

Anderson localization for a generalized Maryland model with potentials given by skew shifts

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Abstract. In this paper, we proved Anderson localization for the following long-range operator

$$H = \tan \pi \left(x_0 + my_0 + \frac{m(m-1)}{2} \omega \right) \delta_{mn} + \epsilon S_\phi,$$

which generalized the Maryland model to potentials given by skew shifts.

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1. Introduction and main result

Spectral theory of Schrödinger operators is an important topic in both physics and mathematics. Let us begin with the Maryland model

$$H = \lambda \tan \pi(x + n\omega) \delta_{nn'} + \Delta, \tag{1.1}$$

where $\lambda > 0$ is the coupling, $x \in \mathbb{T} = \mathbb{R}/\mathbb{Z}$ is the phase, $\omega \in \mathbb{R} \setminus \mathbb{Q}$ is the frequency, Δ is the lattice Laplacian on \mathbb{Z}

$$\Delta(n, n') = \begin{cases} 1 & |n - n'| = 1; \\ 0 & |n - n'| \neq 1. \end{cases}$$

This is an unbounded self adjoint operator on $l^2(\mathbb{Z})$. We assume

$$x + n\omega - \frac{1}{2} \notin \mathbb{Z}, \quad \text{for } n \in \mathbb{Z},$$

to make the operator well defined.

The Maryland model was originally proposed in physics paper [8]. Simon [14] proved Anderson localization for the Maryland model H (1.1) with Diophantine frequencies ω . Anderson localization means that H has pure point spectrum with exponentially decaying eigenfunctions. Recently, using transfer matrix and Lyapunov exponent, Jitomirskaya and Yang [12] developed a constructive method to prove Anderson localization for the Maryland model. More exactly, Jitomirskaya and Liu [10] proved arithmetic spectral transitions for the Maryland model.

Schrödinger operators can be generalized to long range operators. For example, consider the following long-range operator

$$H_\omega(x) = \cos(x + n\omega)\delta_{nn'} + \epsilon S_\phi, \quad (1.2)$$

where S_ϕ is a Toeplitz operator

$$S_\phi(n, n') = \hat{\phi}(n - n')$$

and ϕ is real analytic, $\hat{\phi}(n)$ is the n -th Fourier coefficient of ϕ . Bourgain and Jitomirskaya [5] proved that there is $\epsilon_0 = \epsilon_0(\phi) > 0$, such that if $0 < \epsilon < \epsilon_0$, $H_\omega(x)$, defined in (1.2), satisfies Anderson localization for $(x, \omega) \in \mathbb{T}^2$ in a set of full measure. This result is non-perturbative, since ϵ_0 does not depend on ω . Note that in the long range case, we cannot use the transfer matrix formalism. In this paper, we focus on 1D quasi-periodic operators. For quasi-periodic operators on \mathbb{Z}^d , we only mention the recent progress made by Jitomirskaya, Liu and Shi [11].

The Maryland model has many generalizations. Using KAM methods, Bellissard, Lima, and Scoppola [6] generalized localization results of the Maryland model to v -dimensional incommensurate structures. This result is perturbative. Recently, Kachkovskiy [13] established non-perturbative Anderson localization for a wide class of quasi-periodic Schrödinger operators with unbounded monotone potentials, extending the classical localization results of the Maryland model. The authors [15] gave a non-perturbative proof of Anderson localization for the Maryland model with long range interactions

$$H(x) = \tan \pi(x + n\omega)\delta_{nn'} + \epsilon S_\phi. \quad (1.3)$$

Note that for all operators mentioned above, potentials are given by shifts

$$Tx = x + \omega. \quad (1.4)$$

Now, let T be skew shifts

$$T(x_1, x_2) = (x_1 + x_2, x_2 + \omega), \quad (x_1, x_2) \in \mathbb{T}^2. \quad (1.5)$$

In this case, using transfer matrix and Lyapounov exponent, Bourgain, Goldstein, and Schlag [4] proved Anderson localization for

$$H = \lambda v(T^n x) + \Delta. \quad (1.6)$$

For long range case, Bourgain [1] considered the following operator

$$H(x) = v(T^m x)\delta_{mn} + \phi_{m-n}(T^m x) + \overline{\phi_{n-m}(T^n x)}, \quad (1.7)$$

where v is a real, nonconstant, trigonometric polynomial, ϕ_k are trigonometric polynomials and T is the skew shift on \mathbb{T}^2 . Using multi-scale method, Bourgain proved Anderson localization for the operator (1.7).

In this paper, we will study a generalized Maryland model with potentials given by skew shifts. More precisely, we consider the following operator

$$H(x) = \tan \pi(T^m x)_1 \delta_{mn} + \epsilon S_\phi, \quad (1.8)$$

where T is the skew shift on \mathbb{T}^2 and $(T^m x)_1$ refers to the first coordinate of $T^m x$. To make the operator (1.8) well defined, we will always assume

$$(T^m x)_1 - \frac{1}{2} \notin \mathbb{Z}, \quad \text{for all } m \in \mathbb{Z}. \quad (1.9)$$

We will prove the following result.

Theorem 1.1. *Consider a lattice operator $H_\omega(x)$ associated to the skew shift $T = T_\omega$ of the form (1.8). Assume $\omega \in \text{DC}$ (diophantine condition)*

$$\|k\omega\| > \gamma|k|^{-2} \quad \text{for all } k \in \mathbb{Z} \setminus \{0\} \quad (1.10)$$

and ϕ real analytic satisfying

$$|\hat{\phi}(n)| < e^{-\rho|n|}, \quad \text{for all } n \in \mathbb{Z}, \quad (1.11)$$

for some $\rho > 0$. Fix $x_0 \in \mathbb{T}^2$. Then for almost all $\omega \in \text{DC}$ and ϵ taken sufficiently small (depending on γ, ρ), $H_\omega(x_0)$ satisfies Anderson localization.

In the long range case here, the transfer matrix formalism is not applicable. Unlike shifts cases, we cannot obtain non-perturbative results in skew shifts cases. Our basic strategy is the same as that in [1], but the main difficulty is that the potential \tan has singularity and the operator H is unbounded.

In order to prove Anderson localization, we need Green's function estimates for

$$G_{[0,N]}(x, E) = (R_{[0,N]}(H(x) - E)R_{[0,N]})^{-1}, \quad (1.12)$$

where R_Λ is the restriction operator to $\Lambda \subset \mathbb{Z}$. Note that

$$R_{[0,N]}(H(x) - E)R_{[0,N]} = D(x)B(x), \tag{1.13}$$

where

$$D(x) = \text{diag}\left(\frac{1}{\cos \pi x_1}, \dots, \frac{1}{\cos \pi (T^N x)_1}\right). \tag{1.14}$$

Hence

$$G_{[0,N]}(x, E) = B(x)^{-1}D(x)^{-1}. \tag{1.15}$$

Since in $D(x)^{-1}$, the singularity $\frac{1}{\cos}$ vanishes, we only need Green’s function estimates for $B(x)^{-1}$. We need to point out that $B(x)$ is not self-adjoint. Fortunately, we find that multi-scale analysis still applies to this case. Since the operator H is unbounded and the energy E is unbounded, we use the specific property of trigonometric functions to overcome the difficulty of the unboundedness of the energy E .

We summarize the structure of this paper. First, we will prove Green’s function estimates in Section 2. Then we recall some facts about semi-algebraic sets in Section 3 and give the proof of Anderson localization in Section 4.

We will use the following notations. For positive numbers a, b , $a \lesssim b$ means $Ca \leq b$ for some constant $C > 0$. $a \ll b$ means C is large. $a \sim b$ means $a \lesssim b$ and $b \lesssim a$. N^{1-} means $N^{1-\epsilon}$ with some small $\epsilon > 0$. For $x \in \mathbb{R}$, $\|x\| = \inf_{m \in \mathbb{Z}} |x - m|$, for $x = (x_1, x_2) \in \mathbb{T}^2$, $\|x\| = \|x_1\| + \|x_2\|$.

2. Green’s function estimates

In this section, we will prove the Green’s function estimates using multi-scale analysis in [1].

We need the following lemma.

Lemma 2.1 (Lemma 3.16 in [1]). *Let $A(x) = \{A_{mn}(x)\}_{1 \leq m, n \leq N}$ be a matrix-valued function on \mathbb{T}^d such that*

$$A(x) \text{ is self-adjoint for } x \in \mathbb{T}^d, \tag{2.1}$$

$$A_{mn}(x) \text{ is a trigonometric polynomial of degree } < N^{C_1}, \tag{2.2}$$

$$|A_{mn}(x)| < C_2 e^{-c_2|m-n|}, \tag{2.3}$$

where $c_2, C_1, C_2 > 0$ are constants.

Let $0 < \delta < 1$ be sufficiently small, $M = N^{\delta^6}$, $L_0 = N^{\frac{1}{100}\delta^2}$, $0 < c_3 < \frac{1}{10}c_2$.

Assume that for any interval $I \subset [1, N]$ of size L_0 , except for x in a set of measure $< e^{-L_0^{\delta^3}}$,

$$\|(R_I A(x) R_I)^{-1}\| < e^{L_0^{1-}}, \tag{2.4}$$

$$|(R_I A(x) R_I)^{-1}(m, n)| < e^{-c_3|m-n|}, \quad m, n \in I, |m - n| > \frac{L_0}{10}. \tag{2.5}$$

For fixed $x \in \mathbb{T}^d$, $n_0 \in [1, N]$ is called a good site if $I_0 = [n_0 - \frac{M}{2}, n_0 + \frac{M}{2}] \subset [1, N]$ and

$$\|(R_{I_0} A(x) R_{I_0})^{-1}\| < e^{M^{1-}}, \tag{2.6}$$

$$|(R_{I_0} A(x) R_{I_0})^{-1}(m, n)| < e^{-c_3|m-n|}, \quad m, n \in I_0, |m - n| > \frac{M}{10}. \tag{2.7}$$

Denote $\Omega(x) \subset [1, N]$ the set of bad sites. Assume that for any interval $J \subset [1, N]$ such that $|J| > N^{\frac{\delta}{5}}$, we have

$$|J \cap \Omega(x)| < |J|^{1-\delta}. \tag{2.8}$$

Then

$$\|A(x)^{-1}\| < e^{N^{1-\frac{\delta}{C(d)}}}, \tag{2.9}$$

$$|A(x)^{-1}(m, n)| < e^{-c'_3|m-n|}, \quad |m - n| > \frac{N}{10} \tag{2.10}$$

except for x in a set of measure $< e^{-\frac{N^{\delta^2}}{C(d)}}$, where $C(d)$ is a constant depending on d and $c'_3 > c_3 - (\log N)^{-8}$.

We also need the following ergodic property of skew shifts on \mathbb{T}^2 .

Lemma 2.2 (Lemma 15.21 in [2]). Assume $\omega \in \text{DC}$, $T = T_\omega$ is the skew shift on \mathbb{T}^2 , $\epsilon > L^{-\frac{1}{10}}$. Then

$$\#\{n = 1, \dots, L : \|T^n x - a\| < \epsilon\} < C\epsilon^2 L.$$

Remark 2.3. In the proof of Lemma 2.2, we only need to assume

$$\|k\omega\| > \gamma|k|^{-2}, \quad \text{for all } 0 < |k| \leq L.$$

By Lemma 2.1, Lemma 2.2, we can prove the Green's function estimates.

Proposition 2.4. Let $T = T_\omega$ be the skew shift and

$$H_{mn}(x) = \tan \pi(T^m x)_1 \delta_{mn} + \epsilon S_\phi. \tag{2.11}$$

Assume ϕ real analytic satisfying

$$|\hat{\phi}(n)| < e^{-\rho|n|}, \quad \text{for all } n \in \mathbb{Z} \tag{2.12}$$

for some $\rho > 0$ and ω satisfying

$$\|k\omega\| > \gamma|k|^{-2}, \quad \text{for all } 0 < |k| \leq N, \tag{2.13}$$

ϵ is small (depending on γ, ρ). Then for energy E ,

$$\|G_{[0,N]}(x, E)\| < e^{N^{1-}}, \tag{2.14}$$

$$|G_{[0,N]}(x, E)(m, n)| < e^{-\frac{\rho}{100}|m-n|}, \quad 0 \leq m, n \leq N, |m - n| > \frac{N}{10} \tag{2.15}$$

for $x \notin \Omega_N(E)$, where

$$\text{mes } \Omega_N(E) < e^{-N^\sigma}, \quad \sigma > 0. \tag{2.16}$$

Proof. Write

$$H_{[0,N]}(x) - E = D_{[0,N]}(x)B_{[0,N]}(x), \tag{2.17}$$

where

$$D_{mn}(x) = \frac{\sqrt{1 + (\epsilon\hat{\phi}(0) - E)^2}}{\cos \pi(T^m x)_1} \delta_{mn}, \tag{2.18}$$

$$B_{mm}(x) = \frac{1}{\sqrt{1 + (\epsilon\hat{\phi}(0) - E)^2}} [\sin \pi(T^m x)_1 + (\epsilon\hat{\phi}(0) - E) \cos \pi(T^m x)_1], \tag{2.19}$$

$$B_{mn}(x) = \frac{\epsilon\hat{\phi}(m - n) \cos \pi(T^m x)_1}{\sqrt{1 + (\epsilon\hat{\phi}(0) - E)^2}}, \quad m \neq n. \tag{2.20}$$

We will apply Lemma 2.1 to $B_{[0,N]}(x)$. Note that $B_{[0,N]}(x)$ is not self-adjoint. However, in the proof of Lemma 2.1, we don't need (2.1). Since

$$T^m(x_1, x_2) = \left(x_1 + mx_2 + \frac{m(m - 1)}{2}\omega, x_2 + m\omega\right), \tag{2.21}$$

$B_{mn}(x)$ is a trigonometric polynomial of degree $< |m|$. (2.3) holds with $C_2 = 1$, $c_2 = \rho$.

We need to prove

$$\text{mes}\{x \in \mathbb{T}^2: \text{there exist } m \text{ and } n, \text{ with } 0 \leq m, n \leq N, \text{ such that } |B_{[0,N]}(x)^{-1}(m, n)| > e^{N^{1-c_3|m-n|} \chi_{|m-n| > \frac{N}{10}}}\} < e^{-N^{\delta^3}} \tag{2.22}$$

for some $c_3 > \frac{\rho}{100}$, $0 < \delta < 1$.

By

$$\begin{aligned} & |\sin \pi x + (\epsilon \hat{\phi}(0) - E) \cos \pi x| \\ &= \sqrt{1 + (\epsilon \hat{\phi}(0) - E)^2} |\cos \pi(x - \alpha)|, \quad 0 < \alpha < 1, \end{aligned}$$

using the fact

$$\text{mes}\{x \in [0, 1]: |\cos \pi x| < \eta\} < \eta, \quad \text{for all } 0 < \eta < 1,$$

we have

$$\text{mes} \left[x \in [0, 1] \mid \frac{1}{\sqrt{1 + (\epsilon \hat{\phi}(0) - E)^2}} |\sin \pi x + (\epsilon \hat{\phi}(0) - E) \cos \pi x| < \epsilon_0 \right] < \epsilon_0. \quad (2.23)$$

Since T is a measure-preserving transformation,

$$\begin{aligned} \text{mes} \left[x \in \mathbb{T}^2 \mid \frac{1}{\sqrt{1 + (\epsilon \hat{\phi}(0) - E)^2}} |\sin \pi(T^m x)_1 + (\epsilon \hat{\phi}(0) - E) \cos \pi(T^m x)_1| \right. \\ \left. < \epsilon_0 \right] < \epsilon_0. \end{aligned} \quad (2.24)$$

Hence

$$\text{mes}[x \in \mathbb{T}^2 \mid \min_{0 \leq m \leq N_0} |B_{mm}(x)| < \epsilon_0] < N_0 \epsilon_0. \quad (2.25)$$

If $\min_{0 \leq m \leq N_0} |B_{mm}(x)| > \epsilon_0 > \epsilon$, take $\epsilon_0 = e^{-N_0^{\frac{1}{2}}}$, $\epsilon = e^{-N_0}$, by Neumann expansion and (2.25), we have

$$|B_{[0, N_0]}(x)^{-1}(m, n)| < e^{N_0^{\frac{1}{2}} - \frac{\rho}{2}|m-n|}, \quad m, n \in [0, N_0] \quad (2.26)$$

except for x in a set of measure $< e^{-cN_0^{\frac{1}{2}}}$. So, (2.22) holds for an initial scale N_0 .

Assume (2.22) holds up to scale $L_0 = N^{\frac{1}{100}\delta^2}$, since

$$B_{m+1, n+1}(x) = B_{mn}(Tx), \quad (2.27)$$

(2.4) and (2.5) will hold for x outside a set of measure at most $e^{-L_0^{\delta^3}}$. Denote $\Omega(x) \subset [0, N]$ the set of bad sites with respect to scale $M = N^{\delta^6}$. $n_0 \notin \Omega(x)$ means

$$\begin{aligned} & |B_{[0, M]}(T^{n_0 - \frac{M}{2}} x)^{-1}(m, n)| \\ &= \left| B_{[n_0 - \frac{M}{2}, n_0 + \frac{M}{2}]}(x)^{-1} \left(m + n_0 - \frac{M}{2}, n + n_0 - \frac{M}{2} \right) \right| \\ &< e^{M^{1-c_3|m-n|} \chi_{|m-n| > \frac{M}{10}}}, \quad m, n \in [0, M]. \end{aligned} \quad (2.28)$$

From the inductive hypothesis, we have

$$|B_{[0,M]}(x)^{-1}(m, n)| < e^{M^{1-c_3|m-n|} \chi_{|m-n| > \frac{M}{10}}}, \quad m, n \in [0, M] \tag{2.29}$$

for $x \notin \Omega_0$, $\text{mes } \Omega_0 < e^{-M^{\delta^3}}$. By (2.28), (2.29), Lemma 2.1, we only need to show that for any $x \in \mathbb{T}^2$, $N^{\frac{\delta}{5}} < L < N$,

$$\#\{1 \leq n \leq L \mid T^n x \in \Omega_0\} < L^{1-\delta}. \tag{2.30}$$

Expressing (2.29) as a ratio of determinants and replacing \cos, \sin by truncated power series, Ω_0 may be viewed as a semi-algebraic set of degree at most M^6 . (For properties of semi-algebraic sets, see Section 3.) If $r > e^{-\frac{1}{2}M^{\delta^3}}$, by Proposition 3.2, Ω_0 may be covered by at most $M^C \left(\frac{1}{r}\right) r$ -balls. Choosing $r = L^{-\frac{1}{20}} > N^{-1} > e^{-\frac{1}{2}M^{\delta^3}}$, using Lemma 2.2, Remark 2.3, we have

$$\#\{1 \leq n \leq L \mid T^n x \in \Omega_0\} < M^C \left(\frac{1}{r}\right) r^2 L < L^{C\delta^5 + 1 - \frac{1}{20}} < L^{1-\delta}.$$

This proves (2.30) and (2.22).

By (2.17),

$$G_{[0,N]}(x, E) = (H_{[0,N]}(x) - E)^{-1} = B_{[0,N]}(x)^{-1} D_{[0,N]}(x)^{-1}, \tag{2.31}$$

hence

$$G_{[0,N]}(x, E)(m, n) = \frac{\cos \pi(T^n x)_1}{\sqrt{1 + (\epsilon \hat{\phi}(0) - E)^2}} B_{[0,N]}(x)^{-1}(m, n), \quad m, n \in [0, N]. \tag{2.32}$$

By (2.31) and (2.32),

$$\|G_{[0,N]}(x, E)\| \leq \|B_{[0,N]}(x)^{-1}\|, \tag{2.33}$$

$$|G_{[0,N]}(x, E)(m, n)| \leq |B_{[0,N]}(x)^{-1}(m, n)|, \quad m, n \in [0, N]. \tag{2.34}$$

Proposition 2.4 follows from (2.22), (2.33), and (2.34). □

3. Semi-algebraic sets

We recall some basic facts of semi-algebraic sets in this section, which is needed in Section 4. Let $\mathcal{P} = \{P_1, \dots, P_s\} \subset \mathbb{R}[X_1, \dots, X_n]$ be a family of real polynomials whose degrees are bounded by d . A semi-algebraic set is given by

$$S = \bigcup_j \bigcap_{l \in L_j} \{\mathbb{R}^n \mid P_l s_{jl} 0\}, \tag{3.1}$$

where $L_j \subset \{1, \dots, s\}$, $s_{jl} \in \{\leq, \geq, =\}$ are arbitrary. We say that S has degree at most sd and its degree is the inf of sd over all representations as in (3.1).

We need the following quantitative version of the Tarski-Seidenberg principle.

Proposition 3.1 ([7]). *Let $S \subset \mathbb{R}^n$ be a semi-algebraic set of degree B , then any projection of S is semi-algebraic of degree at most B^C , $C = C(n)$.*

We also need the following fact.

Proposition 3.2 (Corollary 9.6 in [2]). *Let $S \subset [0, 1]^n$ be semi-algebraic of degree B . Let $\epsilon > 0$, $\text{mes}_n S < \epsilon^n$. Then S may be covered by at most $B^C (\frac{1}{\epsilon})^{n-1} \epsilon$ -balls.*

Finally, we will use the following lemma.

Lemma 3.3 (Lemma 15.26 in [2]). *Let $S \subset \mathbb{T}^3$ be a semi-algebraic set of degree B such that*

$$\text{mes } S < e^{-B^\sigma}, \quad \sigma > 0.$$

Let M be an integer satisfying

$$\log \log M \ll \log B \ll \log M.$$

Then for any fixed $x_0 \in \mathbb{T}^2$,

$$\text{mes}[\omega \in \mathbb{T} \mid \text{there exists } j \sim M \text{ such that } (\omega, T_\omega^j x_0) \in S] < M^{-c}, \quad c > 0,$$

where T_ω is the skew shift with frequency ω .

4. Proof of Anderson localization

In this section, we give the proof of Anderson localization as in [3].

By application of the resolvent identity, we have the following

Lemma 4.1. *Let $I \subset \mathbb{Z}$ be an interval of size N and $\{I_\alpha\}$ subintervals of size $M \ll N$, $N = e^{(\log M)^2}$. Assume that, for all $k \in I$, there is some α such that*

$$\left[k - \frac{M}{4}, k + \frac{M}{4} \right] \cap I \subset I_\alpha \tag{4.1}$$

and, for all α ,

$$\|G_{I_\alpha}\| < e^{M^{1-}}, \quad |G_{I_\alpha}(n_1, n_2)| < e^{-\frac{\rho}{100}|n_1 - n_2|}, \tag{4.2}$$

with $n_1, n_2 \in I_\alpha$, $|n_1 - n_2| > \frac{M}{10}$.

Then

$$|G_I(n_1, n_2)| < e^M, \quad n_1, n_2 \in I, \tag{4.3}$$

$$|G_I(n_1, n_2)| < e^{-\frac{\rho}{200}|n_1-n_2|}, \quad n_1, n_2 \in I, |n_1 - n_2| > \frac{N}{10}. \tag{4.4}$$

Proof. For $m, n \in I$, there is some α such that

$$\left[m - \frac{M}{4}, m + \frac{M}{4} \right] \cap I \subset I_\alpha. \tag{4.5}$$

By resolvent identity,

$$|G_I(m, n)| \leq e^{M^{1-}} + \sum_{n_1 \in I_\alpha, n_2 \notin I_\alpha} |G_{I_\alpha}(m, n_1)| e^{-\rho|n_1-n_2|} |G_I(n_2, n)|. \tag{4.6}$$

If $|m - n_1| \leq \frac{M}{8}$, then $|n_1 - n_2| \geq \frac{M}{8}$, hence

$$\sum_{|m-n_1| \leq \frac{M}{8}, n_2 \notin I_\alpha} |G_{I_\alpha}(m, n_1)| e^{-\rho|n_1-n_2|} \leq M e^{M^{1-}} e^{-\rho \frac{M}{8}} < \frac{1}{4}. \tag{4.7}$$

If $|m - n_1| > \frac{M}{8}$, then $|G_{I_\alpha}(m, n_1)| < e^{-\frac{\rho}{100}|m-n_1|}$, hence

$$\sum_{|m-n_1| > \frac{M}{8}, n_2 \notin I_\alpha} |G_{I_\alpha}(m, n_1)| e^{-\rho|n_1-n_2|} < e^{-\frac{\rho}{100}M} < \frac{1}{4}. \tag{4.8}$$

By (4.6), (4.7), and (4.8),

$$\max_{m, n \in I} |G_I(m, n)| < e^{M^{1-}} + \frac{1}{2} \max_{m, n \in I} |G_I(m, n)|. \tag{4.9}$$

(4.3) follows from (4.9).

Take $m, n \in I, |m - n| > \frac{N}{10}$, assume (4.5), by resolvent identity,

$$\begin{aligned} |G_I(m, n)| &\leq \sum_{n_0 \in I_\alpha, n_1 \notin I_\alpha} |G_{I_\alpha}(m, n_0)| e^{-\rho|n_0-n_1|} |G_I(n_1, n)| \\ &\leq M \sum_{|m-n_1| > \frac{M}{4}} e^{-\frac{\rho}{100}|m-n_1|} |G_I(n_1, n)| \\ &\leq M^t \sum_{|m-n_1| > \frac{M}{4}, \dots, |n_{t-1}-n_t| > \frac{M}{4}} e^{-\frac{\rho}{100}(|m-n_1| + \dots + |n_{t-1}-n_t|)} |G_I(n_t, n)| \end{aligned} \tag{4.10}$$

where $t \leq 10 \frac{N}{M}$.

If $|n - n_t| \leq M$, then by (4.3) and (4.10),

$$|G_I(m, n)| \leq M^t N^t e^{M - \frac{\rho}{100}|m-n_t|} \leq e^{20 \frac{N}{M} \log N + 2M - \frac{\rho}{100}|m-n|} < e^{-\frac{\rho}{200}|m-n|}. \tag{4.11}$$

If $t = 10\frac{N}{M}$, then by (4.3) and (4.10),

$$|G_I(m, n)| \leq M^t N^t e^{-\frac{\rho}{100}\frac{10N}{M}\frac{M}{4} + M} \leq e^{20\frac{N}{M} \log N + M - \frac{\rho}{100}2N} < e^{-\frac{\rho}{100}|m-n|}. \tag{4.12}$$

(4.4) follows from (4.11) and (4.12). This proves Lemma 4.1. \square

Now we can prove the main result.

Theorem 4.2. *Consider a lattice operator $H_\omega(x)$ associated to the skew shift $T = T_\omega$ of the form*

$$H_\omega(x) = \tan \pi(T^m x)_1 \delta_{mn} + \epsilon S_\phi. \tag{4.13}$$

Assume $\omega \in \text{DC}$ (diophantine condition)

$$\|k\omega\| > \gamma|k|^{-2}, \quad \text{for all } k \in \mathbb{Z} \setminus \{0\} \tag{4.14}$$

and ϕ real analytic satisfying

$$|\hat{\phi}(n)| < e^{-\rho|n|}, \quad \text{for all } n \in \mathbb{Z} \tag{4.15}$$

for some $\rho > 0$. Fix $x_0 \in \mathbb{T}^2$. Then for almost all $\omega \in \text{DC}$ and ϵ taken sufficiently small(depending on γ, ρ), $H_\omega(x_0)$ satisfies Anderson localization.

Proof. By Shnol’s theorem [9], to establish Anderson localization, it suffices to show that if $\xi = (\xi_n)_{n \in \mathbb{Z}}$, $E \in \mathbb{R}$ satisfy

$$\xi_0 = 1, |\xi_n| < C|n|, \quad |n| \rightarrow \infty, \tag{4.16}$$

$$H(x_0)\xi = E\xi, \tag{4.17}$$

then

$$|\xi_n| < e^{-c|n|}, \quad |n| \rightarrow \infty. \tag{4.18}$$

Denote $\Omega = \Omega(E) \subset \mathbb{T}^2$ the set of x such that

$$|G_{[-N, N]}(x, E)(m, n)| < e^{N^{1-\frac{\rho}{100}}|m-n| \chi_{|m-n| > \frac{N}{10}}} \tag{4.19}$$

fails for some $|m|, |n| \leq N$. Let $N_1 = N^{C_1}$, C_1 is a sufficiently large constant. Then by Proposition 2.4,

$$\text{mes } \Omega(E) < e^{-N^\sigma}, \tag{4.20}$$

$$\#\{|j| \leq N_1 | T^j x_0 \in \Omega\} < N_1^{1-\delta}. \tag{4.21}$$

So, we may find an interval $I \subset [0, N_1]$ of size N such that

$$T^j x_0 \notin \Omega, \quad \text{for all } j \in I \cup (-I). \tag{4.22}$$

Hence

$$|G_{[j-N, j+N]}(x_0, E)(m, n)| < e^{N^{1-\frac{\rho}{100}}|m-n|\chi_{|m-n|>\frac{N}{10}}}, \quad m, n \in [j - N, j + N]. \tag{4.23}$$

By (4.16), (4.17), and (4.23), we have

$$|\xi_j| \leq C \sum_{\substack{n_1 \in [j-N, j+N] \\ n_2 \notin [j-N, j+N]}} e^{N^{1-\frac{\rho}{100}}|j-n_1|\chi_{|j-n_1|>\frac{N}{10}}} e^{-\rho|n_1-n_2|} |n_2| < e^{-\frac{\rho}{200}N}. \tag{4.24}$$

Denoting j_0 the center of I , we have

$$1 = \xi_0 \leq \|G_{[-j_0, j_0]}(x_0, E)\| \|R_{[-j_0, j_0]}H(x_0)R_{Z \setminus [-j_0, j_0]}\xi\|. \tag{4.25}$$

By (4.16) and (4.24), we have for $|n| \leq j_0$,

$$\begin{aligned} & |(R_{[-j_0, j_0]}H(x_0)R_{Z \setminus [-j_0, j_0]}\xi)_n| \\ & \leq \sum_{|n_1|>j_0} e^{-\rho|n-n_1|} |\xi_{n_1}| \\ & \leq \sum_{j_0 < |n_1| \leq j_0 + \frac{N}{2}} e^{-\rho|n-n_1|} e^{-\frac{\rho}{200}N} + C \sum_{|n_1|>j_0 + \frac{N}{2}} e^{-\rho|n-n_1|} |n_1| < e^{-\frac{\rho}{400}N}. \end{aligned} \tag{4.26}$$

By (4.25) and (4.26),

$$\|G_{[-j_0, j_0]}(x_0, E)\| > e^{\frac{\rho}{500}N}, \tag{4.27}$$

hence

$$\text{dist}(E, \text{spec } H_{[-j_0, j_0]}(x_0)) < e^{-\frac{\rho}{500}N}. \tag{4.28}$$

Denote

$$\mathcal{E}_\omega = \bigcup_{|j| \leq N_1} \text{spec } H_{[-j, j]}(x_0). \tag{4.29}$$

It follows from (4.28) that if $x \notin \bigcup_{E' \in \mathcal{E}_\omega} \Omega(E')$, then

$$|G_{[-N, N]}(x, E)(m, n)| < e^{N^{1-\frac{\rho}{100}}|m-n|\chi_{|m-n|>\frac{N}{10}}}, \quad |m|, |n| \leq N. \tag{4.30}$$

Consider the set $S \subset \mathbb{T}^3 \times \mathbb{R}$ of (ω, x, E') , where

$$\|k\omega\| > c|k|^{-2}, \quad \text{for all } 0 < |k| \leq N, \quad (4.31)$$

$$x \in \Omega(E'), \quad (4.32)$$

$$E' \in \mathcal{E}_\omega. \quad (4.33)$$

By Proposition 3.1,

$$\text{Proj}_{\mathbb{T}^3} S \text{ is a semi-algebraic set of degree } < N^C, \quad (4.34)$$

and by (4.20),

$$\text{mes}(\text{Proj}_{\mathbb{T}^3} S) < e^{-\frac{1}{2}N^\sigma}. \quad (4.35)$$

Let $N_2 = e^{(\log N)^2}$,

$$\mathcal{R}_N = \{\omega \in \mathbb{T} \mid (\omega, T^j x_0) \in \text{Proj}_{\mathbb{T}^3} S, \exists |j| \sim N_2\}. \quad (4.36)$$

By (4.34), (4.35), (4.36), using Lemma 3.3, $\text{mes } \mathcal{R}_N < N_2^{-c}$, $c > 0$. Let

$$\mathcal{R} = \bigcap_{N_0 \geq 1} \bigcup_{N \geq N_0} \mathcal{R}_N, \quad (4.37)$$

then, by Borel–Cantelli theorem, $\text{mes } \mathcal{R} = 0$. We restrict $\omega \notin \mathcal{R}$.

If $\omega \notin \mathcal{R}_N$, we have for all $|j| \sim N_2$, $(\omega, T^j x_0) \notin \text{Proj}_{\mathbb{T}^3} S$, by (4.30),

$$|G_{[j-N, j+N]}(x_0, E)(m, n)| < e^{N^{1-\frac{\rho}{100}}|m-n| \chi_{|m-n| > \frac{N}{10}}}, \quad m, n \in [j-N, j+N]. \quad (4.38)$$

Let

$$\Lambda = \bigcup_{\frac{1}{4}N_2 < j < 4N_2} [j-N, j+N] \supset \left[\frac{1}{4}N_2, 4N_2 \right],$$

by Lemma 4.1, we deduce from (4.38) that

$$|G_\Lambda(x_0, E)(m, n)| < e^{-\frac{\rho}{200}|m-n|}, \quad |m-n| > \frac{N_2}{10}, \quad (4.39)$$

and therefore

$$|\xi_j| < e^{-\frac{\rho}{4000}|j|}, \quad \frac{1}{2}N_2 \leq j \leq N_2. \quad (4.40)$$

Since $\omega \notin \mathcal{R}$, by (4.37), there is some $N_0 > 0$, such that for all $N \geq N_0$, $\omega \notin \mathcal{R}_N$. So, (4.40) holds for $j \in \bigcup_{N \geq N_0} [\frac{1}{2}e^{(\log N)^2}, e^{(\log N)^2}] = [\frac{1}{2}e^{(\log N_0)^2}, \infty)$. This proves (4.18) for $j > 0$, similarly for $j < 0$. Hence Theorem 4.2 follows. \square

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References

- [1] J. Bourgain, Estimates on Green's functions, localization and the quantum kicked rotor model. *Ann. of Math. (2)* **156** (2002), no. 1, 249–294. [MR 1935847](#)
[Zbl 1213.82054](#)
- [2] J. Bourgain, *Green's function estimates for lattice Schrödinger operators and applications*. Annals of Mathematics Studies, 158. Princeton University Press, Princeton, N.J., 2005. [MR 2100420](#) [Zbl 1137.35001](#)
- [3] J. Bourgain and M. Goldstein, On nonperturbative localization with quasi-periodic potential. *Ann. of Math. (2)* **152** (2000), no. 3, 835–879. [MR 1815703](#) [Zbl 1053.39035](#)
- [4] J. Bourgain, M. Goldstein, and W. Schlag, Anderson localization for Schrödinger operators on \mathbb{Z} with potentials given by the skew-shift. *Comm. Math. Phys.* **220** (2001), no. 3, 583–621. [MR 1843776](#) [Zbl 0994.82044](#)
- [5] J. Bourgain and S. Jitomirskaya, Absolutely continuous spectrum for 1D quasiperiodic operators. *Invent. Math.* **148** (2002), no. 3, 453–463. [MR 1908056](#)
[Zbl 1036.47019](#)
- [6] J. B ellissard, R. Lima and E. Scoppola, Localization in v -dimensional incommensurate structures. *Comm. Math. Phys.* **88** (1983), no. 4, 465–477. [MR 0702564](#)
[Zbl 0543.35018](#)
- [7] S. Basu, R. Pollack and M.-F Roy, On the combinatorial and algebraic complexity of quantifier elimination. *J. ACM* **43** (1996), no. 6, 1002–1045. [MR 1434910](#)
[Zbl 0885.68070](#)
- [8] D. Grempel, S. Fishman and R. Prange, Localization in an incommensurate potential: an exactly solvable model. *Phys. Rev. Lett.* **49**, no. 11 (1982), 833–836.
- [9] R. Han, Shnol's theorem and the spectrum of long range operators. *Proc. Amer. Math. Soc.* **147** (2019), no. 7, 2887–2897. [MR 3973892](#) [Zbl 07073444](#)
- [10] S. Jitomirskaya and W. Liu, Arithmetic spectral transitions for the Maryland model. *Comm. Pure Appl. Math.* **70** (2017), no. 6, 1025–1051. [MR 3639318](#) [Zbl 06734215](#)
- [11] S. Jitomirskaya, W. Liu, and Y. Shi, Anderson localization for multi-frequency quasi-periodic operators on \mathbb{Z}^d . *Geom. Funct. Anal.* **30** (2020), no. 2, 457–481.
[MR 4108613](#) [Zbl 07219753](#)
- [12] S. Jitomirskaya and F. Yang, Pure point spectrum for the Maryland model: a constructive proof. *Ergodic Theory Dynam. Systems* **41** (2021), no. 1, 283–294. [MR 4190056](#)
[Zbl 1456.37053](#)

- [13] I. Kachkovskiy, Localization for quasiperiodic operators with unbounded monotone potentials. *J. Funct. Anal.* **277** (2019), no. 10, 3467–3490. [MR 4001077](#)
[Zbl 07104034](#)
- [14] B. Simon, Almost periodic Schrödinger operators. IV. The Maryland model. *Ann. Physics* **159** (1985), no. 1, 157–183. [MR 0776654](#) [Zbl 0595.35032](#)
- [15] J. Shi and X. Yuan, Anderson localization for the Maryland model with long range interactions. Preprint, 2019. [arXiv:1909.06542](#)

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