

Finite-Element Methods for Singularly Perturbed Elliptic Boundary Value Problems and its Application to the Stationary Navier-Stokes Equations¹⁾

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Wir betrachten einige Varianten der Methode der finiten Elemente zur Lösung singular gestörter elliptischer Randwertaufgaben 2. Ordnung, wobei das reduzierte Problem von erster Ordnung ist. Stabilitäts- und Konvergenzeigenschaften dieser Methoden und ihre Anwendung auf die stationären Navier-Stokes-Gleichungen für große Reynolds-Zahlen werden untersucht.

Мы рассматриваем некоторые варианты метода конечных элементов для решения сингулярно возмущённых эллиптических граничных задач второго порядка, при которых редуцированная задача — проблема первого порядка. Исследуются сходимость и устойчивость таких методов и их применение на стационарное уравнение Навье—Стокса при больших числах Рейнолдса.

We consider some modifications of finite-element methods for solving singularly perturbed elliptic boundary value problems of second order where the reduced problem is of first order. Stability and convergence properties of such methods and the application to the stationary Navier-Stokes problem for high Reynolds numbers are studied.

1. Introduction

For solving singularly perturbed elliptic boundary value problems, the common finite-element methods using, for instance, piecewise linear functions, are unstable and exhibit spurious oscillations unless the discretization parameter is not sufficiently small. Moreover, the occurrence of boundary layers influences the approximation properties in a negative sense — the usual error estimates become meaningless. For this reason, in the last years many modifications of the standard finite-element method have been developed to overcome these difficulties, i.e. to guarantee stability and to obtain a good approximation of the exact solution. Such proposals are Petrov-Galerkin methods with piecewise polynomial bases [9], special integration rules for the convective terms [5], the utilization of special directional derivatives [3, 25], some mixed finite-element methods [13], the symmetrization of the bilinear form [2], the use of artificial diffusion [15, 16], asymptotically fitted methods [8, 22], hybrid upwind finite-element methods [12, 21, 26] and the streamline diffusion method [10, 11, 19].

Nowadays, from a mathematical point of view the last three methods are the best ones for solving singularly perturbed problems because of its mathematical foundations and its favourable properties. In the following we will give an overview on the main ideas and results of these three classes of methods in the case of the two-dimensional stationary convection-diffusion problem ($\Omega \subset \mathbb{R}^2$, $\Gamma = \partial\Omega$)

$$-\varepsilon \Delta u + b(x) \nabla u + c(x) u = f(x) \text{ in } \Omega, \quad u = 0 \text{ on } \Gamma, \quad (1.1)$$

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and discuss its application to the stationary Navier-Stokes equations

$$-\varepsilon \Delta u + u \nabla u + \nabla p = f \text{ and } \operatorname{div} u = 0 \text{ in } \Omega, \quad u = 0 \text{ on } \Gamma.$$

We will assume that Ω is a polygon, the functions b , c and f are sufficiently regular and ε is a small positive parameter. Furthermore, we denote, by $\|\cdot\|_{k,p}$, $\|\cdot\|_{k,p}$ the usual seminorm and norm in the Sobolev space $W^{k,p}(\Omega)$, by (\cdot, \cdot) the scalar product in the space $L^2(\Omega)$ and by C a generic constant independent of ε and h .

2. Asymptotically fitted finite-element methods

The weak formulation of problem (1.1) reads as follows:

$$\begin{aligned} \text{Find } u \in H_0^1(\Omega) \text{ such that for all } v \in H_0^1(\Omega) \\ B(u, v) \equiv \varepsilon(\nabla u, \nabla v) + (b \nabla u + cu, v) = (f, v). \end{aligned} \quad (2.1)$$

Under the assumption $(c - 2^{-1} \operatorname{div} b)(x) \geq \alpha > 0$ for all x the bilinear form B is H_0^1 -elliptic and by means of Lax-Milgram's theorem we obtain the existence of a unique solution u of the problem (2.1). Choosing a conforming finite-element space $V_h \subset H_0^1(\Omega)$ which we will specify later we get the following discrete problem:

$$\begin{aligned} \text{Find } u_h \in V_h \text{ such that for all } v \in V_h \\ B(u_h, v) = (f, v). \end{aligned} \quad (2.2)$$

Under the above assumption also problem (2.2) has a unique solution u_h . In order to obtain error estimates with ε -independent error constants, we will use the ε -weighted H^1 -norm and the fitted norm defined by

$$\|v\|_\varepsilon = (\varepsilon \|v\|_{1,2}^2 + \|v\|_{0,2}^2)^{1/2} \quad \text{and} \quad \| \|v\| \| = \|v\|_\varepsilon + \sup_{0 \neq v \in V} \frac{B(v, v)}{\|v\|_\varepsilon},$$

respectively, which are equivalent to the H^1 -norm for fixed $\varepsilon > 0$. Moreover, it holds $\|u\|_{0,2} \leq \|u\|_\varepsilon \leq \| \|u\| \| \leq \|u\|_{1,2}$. From the H_0^1 -ellipticity of B in the ε -norm with the constant $c_1 = \min(1, \alpha)$ we have for each $w_h \in V_h$

$$\|w_h - u_h\|_\varepsilon \leq \frac{1}{c_1} \frac{B(w_h - u_h, w_h - u_h)}{\|w_h - u_h\|_\varepsilon} = \frac{1}{c_1} \frac{B(w_h - u, w_h - u_h)}{\|w_h - u_h\|_\varepsilon}$$

and conclude

$$\|u - u_h\|_\varepsilon \leq \|u - w_h\|_\varepsilon + \|w_h - u_h\|_\varepsilon \leq \|u - w_h\|_\varepsilon + \frac{1}{c_1} \sup_{v \in V_h} \frac{B(w_h - u, v)}{\|v\|_\varepsilon}.$$

Thus, the estimation of the error is reduced to the estimation of the approximation error in the fitted norm $\|u - u_h\|_\varepsilon \leq C \inf \{ \| \|u - w_h\| \| : w_h \in V_h \}$. Usually, the approximation error is replaced by the interpolation error such that for instance for spaces of bilinear elements the estimate $\inf \{ \| \|u - w_h\| \| : w_h \in V_h \} \leq Ch |u|_{2,2}$ holds. Now, the occurrence of boundary layers in the solution affects the boundedness of $|u|_{2,2}$ for $\varepsilon \rightarrow 0$, i.e. only $\varepsilon^{3/2} |u|_{2,2}$ is uniformly ε -bounded and the resulting estimate reads

$$\|u - u_h\|_\varepsilon \leq Ch \varepsilon^{-3/2}. \quad (2.3)$$

It is clear that this becomes meaningless provided that $\varepsilon \ll h$.

The principle of asymptotically fitted finite-element methods consists of splitting up the approximation error in two parts

$$\inf_{u_h \in V_h} \| \|u - w_h\| \| \leq \| \|u - u_{as}\| \| + \inf_{w_h \in V_h} \| \|u_{as} - w_h\| \|,$$

namely in the asymptotic error $\| \|u - u_{as}\| \|$ and the approximation error for the asymptotic solution u_{as} . Thus, the investigation consists of two steps:

- (i) Studying the structure of the asymptotic solution u_{as} and estimating the difference to the exact solution in the fitted norm.
- (ii) Fitting the finite-element space V_h according to the structure of u_{as} and estimating the approximation error.

A detailed treatment of the asymptotic behaviour and of the construction of asymptotic solutions u_{as} in the maximum norm can be found in [7]. Let us consider the special case

$$(A1) \quad \Omega = (0, 1) \times (0, 1), b = (b_1, b_2) \quad \text{with } b_i(x) > 0, \quad i = 1, 2$$

in which the asymptotic approximation u_{as} consists of the solution u_0 of the reduced problem ($\Gamma_- = \{x \in \Gamma: b(x) \cdot n(x) < 0, n$ outer normal})

$$b(x) \nabla u_0 + c(x) u_0 = f(x) \text{ in } \Omega, \quad u_0 = 0 \text{ on } \Gamma_-,$$

two ordinary boundary layer terms

$$v_1(x) = -u_0(1, x_2) \left[\exp\left(-\frac{b_1(1, x_2)}{\varepsilon}\right) (1 - x_1) - (1 - x_1) \exp\left(-\frac{b_1(1, x_2)}{\varepsilon}\right) \right], \tag{2.4}$$

$$v_2(x) = -u_0(x_1, 1) \left[\exp\left(-\frac{b_2(x_1, 1)}{\varepsilon}\right) (1 - x_2) - (1 - x_2) \exp\left(-\frac{b_2(x_1, 1)}{\varepsilon}\right) \right] \tag{2.5}$$

and a corner layer term

$$v_3(x) = u_0(1, 1) \prod_{i=1}^2 \left[\exp\left(-\frac{b_i(1, 1)}{\varepsilon}\right) (1 - x_i) - (1 - x_i) \exp\left(-\frac{b_i(1, 1)}{\varepsilon}\right) \right]. \tag{2.6}$$

In [22] the estimate

$$\| \|u - u_{as}\| \| \leq C\varepsilon^{1/2} \tag{2.7}$$

was proven where $u_{as} = u_0 + v_1 + v_2 + v_3$. It should be mentioned that the proof of (2.7) is nonstandard because the solution u_0 of the reduced problem belongs only to the space $C(\bar{\Omega}) \cap W^{1,\infty}(\Omega)$.

Now we have to choose the finite-element space V_h in such a way that the interpolation error becomes small. On a square mesh of fineness $h = 1/N$ we add the common bilinear functions defined by

$$\varphi_{ij}(x) = \varphi_i(x_1) \varphi_j(x_2), \quad \varphi_i(t) = \begin{cases} \frac{t}{h} - i + 1, & t \in [(i - 1)h, ih], \\ i + 1 - \frac{t}{h}, & t \in [ih, (i + 1)h], \\ 0, & \text{otherwise.} \end{cases}$$

The following functions fitting the boundary layer terms (2.4)–(2.6):

$$w_j^1(x) = \begin{cases} \varphi_j(x_2) \left[\exp\left(-\frac{b_1(1, jh)}{\varepsilon} (1 - x_1)\right) - \varphi_{N-1}(x_1) \exp\left(-\frac{b_1(1, jh) h}{\varepsilon}\right) - \varphi_N(x_1) \right], & x_1 \in [1 - h, 1], \\ 0, & \text{otherwise,} \end{cases}$$

$$w_j^2(x) = \begin{cases} \varphi_j(x_1) \left[\exp\left(-\frac{b_2(jh, 1)}{\varepsilon} (1 - x_2)\right) - \varphi_{N-1}(x_2) \exp\left(-\frac{b_2(jh, 1) h}{\varepsilon}\right) - \varphi_N(x_2) \right], & x_2 \in [1 - h, 1], \\ 0, & \text{otherwise,} \end{cases}$$

$$w^3(x) = \begin{cases} \prod_{i=1}^2 \left[\exp\left(-\frac{b_i(1, 1)}{\varepsilon} (1 - x_i)\right) - \varphi_{N-1}(x_i) \exp\left(-\frac{b_i(1, 1) h}{\varepsilon}\right) - \varphi_N(x_i) \right], & x \in [1 - h, 1] \times [1 - h, 1], \\ 0, & \text{otherwise.} \end{cases}$$

Now, the asymptotically fitted finite-element method is characterized by (2.2) with $V_h = \text{span}(\varphi_{ij}, w_j^1, w_j^2, w^3; i, j = 1, \dots, N - 1)$. In [22] the interpolation error was estimated by

$$\inf_{w_h \in V_h} \|u_{as} - w_h\| \leq C(\varepsilon + h^{1/2}) + C(m) (\varepsilon/h)^m \quad (2.8)$$

with $m > 0$, arbitrary. The value 1/2 of the h -exponent is due to the fact that u_0 is not sufficiently regular.

Theorem 1: *Let \tilde{u}_h be the solution of the discrete problem (2.2) without fitting if $\varepsilon \geq ch^{1/2}$ and with the above mentioned fitting if $\varepsilon \leq ch^{1/2}$. Then, under the assumption (A1) we have the error estimate*

$$\|u - \tilde{u}_h\| = Ch^{1/4} \quad (2.9)$$

uniformly with respect to ε .

Proof: Combining the estimates (2.3), (2.7), and (2.8) we obtain (2.9) ■

In [22] asymptotically fitted finite-element methods for some other cases have been derived and error estimates have been given. The main advantage of asymptotically fitted finite-element methods consists of the favourable approximation properties within the layers which is based on the relatively large amount of analytical a-priori knowledge. Thus, this method can be used if some information on the position and the structure of boundary layers are known. In numerical experiments it was demonstrated that the method is stable and does not produce any oscillations in the solution.

3. Hybrid upwind finite-element methods

The main objective of hybrid upwind finite-element methods is to conserve the inverse monotonicity of the continuous problem which can be established if

$$(A2) \quad c(x) \geq c_0 \geq 0 \text{ in } \Omega$$

is fulfilled. As a consequence stability of the discrete problem in L^∞ -norms can be shown.

Let \mathfrak{T}_h be an admissible, regular triangulation of weakly acute type, i.e. all interior angles are smaller or equal to $\pi/2$, and let $\{B_i\}$ be the set of nodal points. On each triangle K we use linear functions, i.e. we set $V_h = \{v_h \in C(\bar{\Omega}) : v_h|_K \in P_1(K), v_h|_r = 0\}$. Let $\{\Phi_i\}$ be a basis of V_h with $\Phi_i(B_j) = \delta_{ij}$. Then, the discretization of $\varepsilon(\nabla u, \nabla v)$ corresponds to an M -matrix but in general we have positive outer diagonal elements from $(b\nabla u + cu, v)$. In order to modify the standard finite-element method we consider a secondary decomposition of $\bar{\Omega}$ into dual domains $\{D_i\}$ where D_i denotes that polygon whose vertices are circumcentres of triangles surrounding B_i . Furthermore, let A_i be the set of indices j such that B_j and B_i are neighbour nodes, B_{ij} the midpoint of the side $B_i B_j$ and Γ_{ij} the side of ∂D_i passing B_{ij} . We also need the characteristic function Φ_i of the domain D_i and introduce the lumping operator from $C(\bar{\Omega})$ into $L^2(\Omega)$ defined by $v \rightarrow \bar{v}_h = \sum_i v(B_i) \bar{\Phi}_i$. We denote by n_{ij} the unit outer normal vector on the part Γ_{ij} of the boundary ∂D_i and by β_{ij} an approximation of $\int b \cdot n_{ij} ds$. Now, the hybrid upwind finite-element method is given as follows:

Find $u_h \in V_h$ such that for all $v_h \in V_h$

$$B_h(u_h, v_h) \equiv \varepsilon(\nabla u_h, \nabla v_h) + b_h(u_h, v_h) + (c\bar{u}_h, \bar{v}_h) = (f, \bar{v}_h), \tag{3.1}$$

where

$$b_h(u_h, v_h) = \sum_i \psi(P_i) \sum_{j \in A_i} \beta_{ij} (\lambda_{ij} - 1) [u_h(B_i) - u_h(B_j)], \quad \lambda_{ij} = \begin{cases} 1, & \beta_{ij} \geq 0, \\ 0, & \beta_{ij} < 0. \end{cases}$$

Using the above mentioned basis $\{\Phi_i\}$ it is easy to see that the matrices corresponding to $b_h(u_h, v_h)$ and $(c\bar{u}_h, \bar{v}_h)$ have nonpositive outer diagonal elements and, by means of a chain property, we can establish that the system matrix of (3.1) is an M -matrix. Consequently, the discrete problem is inverse monotone.

In order to formulate the convergence properties of the method we use the same norms as in Section 2. Let β_{ij} be calculated by the midpoint rule applied to $\int b \cdot n_{ij} ds$.

Then we obtain the V_h -ellipticity of B_h and, by a modification of the first lemma of Strang, the estimate

$$\|u - u_h\|_\varepsilon \leq C \inf_{v_h \in V_h} (\|u - v_h\|) + \sup_{w_h \in V_h} \frac{B(v_h, w_h) - B_h(v_h, w_h)}{\|w_h\|_\varepsilon} + \sup_{w_h \in V_h} \frac{(f, w_h - \bar{w}_h)}{\|w_h\|_\varepsilon}. \tag{3.2}$$

On the basis of this the estimates

$$\|u - u_h\|_\varepsilon \leq \begin{cases} Ch\varepsilon^{-1/2} |u|_{2,2} & \text{for an arbitrary mesh,} \\ Ch |u|_{2,2} & \text{for a regular mesh} \end{cases} \tag{3.3}$$

were proven in [21]. Here a mesh is called regular if it is built from three families of parallels. Using a lemma from Stampaccia, it is also possible to obtain the L^∞ -estimates

$$\|u - u_h\|_{0,\infty} \leq \begin{cases} C(\sigma) h^\sigma \varepsilon^{-1/2} |u|_{2,\infty} & \text{for an arbitrary mesh,} \\ C(\sigma) h^\sigma |u|_{2,\infty} & \text{for a regular mesh} \end{cases} \tag{3.4}$$

with $\sigma \in (0, 1)$, arbitrary. Because the seminorms $|u|_{2,2}$ and $|u|_{2,\infty}$ are not uniformly bounded with respect to ε , the estimates (3.3) and (3.4) become meaningless if ε tends to zero. However, RISCHE [21] was successful in proving local estimates in subdomains where no boundary layers occur (called global domain in the notation of [7]). For the special case

(A3) $\Omega = (0, 1) \times (0, 1)$, $b = (b_1, 0)$ with $b_1 > 0$, \mathfrak{T}_h a regular mesh,

we have boundary layers of ordinary and parabolic type near $x_1 = 1$, $x_2 = 0$, $x_2 = 1$ and corner layers of different type in the vertexes of Ω .

Theorem 2: *Let the assumptions (A2), (A3) be fulfilled and the distances d , d_1 and d_2 be defined by $d = \max(\varepsilon, h)$, $d_1 = O(d^{1/2} |\ln d|)$, $d_2 = O(d |\ln d|)$. Then, in the global domain $\Omega^* = [0, 1 - d_2] \times [d_1, 1 - d_1]$ we have the local estimates*

$$\|u - u_h\|_{\varepsilon, \Omega^*} \leq Ch, \quad \|u - u_h\|_{0, \infty, \Omega^*} \leq C(\sigma) h^\sigma, \quad \sigma \in (0, 1) \text{ arbitrary.}$$

For the proof see the more complex cases considered in [21] ■

The main advantages of the hybrid upwind finite-element methods consist in the good stability properties (inverse monotonicity of the discrete problem) combined with localization properties which do not require any a-priori knowledge on the position of boundary layers. We mention that in [21] also the case of a system of equations of the form (1.1) was studied.

4. Streamline diffusion method

The mathematical foundation of the streamline diffusion method was given by NÄVERT [19]. The method combines high order of convergence with good stability properties. In order to sketch the procedure we start with the variational form of (1.1):

Find $u \in H_0^1(\Omega)$ such that for all $v \in H_0^1(\Omega)$

$$B(u, v) \equiv \varepsilon(\nabla u, \nabla v) + (b\nabla u + cu, v) = (f, v).$$

Provided that the exact solution belongs to $H^2(\Omega)$, for all $v \in H_0^1(\Omega)$ the relation $(-\varepsilon\Delta u + b\nabla u + cu, b\nabla v) = (f, b\nabla v)$ is satisfied such that each solution $u \in H_0^1(\Omega) \cap H^2(\Omega)$ of (2.1) fulfils

$$B_\delta(u, v) \equiv (-\varepsilon\delta\Delta u, b \cdot \nabla v) + \varepsilon(\nabla u, \nabla v) + (b\nabla u + cu, v + \delta b\nabla v) = (f, v + \delta b\nabla v) \quad (4.1)$$

for all $v \in H_0^1(\Omega)$. We use piecewise polynomials of degree k , that means we set $V_h = \{v_h \in C(\bar{\Omega}) : v_h|_K \in P_k(K), v_h|_\Gamma = 0\}$. Then, the following discrete problem is derived from (4.1):

Find $u_h \in V_h$ such that for all $v_h \in V_h$

$$B_h(u_h, v_h) = (f, v_h + \delta b \cdot \nabla v_h), \quad (4.2)$$

where the bilinear form B_h is defined by

$$B_h(u_h, v_h) \equiv -\varepsilon\delta \sum_K (\Delta u_h, b \cdot \nabla v_h)_K + \varepsilon(\nabla u_h, \nabla v_h) + (b\nabla u_h + cu_h, v_h + \delta b \cdot \nabla v_h).$$

We remark that for $\delta = 0$ (4.2) is equal to the standard Galerkin finite-element method.

We introduce a fitted norm in $H_0^1(\Omega)$ defined by $\|u\| = (\varepsilon \|u\|_{1,2}^2 + \delta \|b \cdot \nabla u\|_{0,2}^2 + \alpha \|u\|_{0,2}^2)^{1/2}$ where α fulfils $(c - 2^{-1} \operatorname{div} b)(x) \geq \alpha > 0$ for all $x \in \Omega$. By means of inverse inequalities for estimating of $\sum_K (\Delta u_h, b \cdot \nabla u_h)_K$, we can show the H_0^1 -ellipticity of B_h in the fitted norm provided δ is sufficiently small. To be more specific, we

have for $0 \leq \delta \leq \min(\delta_0, c_1 h^2 \varepsilon^{-1})$ the estimate

$$B_h(u_h, u_h) \geq C_2 \|u_h\|^2 \quad \text{for all } u_h \in H_0^1(\Omega), C_2 > 0. \tag{4.3}$$

Taking into consideration that

$$|(f, v_h + \delta b \cdot \nabla v_h)| \leq \|f\|_{0,2} (\|v_h\|_{0,2} + \delta \|b \cdot \nabla v_h\|_{0,2}) \leq C_3 \|f\|_{0,2} \|v_h\|$$

we obtain from (4.3) a stability result in the form $\|u_h\| \leq C \|f\|_{0,2}$.

The good stability results observed in numerical test problems are due to the term $\delta \|b \cdot \nabla u\|_{0,2}^2$ contained in the fitted norm defined above. As for as the inverse monotonicity of the discrete problem is concerned it is easy to see that the system matrix corresponding to (4.2) in general is not an M -matrix. For example, in the particular case that $b = \text{const}$, $c = 0$, $\Omega = (0, 1) \times (0, 1)$, piecewise linear elements are used and the triangulation is of Friedrichs-Keller type, the nonnegativity of the outer diagonal elements can not be fulfilled for sufficiently small ε compared with h . Moreover, numerical test examples permit the conclusion that the discrete problem is in fact not inverse monoton for ε tending to zero.

By carefully handling in majorizing the bilinear form B_h Nävert was able to prove error estimates.

Theorem 3: *Let $\varepsilon < h$ and $\delta = c_1 h$. Then we have the error estimate*

$$\|u - u_h\| = Ch^{k+1/2} |u|_{k+1,2} \tag{4.4}$$

for piecewise polynomials of degree k .

However, the seminorm $|u|_{k+1,2}$ on the right-hand side of (4.4) is not uniformly bounded with respect to ε . Therefore, local estimates in domains without boundary layers are more important.

For the case (A3) already studied in Theorem 2 in frame of hybrid upwind finite-element methods, the results of [19] yield the estimate $\|u - u_h\|_{\Omega^*} \leq Ch^{k+1/2}$, $\varepsilon < h$.

Comparing the streamline diffusion method with the asymptotically fitted method and with the hybrid upwind method we observe that it represents an intermediate position in some sense. The asymptotically fitted methods guarantee ε -uniform convergence but stability is only obtained in a rather weak sense. The streamline diffusion method on the one hand yields better stability properties but on the other hand it does not guarantee ε -uniform convergence. However, the local estimates of high order of convergence show the capability of the method. Finally, the hybrid upwind finite-element method gives stability in L^∞ -norms and preserves the inverse monotonicity of the problem to be solved. But, in contrary to the streamline diffusion method, it only works for linear elements.

5. Application to the stationary Navier-Stokes equations

The stationary Navier-Stokes problem consists in determining the velocity u and the pressure p of a fluid which are solutions of the system of equations ($\Omega \subset \mathbb{R}^2$, $\Gamma = \partial\Omega$)

$$-\varepsilon \Delta u + (u \nabla) u + \nabla p = f \quad \text{and} \quad \text{div } u = 0 \quad \text{in } \Omega, \quad u = 0 \quad \text{on } \Gamma.$$

Multiplying these equations by functions belonging to $V = H_0^1(\Omega)^2$ and $Q = L_0^2(\Omega) = \left\{ q \in L^2(\Omega) : \int_{\Omega} q dx = 0 \right\}$, respectively, and integrating by parts we obtain the variational formulation in the primitive variables:

Find $(u, p) \in V \times Q$ such that

$$\begin{aligned} \varepsilon a(u, v) + b(u, u, v) - (p, \operatorname{div} v) &= (f, v) & \forall v \in V, \\ (q, \operatorname{div} u) &= 0 & \forall q \in Q, \end{aligned} \quad (5.1)$$

where we used the notations

$$a(u, v) = \int_{\Omega} \nabla u \nabla v \, dx, \quad b(u, v, w) = \frac{1}{2} \int_{\Omega} (u \nabla v w - u \nabla w v) \, dx.$$

It is well known that the problem (5.1) admits at least one solution and that this solution is unique provided that $\varepsilon > \varepsilon_0(\|\cdot\|_{0,2}, \Omega)$ [6]. It should be mentioned that stability problems which are similar to the case of convection-diffusion equations arise for small ε (large Reynold numbers) already before the bound ε_0 is attained.

Because of the lack of sufficient a-priori knowledge concerning the asymptotic behaviour of the solutions of (5.1) for small values of ε , we shall not discuss asymptotically fitted methods in this section.

At first we consider a finite-element method of hybrid upwind type. To this end we start with a pair of finite-element spaces V_h, Q_h satisfying the discrete version of the Ladyzhenskaya-Babuska-Brezzi condition

$$\sup_{v_h \in V_h} (q_h, \operatorname{div} v_h) / \|v_h\|_{1,2} \geq \beta \|q_h\|_{0,2} \quad \text{for all } q_h \in Q_h, \beta > 0$$

which is very important for deriving convergence results. Let Ω be a polygon divided into triangles K . We denote by B_j the midpoints of edges and define the finite-element spaces by

$$V_h = \{v_h: v_h|_K \in P_1(K)^2, v_h \text{ cont. in } B_j, v_h(B_j) = 0 \text{ if } B_j \in \Gamma\}, \quad (5.2)$$

$$Q_h = \{q_h \in Q: q_h|_K \in P_0(K)\}. \quad (5.3)$$

Because the discrete velocity space V_h consists of piecewise linear functions which on the edges are continuous only in the midpoints, V_h is not contained in K (nonconforming finite-element methods) and we have to extend the bi- and trilinear form in (5.1). This can be done in a natural way by an elementwise calculation of the corresponding integrals. Let us introduce the meshdependent norm

$$\|u\|_h = \left(\sum_K \int_K (\nabla u)^2 \, dx \right)^{1/2} \quad \text{for } u \in V + V_h.$$

The standard finite-element method studied in [4, 27] reads as follows:

Find $(u_h, p_h) \in V_h \times Q_h$ such that

$$\begin{aligned} \varepsilon a(u_h, v_h) + b(u_h, u_h, v_h) - (p_h, \operatorname{div} v_h) &= (f, v_h) & \forall v_h \in V_h, \\ (q_h, \operatorname{div} u_h) &= 0 & \forall q_h \in Q_h. \end{aligned} \quad (5.4)$$

It converges of order one. Applying a hybrid upwind method to (5.1), we have only to change the discretization of the convective term with the aim of a better reflection of the dominate influence of the convective term for small values of ε . For this, following the idea of OHMORI and USHIJIMA [20] we define a secondary decomposition of Ω into domains D_i . Each inner node $B_i, i = 1, \dots, N$, corresponds to a dual domain D_i which is defined by the barycenters S_1, S_2 of the neighbouring triangles K_1, K_2 of B_i . Furthermore, let A_i be the set of indices j such that B_i and B_j are neighbour nodes and n_{ij} be the unit outer normal vector with respect to D_i along the part Γ_{ij}

of ∂D_i between B_i and B_j . Now, we can derive the following upwind discretization b_h of the trilinear form b :

$$b_h(u, v, w) = \sum_{i=1}^N \sum_{j \in A_i} \int_{\Gamma_{ij}} u n_{ij} ds (1 - \lambda_{ij}(u)) [v(B_j) - v(B_i)] w(B_i), \tag{5.5}$$

where $\lambda_{ij}(u)$ depends on the flux through Γ_{ij} according to

$$\lambda_{ij}(u) = \begin{cases} 1 & \text{if } \int_{\Gamma_{ij}} u n_{ij} ds \geq 0 \\ 0 & \text{otherwise.} \end{cases}$$

Our hybrid upwind finite-element method for solving the stationary Navier-Stokes problem (5.1) is characterized as follows:

Find $(u_h, p_h) \in V_h \times Q_h$ such that

$$\begin{aligned} \varepsilon a(u_h, v_h) + b_h(u_h, u_h, v_h) - (p_h, \operatorname{div} v_h) &= (f, v_h), & \forall v_h \in V_h, \\ (q_h, \operatorname{div} u_h) &= 0 & \forall q_h \in Q_h. \end{aligned} \tag{5.6}$$

Theorem 4: *The discrete problem (5.6) admits at least one solution $(u_h, p_h) \in V_h \times Q_h$ which is unique, provided $\varepsilon > \varepsilon_0(h)$ where $\varepsilon_0(h) \rightarrow \varepsilon_0$ as $h \rightarrow 0$. Moreover, for $\varepsilon > \varepsilon_0(h)$ the error estimate*

$$\|u - u_h\|_h + \|p - p_h\|_{0,2} \leq C(\sigma) h^{1-\sigma}, \quad \sigma > 0 \text{ arbitrary,}$$

is satisfied if $(u, p) \in H^2(\Omega)^2 \times H^1(\Omega)$.

The crucial point in proving Theorem 4 is that the function u in (5.5) (contrary to the function b in (1.1)) is not sufficiently regular since it only belongs to the space V_h . For details we refer to [23], numerical test examples can be found in [24]. A further advantage of the proposed methods consists in a favourable property of the linear systems of equations if they are generated by a fixed point procedure. Namely, let $\{(\Phi_i, 0)\}, \{(0, \Phi_i)\}$ be a basis of V_h defined by $\Phi_i(B_j) = \delta_{ij}$. Then, for fixed $z \in V_h$ the matrix corresponding to $\varepsilon a(u_h, v_h) + b_h(z, u_h, v_h)$ is an M -matrix, provided the mesh is of weakly acute type [23].

Finally, we discuss a nonconforming streamline diffusion method for the stationary Navier-Stokes problem. We will use the same finite-element spaces V_h, Q_h defined by (5.2), (5.3) and start with the standard finite-element method (5.4). Provided that the exact solution (u, p) belongs to the space $H^2(\Omega)^2 \times H^1(\Omega)$ we obtain by testing the relation $-\varepsilon \Delta u + (u \nabla) u + \nabla p = f$ on each element with $\delta u_h \nabla v$

$$-\varepsilon \delta \sum_K (\Delta u, u_h \nabla v)_K + \delta \sum_K (u \nabla u, u_h \nabla v)_K + \delta \sum_K (\nabla p, u_h \nabla v)_K = \delta \sum_K (f, u_h \nabla v)_K.$$

Since we have piecewise linear elements for the velocity and piecewise constant elements for the pressure, the discrete problem reduces to the following form:

Find $(u_h, p_h) \in V_h \times Q_h$ such that

$$\begin{aligned} \varepsilon a(u_h, v_h) + b(u_h, u_h, v_h) + \delta \sum_K (u_h \nabla u_h, u_h \nabla v_h)_K - (p_h, \operatorname{div} v_h) \\ = \sum_K (f, v_h + \delta u_h \nabla v_h)_K \quad \text{for all } v_h \in V_h, \end{aligned} \tag{5.7}$$

$$(q_h, \operatorname{div} u_h) = 0 \quad \text{for all } q_h \in Q_h.$$

The additional term $\delta \sum (u_h \nabla u_h, u_h \nabla v_h)_K$ on the left-hand side of (5.7) has a stabilizing effect on the discrete problem for small ε .

Theorem 5: *Let δ satisfy $0 \leq \delta \leq C_1 h^{1+\sigma}$, $\sigma > 0$ arbitrary, and let the exact solution (u, p) belong to $(W^{1,\infty}(\Omega) \times H^2(\Omega))^2 \times H^1(\Omega)$. Then there are constants ε_0 and h_0 such that for $\varepsilon > \varepsilon_0$ and $h \leq h_0$ the problems (5.1) and (5.7) have unique solutions which satisfy the error estimate*

$$\varepsilon \|u - u_h\|_h^2 + \sum_K \|u_h \cdot \nabla(u - u_h)\|_{0,2,K}^2 \leq Ch^2, \quad \|p - p_h\|_{0,2} \leq Ch.$$

For the proof and further results concerning the case $\varepsilon > \varepsilon_0$ we refer to [18]. ■

Comparing the streamline diffusion method with the hybrid upwind finite-element method we observe that, contrary to the linear case, the strong smoothness assumption $u \in W^{1,\infty}(\Omega)^2$ becomes necessary. First order of convergence can be established whereas the hybrid upwind method converges almost of first order. As far as the stability is concerned both methods achieve a stabilizing effect by the additional terms $b_h(u_h, u_h, u_h)$ and $\delta \sum_K \|u_h \nabla u_h\|_{0,2,K}^2$ which are not identically zero on V_h as in the standard finite-element method.

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