) and (2.3) is defined as

$$m_{\alpha,\beta}(z) := \frac{\cos(\alpha)u_z'(0) + \sin(\alpha)u_z(0)}{-\sin(\alpha)u_z'(0) + \cos(\alpha)u_z(0)}$$

where $\alpha, \beta \in [0, \pi)$.

It is well known that Weyl *m*-function $m_{\alpha,\beta}$ is a meromorphic Herglotz function. The definition of a Herglotz function and other definitions and results from complex function theory used in this paper can be found in Appendix A. Everitt [18] proved that the Weyl *m*-function has the asymptotic

$$m_{0,\beta}(z) = i\sqrt{z} + o(1)$$

for $\alpha = 0$, and

$$m_{\alpha,\beta}(z) = \frac{\cos \alpha}{\sin \alpha} + \frac{1}{\sin^2 \alpha} \frac{i}{\sqrt{z}} + O\left(\frac{1}{|z|}\right)$$

for $\alpha \in (0, \pi)$ as *z* goes to infinity in the upper half plane. Asymptotics of Weyl *m*-function and Herglotz representation theorem imply that $m_{\alpha,\beta}$ is represented as the Herglotz integral of a discrete positive Poisson-finite measure supported on the spectrum $\sigma_{\alpha,\beta}$:

$$m_{\alpha,\beta}(z) = a + \int_{\mathbb{R}} \left[\frac{1}{t-z} - \frac{t}{1+t^2} \right] d\mu_{\alpha,\beta}(t), \qquad (2.8)$$

where $a = \operatorname{Re}(m_{\alpha,\beta}(i))$, $\sigma_{\alpha,\beta} = \{a_n\}_{n \in \mathbb{N}}$ and $\mu_{\alpha,\beta} = \sum_{n \in \mathbb{N}} \gamma_n \delta_{a_n}$. The measure $\mu_{\alpha,\beta}$ is the spectral measure of the Schrödinger operator *L* corresponding to the *m*-function $m_{\alpha,\beta}$. The point masses of the spectral measure is represented in terms of norming constants as $\gamma_n = (\tau_\alpha(a_n))^{-1}$.

Definition 2.3. The spectral measure of the Schrödinger operator *L* corresponding to the *m*-function $m_{\alpha,\beta}$ (or the boundary conditions (2.2), (2.3)) is defined as

$$\mu_{\alpha,\beta} := \sum_{n \in \mathbb{N}} \frac{\delta_{a_n}}{\tau_{\alpha}(a_n)},$$

where $\alpha, \beta \in [0, \pi)$ and $\sigma_{\alpha, \beta} = \{a_n\}_{n \in \mathbb{N}}$.

Since $\mu_{\alpha,\beta}$ is a Poisson-finite measure, the spectrum and the point masses of the spectral measure satisfy

$$\sum_{n\in\mathbb{N}}\frac{\gamma_n}{1+a_n^2}<\infty.$$

These properties of the *m*-function, the spectral measure and a detailed discussion of one dimensional Schrödinger operators appear in Chapter 9 of [43].

In order to illustrate what we have discussed so far, let us consider the free operator $(q \equiv 0)$ with Dirichlet (u = 0) and Neumann (u' = 0) boundary conditions.

Example 2.4. The spectra, the *m*-function and the spectral measure for $q \equiv 0$ on $(0, \pi)$ with Dirichlet–Dirichlet, Neumann–Dirichlet, and Neumann–Neumann boundary conditions are as follows:

$$\sigma_{\text{DD}} := \sigma_{0,0} = \{n^2\}_{n \in \mathbb{N}},$$

$$m_{0,0} = -\sqrt{z} \cot(\sqrt{z}\pi),$$

$$\mu_{0,0} = \frac{2}{\pi} \sum_{n=1}^{\infty} n^2 \delta_{n^2};$$

$$\sigma_{\rm ND} := \sigma_{\pi/2,0} = \left\{ \left(n - \frac{1}{2} \right)^2 \right\}_{n \in \mathbb{N}},$$

$$m_{\pi/2,0} = \frac{\tan(\sqrt{z\pi})}{\sqrt{z}},$$

$$\mu_{\pi/2,0} = \frac{2}{\pi} \sum_{n=1}^{\infty} \delta_{(n-1/2)^2};$$

$$\sigma_{\rm NN} := \sigma_{\pi/2,\pi/2} = \{ (n-1)^2 \}_{n \in \mathbb{N}},$$

$$m_{\pi/2,\pi/2} = \frac{\cot(\sqrt{z\pi})}{\sqrt{z}},$$

$$\mu_{\pi/2,\pi/2} = \frac{2}{\pi} \sum_{n=1}^{\infty} \delta_{(n-1)^2}.$$

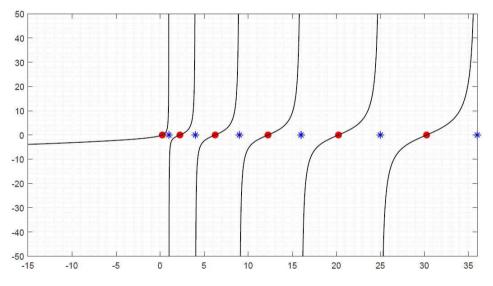


Figure 1. The graph of Weyl *m*-function $m_{0,0}$ on \mathbb{R} , Neumann–Dirichlet spectrum σ_{ND} (•) and Dirichlet–Dirichlet spectrum σ_{DD} (*) for the free operator ($q \equiv 0$).

3. Inverse spectral theory of regular Schrödinger operators

3.1. Classical results. The first inverse spectral result on Schrödinger operators was given by Ambarzumian.

Theorem 3.1 (Ambarzumian [1]; see Horváth [28]). Let $q \in C[0, \pi]$ and let $\sigma_{\pi/2,\pi/2} = \{n^2\}_{n=0}^{\infty}$. Then $q \equiv 0$.

Later Borg found that in most cases two spectra is the required spectral information to recover the operator uniquely.

Theorem 3.2 (Borg [12]; see Horváth [28]). Let $q \in L^1(0, \pi)$, $\sigma_1 = \sigma_{0,\beta}$, $\sigma_2 = \sigma_{\alpha_2,\beta}$, $\sin \alpha_2 \neq 0$ and

$$\tilde{\sigma}_2 = \begin{cases} \sigma_2 & \text{if } \sin \beta = 0, \\ \sigma_2 \setminus \{a_1\} & \text{if } \sin \beta \neq 0. \end{cases}$$

Then $\sigma_1 \cup \tilde{\sigma}_2$ determines the potential and no proper subset has the same property.

A Schrödinger operator (or a potential) is said *determined* (or *recovered*) by its spectral data, if any other operator with the same data must have the same potential a.e. on $(0, \pi)$. Levinson extended Borg's result by removing the Dirichlet boundary condition restriction from the first spectrum.

Theorem 3.3 (Levinson [34]; see Horváth [28]). Let $q \in L^1(0, \pi)$ and suppose $\sin(\alpha_1 - \alpha_2) \neq 0$. Then $\sigma_{\alpha_1,\beta}$ and $\sigma_{\alpha_2,\beta}$ determine the potential.

Marchenko showed that the spectral measure or the corresponding Weyl *m*-function provides sufficient spectral data to recover the potential uniquely.

Theorem 3.4 (Marchenko [38]; see Teschl [43], Section 9.4). Let $q \in L^1(0, \pi)$. Then $\mu_{\alpha,\beta}$ or $m_{\alpha,\beta}$ determines the potential.

In the notations of Section 2.2, Marchenko's theorem says that the spectrum $\sigma_{\alpha,\beta} = \{a_n\}_{n \in \mathbb{N}}$ and the point masses $\{\gamma_n\}_{n \in \mathbb{N}}$ of the corresponding spectral measure (or the norming constants $\{\tau_{\alpha}(a_n)\}_{n \in \mathbb{N}}$) provide sufficient spectral data to recover the operator uniquely.

Hochstadt and Lieberman observed that one spectrum recovers the potential if the first half of it is known.

Theorem 3.5 (Hochstadt and Lieberman [27]). Let $q \in L^1(0, \pi)$. Then q on $(0, \pi/2)$ and $\sigma_{\alpha,\beta}$ determine the potential.

3.2. Some recent results in the finite interval case. For any discrete real sequence $A = \{x_n\}_{n \in \mathbb{N}}, x_n \to \infty$ the counting function is defined as

$$n_A(t) := \sum_{x_n \le t} 1.$$

Gesztesy, Simon and del Rio generalized Levinson's theorem to three spectra.

Theorem 3.6 (del Rio, Gesztesy, and Simon [13]). Let $q \in L^1(0, \pi)$. Then $S \subset \sigma_{\alpha_1,\beta} \cup \sigma_{\alpha_2,\beta} \cup \sigma_{\alpha_3,\beta}$ satisfying

$$n_{S}(t) \ge (2/3)n_{(\sigma_{\alpha_{1},\beta}\cup\sigma_{\alpha_{2},\beta}\cup\sigma_{\alpha_{3},\beta})}(t)$$

for sufficiently large t > 0, determine the potential.

Gesztesy and Simon observed that the knowledge of the eigenvalues can be replaced by information on the derivatives of the potential around the midpoint of the interval.

Theorem 3.7 (Gesztesy and Simon [21]). Let $q \in L^1(0, \pi)$, $\alpha, \beta \neq 0$ and $q \in C^{2k}(\pi/2 - \epsilon, \pi/2 + \epsilon)$ for some $k \in \mathbb{N}$ and $\epsilon > 0$. Then q on $(0, \pi/2)$ and $\sigma_{\alpha,\beta}$ except for k + 1 eigenvalues determine the potential.

In the same paper, they generalized Hochstadt-Lieberman theorem.

Theorem 3.8 (Gesztesy and Simon [21]). Let $q \in L^1(0, \pi)$ and $\pi/2 < a < \pi$. Then q on (0, a) and $S \subset \sigma_{\alpha,\beta}$ satisfying

$$n_S(t) \ge 2(1 - a/\pi)n_{\sigma_{\alpha,\beta}}(t) + a/\pi - 1/2$$

for sufficiently large t > 0, determine the potential.

Amour, Raoux, and Faupin proved similar results using extra information on smoothness of the potential.

Theorem 3.9 (Amour and Raoux [4]). Let α , β_1 , $\beta_2 \neq 0$, $p \in [1, \infty)$, $q_1, q_2 \in L^1(0, \pi)$, $q_1 - q_2 \in L^p(a, \pi)$ and $\pi/2 < a < \pi$. If $q_1 = q_2$ a.e. on (0, a) and $S \subset \sigma_{\alpha,\beta_1}(q_1) \cap \sigma_{\alpha,\beta_2}(q_2)$ satisfies

$$2(1 - a/\pi)n_{\sigma}(t) + C \ge n_{S}(t) \ge 2(1 - a/\pi)n_{\sigma}(t) + 1/(2p) + 2a/\pi - 2$$

for a real number C and sufficiently large t > 0, where σ denotes either of $\sigma_{\alpha,\beta_k}(q_k)$, then $q_1 = q_2$ a.e. on $(0, \pi)$.

Theorem 3.10 (Amour, Faupin, and Raoux [3]). Let α , β_1 , $\beta_2 \neq 0$, $k \in \{0, 1, 2\}$, $p \in [1, \infty)$, $q_1, q_2 \in W^{k,1}(0, \pi)$, $q_1 - q_2 \in W^{k,p}(a, \pi)$ and $\pi/2 < a < \pi$. If $q_1 = q_2$ on (0, a) and $S \subset \sigma_{\alpha,\beta_1}(q_1) \cap \sigma_{\alpha,\beta_2}(q_2)$ satisfying

$$n_{S}(t) \ge 2(1 - a/\pi)n_{\sigma}(t) - k/2 + 1/(2p) + a/\pi - 3/2$$

for sufficiently large t > 0, where σ denotes either of $\sigma_{\alpha,\beta_k}(q_k)$, then $q_1 = q_2$ a.e. on $(0, \pi)$.

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Theorem 3.11 (Amour, Faupin, and Raoux [3]). Let α , β_1 , $\beta_2 \neq 0$, $k \in \{0, 1, 2\}$, $p \in [1, \infty)$, $q_1, q_2 \in W^{k,1}(0, \pi)$, $q_1 - q_2 \in W^{k,p}(a, \pi)$ and $\pi/2 < a < \pi$. If $q_1 = q_2$ on (0, a) and $S \subset \sigma_{\alpha,\beta_1}(q_1) \cap \sigma_{\alpha,\beta_2}(q_2)$ satisfying

$$2(1 - a/\pi)n_{\sigma}(t) + C \ge n_{S}(t) \ge 2(1 - a/\pi)n_{\sigma}(t) - k/2 + 1/(2p) + 2a/\pi - 2$$

for sufficiently large t > 0, where σ denotes either of $\sigma_{\alpha,\beta_k}(q_k)$, then $q_1 = q_2$ a.e. on $(0, \pi)$.

Horváth proved a remarkable characterization theorem, which represents a connection between inverse spectral theory and completeness of exponential systems.

Theorem 3.12 (Horváth [28]). Let $1 \le p \le \infty$, $q \in L^p(0, \pi)$, $0 \le a < \pi$ and $\lambda_n \in \sigma_{\alpha,0}$. Then q on (0, a) and the eigenvalues λ_n determine q if and only if the system

$$e(\Lambda) = \{ e^{\pm 2i\mu x}, e^{\pm 2i\sqrt{\lambda_n}x} : n \ge 1 \}$$

is complete in $L^p(a - \pi, \pi - a)$ for some $\mu \neq \pm \sqrt{\lambda_n}$.

Horváth and Sáfár proved similar results for the norming constants in terms of a cosine system. For a sequence $\Lambda = \{\lambda_1, \lambda_2, ...\} \subset \mathbb{R}$ and a subset $S \subset \Lambda$ they considered the following cosine system:

$$C(\Lambda, S) = \{\cos(2\sqrt{\lambda_n}t) : n \in \mathbb{N}\} \cup \{t \cos(2\sqrt{\lambda_n}t) : \lambda_n \in S\}.$$

Theorem 3.13 (Horváth and Sáfár [31]). Let $\beta \neq 0, 1 \leq p \leq \infty, q \in L^1(0, \pi)$, $q \in L^p(a, \pi), 0 \leq a < \pi$ and

$$\Lambda = \{\lambda_n : \lambda_n \in \sigma_{\alpha_n,\beta}, n \in \mathbb{N}\}$$

be a subset of eigenvalues such that $\lambda_n \not\rightarrow -\infty$ are different real numbers and $S \subset \Lambda$. Then q on (0, a), Λ and $\{\tau_{\alpha_n}(\lambda_n)\}_{\lambda_n \in S}$ determine q if the system $C(\Lambda, S)$ is complete in $L^p(0, \pi - a)$.

For Dirichlet boundary condition Horváth and Sáfár obtained an optimal condition.

Theorem 3.14 (Horváth and Sáfár [31]). Let us have the assumptions of Theorem 3.13, but $\beta = 0$. Let $\mu \neq \pm \sqrt{\lambda_n}$, $\mu \in \mathbb{R}$. Then the system $C(\Lambda, S) \cup \{\cos(2\sqrt{\mu}t)\}$ is complete in $L^p(0, \pi - a)$ if and only if q on (0, a), Λ and $\{\tau_{\alpha_n}(\lambda_n)\}_{\lambda_n \in S}$ determine q.

Makarov and Poltoratski gave a characterization theorem in terms of exterior Beurling–Malliavin density as a corollary of Horváth's result [28] (Theorem 3.12 above) and the Beurling–Malliavin theorem [6, 7].

If $\{I_n\}_{n \in \mathbb{N}}$ is a sequence of disjoint intervals on the real line, it is called *short* if

$$\sum_{n\in\mathbb{N}}\frac{|I_n|^2}{1+\operatorname{dist}^2(0,I_n)}<\infty$$

and long otherwise.

If Λ is a sequence of real points, its exterior (effective) Beurling–Malliavin density is defined as

$$D^*(\Lambda) = \sup\{d: \text{ there exists a long } \{I_n\} \text{ such that}$$

 $\#(\Lambda \cap I_n) \ge d|I_n|, \text{ for all } n \in \mathbb{N}\}.$

For a non-real sequence its density is defined as $D^*(\Lambda) = D^*(\Lambda')$, where Λ' is a real sequence $\lambda'_n = \left(\operatorname{Re} \frac{1}{\lambda_n}\right)^{-1}$, if Λ has no imaginary points, and as $D^*(\Lambda) = D^*((\Lambda + c)')$ otherwise.

For any complex sequence Λ its radius of completeness is defined as

 $R(\Lambda) = \sup\{a: \{e^{i\lambda z}\}_{\lambda \in \Lambda} \text{ is complete in } L^2(0, a)\}.$

Now we are ready to state one of the fundamental results of Harmonic Analysis.

Theorem 3.15 (Beurling–Malliavin theorem [6, 7]). *Let* Λ *be a discrete sequence. Then*

$$R(\Lambda) = 2\pi D^*(\Lambda).$$

Let us note that Makarov and Poltoratski considered the Schrödinger equation $Lu = -u'' + qu = z^2u$ and the *m*-function corresponding to this equation, which is obtained by applying the square root transform to the *m*-function we have discussed so far. Let us denote their *m*-function by \tilde{m} .

Theorem 3.16 (Makarov, Poltoratski [37]). Let $\Lambda = {\lambda_n}_{n \in \mathbb{N}}$ be a sequence of discrete non-zero complex numbers, $q \in L^2(0, \pi)$ and $0 \le a \le 1$. The following statements are equivalent:

- (1) q on (0, d) for some d > a and $\{\tilde{m}(\lambda_n)\}_{n \in \mathbb{N}}$ determine q;
- (2) $\pi D^*(\Lambda) \ge 1-a$.

Makarov and Poltoratski's observation shows that Horváth's theorem establishes equivalence between mixed spectral problems for Schrödinger operators and the Beurling–Malliavin problem on completeness of exponentials in L^2 spaces.

In the same paper they obtained an uncertainty version of Borg's theorem.

Theorem 3.17 (Makarov, Poltoratski [37]). Let $\{I_n\}_{n \in \mathbb{N}}$ be a sequence of intervals on \mathbb{R} and $q \in L^2(0, \pi)$. The following statements are equivalent:

- (1) the condition $\sigma_{\text{DD}} \cup \sigma_{\text{ND}} \subset \bigcup_{n \in \mathbb{N}} I_n$ and q on $(0, \epsilon)$ for some $\epsilon > 0$ determine the potential q;
- (2) for any long sequence of intervals $\{J_n\}_{n \in \mathbb{N}}$,

$$\frac{\sum_{I_n \cap J_n} \log_{-} |I_n|}{|J_n|} \not\to 0$$

as $n \to \infty$.

4. An inverse spectral problem with mixed data

4.1. The main result with Dirichlet–Dirichlet and Neumann–Dirichlet boundary conditions. We prove our main result, Theorem 4.2, by representing the Weyl–Titchmarsh *m*-function as an infinite product in terms of Dirichlet–Dirichlet $(\alpha = 0, \beta = 0)$ and Neumann–Dirichlet $(\alpha = \pi/2, \beta = 0)$ spectra. We follow the notations introduced in Example 2.4 for these two spectra, i.e. $\sigma_{DD} := \sigma_{0,0}$ and $\sigma_{ND} := \sigma_{\pi/2,0}$. For simplicity, let us also denote $m_{0,0}$ by *m*. For any infinite product (or sum) defined on an open set $\Omega \subset \mathbb{C}$, normal convergence means that the product (or the sum) converges uniformly on every compact subset of Ω .

Lemma 4.1. The *m*-function of a regular Schrödinger operator $(q \in L^1(0, \pi))$ for Dirichlet–Dirichlet boundary conditions $(\alpha = 0, \beta = 0)$ has representations in terms of Dirichlet–Dirichlet and Neumann–Dirichlet spectra:

$$m(z) = C\left(\frac{z}{b_1} - 1\right) \prod_{n \in \mathbb{N}} \left(\frac{z}{b_{n+1}} - 1\right) \left(\frac{z}{a_n} - 1\right)^{-1},$$
(4.1)

and

$$m(z) = -C \prod_{n \in \mathbb{N}} \left(\frac{z}{b_n} - 1\right) \left(\frac{z}{a_n} - 1\right)^{-1},$$
(4.2)

where C > 0, $\sigma_{DD} = \{a_n\}_{n \in \mathbb{N}}$, $\sigma_{ND} = \{b_n\}_{n \in \mathbb{N}}$ and the product converges normally on $\mathbb{C} \setminus \bigcup_{n \in \mathbb{N}} a_n$.

Proof. Let $m = u'_z(0)/u_z(0)$ be the Weyl *m*-function with boundary conditions $u(\pi) = 0, u'(\pi) = -1$. Since *m* is a meromorphic Herglotz function, $\Theta := \frac{m-i}{m+i}$ is the corresponding meromorphic inner function. See Appendix A for the definition of a meromorphic inner function and the relation between Herglotz and inner functions.

Let us define the $E := \{z \in \mathbb{R} : \operatorname{Im} \Theta(z) > 0\}$. The set E is given in terms of $\sigma_{\text{DD}} = \{a_n\}_{n \in \mathbb{N}} \text{ and } \sigma_{\text{ND}} = \{b_n\}_{n \in \mathbb{N}}, \text{ namely } E = (-\infty, b_1) \cup \bigcup_{n \in \mathbb{N}} (a_n, b_{n+1}).$

The characteristic function of E coincides with the real part of the function $\frac{1}{i\pi} \log\left(i \frac{1+\Theta}{1-\Theta}\right)$ a.e. on \mathbb{R} . Since *m* is a meromorphic Herglotz function mapping \mathbb{R} to \mathbb{R} a.e., $\log(m) = \log(i\frac{1+\Theta}{1-\Theta})$ is a well-defined holomorphic function on \mathbb{C}_+ and its imaginary part takes values 0 and π a.e. on \mathbb{R} . Therefore $\frac{1}{i\pi}\log(m) =$ $\frac{1}{i\pi} \log \left(i \frac{1+\Theta}{1-\Theta} \right)$ and the Schwarz integral of χ_E , S_{χ_E} differ by a purely imaginary number on a.e. \mathbb{R} , i.e.

$$\frac{1}{i\pi}\log\left(i\frac{1+\Theta}{1-\Theta}\right) = S_{\chi_E} + ic = P_{\chi_E} + iQ_{\chi_E} + ic, \quad c \in \mathbb{R},$$

where P and Q are Poisson and conjugate Poisson integrals of χ_E , respectively. Definitions of S, P and O appear in the appendix. Therefore

$$i\frac{1+\Theta}{1-\Theta} = \exp(i\pi S_{\chi_E} - \pi c) = \exp(i\pi P_{\chi_E} - \pi Q_{\chi_E} - \pi c), \quad c \in \mathbb{R}.$$

On the real line, $\exp(S_h) = \exp(h + i\tilde{h})$ for any Poisson-summable function h, where \tilde{h} is the Hilbert transform of *h*. If we let $h := \chi_E$, then

$$\tilde{h}(x) = \frac{1}{\pi} \Big[\log \Big(\frac{\sqrt{1+b_1^2}}{|x-b_1|} \Big) + \sum_{n \in \mathbb{N}} \log \Big(\frac{|x-a_n|}{|x-b_{n+1}|} \Big) + \frac{1}{2} \sum_{n \in \mathbb{N}} \log \Big(\frac{1+b_{n+1}^2}{1+a_n^2} \Big) \Big].$$

Therefore

$$\exp(-\pi\tilde{h}(x)) = \frac{|x-b_1|}{\sqrt{1+b_1^2}} \prod_{n \in \mathbb{N}} \frac{|x-b_{n+1}|}{|x-a_n|} \prod_{n \in \mathbb{N}} \left(\frac{1+a_n^2}{1+b_{n+1}^2}\right)^{1/2}.$$

Noting that $\exp(i\pi h)$ is -1 on E and 1 on $\mathbb{R} \setminus E$, the Weyl *m*-function can be given in terms of σ_{DD} and σ_{ND} a.e. on \mathbb{R} :

$$\begin{split} m(x) &= i \frac{1 + \Theta(x)}{1 - \Theta(x)} \\ &= \exp(i \pi S_{\chi E} - \pi c) \\ &= \frac{x - b_1}{\sqrt{1 + b_1^2}} \prod_{n \in \mathbb{N}} \frac{x - b_{n+1}}{x - a_n} \prod_{n \in \mathbb{N}} \left(\frac{1 + a_n^2}{1 + b_{n+1}^2}\right)^{1/2} \exp(-\pi c) \\ &= C\left(\frac{x}{b_1} - 1\right) \prod_{n \in \mathbb{N}} \left(\frac{x}{b_{n+1}} - 1\right) \left(\frac{x}{a_n} - 1\right)^{-1}, \\ C &= \exp(-\pi c) \prod_{n \in \mathbb{N}} \frac{\sqrt{1 + a_n^2}}{a_n} \frac{b_n}{\sqrt{1 + a_n^2}}. \end{split}$$

where ($\sqrt{1+b_n^2}$ Since m(z) and $C(\frac{z}{b_1}-1)\prod_{n\in\mathbb{N}}(\frac{z}{b_{n+1}}-1)(\frac{z}{a_n}-1)^{-1}$ are meromorphic functions that agree a.e. on \mathbb{R} , they are identical by the identity theorem for meromorphic functions. This gives the first representation (4.1). The second representation (4.2) follows from normal convergence of $\{z/b_n - 1\}_{n\in\mathbb{N}}$ to -1 in \mathbb{C} .

Using this representation of the *m*-function, we prove our main result. At this point let us note that the points in a spectrum are enumerated in increasing order, which is done following the asymptotics (2.4)-(2.7).

Theorem 4.2. Let $q \in L^1(0, \pi)$ and $A \subseteq \mathbb{N}$. Then $\{a_n\}_{n \in \mathbb{N}}$, $\{b_n\}_{n \in \mathbb{N} \setminus A}$ and $\{\gamma_n\}_{n \in A}$ determine the potential q, where $\sigma_{DD} = \{a_n\}_{n \in \mathbb{N}}$, $\sigma_{ND} = \{b_n\}_{n \in \mathbb{N}}$ are Dirichlet–Dirichlet and Neumann–Dirichlet spectra and $\{\gamma_n\}_{n \in \mathbb{N}}$ are point masses of the spectral measure $\mu_{0,0} = \sum_{n \in \mathbb{N}} \gamma_n \delta_{a_n}$.

Proof. By representation (2.8) of the *m*-function as a Herglotz integral of the spectral measure, knowing γ_n means knowing $\text{Res}(m, a_n)$. Therefore, in terms of the *m*-function our claim says that the set of poles, $\{a_n\}_{n \in \mathbb{N}}$, the set of zeros except the index set A, $\{b_n\}_{n \in \mathbb{N} \setminus A}$, and the residues with the same index set A, $\{\text{Res}(m, a_n)\}_{n \in A}$ determine the *m*-function uniquely. Before starting to prove this claim let us briefly list the main steps of the proof. We will use similar ideas to prove our results in non-matching index sets case and for general boundary conditions.

Step 1. Reduce the claim to the problem of unique recovery of the infinite product

$$G(z) := -C \prod_{n \in A} \left(\frac{z}{b_n} - 1\right) \left(\frac{z}{a_n} - 1\right)^{-1}$$

from its sets of poles and residues.

- Step 2. Observe that G(z) is a meromorphic Herglotz function and has a representation in terms of its poles, residues and a linear polynomial dz + e.
- Step 3. Show uniqueness of d.
- Step 4. Show uniqueness of *e*.
- Step 5. Use the representation from Step 2 to get uniqueness of the two spectra and prove the claim by Borg's theorem.

Step 1. From Lemma 4.1, the Weyl *m*-function can be represented in terms of σ_{DD} and σ_{ND} ,

$$m(z) = -C \prod_{n \in \mathbb{N}} \left(\frac{z}{b_n} - 1\right) \left(\frac{z}{a_n} - 1\right)^{-1}.$$

Note that for any $k \in A$, we know

$$\operatorname{Res}(m, a_k) = C(b_k - a_k) \frac{a_k}{b_k} \prod_{\substack{n \in \mathbb{N} \\ n \neq k}} \left(\frac{a_k}{b_n} - 1\right) \left(\frac{a_k}{a_n} - 1\right)^{-1}.$$
 (4.3)

Let m(z) = F(z)G(z), where F and G are two infinite products defined as

$$G(z) := -C \prod_{n \in A} \left(\frac{z}{b_n} - 1\right) \left(\frac{z}{a_n} - 1\right)^{-1}, \quad F(z) := \prod_{n \in \mathbb{N} \setminus A} \left(\frac{z}{b_n} - 1\right) \left(\frac{z}{a_n} - 1\right)^{-1}.$$

Also note that at any point of $\{a_n\}_{n \in A}$, the infinite product

$$F(z) = \prod_{n \in \mathbb{N} \setminus A} \left(\frac{z}{b_n} - 1\right) \left(\frac{z}{a_n} - 1\right)^{-1}$$
(4.4)

is known.

Conditions (4.3) and (4.4) imply that for any $k \in A$, we know

$$\operatorname{Res}(G, a_k) = \frac{\operatorname{Res}(m, a_k)}{F(a_k)},$$

i.e. we know all of the poles and residues of G(z), but none of its zeros. We claim that G(z) can be uniquely recovered from this data set.

Step 2. Let us observe that $\arg(G(z)) = \pi - \sum_{n \in A} [\arg(z - b_n) - \arg(z - a_n)]$. Since zeros and poles of G(z) are real and interlacing, $0 < \arg(G(z)) < \pi$ for any z in the upper half plane, i.e. G(z) is a meromorphic Herglotz function. Therefore by Čebotarev's theorem, see Theorem A.1, G(z) has the representation

$$G(z) = dz + e + \sum_{n \in A} A_n \left(\frac{1}{a_n - z} - \frac{1}{a_n} \right),$$
(4.5)

where $d \ge 0$, $e \in \mathbb{R}$ and $\sum_{n \in A} A_n / a_n^2$ is absolutely convergent.

Note that $A_k = -\text{Res}(G(z), a_k)$ for any $k \in A$, which means there are only two unknowns on the right hand side of (4.5), namely constants d and e.

Step 3. Now let us show uniqueness of G(z) by showing uniqueness of dz + e. Let $\tilde{G}(z)$ be another infinite product sharing same properties with G(z), namely:

• the infinite product $\tilde{G}(z)$ is defined as

$$\widetilde{G}(z) := -\widetilde{C} \prod_{n \in A} \left(\frac{z}{\widetilde{b}_n} - 1 \right) \left(\frac{z}{\widetilde{a}_n} - 1 \right)^{-1},$$

where $\tilde{C} > 0$, the set of poles $\{\tilde{a}_n\}_{n \in A}$ satisfies asymptotics (2.5) and the set of zeros $\{\tilde{b}_n\}_{n \in A}$ satisfies asymptotics (2.6);

- G(z) and G̃(z) share same set of poles with equivalent residues at the corresponding poles, i.e. ã_k = a_k and Res(G̃, a_k) = Res(G, a_k) for any k ∈ A;
- by the equivalence of poles and residues of G(z) and $\tilde{G}(z)$ and Čebotarev's theorem, $\tilde{G}(z)$ has the representation

$$\widetilde{G}(z) = \widetilde{d}z + \widetilde{e} + \sum_{n \in A} A_n \Big(\frac{1}{a_n - z} - \frac{1}{a_n} \Big), \tag{4.6}$$

where $\tilde{d} \ge 0, \tilde{e} \in \mathbb{R}$.

Note that we defined \tilde{a}_n and \tilde{b}_n only for $n \in A$. Let $\tilde{a}_n := a_n$ and $\tilde{b}_n := b_n$ for every $n \in \mathbb{N} \setminus A$. Let us also note that $\{\tilde{a}_n\}_{n \in \mathbb{N}}$ and $\{\tilde{b}_n\}_{n \in \mathbb{N}}$ are interlacing sequences so that they represent two spectra of a potential function $\tilde{q}(x)$.

Let $k \in A$ and $b_k \neq \tilde{b}_k$. Since $G(b_k) = 0$ and $\tilde{G}(\tilde{b}_k) = 0$, using representations (4.5) and (4.6) we get

$$-db_k - e = \sum_{n \in A} A_n \Big(\frac{1}{a_n - b_k} - \frac{1}{a_n} \Big), \tag{4.7}$$

$$-\tilde{d}\tilde{b}_k - \tilde{e} = \sum_{n \in A} A_n \Big(\frac{1}{a_n - \tilde{b}_k} - \frac{1}{a_n} \Big), \tag{4.8}$$

$$G(\tilde{b}_k) = G(\tilde{b}_k) - \tilde{G}(\tilde{b}_k) = (d - \tilde{d})\tilde{b}_k + e - \tilde{e}.$$
(4.9)

Replacing $e - \tilde{e}$ by $G(\tilde{b}_k) - (d - \tilde{d})\tilde{b}_k$ and taking difference of (4.7) and (4.8) we get

$$db_k - \tilde{d}\tilde{b}_k - d\tilde{b}_k + \tilde{d}\tilde{b}_k + G(\tilde{b}_k) = \sum_{n \in A} A_n \Big(\frac{\tilde{b}_k - b_k}{(a_n - \tilde{b}_k)(a_n - b_k)}\Big).$$

Dividing both sides by $\tilde{b}_k(\tilde{b}_k - b_k)$ we get

$$\frac{-d}{\tilde{b}_k} + \frac{G(b_k)}{\tilde{b}_k(\tilde{b}_k - b_k)} = \sum_{n \in A} \left(\frac{A_n}{\tilde{b}_k(a_n - \tilde{b}_k)(a_n - b_k)} \right).$$
(4.10)

Note that since $\{a_n\}_{n \in A}$ satisfies asymptotics (2.5) and $\{b_n\}_{n \in A}$, $\{\tilde{b}_n\}_{n \in A}$ satisfy asymptotics (2.6), the inequalities

$$|\tilde{b}_k(a_n - b_k)(a_n - \tilde{b}_k)|^{-1} \le |\tilde{b}_n(a_n - b_n)(a_n - \tilde{b}_n)|^{-1} \le 2/a_n^2$$
(4.11)

are valid for any $k \in A$, for sufficiently large $n \in A$. In addition, $\sum_{n \in A} A_n/a_n^2$ is absolutely convergent. Therefore right hand side of (4.10) converges to 0 as k goes to ∞ . Also note that by (4.9), left hand side of (4.10) is

$$\frac{-d}{\tilde{b}_k} + \frac{G(b_k)}{\tilde{b}_k(\tilde{b}_k - b_k)} = \frac{-d}{\tilde{b}_k} + \frac{G(b_k) - \tilde{G}(b_k)}{\tilde{b}_k(\tilde{b}_k - b_k)}$$

$$= -\frac{d}{\tilde{b}_k} + \frac{1}{\tilde{b}_k - b_k} \Big[d - \tilde{d} + \frac{e - \tilde{e}}{\tilde{b}_k} \Big].$$
(4.12)

Now let us show $\tilde{b}_k - b_k$ converges to 0 as k goes to ∞ . Recall that poles of G and \tilde{G} satisfy asymptotics

$$n^2 + \frac{1}{\pi} \int_0^{\pi} q(x)dx + \alpha_n$$
 and $n^2 + \frac{1}{\pi} \int_0^{\pi} \tilde{q}(x)dx + \tilde{\alpha}_n$,

respectively, where $\alpha_n = o(1)$ and $\tilde{\alpha}_n = o(1)$ as $n \to \infty$. Equivalance of poles of G and \tilde{G} imply equivalence of $\int_0^{\pi} q(x) dx$ and $\int_0^{\pi} \tilde{q}(x) dx$. Therefore b_n and \tilde{b}_n satisfy asymptotics

$$\left(n-\frac{1}{2}\right)^2 + \frac{1}{\pi}\int_0^{\pi} q(x)dx + \beta_n \text{ and } \left(n-\frac{1}{2}\right)^2 + \frac{1}{\pi}\int_0^{\pi} q(x)dx + \tilde{\beta}_n,$$

where $\beta_n = o(1)$ and $\tilde{\beta}_n = o(1)$ as $n \to \infty$. Hence $\tilde{b}_k - b_k = o(1)$ as k goes to ∞ . Therefore by (4.12), left hand side of (4.10) goes to ∞ if $d - \tilde{d} \neq 0$, so we get a contradiction unless $d = \tilde{d}$. This implies that $G(z) - \tilde{G}(z)$ is a real constant, which is $G(0) - \tilde{G}(0) = \tilde{C} - C$.

Step 4. Now let us show $\tilde{C} - C = 0$. Positivity of $(\tilde{b}_k - b_n)/(\tilde{b}_k - a_n)$ for all $n \neq k$, which follows from interlacing property of $\{a_n\}_{n \in \mathbb{N}}$ and $\{b_n\}_{n \in \mathbb{N}}$ and interlacing property of $\{a_n\}_{n \in \mathbb{N}}$ and $\{\tilde{b}_n\}_{n \in \mathbb{N}}$, implies $\operatorname{sgn}(\tilde{C} - C) = \operatorname{sgn}(\tilde{\beta}_k - \beta_k)$ for all $k \in \mathbb{N}$, i.e. $\{b_n\}_{n \in A}$ and $\{\tilde{b}_n\}_{n \in A}$ are interlacing sequences.

Let us assume $\tilde{C} > C$ and wlog the two spectra lie on the positive real line. This implies $\tilde{b}_n > b_n$ for all $n \in A$. Observe that $\prod_{n \in A} \tilde{b}_n / b_n$ is finite, since

$$\sum_{n \in A} \frac{\tilde{b}_n - b_n}{b_n} = \sum_{n \in A} \frac{\tilde{\beta}_n - \beta_n}{b_n} \le \max_{n \in A} (\tilde{\beta}_n - \beta_n) \sum_{n \in A} \frac{1}{b_n} < \infty$$

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Therefore the infinite product $H(z) := G(z)/\tilde{G}(z)$ is represented as

$$H(z) := \frac{G(z)}{\widetilde{G}(z)} = \frac{C}{\widetilde{C}} \prod_{n \in A} \frac{z - b_n}{b_n} \frac{\widetilde{b}_n}{z - \widetilde{b}_n} = \frac{C}{\widetilde{C}} \prod_{n \in A} \frac{\widetilde{b}_n}{b_n} \prod_{n \in A} \frac{z - b_n}{z - \widetilde{b}_n}$$

Let us denote the constant factor of H(z) by $N := (C/\tilde{C}) \prod_{n \in A} \tilde{b}_n/b_n$. Then by interlacing property of $\{b_n\}_{n \in A}$ and $\{\tilde{b}_n\}_{n \in A}$, the infinite product -H is a meromorphic Herglotz function and hence by Theorem A.1 it is represented as

$$-H(z) = -N \prod_{n \in A} \frac{z - b_n}{z - \tilde{b}_n} = Dz + E + \sum_{n \in A} B_n \left(\frac{1}{z - \tilde{b}_n} + \frac{1}{\tilde{b}_n}\right), \quad (4.13)$$

where $B_k = -\operatorname{Res}(H, \tilde{b}_k)$ and $D, E \in \mathbb{R}$.

Now let us show that $\{B_k/\tilde{b}_k\}_{k \in A}$ is summable:

$$\begin{split} \frac{B_k}{\tilde{b}_k} \bigg| &= N \; \frac{\tilde{b}_k - b_k}{b_k} \prod_{\substack{n \in A}} \frac{\tilde{b}_k - b_n}{\tilde{b}_k - \tilde{b}_n} \\ &\leq N \; \frac{\tilde{b}_k - b_k}{\tilde{b}_k} \prod_{\substack{n \in A}} \frac{\tilde{b}_k - b_n}{\tilde{b}_k - \tilde{b}_n} \\ &= N \; \frac{\tilde{b}_k - b_k}{\tilde{b}_k} \prod_{\substack{n \in A \\ 1 \le n \le k - 1}} \left(1 + \frac{\tilde{b}_n - b_n}{\tilde{b}_k - \tilde{b}_n} \right) \\ &= N \; \frac{\tilde{b}_k - b_k}{\tilde{b}_k} \prod_{\substack{n \in A \\ 1 \le n \le k - 1}} \left(1 + \frac{\tilde{\beta}_n - \beta_n}{(k - 1/2)^2 - (n - 1/2)^2 + \tilde{\beta}_k - \tilde{\beta}_n} \right) \\ &\leq N \; \frac{\tilde{b}_k - b_k}{\tilde{b}_k} \prod_{\substack{n \in A \\ 1 \le n \le k - 1}} \left(1 + \frac{\tilde{\beta}_n - \beta_n}{(n + 1 - 1/2)^2 - (n - 1/2)^2 + \tilde{\beta}_k - \tilde{\beta}_n} \right) \\ &\leq N \; \frac{\tilde{b}_k - b_k}{\tilde{b}_k} \prod_{\substack{n \in A \\ 1 \le n \le k - 1}} \left(1 + \frac{\tilde{\beta}_n - \beta_n}{(n + 1 - 1/2)^2 - (n - 1/2)^2 + \tilde{\beta}_k - \tilde{\beta}_n} \right) \\ &\leq N \; \frac{\tilde{b}_k - b_k}{\tilde{b}_k} M \; \prod_{n=1}^{n \in A} \left(1 + \frac{1}{2n} \right), \end{split}$$

for sufficiently large k, where M is a real constant independent of k. Since $\tilde{b}_k - b_k = o(1)$, $\tilde{b}_k = O(k^2)$ and $\prod_{n=1}^{k-1}(1+1/2n) = O(\sqrt{k})$ as k goes to ∞ , we get the asymptotics $B_k/\tilde{b}_k = o(1/k^{3/2})$ as k goes to ∞ and hence $\sum_{n \in A} B_n/\tilde{b}_n$ is absolutely convergent. Then by letting z tend to $-\infty$ in (4.13) we get

$$-N = \lim_{t \to -\infty} \left(Dt + E + \sum_{n \in A} \frac{B_n}{\tilde{b}_n} + \sum_{n \in A} \frac{B_n}{t - \tilde{b}_n} \right)$$

and hence D = 0 and $-N = E + \sum_{n \in A} B_n / \tilde{b}_n$, i.e. -H(z) has the representation

$$-H(z) = N - \sum_{n \in A} \frac{B_n}{z - \tilde{b}_n}.$$
(4.14)

Noting that $H(b_k) = 0$ and $\operatorname{Res}(G, a_k) = \operatorname{Res}(\widetilde{G}, a_k)$, i.e. $H(a_k) = 1$ for all $k \in A$, we get

$$1 = H(a_k) - H(b_k) = -N + \sum_{n \in A} \frac{B_n}{a_k - \tilde{b}_n} + N - \sum_{n \in A} \frac{B_n}{b_k - \tilde{b}_n}$$
$$= \sum_{n \in A} B_n \frac{(b_k - a_k)}{(a_k - \tilde{b}_n)(b_k - \tilde{b}_n)}.$$

Each term of the infinite sum on the right end is positive, so by letting k go to ∞ we get the following contradiction.

$$1 = \lim_{k \to \infty} \sum_{n \in A} B_n \frac{(b_k - a_k)}{(a_k - \tilde{b}_n)(b_k - \tilde{b}_n)} = \sum_{n \in A} B_n \lim_{k \to \infty} \frac{(b_k - a_k)}{(a_k - \tilde{b}_n)(b_k - \tilde{b}_n)} = 0.$$

Similar arguments give another contradiction when $\tilde{C} < C$, so $C = \tilde{C}$.

Step 5. Step 4 implies uniqueness of dz + e, i.e. uniqueness of G(z) and hence uniqueness of $\{b_n\}_{n \in A}$. After unique recovery of the two spectra $\sigma_{DD} = \{a_n\}_{n \in \mathbb{N}}$ and $\sigma_{ND} = \{b_n\}_{n \in \mathbb{N}}$, the potential is uniquely determined by Borg's theorem. \Box

Remark 4.3. If we let $A = \mathbb{N}$, Theorem 4.2 gives Marchenko's theorem with Dirichlet–Dirichlet, Neumann–Dirichlet boundary conditions as a corollary. By letting $A = \emptyset$, we get the statement of Borg's theorem with Dirichlet–Dirichlet, Neumann–Dirichlet boundary conditions.

Remark 4.4. Spectral data of Theorem 4.2 can be seen as $\{a_n\}_{n \in \mathbb{N}}, \{b_n\}_{n \in \mathbb{N} \setminus A}$ and $\{\tau_{\alpha}(a_n)\}_{n \in A}$, where $\{\tau_{\alpha}(a_n)\}_{n \in A}$ is the set of norming constants for $\sigma_{DD} = \{a_n\}_{n \in \mathbb{N}}$.

4.2. Non-matching index sets. If the known point masses of the spectral measure and unknown eigenvalues of the Neumann–Dirichlet spectrum have different index sets, one needs some control over eigenvalues of the Dirichlet– Dirichlet spectrum corresponding to known point masses and unknown part of the Neumann–Dirichlet spectrum. In this case we get a Čebotarev type representation result. Before the statement, let us clarify the notations we use. For any subsequence $\{a_{k_n}\}_{n \in \mathbb{N}} \subset \sigma_{\text{DD}}$ and $\{b_{l_n}\}_{n \in \mathbb{N}} \subset \sigma_{\text{ND}}$, by $A_{k_n,m}$ and A_{k_n} we denote

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the residues at a_{k_n} of partial and infinite products, respectively, consisting of these subsequences:

$$A_{k_n,m} := \operatorname{Res}(G_m, a_{k_n}) = \frac{a_{k_n}}{b_{l_n}} (a_{k_n} - b_{l_n}) \prod_{\substack{1 \le j \le m \\ j \ne n}} \frac{a_{k_j}}{a_{k_n}} \frac{a_{k_n} - b_{l_j}}{a_{k_n} - a_{k_j}},$$
$$A_{k_n} := \operatorname{Res}(G, a_{k_n}) = \frac{a_{k_n}}{b_{l_n}} (a_{k_n} - b_{l_n}) \prod_{\substack{j \in \mathbb{N} \\ j \ne n}} \frac{a_{k_j}}{b_{l_j}} \frac{a_{k_n} - b_{l_j}}{a_{k_n} - a_{k_j}},$$

where

$$G_m(z) := \prod_{n=1}^m \left(\frac{z}{b_{l_n}} - 1\right) \left(\frac{z}{a_{k_n}} - 1\right)^{-1}, \quad G(z) := \prod_{n \in \mathbb{N}} \left(\frac{z}{b_{l_n}} - 1\right) \left(\frac{z}{a_{k_n}} - 1\right)^{-1}.$$

Note that these subsequences are ordered according to their indices, i.e. $a_{k_n} < a_{k_{n+1}}$ and $b_{l_n} < b_{l_{n+1}}$ for any $n \in \mathbb{N}$. This follows from the asymptotics of the spectra.

Lemma 4.5. Let $\{a_{k_n}\}_{n \in \mathbb{N}} \subset \sigma_{DD}$ and $\{b_{l_n}\}_{n \in \mathbb{N}} \subset \sigma_{ND}$ satisfy following properties:

• $\lim_{m \to \infty} \sum_{n=1}^{m} (|A_{k_n,m} - A_{k_n}|/a_{k_n}^2) < \infty;$ • $\{A_{k_n}/a_{k_n}^2\}_{m \in \mathbb{N}} \in l^1.$

•
$$\{A_{k_n}/a_{k_n}^-\}_{n\in\mathbb{N}}\in$$

Then

$$G(z) = cz^{2} + dz + e + \sum_{n \in \mathbb{N}} A_{k_{n}} \left(\frac{1}{z - a_{k_{n}}} + \frac{1}{a_{k_{n}}} \right),$$
(4.15)

where c, d, e are real numbers, A_{k_n} is the residue of G(z) at the point $z = a_{k_n}$ and the sum converges normally on $\mathbb{C} \setminus \bigcup_{n \in \mathbb{N}} a_{k_n}$.

Proof. Let p(z) be the difference of G(z) and the infinite sum on the right hand side of (4.15). Then, p(z) is an entire function, since the infinite product and the infinite sum share the same set of poles with equivalent degrees and residues. We represent $G_m(z)$ as partial sums:

$$\prod_{n=1}^{m} \left(\frac{z}{b_{l_n}} - 1\right) \left(\frac{z}{a_{k_n}} - 1\right)^{-1} = \sum_{n=1}^{m} A_{k_n,m} \left(\frac{1}{z - a_{k_n}} + \frac{1}{a_{k_n}}\right) + 1$$

where $A_{k_n,m} = \operatorname{Res}(G_m, a_{k_n})$.

Let C_n be the circle with radius b_{l_n} centered at the origin. This sequence of circles satisfy following properties:

- C_n omits all the poles a_{k_n} ;
- each C_n lies inside C_{n+1} ;
- the radius of C_n , b_{l_n} diverges to infinity as *n* goes to infinity.

Then,

$$\begin{split} \max_{z \in C_{t}} \left| \frac{p(z) - 1}{b_{l_{t}}^{2}} \right| &= \max_{z \in C_{t}} \left| \frac{G(z) - 1 - \sum_{n \in \mathbb{N}} A_{k_{n}} \left(\frac{1}{z - a_{k_{n}}} + \frac{1}{a_{k_{n}}} \right)}{b_{l_{t}}^{2}} \right| \\ &= \frac{1}{b_{l_{t}}^{2}} \max_{z \in C_{t}} \lim_{m \to \infty} \left| \sum_{n=1}^{m} A_{k_{n},m} \left(\frac{1}{z - a_{k_{n}}} + \frac{1}{a_{k_{n}}} \right) \right| \\ &- \sum_{n=1}^{m} A_{k_{n}} \left(\frac{1}{z - a_{k_{n}}} + \frac{1}{a_{k_{n}}} \right) \right| \\ &= \lim_{m \to \infty} \frac{1}{b_{l_{t}}^{2}} \max_{z \in C_{t}} \left| \sum_{n=1}^{m} (A_{k_{n},m} - A_{k_{n}}) \frac{z}{a_{k_{n}}(z - a_{k_{n}})} \right| \\ &\leq \lim_{m \to \infty} \frac{1}{b_{l_{t}}^{2}} \sum_{n=1}^{m} |A_{k_{n},m} - A_{k_{n}}| \frac{b_{l_{t}}}{a_{k_{n}}|b_{l_{t}} - a_{k_{n}}|} \\ &= \lim_{m \to \infty} \sum_{n=1}^{m} |A_{k_{n},m} - A_{k_{n}}| \frac{1}{a_{k_{n}}b_{l_{t}}|b_{l_{t}} - a_{k_{n}}|} \\ &\leq \lim_{m \to \infty} \sum_{n=1}^{m} |A_{k_{n},m} - A_{k_{n}}| \frac{1}{a_{k_{n}}b_{l_{t}}|b_{l_{t}} - a_{k_{n}}|} \\ &\leq \lim_{m \to \infty} \sum_{n=1}^{m} |A_{k_{n},m} - A_{k_{n}}| \frac{1}{a_{k_{n}}b_{l_{t}}|b_{l_{t}} - a_{k_{n}}|} \\ &\leq \lim_{m \to \infty} \sum_{n=1}^{m} |A_{k_{n},m} - A_{k_{n}}| \frac{1}{a_{k_{n}}b_{l_{t}}|b_{l_{t}} - a_{k_{n}}|} \\ &\leq \lim_{m \to \infty} \sum_{n=1}^{m} |A_{k_{n},m} - A_{k_{n}}| \frac{1}{a_{k_{n}}b_{l_{t}}|b_{l_{t}} - a_{k_{n}}|} \\ &\leq \lim_{m \to \infty} C' \sum_{n=1}^{m} \frac{|A_{k_{n},m} - A_{k_{n}}|}{a_{k_{n}}^{2}} < \infty. \end{split}$$

Note that the second inequality is a consequence of

$$\sup_{t \in \mathbb{N}} (b_{l_t} | b_{l_t} - a_{k_n} |)^{-1} \le (b_{l_1} | b_{l_1} - a_{k_n} |)^{-1},$$

which follows from asymptotics of $\{a_n\}_{n \in \mathbb{N}}$ and $\{b_n\}_{n \in \mathbb{N}}$. Therefore $|p(z) - 1| \leq C''|z|^2$ on the circle C_t for any $t \in \mathbb{N}$, where C' and C'' are real numbers. By the maximum modulus theorem and the entireness of p(z), we conclude that p(z) is a polynomial of at most second degree. Since G(0), G'(0) and G''(0) are real numbers, $c, d, e \in \mathbb{R}$.

Using this Čebotarev type representation we prove our main result in nonmatching index sets case with Dirichlet–Dirichlet, Neumann–Dirichlet boundary conditions. However, we need extra information of an eigenvalue from $\{b_{l_n}\}_{n \in \mathbb{N}}$.

Theorem 4.6. Let $q \in L^1(0, \pi)$, and $\{a_{k_n}\}_{n \in \mathbb{N}} \subset \sigma_{\text{DD}}, \{b_{l_n}\}_{n \in \mathbb{N}} \subset \sigma_{\text{ND}}$ satisfy following properties:

- $\lim_{m\to\infty} \sum_{n=1}^m (|A_{k_n,m} A_{k_n}|/a_{k_n}^2) < \infty;$
- $\{A_{k_n}/a_{k_n}^2\}_{n \in \mathbb{N}} \in l^1$.

Then $\{a_n\}_{n\in\mathbb{N}}, \{b_n\}_{n\in\mathbb{N}}\setminus\{b_{l_n}\}_{n\in\mathbb{N}\setminus\{s\}}$ and $\{\gamma_{k_n}\}_{n\in\mathbb{N}}$ determine the potential q for any $s\in\mathbb{N}$, where $\sigma_{\text{DD}}=\{a_n\}_{n\in\mathbb{N}}, \sigma_{\text{ND}}=\{b_n\}_{n\in\mathbb{N}}$ are Dirichlet–Dirichlet and Neumann–Dirichlet spectra and $\{\gamma_n\}_{n\in\mathbb{N}}$ are point masses of the spectral measure $\mu_{0,0}=\sum_{n\in\mathbb{N}}\gamma_n\delta_{a_n}$.

Proof. By representation of the *m*-function as the Herglotz integral of the spectral measure, from γ_n , we know $\text{Res}(m, a_n)$. Therefore, in terms of the *m*-function our claim says that the set of poles, $\{a_n\}_{n \in \mathbb{N}}$, the set of zeros except the index set $\{l_n\}_{n \in \mathbb{N} \setminus \{s\}}, \{b_{l_s}\} \cup \{b_n\}_{n \in \mathbb{N} \setminus \{l_n\}_{n \in \mathbb{N}}}$, and the residues with the index set $\{k_n\}_{n \in \mathbb{N}}, \{\text{Res}(m, a_{k_n})\}_{n \in \mathbb{N}}$ determine the *m*-function uniquely.

From Lemma 4.1, the Weyl *m*-function can be represented in terms of σ_{DD} and σ_{ND} ,

$$m(z) = -C \prod_{n \in \mathbb{N}} \left(\frac{z}{b_n} - 1\right) \left(\frac{z}{a_n} - 1\right)^{-1}.$$

Note that for any $n \in \mathbb{N}$, we know

$$\operatorname{Res}(m, a_{k_n}) = C(b_{k_n} - a_{k_n}) \frac{a_{k_n}}{b_{k_n}} \prod_{\substack{j \in \mathbb{N} \\ j \neq k_n}} \left(\frac{a_{k_n}}{b_j} - 1 \right) \left(\frac{a_{k_n}}{a_j} - 1 \right)^{-1}.$$
 (4.16)

Let m(z) = F(z)G(z), where F and G are two infinite products defined as

$$G(z) := -C \prod_{n \in \mathbb{N}} \left(\frac{z}{b_{l_n}} - 1\right) \left(\frac{z}{a_{k_n}} - 1\right)^{-1},$$

$$F(z) := \prod_{n \in \mathbb{N} \setminus \{l_n\}_{n \in \mathbb{N}}} \left(\frac{z}{b_n} - 1\right) \prod_{n \in \mathbb{N} \setminus \{k_n\}_{n \in \mathbb{N}}} \left(\frac{z}{a_n} - 1\right)^{-1},$$

Also note that $F(a_{k_n})$ is known for any $n \in \mathbb{N}$. This condition and (4.16) imply that for any $n \in \mathbb{N}$, we know

$$\operatorname{Res}(G, a_{k_n}) = \frac{\operatorname{Res}(m, a_{k_n})}{F(a_{k_n})}.$$

By Lemma 4.5, G(z) has the following representation

$$G(z) = cz^{2} + dz + e + \sum_{n \in \mathbb{N}} A_{k_{n}} \left(\frac{1}{z - a_{k_{n}}} + \frac{1}{a_{k_{n}}} \right),$$
(4.17)

where $A_{k_n} = \text{Res}(G, a_{k_n})$. In order to show uniqueness of G(z), let us consider $\tilde{G}(z)$ similar to the proof of Theorem 4.2, i.e. $\tilde{G}(z)$ has the following properties.

• The infinite product $\tilde{G}(z)$ is defined as

$$\widetilde{G}(z) := -\widetilde{C} \prod_{n \in \mathbb{N}} \left(\frac{z}{\widetilde{b}_{l_n}} - 1 \right) \left(\frac{z}{\widetilde{a}_{k_n}} - 1 \right)^{-1},$$

where $\tilde{C} > 0$, the set of poles $\{\tilde{a}_{k_n}\}_{n \in \mathbb{N}}$ satisfies asymptotics (2.5) and the set of zeros $\{\tilde{b}_{l_n}\}_{n \in \mathbb{N}}$ satisfies asymptotics (2.6). For the given eigenvalues from $\sigma_{\text{ND}} = \{b_n\}_{n \in \mathbb{N}}$, let \tilde{b}_n be defined as b_n , i.e. $\tilde{b}_j := b_j$ for $j \in \mathbb{N} \setminus \{l_n\}_{n \in \mathbb{N}}$. Similarly let $\tilde{a}_j := a_j$ for $j \in \mathbb{N} \setminus \{k_n\}_{n \in \mathbb{N}}$.

- G(z) and G̃(z) share same set of poles with equivalent residues at the corresponding poles, i.e. ã_{kn} = a_{kn} and Res(G̃, a_{kn}) = Res(G, a_{kn}) for any n ∈ N.
- G(z) and $\tilde{G}(z)$ share one zero, namely $b_{l_s} = \tilde{b}_{l_s}$.
- By the equivalence of poles and residues of G(z) and $\tilde{G}(z)$ and Lemma 4.5, $\tilde{G}(z)$ has the representation

$$\tilde{G}(z) = \tilde{c}z^2 + \tilde{d}z + \tilde{e} + \sum_{n \in \mathbb{N}} A_{k_n} \Big(\frac{1}{a_{k_n} - z} - \frac{1}{a_{k_n}} \Big),$$
(4.18)

where $\tilde{c}, \tilde{d}, \tilde{e} \in \mathbb{R}$.

Let $m \in \mathbb{N} \setminus \{s\}$ and $b_{l_m} \neq \tilde{b}_{l_m}$. Since $G(b_{l_m}) = 0$ and $\tilde{G}(\tilde{b}_{l_m}) = 0$, using representations (4.17) and (4.18) we get

$$-cb_{l_m}^2 - db_{l_m} - e = \sum_{n \in \mathbb{N}} A_{k_n} \left(\frac{1}{a_{k_n} - b_{l_m}} - \frac{1}{a_{k_n}} \right), \tag{4.19}$$

$$-\tilde{c}\tilde{b}_{l_m}^2 - \tilde{d}\tilde{b}_{l_m} - \tilde{e} = \sum_{n \in \mathbb{N}} A_{k_n} \Big(\frac{1}{a_{k_n} - \tilde{b}_{l_m}} - \frac{1}{a_{k_n}} \Big),$$
(4.20)

$$G(\tilde{b}_{lm}) = G(\tilde{b}_{lm}) - \tilde{G}(\tilde{b}_{lm}) = (c - \tilde{c})\tilde{b}_{lm}^2 + (d - \tilde{d})\tilde{b}_{lm} + e - \tilde{e}.$$
 (4.21)

Taking difference of (4.19) and (4.20) and replacing $e - \tilde{e}$ by

$$G(\tilde{b}_{l_m}) - (c - \tilde{c})\tilde{b}_{l_m}^2 - (d - \tilde{d})\tilde{b}_{l_m}$$

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we get

$$cb_{l_m}^2 - c\tilde{b}_{l_m}^2 + db_{l_m} - d\tilde{b}_{l_m} + G(\tilde{b}_{l_m}) = \sum_{n \in \mathbb{N}} A_{k_n} \Big(\frac{b_{l_m} - b_{l_m}}{(a_{k_n} - \tilde{b}_{l_m})(a_{k_n} - b_{l_m})} \Big).$$

Dividing both sides by $\tilde{b}_{lm}(\tilde{b}_{lm} - b_{lm})$ we get

$$\frac{-c(b_{l_m} + \tilde{b}_{l_m})}{\tilde{b}_{l_m}} + \frac{-d}{\tilde{b}_{l_m}} + \frac{G(\tilde{b}_{l_m})}{\tilde{b}_{l_m}(\tilde{b}_{l_m} - b_{l_m})} = \sum_{n \in \mathbb{N}} \left(\frac{A_{k_n}}{\tilde{b}_{l_m}(a_{k_n} - \tilde{b}_{l_m})(a_{k_n} - b_{l_m})} \right).$$
(4.22)

Note that since $\{a_n\}_{n \in \mathbb{N}}$ satisfies asymptotics (2.5) and $\{b_n\}_{n \in \mathbb{N}}$, $\{\tilde{b}_n\}_{n \in \mathbb{N}}$ satisfy asymptotics (2.6), the inequalities

$$|\tilde{b}_{l_m}(a_{k_n} - b_{l_m})(a_{k_n} - \tilde{b}_{l_m})|^{-1} \le |\tilde{b}_{k_n}(a_{k_n} - b_{k_n})(a_{k_n} - \tilde{b}_{k_n})|^{-1} \le 2/a_{k_n}^2$$

are valid for any $m \in \mathbb{N} \setminus \{s\}$ and for sufficiently large $n \in \mathbb{N}$. Recall that $\tilde{b}_{k_j} := b_{k_j}$ if $k_j \notin \{l_n\}_{n \in \mathbb{N}}$. In addition, $\sum_{n \in \mathbb{N}} A_{k_n} / a_{k_n}^2$ is absolutely convergent. Therefore right hand side of (4.22) converges to 0 as m goes to ∞ . Also note that by (4.21), left hand side of (4.22) is

$$\frac{-c(b_{l_m} + \tilde{b}_{l_m}) - d}{\tilde{b}_{l_m}} + \frac{G(\tilde{b}_{l_m})}{\tilde{b}_{l_m}(\tilde{b}_{l_m} - b_{l_m})} \\
= \frac{-c(b_{l_m} + \tilde{b}_{l_m}) - d}{\tilde{b}_{l_m}} + \frac{G(\tilde{b}_{l_m}) - \tilde{G}(\tilde{b}_{l_m})}{\tilde{b}_{l_m}(\tilde{b}_{l_m} - b_{l_m})} \\
= \frac{-c(b_{l_m} + \tilde{b}_{l_m}) - d}{\tilde{b}_{l_m}} + \frac{1}{\tilde{b}_{l_m} - b_{l_m}} \Big[(c - \tilde{c}) \tilde{b}_{l_m} + d - \tilde{d} + \frac{e - \tilde{e}}{\tilde{b}_{l_m}} \Big].$$

Let us observe that

$$\lim_{m \to \infty} \frac{-c(b_{l_m} + \tilde{b}_{l_m}) - d}{\tilde{b}_{l_m}} = -2c.$$

Now let us show $\tilde{b}_{l_m} - b_{l_m}$ converges to 0 as *m* goes to ∞ . Recall that poles of *G* and \tilde{G} satisfy asymptotics

$$k_n^2 + \frac{1}{\pi} \int_0^{\pi} q(x)dx + \alpha_{k_n}$$
 and $k_n^2 + \frac{1}{\pi} \int_0^{\pi} \tilde{q}(x)dx + \tilde{\alpha}_{k_n}$

respectively, where $\alpha_n = o(1)$ and $\tilde{\alpha}_n = o(1)$ as $n \to \infty$. Equivalance of poles of G and \tilde{G} imply equivalence of $\int_0^{\pi} q(x) dx$ and $\int_0^{\pi} \tilde{q}(x) dx$. Therefore b_{l_m} and \tilde{b}_{l_m}

satisfy asymptotics

$$\left(l_m - \frac{1}{2}\right)^2 + \frac{1}{\pi} \int_0^{\pi} q(x) dx + \beta_{l_m} \text{ and } \left(l_m - \frac{1}{2}\right)^2 + \frac{1}{\pi} \int_0^{\pi} q(x) dx + \tilde{\beta}_{l_m},$$

where $\beta_m = o(1)$ and $\tilde{\beta}_m = o(1)$ as $m \to \infty$. Hence $\tilde{b}_{lm} - b_{lm} = o(1)$ as m goes to ∞ . Therefore left hand side of (4.22) goes to ∞ if $c - \tilde{c} \neq 0$ or $d - \tilde{d} \neq 0$, so we get a contradiction unless $c = \tilde{c}$ and $d = \tilde{d}$. This implies that $G(z) - \tilde{G}(z)$ is a real constant. However, G(z) and $\tilde{G}(z)$ share the zero b_{ls} . This implies uniqueness of G(z) and hence uniqueness of $\{b_{l_n}\}_{n \in \mathbb{N}}$. After unique recovery of the two spectra σ_{DD} and σ_{ND} , the potential is uniquely determined by Borg's theorem. \Box

We also get the uniqueness result without knowing any point from $\{b_{l_n}\}_{n \in \mathbb{N}}$, but this requires absolute convergence of $\prod_{n \in \mathbb{N}} a_{k_n}/b_{l_n}$. By absolute convergence of $\prod_{n \in \mathbb{N}} a_{k_n}/b_{l_n}$ we mean absolute convergence of $\sum_{n \in \mathbb{N}} (a_{k_n}/b_{l_n} - 1)$. Note that Limit Comparison Test implies that $\prod_{n \in \mathbb{N}} a_{k_n}/b_{l_n}$ is absolutely convergent if and only if $\prod_{n \in \mathbb{N}} b_{l_n}/a_{k_n}$ is absolutely convergent. Absolute convergence of $\prod_{n \in \mathbb{N}} a_{k_n}/b_{l_n}$ also implies the two conditions in Lemma 4.5, so in this case Lemma 4.5 can be written in the following form.

Lemma 4.7. Let $\{a_{k_n}\}_{n \in \mathbb{N}} \subset \sigma_{\text{DD}}$ and $\{b_{l_n}\}_{n \in \mathbb{N}} \subset \sigma_{\text{ND}}$ such that $\prod_{n \in \mathbb{N}} (a_{k_n}/b_{l_n})$ is absolutely convergent. Then

$$G(z) = cz^{2} + dz + e + \sum_{n \in \mathbb{N}} A_{k_{n}} \Big(\frac{1}{z - a_{k_{n}}} + \frac{1}{a_{k_{n}}} \Big),$$

where c, d, e are real numbers, A_{k_n} is the residue of G(z) at the point $z = a_{k_n}$ and the sum converges normally on $\mathbb{C} \setminus \bigcup_{n \in \mathbb{N}} a_{k_n}$.

Proof. We will show that absolute convergence of $\prod_{n \in \mathbb{N}} (a_{k_n}/b_{l_n})$ implies the two conditions in Lemma 4.5, but first we begin by showing that absolute convergence of $\prod_{n \in \mathbb{N}} (a_{k_n}/b_{l_n})$ implies $\{1/(a_{k_n} - b_{l_n})\}_{n \in \mathbb{N}} \in l^1$. Since $\prod_{n \in \mathbb{N}} (a_{k_n}/b_{l_n})$ is absolutely convergent,

$$\sum_{n \in \mathbb{N}} \left| \frac{a_{k_n} - b_{l_n}}{b_{l_n}} \right| = \sum_{n \in \mathbb{N}} \left| \frac{k_n^2 - (l_n - 1/2)^2 + \alpha_{k_n} - \beta_{l_n}}{(l_n - 1/2)^2 + (1/\pi) \int_0^\pi q(x) dx + \beta_{l_n}} \right| < \infty,$$

i.e. $\{(k_n^2 - l_n^2 + l_n)/l_n^2\}_{n \in \mathbb{N}} \in l^1$. Note that $\lim_{n \to \infty} a_{k_n}/b_{l_n} = 1$ implies $\lim_{n \to \infty} k_n/l_n = 1$. Therefore

$$\infty > \sum_{n \in \mathbb{N}} \left| \frac{k_n^2 - l_n^2 + l_n}{l_n^2} \right|$$

= $\sum_{n \in \mathbb{N}} \frac{k_n + l_n}{l_n} \left| \frac{k_n - l_n + l_n / (k_n + l_n)}{l_n} \right|$
 $\ge \sum_{n=1}^N \left| \frac{k_n - l_n + l_n / (k_n + l_n)}{l_n} \right| + \sum_{n=N+1}^\infty \left| \frac{1/4}{l_n} \right|$
 $\ge c_1 \sum_{n \in \mathbb{N}} \frac{1}{l_n},$

where $N \in \mathbb{N}$ and $c_1 > 0$, i.e. $\{1/l_n\}_{n \in \mathbb{N}} \in l^1$ and by Limit Comparison Test $\{1/k_n\}_{n \in \mathbb{N}} \in l^1$. Therefore $\{1/(a_{k_n} - b_{l_n})\}_{n \in \mathbb{N}} \in l^1$, since $1/|a_{k_n} - b_{l_n}| \le 1/|a_{k_n} - b_{k_n}| = O(1/k_n)$ as n goes to ∞ .

The partial product G_N defined in the beginning of Section 4.2 can be represented as

$$G_N(z) = \sum_{n=1}^N \frac{A_{k_n,N}}{z - a_{k_n}} + \prod_{n=1}^N \frac{a_{k_n}}{b_{l_n}}$$

and hence

$$\lim_{N \to \infty} \sum_{n=1}^{N} \frac{A_{k_n,N}}{a_{k_n}} = \lim_{N \to \infty} \left[\prod_{n=1}^{N} \frac{a_{k_n}}{b_{l_n}} - G_N(0) \right] = \prod_{n \in \mathbb{N}} \frac{a_{k_n}}{b_{l_n}} - 1 \in \mathbb{R}.$$
 (4.23)

Since $\{1/a_{k_n}\}_{n \in \mathbb{N}} \in l^1$, existence of this limit implies $\lim_{N \to \infty} \sum_{n=1}^N |A_{k_n,N}/a_{k_n}^2|$ exists.

Now we are ready to prove the first assumption in Lemma 4.5, i.e.

$$\lim_{N \to \infty} \sum_{n=1}^{N} (|A_{k_{n,N}} - A_{k_n}|/a_{k_n}^2) < \infty.$$

For n < N, let us define

$$P_{k_n,N} := \prod_{m=N+1}^{\infty} \frac{a_{k_m}}{b_{l_m}} \frac{a_{k_n} - b_{l_m}}{a_{k_n} - a_{k_m}}.$$

Then

$$\frac{|A_{k_{n,N}} - A_{k_{n}}|}{a_{k_{n}}^{2}} = \left| \left(\frac{A_{k_{n,N}}}{a_{k_{n}}} \right) \left(\frac{a_{k_{n}} - b_{l_{n}}}{a_{k_{n}}} [1 - P_{k_{n,N}}] \right) \left(\frac{1}{a_{k_{n}} - b_{l_{n}}} \right) \right|.$$
(4.24)

Using (4.23), absoult convergence of $\prod_{n \in \mathbb{N}} (a_{k_n}/b_{l_n})$ and hence absolute convergence of $\sum_{n \in \mathbb{N}} [(a_{k_n} - b_{l_n})/a_{k_n}]$ we get that the limits

$$\lim_{N \to \infty} \sum_{n=1}^{N} \frac{A_{k_{n,N}}}{a_{k_n}} \quad \text{and} \quad \lim_{N \to \infty} \sum_{n=1}^{N} \left(\frac{a_{k_n} - b_{l_n}}{a_{k_n}} [1 - P_{k_n,N}] \right)$$

converge.

Recall that we have also showed $\{1/(a_{k_n} - b_{l_n})\}_{n \in \mathbb{N}} \in l^1$. Therefore by (4.24) we get the first assumption in Lemma 4.5,

$$\lim_{N \to \infty} \sum_{n=1}^{N} \frac{|A_{k_{n,N}} - A_{k_n}|}{a_{k_n}^2} < \infty.$$

After recalling that we showed existence of $\lim_{N\to\infty} \sum_{n=1}^{N} |A_{k_n,N}/a_{k_n}^2|$, we get the second assumption in Lemma 4.5, i.e. $\{A_{k_n}/a_{k_n}^2\}_{n\in\mathbb{N}} \in l^1$ as follows:

$$\lim_{N \to \infty} \sum_{n=1}^{N} \frac{|A_{k_n}|}{a_{k_n}^2} \le \lim_{N \to \infty} \sum_{n=1}^{N} \frac{|A_{k_n} - A_{k_n,N}|}{a_{k_n}^2} + \lim_{N \to \infty} \sum_{n=1}^{N} \frac{|A_{k_n,N}|}{a_{k_n}^2} < \infty.$$

Now using Lemma 4.5 we get the desired result.

Theorem 4.8. Let $q \in L^1(0, \pi)$ and $\{a_{k_n}\}_{n \in \mathbb{N}} \subset \sigma_{\text{DD}}, \{b_{l_n}\}_{n \in \mathbb{N}} \subset \sigma_{\text{ND}}$ such that $\prod_{n \in \mathbb{N}} (a_{k_n}/b_{l_n})$ is absolutely convergent. Then $\{a_n\}_{n \in \mathbb{N}}, \{b_n\}_{n \in \mathbb{N}} \setminus \{b_{l_n}\}_{n \in \mathbb{N}}$ and $\{\gamma_{k_n}\}_{n \in \mathbb{N}}$ determine the potential q, where $\sigma_{\text{DD}} = \{a_n\}_{n \in \mathbb{N}}, \sigma_{\text{ND}} = \{b_n\}_{n \in \mathbb{N}}$ are Dirichlet–Dirichlet and Neumann–Dirichlet spectra and $\{\gamma_n\}_{n \in \mathbb{N}}$ are point masses of the spectral measure $\mu_{0,0} = \sum_{n \in \mathbb{N}} \gamma_n \delta_{a_n}$.

Proof. One can use Lemma 4.7 and follow the proof of Theorem 4.6 until the last step, i.e. showing uniqueness of the two spectra after obtaining that $G(z) - \tilde{G}(z)$ is a real constant, so let us show $G(z) - \tilde{G}(z) = 0$. The main differences in this case are that G and \tilde{G} do not share any zero and the infinite products $\prod_{n \in \mathbb{N}} (a_{k_n}/b_{l_n})$ and $\prod_{n \in \mathbb{N}} (a_{k_n}/\tilde{b}_{l_n})$ are absolutely convergent. Let us recall that the infinite products G and \tilde{G} have the following representations:

$$G(z) = cz^{2} + dz - C + \sum_{n \in \mathbb{N}} A_{k_{n}} \left(\frac{1}{z - a_{k_{n}}} + \frac{1}{a_{k_{n}}} \right)$$
$$\tilde{G}(z) = cz^{2} + dz - \tilde{C} + \sum_{n \in \mathbb{N}} A_{k_{n}} \left(\frac{1}{z - a_{k_{n}}} + \frac{1}{a_{k_{n}}} \right)$$

Therefore by taking difference of G(z) and $\tilde{G}(z)$ we get $G(z) + C = \tilde{G}(z) + \tilde{C}$, i.e.

$$-C\prod_{n\in\mathbb{N}} \left(\frac{z}{b_{l_n}} - 1\right) \left(\frac{z}{a_{k_n}} - 1\right)^{-1} + C = -\widetilde{C}\prod_{n\in\mathbb{N}} \left(\frac{z}{\widetilde{b}_{l_n}} - 1\right) \left(\frac{z}{a_{k_n}} - 1\right)^{-1} + \widetilde{C}.$$
 (4.25)

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Note that since the infinite products $\prod_{n \in \mathbb{N}} (a_{k_n}/b_{l_n})$, $\prod_{n \in \mathbb{N}} (a_{k_n}/\tilde{b}_{l_n})$ are absolutely convergent and the two spectra $\{a_n\}_{n \in \mathbb{N}}$, $\{b_n\}_{n \in \mathbb{N}}$ lie on the positive real line, the infinite products on the two sides of (4.25) are uniformly convergent on the second quadrant. Hence by letting *z* go to infinity on the second quadrant we get

$$-C\prod_{n\in\mathbb{N}}\frac{a_{k_n}}{b_{l_n}}+C=-\widetilde{C}\prod_{n\in\mathbb{N}}\frac{a_{k_n}}{\widetilde{b}_{l_n}}+\widetilde{C}.$$
(4.26)

Recall that $\prod_{n \in \mathbb{N}} \tilde{b}_{l_n} / b_{l_n}$ is finite, since

$$\sum_{n\in\mathbb{N}}\frac{|\tilde{b}_{l_n}-b_{l_n}|}{b_{l_n}}=\sum_{n\in\mathbb{N}}\frac{|\tilde{\beta}_{l_n}-\beta_{l_n}|}{b_{l_n}}\leq \max_{n\in\mathbb{N}}|\tilde{\beta}_{l_n}-\beta_{l_n}|\sum_{n\in\mathbb{N}}\frac{1}{b_{l_n}}<\infty.$$

Therefore the infinite product $H(z) := G(z)/\widetilde{G}(z)$ is represented as

$$H(z) := \frac{G(z)}{\widetilde{G}(z)} = \frac{C}{\widetilde{C}} \prod_{n \in \mathbb{N}} \frac{z - b_{l_n}}{b_{l_n}} \frac{\widetilde{b}_{l_n}}{z - \widetilde{b}_{l_n}} = \frac{C}{\widetilde{C}} \prod_{n \in \mathbb{N}} \frac{\widetilde{b}_{l_n}}{b_{l_n}} \prod_{n \in \mathbb{N}} \frac{z - b_{l_n}}{z - \widetilde{b}_{l_n}}$$

We know that G and \tilde{G} share same poles with equivalent residues at the corresponding poles. Therefore for any $m \in \mathbb{N}$

$$1 = H(a_{k_m}) = \frac{C}{\widetilde{C}} \prod_{n \in \mathbb{N}} \frac{\widetilde{b}_{l_n}}{b_{l_n}} \prod_{n \in \mathbb{N}} \frac{a_{k_m} - b_{l_n}}{a_{k_m} - \widetilde{b}_{l_n}}.$$
(4.27)

Now let us find the limit of the infinite product on the right end of (4.27) as *m* goes to ∞ . This infinite product is uniformly convergent if and only if the infinite sum

$$\sum_{n \in \mathbb{N}} \left(\frac{a_{k_m} - b_{l_n}}{a_{k_m} - \tilde{b}_{l_n}} - 1 \right) = \sum_{n \in \mathbb{N}} \frac{\tilde{b}_{l_n} - b_{l_n}}{a_{k_m} - \tilde{b}_{l_n}}$$
(4.28)

is uniformly convergent. Note that asymptotics of the two spectra imply

$$\tilde{b}_{l_j} - b_{l_j} = o(1)$$

as j goes to infinity. Then the asymptotics of $\{a_{k_n}\}_{n \in \mathbb{N}}$, $\{b_{l_n}\}_{n \in \mathbb{N}}$ and $\{\tilde{b}_{l_n}\}_{n \in \mathbb{N}}$ together with absolute convergence of the infinite products $\prod_{n \in \mathbb{N}} (a_{k_n}/b_{l_n})$, $\prod_{n \in \mathbb{N}} (a_{k_n}/\tilde{b}_{l_n})$ imply that

$$\sum_{n\in\mathbb{N}} \left| \frac{b_{l_n} - b_{l_n}}{a_{k_m} - \tilde{b}_{l_n}} \right| \le \sum_{n\in\mathbb{N}} \left| \frac{b_{l_n} - b_{l_n}}{a_{k_n} - \tilde{b}_{l_n}} \right| < \infty, \tag{4.29}$$

since $\{1/(a_{k_n} - \tilde{b}_{l_n})\}_{n \in \mathbb{N}} \in l^1$ as we discussed in the proof of Lemma 4.7.

Therefore by letting *m* go to ∞ in (4.27) we get $\tilde{C}/C = \prod_{j \in \mathbb{N}} \tilde{b}_{l_j}/b_{l_j}$. If we define $\gamma := \prod_{n \in \mathbb{N}} a_{k_n}/b_{l_n}$ and $\tilde{\gamma} := \prod_{n \in \mathbb{N}} a_{k_n}/\tilde{b}_{l_n}$, we get $\tilde{C}/C = \gamma/\tilde{\gamma}$. This identity and (4.26) imply $\frac{\gamma - 1}{\tilde{\gamma} - 1} = \frac{\gamma}{\tilde{\gamma}}$ and hence $\gamma = \tilde{\gamma}$. Therefore $C = \tilde{C}$. This implies uniqueness of G(z) and hence uniqueness of $\{b_{l_n}\}_{n \in \mathbb{N}}$. After unique recovery of the two spectra σ_{DD} and σ_{ND} , the potential is uniquely determined by Borg's theorem.

4.3. General boundary conditions. As discussed in Section 2.2, the Weyl *m*-function for the Schrödinger equation

$$Lu = -u'' + qu = zu$$
 (4.30)

with boundary conditions

$$u(0)\cos\alpha - u'(0)\sin\alpha = 0,$$
 (4.31)

$$u(\pi)\cos\beta + u'(\pi)\sin\beta = 0, \qquad (4.32)$$

is defined as

$$m_{\alpha,\beta}(z) = \frac{\cos(\alpha)u_z'(0) + \sin(\alpha)u_z(0)}{-\sin(\alpha)u_z'(0) + \cos(\alpha)u_z(0)},$$

where $u_z(t)$ is a solution of (4.30) satisfying (4.32) and $\alpha, \beta \in [0, \pi)$. In order to prove our result with boundary conditions (4.31) and (4.32) we need to consider more general *m*-functions. Recall that we have defined the *m*-function in Section 2.2 by introducing two solutions $s_z(t)$ and $c_z(t)$ of (4.30) satisfying the initial conditions

$$s_z(0) = \sin(\alpha), \quad s'_z(0) = \cos(\alpha),$$
$$c_z(0) = \cos(\alpha), \quad c'_z(0) = -\sin(\alpha)$$

and $u_z(t)$, a solution of (4.30) with boundary conditions $u_z(\pi) = \sin \beta$, $u'_z(\pi) = -\cos \beta$. The same steps to define the *m*-function as in Section 2.2 can be followed if $c_z(t)$ is a linearly independent solution with $W(c_z, s_z) = 1$. Therefore we introduce two solutions $s_z(t)$ and $c_z(t)$ of (4.30) satisfying the initial conditions

$$s_z(0) = \sin(\alpha_2), \qquad s'_z(0) = \cos(\alpha_2),$$

$$c_z(0) = \frac{\sin(\alpha_1)}{\sin(\alpha_1 - \alpha_2)}, \quad c'_z(0) = \frac{\cos(\alpha_1)}{\sin(\alpha_1 - \alpha_2)}$$

for $\alpha_1, \alpha_2 \in [0, \pi)$, $\sin(\alpha_1 - \alpha_2) \neq 0$ and same $u_z(t)$. Then we can define the *m*-function $m_{\alpha_1,\alpha_2,\beta}$.

Definition 4.9. The *m*-function $m_{\alpha_1,\alpha_2,\beta}$ is defined as

$$m_{\alpha_1,\alpha_2,\beta}(z) := \frac{1}{\sin(\alpha_2 - \alpha_1)} \left[\frac{-\sin(\alpha_1)u_z'(0) + \cos(\alpha_1)u_z(0)}{-\sin(\alpha_2)u_z'(0) + \cos(\alpha_2)u_z(0)} \right]$$

where $\alpha_1, \alpha_2, \beta \in [0, \pi)$, $\sin(\alpha_2 - \alpha_1) \neq 0$ and $u_z(t)$ is a solution of (4.30) with boundary conditions $u_z(\pi) = \sin \beta$, $u'_z(\pi) = -\cos \beta$.

Remark 4.10. The *m*-function $m_{\alpha,\beta}$ we discussed in Section 2.2 is obtained by letting $\alpha_1 = \alpha - \pi/2$ and $\alpha_2 = \alpha$, i.e. $m_{\alpha - \frac{\pi}{2}, \alpha, \beta}(z) = m_{\alpha,\beta}(z)$.

The *m*-function $m_{\alpha_1,\alpha_2,\beta}(z)$ is a meromorphic Herglotz function having real zeros on $\sigma_{\alpha_1,\beta}$ and real poles on $\sigma_{\alpha_2,\beta}$, which are interlacing. It is a meromorphic Herglotz function, since $m_{0,\beta}(z) = u'_z(0)/u_z(0)$ is a meromorphic Herglotz function and sgn[Im $(m_{\alpha_1,\alpha_2,\beta}(z))$] = sgn[Im $(m_{0,\beta}(z))$]. Therefore Herglotz representation theorem implies

$$m_{\alpha_1,\alpha_2,\beta}(z) = az + b + \int \left[\frac{1}{t-z} - \frac{t}{1+t^2}\right] d\mu_{\alpha_1,\alpha_2,\beta}(t),$$

where $a, b \in \mathbb{R}$ and $\mu_{\alpha_1,\alpha_2,\beta}$ is a positive discrete Poisson-summable measure supported on the spectrum $\sigma_{\alpha_2,\beta}$. Let us call $\mu_{\alpha_1,\alpha_2,\beta}$ the spectral measure corresponding to $(\alpha_1, \alpha_2, \beta)$. Now we prove our results with general boundary conditions.

Theorem 4.11. Let $q \in L^1(0, \pi)$, $A \subset \mathbb{N}$, $\sin(\alpha_2 - \alpha_1) \neq 0$ and $\alpha_1, \alpha_2, \beta \in [0, \pi)$. Then $\{a_n\}_{n \in \mathbb{N}}$, $\{b_n\}_{n \in \mathbb{N} \setminus A}$ and $\{\gamma_n\}_{n \in A}$ determine the potential q, where $\sigma_{\alpha_2,\beta} = \{a_n\}_{n \in \mathbb{N}}$, $\sigma_{\alpha_1,\beta} = \{b_n\}_{n \in \mathbb{N}}$ are two spectra and $\{\gamma_n\}_{n \in \mathbb{N}}$ are point masses of the spectral measure $\mu_{\alpha_1,\alpha_2,\beta} = \sum_{n \in \mathbb{N}} \gamma_n \delta_{a_n}$.

Proof. Wlog let a_n and b_n be positive for all $n \in \mathbb{N}$. We follow the arguments we used in the proofs of Lemma 4.1 and Theorem 4.2, but there are two differences: asymptotics of the two spectra, depending on $\alpha_1, \alpha_2, \beta$ and hence the order relation betweeen a_n and b_n . Thus, we consider the following cases.

(i) $\alpha_1 \neq 0, \alpha_2 \neq 0, \alpha_1 > \alpha_2$. When $\beta \neq 0$, the two spectra $\sigma_{\alpha_2,\beta} = \{a_n\}_{n \in \mathbb{N}}$ and $\sigma_{\alpha_1,\beta} = \{b_n\}_{n \in \mathbb{N}}$ satisfy the asymptotics (2.4) and hence $a_n > b_n$ for all $n \in \mathbb{N}$. Therefore using the proof of Lemma 4.1, $m_{\alpha_1,\alpha_2,\beta}(z)$ can be represented as (4.2). Using this representation and Čebotarev's theorem as we discussed in the proof of Theorem 4.2, the meromorphic Herglotz function G(z) defined as

$$G(z) := -C \prod_{n \in A} \left(\frac{z}{b_n} - 1\right) \left(\frac{z}{a_n} - 1\right)^{-1}$$
(4.33)

has the following representation:

$$G(z) = dz + e + \sum_{n \in A} A_n \left(\frac{1}{z - a_n} + \frac{1}{a_n}\right).$$

Only unknown constants on the right hand side are d and e. In order to show uniqueness of the linear term dz + e, let us introduce $\tilde{G}(z)$ as we did in the proof of Theorem 4.2:

• the infinite product \tilde{G} is defined as

$$\widetilde{G}(z) := -\widetilde{C} \prod_{n \in A} \left(\frac{z}{\widetilde{b}_n} - 1 \right) \left(\frac{z}{\widetilde{a}_n} - 1 \right)^{-1},$$

where $\tilde{C} > 0$, the set of poles $\{\tilde{a}_n\}_{n \in A}$ and the set of zeros $\{\tilde{b}_n\}_{n \in A}$ satisfy asymptotics (2.4). Let $\tilde{a}_k := a_k$ and $\tilde{b}_k := b_k$ for $k \in \mathbb{N} \setminus A$;

- G and G̃ share same set of poles with equivalent residues at the corresponding poles, i.e. ã_k = a_k and Res(G̃, a_k) = Res(G, a_k) for any k ∈ A;
- by the equivalence of poles and residues of G and G and Cebotarev's theorem, G(z) has the representation

$$\widetilde{G}(z) = \widetilde{d}z + \widetilde{e} + \sum_{n \in A} A_n \Big(\frac{1}{z - a_n} + \frac{1}{a_n} \Big),$$

where $\tilde{d} \ge 0$, $\tilde{e} \in \mathbb{R}$.

Therefore difference of G and \tilde{G} is a linear polynomial, i.e.

$$G(z) - \tilde{G}(z) = (d - \tilde{d})z + e - \tilde{e}$$
(4.34)

Note that since $\{a_n\}_{n \in \mathbb{N}}$, $\{b_n\}_{n \in \mathbb{N}}$ and $\{\tilde{b}_n\}_{n \in \mathbb{N}}$ are subsets of $(0, \infty)$ and satisfy asymptotics (2.4), for any $x \in (-\infty, 0)$ we get

$$|G(x) - \widetilde{G}(x)| \le \left| C \prod_{n \in A} \left(\frac{x}{b_n} - 1 \right) \left(\frac{x}{a_n} - 1 \right)^{-1} \right| + \left| \widetilde{C} \prod_{n \in A} \left(\frac{x}{\widetilde{b}_n} - 1 \right) \left(\frac{x}{a_n} - 1 \right)^{-1} \right|$$
$$\le C \prod_{n \in A} \frac{a_n}{b_n} + \widetilde{C} \prod_{n \in A} \frac{a_n}{\widetilde{b}_n} < \infty.$$

Convergence of the infinite product $\prod_{n \in A} a_n / b_n$ follows from the fact that

$$\sum_{n\in A}\frac{a_n-b_n}{b_n}\leq M\sum_{n\in A}\frac{1}{n^2},$$

for some $M < \infty$, since asymptotics (2.4) imply $|a_n - b_n| \le M_1$ for some $M_1 < \infty$ independent of *n* and $a_n = n^2 + o(n^2)$, $b_n = n^2 + o(n^2)$ as *n* goes to infinity. Therefore

$$\lim_{x \to -\infty} |(d - \tilde{d})x + e - \tilde{e}| = \lim_{x \to -\infty} |G(x) - \tilde{G}(x)| = \left| C \prod_{n \in A} \frac{a_n}{b_n} - \tilde{C} \prod_{n \in A} \frac{a_n}{\tilde{b}_n} \right| < \infty,$$

so we get a contradiction unless $d = \tilde{d}$. This implies that $G(z) - \tilde{G}(z)$ is a real constant, which is $G(0) - \tilde{G}(0) = \tilde{C} - C$. In order to show $\tilde{C} = C$, we follow exactly the same arguments used in the proof of Theorem 4.2.

This gives uniqueness of G(z) and hence uniqueness of $\{b_n\}_{n \in A}$. After unique recovery of the two spectra $\sigma_{\alpha_2,\beta}$ and $\sigma_{\alpha_1,\beta}$, Levinson's theorem uniquely determines the potential.

When $\beta = 0$, one can apply same arguments. The only difference appears in asymptotics of $\sigma_{\alpha_2,\beta} = \{a_n\}_{n \in \mathbb{N}}$ and $\sigma_{\alpha_1,\beta} = \{b_n\}_{n \in \mathbb{N}}$, which does not affect the result.

(ii) $\alpha_1 \neq 0$, $\alpha_2 = 0$, $\beta = 0$. The two spectra $\sigma_{\alpha_2,\beta} = \{a_n\}_{n \in \mathbb{N}}$ and $\sigma_{\alpha_1,\beta} = \{b_n\}_{n \in \mathbb{N}}$ satisfy the asymptotics (2.5) and (2.6) respectively. One then obtains the result by following the proofs of Lemma 4.1 and Theorem 4.2.

(iii) $\alpha_1 \neq 0, \alpha_2 = 0, \beta \neq 0$. The two spectra $\sigma_{\alpha_2,\beta} = \{a_n\}_{n \in \mathbb{N}}$ and $\sigma_{\alpha_1,\beta} = \{b_n\}_{n \in \mathbb{N}}$ satisfy the asymptotics (2.7) and (2.4) respectively, which is similar to the previous case.

(iv) $\alpha_1 \neq 0, \alpha_2 \neq 0, \alpha_1 < \alpha_2$ or $\alpha_1 = 0, \alpha_2 \neq 0, \beta \neq 0$ or $\alpha_1 = 0, \alpha_2 \neq 0, \beta = 0$. In all of these three cases, $a_n < b_n$ for all $n \in \mathbb{N}$. Therefore using the proof of Lemma 4.1, $m_{\alpha_1,\alpha_2,\beta}(z)$ can be represented as

$$m_{\alpha_1,\alpha_2,\beta}(z) = C \prod_{n \in \mathbb{N}} \left(\frac{z}{b_n} - 1\right) \left(\frac{z}{a_n} - 1\right)^{-1}.$$

In order to represent G(z) as (4.33), an extra factor is required, so we shift indices of b_n up by one inside A and let b_1 be a positive real number less than a_1 , assuming wlog $1 \in A$. Then $\frac{z-b_1}{b_1}G(z)$ can be represented as (4.33). Using this representation and Čebotarev's theorem, the meromorphic Herglotz function $\frac{z-b_1}{b_1}G(z)$ has the following representation:

$$\left(\frac{z-b_1}{b_1}\right)G(z) = az+b+\sum_{n\in A}A_n\left(\frac{1}{a_n-z}-\frac{1}{a_n}\right).$$

Therefore if we introduce $\tilde{G}(z)$ similar to the previous cases, then $\frac{z-b_1}{b_1}G(z)$ and $\frac{z-b_1}{b_1}\tilde{G}(z)$ share the same set of poles $\{a_n\}_{n\in A}$ with the same residues $\{-A_n\}_{n\in A}$ and have the sets of zeros $\{b_n\}_{n\in A}$ and $\{b_1\} \cup \{\tilde{b}_n\}_{n\in A\setminus\{1\}}$ respectively, so the difference of $\frac{z-b_1}{b_1}G(z)$ and $\frac{z-b_1}{b_1}\tilde{G}(z)$ is a linear polynomial with real coefficients and hence $G(z) - \tilde{G}(z)$ is a real constant, which is $G(0) - \tilde{G}(0) = \tilde{C} - C$. In order to show $\tilde{C} = C$, we follow exactly the same arguments used in the proof of Theorem 4.2.

This implies uniqueness of G(z) and hence uniqueness of $\{b_n\}_{n \in A}$. After unique recovery of the two spectra $\sigma_{\alpha_2,\beta}$ and $\sigma_{\alpha_1,\beta}$, Levinson's theorem uniquely determines the potential.

Remark 4.12. Theorem 4.11 gives Marchenko's theorem with the *m*-function $m_{\alpha_1,\alpha_2,\beta}$ as a corollary if we let $A = \mathbb{N}$. By letting $A = \emptyset$, we get the statement of Levinson's theorem.

For the non-matching index sets case, let us recall the definitions of $A_{k_n,m}$ and A_{k_n} :

$$A_{k_n,m} := \frac{a_{k_n}}{b_{l_n}} (a_{k_n} - b_{k_n}) \prod_{\substack{j=1\\j \neq n}}^m \frac{a_{k_j}}{b_{l_j}} \frac{a_{k_n} - b_{l_j}}{a_{k_n} - a_{k_j}},$$
$$A_{k_n} := \frac{a_{k_n}}{b_{l_n}} (a_{k_n} - b_{k_n}) \prod_{\substack{j=1\\j \neq n}}^\infty \frac{a_{k_j}}{b_{l_j}} \frac{a_{k_n} - b_{l_j}}{a_{k_n} - a_{k_j}}.$$

We can prove Theorem 4.6 and Theorem 4.8 with general boundary conditions following the same proofs. However, if boundary conditions α_1 and α_2 are nonzero, then we need that eventually the two index sets $\{k_n\}_{n \in \mathbb{N}}$ and $\{l_n\}_{n \in \mathbb{N}}$ have no common element.

Theorem 4.13. Let $q \in L^1(0, \pi)$, $\sin(\alpha_2 - \alpha_1) \neq 0$, $\alpha_1, \alpha_2, \beta \in [0, \pi)$ and $\{a_{k_n}\}_{n \in \mathbb{N}} \subset \sigma_{\alpha_2,\beta}$, $\{b_{l_n}\}_{n \in \mathbb{N}} \subset \sigma_{\alpha_1,\beta}$ satisfy following properties:

- $\lim_{m\to\infty} \sum_{n=1}^{m} (|A_{k_n,m} A_{k_n}|/a_{k_n}^2) < \infty;$
- $\{A_{k_n}/a_{k_n}^2\}_{n\in\mathbb{N}}\in l^1$.

(i) If $\alpha_1 = 0$ or $\alpha_2 = 0$, then $\{a_n\}_{n \in \mathbb{N}}$, $\{b_n\}_{n \in \mathbb{N}} \setminus \{b_{l_n}\}_{n \in \mathbb{N} \setminus \{s\}}$ and $\{\gamma_{k_n}\}_{n \in \mathbb{N}}$ determine the potential q for any $s \in \mathbb{N}$, where $\sigma_{\alpha_2,\beta} = \{a_n\}_{n \in \mathbb{N}}$, $\sigma_{\alpha_1,\beta} = \{b_n\}_{n \in \mathbb{N}}$ are two spectra and $\{\gamma_n\}_{n \in \mathbb{N}}$ are point masses of the spectral measure $\mu_{\alpha_1,\alpha_2,\beta} = \sum_{n \in \mathbb{N}} \gamma_n \delta_{a_n}$. (ii) If $\alpha_1 \neq 0$, $\alpha_2 \neq 0$ and there exists $N \in \mathbb{N}$ such that $k_n \neq l_n$ for all n > N, then $\{a_n\}_{n \in \mathbb{N}}$, $\{b_n\}_{n \in \mathbb{N}} \setminus \{b_{l_n}\}_{n \in \mathbb{N} \setminus \{s\}}$ and $\{\gamma_{k_n}\}_{n \in \mathbb{N}}$ determine the potential q for any $s \in \mathbb{N}$, where $\sigma_{\alpha_2,\beta} = \{a_n\}_{n \in \mathbb{N}}$, $\sigma_{\alpha_1,\beta} = \{b_n\}_{n \in \mathbb{N}}$ are two spectra and $\{\gamma_n\}_{n \in \mathbb{N}}$ are point masses of the spectral measure $\mu_{\alpha_1,\alpha_2,\beta} = \sum_{n \in \mathbb{N}} \gamma_n \delta_{a_n}$.

Proof. In the proof of Theorem 4.6 we used the inequalities (4.11), namely

$$|\tilde{b}_{l_m}(a_{k_n} - b_{l_m})(a_{k_n} - \tilde{b}_{l_m})|^{-1} \le |\tilde{b}_{k_n}(a_{k_n} - b_{k_n})(a_{k_n} - \tilde{b}_{k_n})|^{-1} \le 2/a_{k_n}^2.$$

If $\alpha_1 = 0$ or $\alpha_2 = 0$, these inequalities are still valid for any $m \in \mathbb{N} \setminus \{s\}$ and for sufficiently large $n \in \mathbb{N}$. Recall that $\tilde{b}_{k_j} := b_{k_j}$ if $k_j \notin \{l_n\}_{n \in \mathbb{N}}$.

If $\alpha_1 \neq 0$, $\alpha_2 \neq 0$ and there exists $N \in \mathbb{N}$ such that $k_n \neq l_n$ for all n > N, we modify these inequalities as follows:

$$|\tilde{b}_{l_m}(a_{k_n} - b_{l_m})(a_{k_n} - \tilde{b}_{l_m})|^{-1} \le |\tilde{b}_{k_n+1}(a_{k_n} - b_{k_n+1})(a_{k_n} - \tilde{b}_{k_{n+1}})|^{-1} \le 2/a_{k_n}^2,$$

which are valid for any $m \in \mathbb{N} \setminus \{s\}$ and for sufficiently large $n \in \mathbb{N}$.

After getting these inequalities we apply proofs of Lemma 4.5 and Theorem 4.6 with the *m*-function $m_{\alpha_1,\alpha_2,\beta}$ and the spectral measure $\mu_{\alpha_1,\alpha_2,\beta}$ and obtain uniqueness of $\{b_{l_n}\}_{n \in \mathbb{N}}$. Even though asymptotics of the spectra may be different than Dirichlet–Dirichlet, Neumann–Dirichlet case, the same arguments can be used. After unique recovery of the two spectra $\sigma_{\alpha_2,\beta}$ and $\sigma_{\alpha_1,\beta}$, Levinson's theorem uniquely determines the potential.

Theorem 4.14. Let $q \in L^1(0, \pi)$, $\sin(\alpha_2 - \alpha_1) \neq 0$, $\alpha_1, \alpha_2, \beta \in [0, \pi)$ and $\prod_{n \in \mathbb{N}} a_{k_n}/b_{l_n}$ be absolutely convergent, where $\{a_{k_n}\}_{n \in \mathbb{N}} \subset \sigma_{\alpha_2,\beta}$, $\{b_{l_n}\}_{n \in \mathbb{N}} \subset \sigma_{\alpha_1,\beta}$.

(i) If $\alpha_1 = 0$ or $\alpha_2 = 0$, then $\{a_n\}_{n \in \mathbb{N}}$, $\{b_n\}_{n \in \mathbb{N}} \setminus \{b_{l_n}\}_{n \in \mathbb{N}}$ and $\{\gamma_{k_n}\}_{n \in \mathbb{N}}$ determine the potential q, where $\sigma_{\alpha_2,\beta} = \{a_n\}_{n \in \mathbb{N}}$, $\sigma_{\alpha_1,\beta} = \{b_n\}_{n \in \mathbb{N}}$ are two spectra and $\{\gamma_n\}_{n \in \mathbb{N}}$ are point masses of the spectral measure $\mu_{\alpha_1,\alpha_2,\beta} = \sum_{n \in \mathbb{N}} \gamma_n \delta_{a_n}$.

(ii) If $\alpha_1 \neq 0$, $\alpha_2 \neq 0$ and there exists $N \in \mathbb{N}$ such that $k_n \neq l_n$ for all n > N, then $\{a_n\}_{n \in \mathbb{N}}, \{b_n\}_{n \in \mathbb{N}} \setminus \{b_{l_n}\}_{n \in \mathbb{N}}$ and $\{\gamma_{k_n}\}_{n \in \mathbb{N}}$ determine the potential q, where $\sigma_{\alpha_2,\beta} = \{a_n\}_{n \in \mathbb{N}}, \sigma_{\alpha_1,\beta} = \{b_n\}_{n \in \mathbb{N}}$ are two spectra and $\{\gamma_n\}_{n \in \mathbb{N}}$ are point masses of the spectral measure $\mu_{\alpha_1,\alpha_2,\beta} = \sum_{n \in \mathbb{N}} \gamma_n \delta_{a_n}$.

Proof. If $\alpha_1 = 0$ or $\alpha_2 = 0$, we follow the proofs of Lemma 4.7 and Theorem 4.8 with the *m*-function $m_{\alpha_1,\alpha_2,\beta}$ and the spectral measure $\mu_{\alpha_1,\alpha_2,\beta}$ and obtain uniqueness of $\{b_{l_n}\}_{n \in \mathbb{N}}$. After unique recovery of the two spectra $\sigma_{\alpha_2,\beta}$ and $\sigma_{\alpha_1,\beta}$, Levinson's theorem uniquely determines the potential.

If $\alpha_1 \neq 0$, $\alpha_2 \neq 0$ and there exists $N \in \mathbb{N}$ such that $k_n \neq l_n$ for all n > N, then the only difference appears in showing $\{1/(a_{k_n} - b_{l_n})_{n \in \mathbb{N}} \in l^1$, so let us show that absolute convergence of $\prod_{n \in \mathbb{N}} (a_{k_n}/b_{l_n})$ implies $\{1/(a_{k_n} - b_{l_n})\}_{n \in \mathbb{N}} \in l^1$. Since $\prod_{n \in \mathbb{N}} (a_{k_n}/b_{l_n})$ is absolutely convergent,

$$\sum_{n \in \mathbb{N}} \left| \frac{a_{k_n} - b_{l_n}}{b_{l_n}} \right| = \sum_{n \in \mathbb{N}} \left| \frac{(k_n - 1)^2 - (l_n - 1)^2 + \gamma_1 + \alpha_{k_n} - \beta_{l_n}}{(l_n - 1)^2 + \gamma_2 + (2/\pi) \int_0^{\pi} q(x) dx + \beta_{l_n}} \right| < \infty,$$

i.e. $\{(k_n^2 - l_n^2 - 2k_n + 2l_n)/l_n^2\}_{n \in \mathbb{N}} \in l^1$. Here $\gamma_1 = 2[\cot(\alpha_2) - \cot(\alpha_1)]/\pi$, $\gamma_2 = 2[\cot(\beta) + \cot(\alpha_1)]/\pi$ and wlog we assume $\beta \neq 0$. Note that $\lim_{n \to \infty} a_{k_n}/b_{l_n} = 1$ implies $\lim_{n \to \infty} k_n/l_n = 1$. Therefore

$$\infty > \sum_{n \in \mathbb{N}} \left| \frac{k_n^2 - l_n^2 - 2(k_n - l_n)}{l_n^2} \right|$$
$$= \sum_{n \in \mathbb{N}} \frac{k_n + l_n - 2}{l_n} \left| \frac{k_n - l_n}{l_n} \right|$$
$$\ge \sum_{n=1}^N \left| \frac{k_n - l_n}{l_n} \right| + \sum_{n=N+1}^\infty \frac{1}{l_n}$$
$$\ge c_1 \sum_{n \in \mathbb{N}} \frac{1}{l_n}$$

where $N \in \mathbb{N}$ and $c_1 > 0$, so $\{1/l_n\}_{n \in \mathbb{N}} \in l^1$ and hence by Limit Comparison Test $\{1/k_n\}_{n \in \mathbb{N}} \in l^1$. Therefore $\{1/(a_{k_n} - b_{l_n})\}_{n \in \mathbb{N}} \in l^1$, since for n > N, $1/|a_{k_n} - b_{l_n}| \le 1/|a_{k_n} - b_{k_n+1}| = O(1/k_n)$ as n goes to ∞ . Now we apply proofs of Lemma 4.7 and Theorem 4.8 with the m-function $m_{\alpha_1,\alpha_2,\beta}$ and the spectral measure $\mu_{\alpha_1,\alpha_2,\beta}$ and obtain uniqueness of $\{b_{l_n}\}_{n \in \mathbb{N}}$. After unique recovery of the two spectra $\sigma_{\alpha_2,\beta}$ and $\sigma_{\alpha_1,\beta}$, Levinson's theorem uniquely determines the potential.

Appendix A. Complex function theoretical tools

In this section we recall some definitions and theorems from complex function theory used in our discussions. We follow [40].

A function on \mathbb{R} is Poisson-summable if it is summable with respect to the Poisson measure Π , defined as $d\Pi := dx/(1 + x^2)$. The space of Poisson-summable functions on \mathbb{R} is denoted by L^1_{Π} .

The Schwarz integral of a Poisson-summable function f is

$$Sf(z) = \frac{1}{i\pi} \int_{\mathbb{R}} \left(\frac{1}{t-z} - \frac{t}{1+t^2} \right) f(t) dt.$$

The Schwarz integral of a real valued Poisson-summable function is given in terms of its Poisson and conjugate Poisson integrals:

$$Sf = Pf + iQf,$$

$$Pf(z) = \frac{1}{\pi} \int \frac{y}{(t-x)^2 + y^2} f(t)dt,$$

$$Qf(z) = \frac{1}{\pi} \int \left(\frac{x-t}{(x-t)^2 + y^2} + \frac{1}{1+t^2}\right) f(t)dt$$

A measure μ on \mathbb{R} is Poisson-finite if $\int \frac{1}{1+t^2} d|\mu|(t) < \infty$. The Schwarz integral of a Poisson-finite measure μ , defined as

$$S\mu(z) = \frac{1}{i\pi} \int_{\mathbb{R}} \left(\frac{1}{t-z} - \frac{t}{1+t^2} \right) d\mu(t),$$

is analytic in the upper half-plane \mathbb{C}_+ .

Outer functions in \mathbb{C}_+ are analytic functions of the form e^{Sf} for $f \in L^1_{\Pi}$.

Inner functions in \mathbb{C}_+ are bounded analytic functions with non-tangential boundary values, equal to 1 in modulus, almost everywhere on \mathbb{R} . If an inner function extends to \mathbb{C} meromorphically, it is called *meromorphic inner function*, usually denoted by Θ .

Hilbert transform of $f \in L^1_{\Pi}$, denoted by \tilde{f} , is defined as the singular integral

$$\tilde{f}(x) = \frac{1}{\pi} \text{ p.v.} \int \left[\frac{1}{x-t} + \frac{t}{1+t^2}\right] f(t) dt,$$

where p.v. denotes the Cauchy principal value. It is the angular limit of Qf = Im Sf, hence the outer function e^{Sf} coincides with $e^{f+i\tilde{f}}$ on \mathbb{R} .

A meromorphic function is said *real* if it maps real numbers to real numbers on its domain. A meromorphic Herglotz function *m* is a real meromorphic function with positive imaginary part on \mathbb{C}_+ . It has negative imaginary part on \mathbb{C}_- via the relation $m(\bar{z}) = \overline{m(z)}$.

There is a one-to-one correspondence between meromorphic inner functions and meromorphic Herglotz functions via equations

$$m = i \frac{1+\Theta}{1-\Theta}, \quad \Theta = \frac{m-i}{m+i}.$$

A meromorphic Herglotz function can be described as the Schwarz integral of a positive discrete Poisson-finite measure:

$$m(z) = az + b + iS\mu,$$

where $a \ge 0, b \in \mathbb{R}$. The term *iS* is also called *Herglotz integral* and usually denoted by *H*. This representation is valid even if the Herglotz function can not be extended meromorphically to \mathbb{C} , in which case μ may not be discrete. It is called *Herglotz representation theorem*. Čebotarev proved a similar result.

Theorem A.1 (Čebotarev [33]). If the real meromorphic function m maps \mathbb{C}_+ onto \mathbb{C}_+ , then its poles $\{a_k\}_{k \in \mathbb{Z}}$ are all real and simple, and it may be represented in the form

$$m(z) = az + b + \sum_{k=N}^{M} A_k \Big(\frac{1}{a_k - z} - \frac{1}{a_k} \Big),$$

where $a \ge 0$, $b \in \mathbb{R}$, $-\infty \le N < M \le \infty$, $A_k \ge 0$ and the sum $\sum_{k=N}^{M} A_k/a_k^2$ converges.

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References

- V. Ambarzumian, Über eine Frage der Eigenwerttheorie. Z. f. Physik 53 (1929), 690–695. JFM 55.0868.01
- [2] L. Amour and J. Faupin, Inverse spectral results for the Schrödinger operator in Sobolev spaces. *Int. Math. Res. Not. IMRN* 2010, no. 22, 4319–4333. MR 2737773 Zbl 1214.34013
- [3] L. Amour, J. Faupin, and T. Raoux, Inverse spectral results for Schrödinger operator on the unit interval with partial information given on the potentials. *J. Math. Phys.* 50 (2009), no. 3, article id. 033505, 14 pp. MR 2510910
- [4] L. Amour and T. Raoux, Inverse spectral results for Schrödinger operators on the unit interval with potentials in *l^p* spaces. *Inverse Problems* 23 (2007), no. 6, 2367–2373. MR 2441008 Zbl 1136.34008
- [5] A. Baranov, Y. Belov, and A. Poltoratski, De Branges functions of Schroedinger equations. *Collect. Math.* 68 (2017), no. 2, 251–263. MR 3633061 Zbl 06748545
- [6] A. Beurling and P. Malliavin, On Fourier transforms of measures with compact support. Acta Math. 107 (1962), 291–309. MR 0147848 Zbl 0127.32601
- [7] A. Beurling and P. Malliavin, On the closure of characters and the zeros of entire functions. Acta Math. 118 (1967), 79–93. MR 0209758 Zbl 0171.11901
- [8] N. P. Bondarenko, A partial inverse problem for the Sturm–Liouville operator on a star-shaped graph. *Anal. Math. Phys.* 8 (2018), no. 1, 155–168. MR 3766644
 Zbl 1391.34037

- [9] N. P. Bondarenko, Partial inverse problems for the Sturm–Liouville operator on a starshaped graph with mixed boundary conditions. *J. Inverse Ill-Posed Probl.* 26 (2018), no. 1, 1–12. MR 3757481 Zbl 1407.34043
- [10] N. P. Bondarenko and C. T. Shieh, Partial inverse problems for Sturm-Liouville operators on trees. *Proc. Roy. Soc. Edinburgh Sect.* A 147 (2017), no. 5, 917–933. MR 3705313 Zbl 1386.34027
- [11] N. P. Bondarenko and C. F. Yang, Partial inverse problems for the Sturm–Liouville operator on a star-shaped graph with different edge lengths. *Results Math.* **73** (2018), no. 2, Paper No. 56, 17 pp. MR 3778931 Zbl 1396.34005
- [12] G. Borg, Eine Umkehrung der Sturm-Liouvilleschen Eigenwertaufgabe. Acta Math. 78 (1946), 1–96. MR 0015185 Zbl 0063.00523
- [13] R. del Rio, F. Gesztesy, and B. Simon, Inverse spectral analysis with partial information on the potential. III. Updating boundary conditions. *Internat. Math. Res. Notices* 1997, no. 15, 751–758. MR 1470376 Zbl 0898.34075
- [14] J. Eckhardt, Direct and inverse spectral theory of singular left-definite Sturm-Liouville operators. J. Differential Equations 253 (2012), no. 2, 604–634. MR 2921208 Zbl 1250.34021
- [15] J. Eckhardt, Inverse uniqueness results for Schrödinger operators using de Branges theory. *Complex Anal. Oper. Theory* 8 (2014), no. 1, 37–50. MR 3147711 Zbl 1307.34029
- [16] J. Eckhardt, F. Gesztesy, N. Roger, and G. Teschl, Inverse spectral theory for Sturm– Liouville operators with distributional potentials. *J. Lond. Math. Soc.* (2) 88 (2013), no. 3, 801–828. MR 3145132 Zbl 1302.34127
- [17] J. Eckhardt and G. Teschl, Uniqueness results for one-dimensional Schrödinger operators with purely discrete spectra. *Trans. Amer. Math. Soc.* 365 (2013), no. 7, 3923–3942. MR 3042609 Zbl 1286.34022
- [18] W. N. Everitt, On a property of the *m*-coefficient of a second order linear differential equation. J. Lond. Math. Soc. (2) 4 (1972), 443–457. Zbl 0262.34012
- [19] F. Gesztesy, Inverse spectral theory as influenced by Barry Simon. In F. Gesztesy, P. Deift, Ch. Galvez, P. Perry and W. Schlag (eds.), *Spectral theory and mathematical physics: a Festschrift in honor of Barry Simon's 60th birthday*. Ergodic Schrödinger operators, singular spectrum, orthogonal polynomials, and inverse spectral theory. Papers from the conference held at the California Institute of Technology, Pasadena, CA, March 27–31, 2006. Proceedings of Symposia in Pure Mathematics, 76, Part 2. American Mathematical Society, Providence, R.I., 2007, 741–820. MR 2307754 Zbl 1135.34006
- [20] F. Gesztesy and B. Simon, A new approach to inverse spectral theory. II. General real potentials and the connection to the spectral measure. *Ann. of Math.* (2) **152** (2000), no. 2, 593–643. MR 1804532 Zbl 0983.34013

- [21] F. Gesztesy and B. Simon, Inverse spectral analysis with partial information on the potential. II. The case of discrete spectrum. *Trans. Amer. Math. Soc.* 352 (2000), no. 6, 2765–2787. MR 1694291 Zbl 0948.34060
- [22] F. Gesztesy and B. Simon, Uniqueness theorems in inverse spectral theory for onedimensional Schrödinger operators. *Trans. Amer. Math. Soc.* 348 (1996), no. 1, 349–373. MR 1329533 Zbl 0846.34090
- [23] N. J. Guliyev, Essentially isospectral transformations and their applications. *Ann. Mat. Pura Appl.* (4) **199** (2020), no. 4, 1621–1648. MR 4117511 Zbl 1448.34064
- [24] N. J. Guliyev, Schrödinger operators with distributional potentials and boundary conditions dependent on the eigenvalue parameter. J. Math. Phys. 60 (2019), no. 6, 063501, 23 pp. MR 3959138 Zbl 1416.81056
- [25] Y. Guo and G. Wei, Inverse Sturm–Liouville problems with the potential known on an interior subinterval. *Appl. Anal.* 94 (2015), no. 5, 1025–1031. MR 3318324 Zbl 1368.34026
- Y. Guo and G. Wei, Inverse problems: dense nodal subset on an interior subinterval. J. Differential Equations 255 (2013), no. 7, 2002–2017. MR 3072679 Zbl 1288.34013
- [27] H. Hochstadt and B. Lieberman, An inverse Sturm–Liouville problem with mixed given data. SIAM J. Appl. Math. 34 (1978), no. 4, 676–680. MR 0470319
 Zbl 0418.34032
- [28] M. Horváth, Inverse spectral problems and closed exponential systems. Ann. of Math. (2) 162 (2005), no. 2, 885–918. MR 2183284 Zbl 1102.34005
- [29] M. Horváth, On the inverse spectral theory of Schrödinger and Dirac operators. *Trans. Amer. Math. Soc.* 353 (2001), no. 10, 4155–4171. MR 1837225 Zbl 0977.34018
- [30] R. O. Hryniv and Y. V. Mykytyuk, Half-inverse spectral problems for Sturm–Liouville operators with singular potentials. *Inverse Problems* 20 (2004), no. 5, 1423–1444. MR 2109127 Zbl 1074.34007
- [31] M. Horváth and O. Sáfár, Inverse eigenvalue problems. J. Math. Phys. 57 (2016), no. 11, artecle id 112102, 16 pp. MR 3574178 Zbl 1351.81047
- [32] A. Kostenko, A. Sakhnovich, and G. Teschl, Weyl–Titchmarsh theory for Schrödinger operators with strongly singular potentials. *Int. Math. Res. Not. IMRN* 2012, no. 8, 1699–1747. MR 2920828 Zbl 1248.34027
- [33] B. Ja. Levin, *Distribution of zeros of entire functions*. Translated from the Russian by R. P. Boas, J. M. Danskin, F. M. Goodspeed, J. Korevaar, A. L. Shields, and H. P. Thielman. Revised edition. Translations of Mathematical Monographs, 5. American Mathematical Society, Providence, R.I., 1980. MR 0589888
- [34] N. Levinson, The inverse Sturm–Liouville problem. *Mat. Tidsskr.* B 1949 (1949), 25–30. MR 0032067 Zbl 0041.42310 Zbl 0045.36402

- [35] B. M. Levitan and M. G. Gasymov, Determination of a differential equation by two of its spectra. *Russ. Math. Surveys* 19 (1964), no. 2, 1–63. English translation of *Uspehi Mat. Nauk* 19 (1964), no. 2 (116), 3–63, in Russian. MR 0162996 Zbl 0145.10903
- [36] B. M. Levitan and I. S. Sargsjan, *Introduction to spectral theory: selfadjoint ordinary differential operators*. Translated from the Russian by A. Feinstein. Translations of Mathematical Monographs, 39. American Mathematical Society, Providence, R.I., 1975. MR 0369797 Zbl 0302.47036
- [37] N. Makarov and A. Poltoratski, Two spectra theorem with uncertainty. J. Spectr. Theory 9 (2019), no. 4, 1249–1285. MR 4033521 Zbl 1433.42024
- [38] V. A. Marchenko, Some questions in the theory of one-dimensional linear differential operators of the second order. I. Am. Math. Soc., Translat. (2) 101 (1973), 1–104. English translation of Trudy Moskov. Mat. Obšč. 1 (1952), 327–420, in Russian. MR 0058065 Zbl 0273.34014
- [39] W. Y. Ping and C. T. Shieh, Inverse problems for Sturm–Liouville equations with boundary conditions linearly dependent on the spectral parameter from partial information. *Results Math.* 65 (2014), no. 1–2, 105–119. MR 3162432 Zbl 1309.34014
- [40] A. Poltoratski, *Toeplitz approach to problems of the uncertainty principle*. CBMS Regional Conference Series in Mathematics, 121. Published for the Conference Board of the Mathematical Sciences, Washington, D.C., by the American Mathematical Society, Providence, R.I., 2015. MR 3309830 Zbl 1317.47003
- [41] M. Sat, Inverse problems for Sturm–Liouville operators with boundary conditions depending on a spectral parameter. *Electron. J. Differential Equations* 2017, Paper No. 26, 7 pp. MR 3609154 Zbl 1360.34028
- [42] B. Simon, A new approach to inverse spectral theory. I. Fundamental formalism. Ann. of Math. (2) 150 (1999), no. 3, 1029–1057. MR 1740987 Zbl 0945.34013
- [43] G. Teschl, Mathematical methods in quantum mechanics. With applications to Schrödinger operators. Graduate Studies in Mathematics, 99. American Mathematical Society, Providence, R.I., 2009. MR 2499016 Zbl 1166.81004
- [44] G. Teschl and K. Unterkofler, Spectral theory as influenced by Fritz Gesztesy. In H. Holden, B. Simon, and G. Teschl (eds.), *Spectral analysis, differential equations* and mathematical physics: a festschrift in honor of Fritz Gesztesy's 60th birthday. Proceedings of Symposia in Pure Mathematics, 87. American Mathematical Society, Providence, R.I., 2013, 341–363. MR 3087914 Zbl 1345.35066
- [45] Y. P. Wang, Uniqueness theorems for Sturm–Liouville operators with interior twindense nodal set. *Electron. J. Differential Equations* 2017, Paper No. 226, 11 pp. MR 3711179 Zbl 1371.34034
- [46] Y. P. Wang and H. Koyunbakan, On the Hochstadt–Lieberman theorem for discontinuous boundary-valued problems. *Acta Math. Sin. (Engl. Ser.)* **30** (2014), no. 6, 985–992. MR 3200446 Zbl 1301.34019

- [47] Y. P. Wang, C. T. Shieh, and Y. T. Ma, Inverse spectral problems for Sturm–Liouville operators with partial information. *Appl. Math. Lett.* 26 (2013), no. 12, 1175–1181. MR 3250430 Zbl 1315.34026
- [48] Y. P. Wang and V. A. Yurko, On the inverse nodal problems for discontinuous Sturm–Liouville operators. J. Differential Equations 260 (2016), no. 5, 4086–4109. MR 3437580 Zbl 1342.34028
- [49] Z. Wei and G. Wei, Uniqueness results for inverse Sturm–Liouville problems with partial information given on the potential and spectral data. *Bound. Value Probl.* 2016, Paper No. 200, 16 pp. MR 3571341 Zbl 1410.34065
- [50] Z. Wei and G. Wei, On the uniqueness of inverse spectral problems associated with incomplete spectral data. J. Math. Anal. Appl. 462 (2018), no. 1, 697–711. MR 3771270 Zbl 1407.34033
- [51] G. S. Wei and H. K. Xu, On the missing eigenvalue problem for an inverse Sturm– Liouville problem. J. Math. Pures Appl. (9) 91 (2009), no. 5, 468–475. MR 2517783 Zbl 1176.34012
- [52] G. S. Wei and H. K. Xu, Inverse spectral problem with partial information given on the potential and norming constants. *Trans. Amer. Math. Soc.* 364 (2012), no. 6, 3265–3288. MR 2888245 Zbl 1255.34017
- [53] C. F. Yang, Inverse spectral problems for the Sturm–Liouville operator on a *d*-star graph. J. Math. Anal. Appl. 365 (2010), no. 2, 742–749. MR 2587077 Zbl 1195.34023

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