

## Elliptic Oscillation Theory

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Wir gebrauchen eine geeignete Darstellung von Gårdings Ungleichung, um Nichtoszillationskriterien für die allgemeine elliptische Differentialgleichung gerader Ordnung

$$\sum_{|\alpha|, |\beta|=0}^m (-1)^{|\alpha|} D^\alpha [A_{\alpha\beta}(x) D^\beta u] = 0 \quad (x \in \Omega)$$

von bekannten Nichtoszillationskriterien für die häufiger betrachtete Gleichung  $(-1)^m \Delta^m u + h(x) u = 0$  ( $x \in \Omega$ ) abzuleiten, wo  $\Omega$  ein unbeschränktes offenes Gebiet des  $\mathbb{R}^n$  ist.

Употребляя один из подходящих вариантов неравенства Гординга (Gårding) мы покажем, как можно вывести новые теоремы неосцилляции для общих эллиптических уравнений четного порядка

$$\sum_{|\alpha|, |\beta|=0}^m (-1)^{|\alpha|} D^\alpha [A_{\alpha\beta}(x) D^\beta u] = 0 \quad (x \in \Omega)$$

из известных теорем неосцилляции для более часто изучаемого уравнения  $(-1)^m \Delta^m u + h(x) u = 0$  ( $x \in \Omega$ ), где  $\Omega$  является неограниченным открытым подмножеством  $\mathbb{R}^n$ .

Using an appropriate version of Gårding's inequality, we show how to deduce new non-oscillation theorems for the general even-order elliptic equation

$$\sum_{|\alpha|, |\beta|=0}^m (-1)^{|\alpha|} D^\alpha [A_{\alpha\beta}(x) D^\beta u] = 0 \quad (x \in \Omega)$$

from known non-oscillation theorems for the more frequently studied equation  $(-1)^m \Delta^m u + h(x) u = 0$  ( $x \in \Omega$ ), where  $\Omega$  is an unbounded, open subset of  $\mathbb{R}^n$ .

**1. Introduction.** Several writers (see, e.g., [2, 3, 5, 8, 9]), have obtained non-oscillation theorems for various forms of the elliptic partial differential equation

$$(-1)^m \sum_{|\alpha|, |\beta|=m} D^\alpha [a_{\alpha\beta}(x) D^\beta u] + a_0(x) u = 0 \quad (x \in \Omega \subseteq \mathbb{R}^n)$$

in an unbounded open set  $\Omega$ . In a recent paper [6], by using a version of Poincaré's inequality, the author obtained non-oscillation theorems for the more general equation

$$(-1)^m \sum_{|\alpha|, |\beta|=m} D^\alpha [A_{\alpha\beta}(x) D^\beta u] + \sum_{|\alpha| \leq m} B_\alpha(x) D^\alpha u = 0.$$

In the present paper, by using an appropriate version of Gårding's inequality, we will extend the results in [6] to the equation

$$Lu := \sum_{|\alpha|, |\beta|=0}^m (-1)^{|\alpha|} D^\alpha [A_{\alpha\beta}(x) D^\beta u] = 0 \quad (x \in \Omega \subseteq \mathbb{R}^n); \quad (1)$$

where the coefficient functions  $A_{\alpha\beta}$  are real-valued and sufficiently smooth. (The multi-index notation employed here is the same as in [1].) Our main result is a comparison theorem, whose proof, based on a suitable version of Gårding's inequality, will show that every known non-oscillation theorem for the equation

$$(-1)^m \Delta^m u + h(x) u = 0 \quad (x \in \Omega \subseteq \mathbb{R}^n) \quad (2)$$

gives rise to a new non-oscillation theorem for (1).

**2. Definitions and preliminary results.** Throughout this paper,  $G$  will denote any non-empty, open (possibly unbounded) subset of  $\Omega$ . If  $k$  is any non-negative integer, we define the seminorm  $|\cdot|_{k,G}$ , the *weighted* seminorm  $|\cdot|_{k,G,w}$ , and the norm  $\|\cdot\|_{k,G}$  as follows:

$$|u|_{k,G} = \left[ \sum_{|\alpha|=k} \int_G |D^\alpha u|^2 dx \right]^{1/2}, \quad (3)$$

$$|u|_{k,G,w} = \left[ \sum_{|\alpha|=k} \int_G (k!/\alpha!) |D^\alpha u|^2 dx \right]^{1/2}, \quad (4)$$

$$\|u\|_{k,G} = \left[ \sum_{j=0}^k |u|_{j,G}^2 \right]^{1/2}. \quad (5)$$

The definition of  $|\cdot|_{k,G,w}$  is motivated by the following formula, which is valid for all real-valued  $\Phi$  in  $C_0^\infty(G)$ :

$$\begin{aligned} (-1)^m \int_G \Phi \Delta^m \Phi dx &= (-1)^m \int_G \Phi \left( \sum_{k=1}^n D_k^2 \right)^m \Phi dx \\ &= (-1)^m \int_G \Phi \sum_{|\alpha|=m} (m!/\alpha!) D^\alpha \Phi dx \\ &= (-1)^m \sum_{|\alpha|=m} \int_G \Phi D^\alpha [(m!/\alpha!) D^\alpha \Phi] dx \\ &= \sum_{|\alpha|=m} \int_G (m!/\alpha!) |D^\alpha \Phi|^2 dx. \end{aligned}$$

To compare the seminorms  $|\cdot|_{m,G}$  and  $|\cdot|_{m,G,w}$ , we let

$$c_0 = \max \{m!/\alpha! : |\alpha| = m\}. \quad (6)$$

Then it is easily seen that

$$|u|_{m,G} \leq |u|_{m,G,w} \leq c_0^{1/2} |u|_{m,G}. \quad (7)$$

We also note that, in (3) and (5), when there is no danger of confusion, we omit the subscript  $G$ . Let  $C_B^k(G) = \{u \in C^k(G) : \|u\|_{k,G} < \infty\}$ , and let  $H_k(G)$  and  $H_k^0(G)$  denote the completions of  $C_B^k(G)$  and  $C_0^\infty(G)$ , respectively, with respect to the norm  $\|\cdot\|_{k,G}$ .

If  $G$  is bounded, and if there exists a non-trivial function  $u$  in  $H_m^0(G) \cap C_B^{2m}(G)$  such that (1) holds, then  $G$  is called a *nodal domain* for  $L$  or a nodal domain for (1). If for all positive  $r$  the region  $\{x \in \Omega : |x| > r\}$  contains a nodal domain for  $L$ , then (1) is said to be *nodally oscillatory* (or *strongly oscillatory*) in  $\Omega$ .

Using integration by parts, we can easily show that if  $G$  is any non-empty, open (possibly unbounded) subset of  $\Omega$ , then for every real-valued  $\Phi$  in  $C_0^\infty(G)$  we have

$$\begin{aligned} \int_G \Phi L \Phi dx &= \sum_{|\alpha|=|\beta|=m} \int_G A_{\alpha\beta}(x) D^\alpha \Phi D^\beta \Phi dx + \int_G \Phi^2 A_{00}(x) dx \\ &\quad + \sum_{|\alpha|+|\beta|=1}^{2m-1} \int_G A_{\alpha\beta} D^\alpha \Phi D^\beta \Phi dx. \end{aligned} \quad (8)$$

The standard proof of the global version of Gårding's inequality [1: Theorem 7.6] now yields the following result.

Lemma 2.1: Let  $E_0$  denote the ellipticity constant of the differential operator  $L$ ; in other words, let

$$E_0 = \inf \left\{ \sum_{|\alpha|=|\beta|=m} A_{\alpha\beta}(x) \xi^{\alpha+\beta} |\xi|^{-2m} : 0 \neq \xi \in \mathbb{R}^n, x \in \Omega \right\}.$$

Suppose that the principal coefficients  $A_{\alpha\beta}$  ( $|\alpha| = |\beta| = m$ ) are uniformly continuous on  $\Omega$ , and that the remaining coefficients  $A_{\alpha\beta}$  ( $|\alpha| + |\beta| \leq 2m - 1$ ) are bounded and measurable on  $\Omega$ . Let  $G$  be any non-empty, open subset of  $\Omega$ . Then there exist constants  $c_1 \in (0, \infty)$  and  $c_2 \in [0, \infty)$  such that, for every real-valued  $\Phi$  in  $C_0^\infty(G)$ , we have

$$\sum_{|\alpha|=|\beta|=m} \int_G A_{\alpha\beta}(x) D^\alpha \Phi D^\beta \Phi dx + \sum_{|\alpha|+|\beta|=1}^{2m-1} \int_G A_{\alpha\beta}(x) D^\alpha \Phi D^\beta \Phi dx \geq c_1 E_0 \|\Phi\|_{m,G}^2 - c_2 \|\Phi\|_{0,G}^2.$$

The constant  $c_1$  depends only on  $m$  and  $n$ ; the constant  $c_2$  depends only on  $m$ ,  $n$ ,  $E_0$ ,  $\sup \{|A_{\alpha\beta}(x)| : x \in \Omega; 1 \leq |\alpha| + |\beta| \leq 2m - 1\}$  and the modulus of continuity for the principal coefficients.

3. The main results. Using Lemma 2.1, we will first obtain a comparison theorem, which we can then employ to obtain new non-oscillation theorems for (1) from all known non-oscillation theorems for (2).

Theorem 3.1: Let  $M$  be the differential operator defined by

$$Mu = (-1)^m c_4 \Delta^m u + [A_{00}(x) - c_2] u, \tag{9}$$

where

$$c_4 = c_1 E_0 / c_0, \tag{10}$$

and  $c_0$  is defined by (6). If (1) is nodally oscillatory in  $\Omega$ , then the differential equation

$$Mu = 0 \tag{11}$$

is also nodally oscillatory in  $\Omega$ .

Proof: If (1) is nodally oscillatory in  $\Omega$ , then for every positive number  $r$  the region  $\{x \in \Omega : |x| > r\}$  contains a nodal domain  $G$  for the differential operator  $L$ . Thus, there exists a non-trivial function  $u$  in  $H_m^0(G) \cap C^{2m}(G)$  such that (1) holds. Furthermore, (8), Lemma 2.1, (9), integration by parts, (4), (7) and (10) imply that, for every  $\Phi$  in  $C_0^\infty(G)$ , we have

$$\begin{aligned} & \int_G \Phi L \Phi dx - \int_G \Phi M \Phi dx \\ & \geq [c_1 E_0 \|\Phi\|_{m,G}^2 - c_2 \|\Phi\|_{0,G}^2 + A_{00} \|\Phi\|_{0,G}^2] - [c_4 \|\Phi\|_{m,G,w}^2 + (A_{00} - c_2) \|\Phi\|_{0,G}^2] \\ & = c_1 E_0 \|\Phi\|_{m,G}^2 - c_4 \|\Phi\|_{m,G,w}^2 \geq c_1 E_0 \|\Phi\|_{m,G}^2 - c_4 \|\Phi\|_{m,G,w}^2 \\ & \geq [(c_1 E_0 / c_0) - c_4] \|\Phi\|_{m,G,w}^2 = 0. \end{aligned} \tag{12}$$

Using (1), (12) and a continuity argument, we obtain  $0 = \int_G u Lu dx \geq \int_G u Mu dx$ .

Therefore, the smallest eigenvalue of the eigenvalue problem  $Mv = \lambda v$ ,  $v \in H_m^0(G) \cap C^{2m}(G)$  is non-positive. Hence, we can apply a known monotonicity principle [4] to show that  $G$  has a non-empty open subset  $G'$  such that zero is the smallest eigenvalue of the eigenvalue problem  $Mw = \mu w$ ,  $w \in H_m^0(G') \cap C^{2m}(G')$ . Thus, we have shown that, for every positive number  $r$ , the equation  $Mw = 0$  has a non-trivial solution  $w$ , with a nodal domain  $G' \subset G \subset \{x \in \Omega : |x| > r\}$ . ■

To illustrate how Theorem 3.1 may be employed to obtain new non-oscillation theorems for (1) from known non-oscillation theorems for (2), we now generalize the non-oscillation portion of [8: Theorem 1]. (In [7] we showed how to obtain new oscillation theorems for (1) from known oscillation theorems for (2).)

**Theorem 3.2:** *Consider the polynomial*

$$\prod_{j=0}^{m-1} [r + [(n - 2m + 4j)/2]^2] =: \sum_{k=0}^m b_k r^k.$$

*If  $n \geq 2m$ , or if  $n < 2m$  and  $n$  is odd, then (1) is nodally non-oscillatory in  $\Omega$  if there exists a positive number  $r_0$  such that for every  $x$  in the region  $\{x \in \Omega: |x| > r_0\}$  we have*

$$[A_{00}(x) - c_2]/c_4 > -|x|^{-2m} \sum_{k=0}^m [(2k - 1)!!] b_k / 4^k \log^{2k} |x|. \quad (13)$$

**Proof:** Suppose to the contrary that (1) is nodally oscillatory in  $\Omega$ . Then it follows from Theorem 3.1 that (11) is nodally oscillatory in  $\Omega$ , contrary to the fact, proved in [8: Theorem 1], that (11) is nodally non-oscillatory in  $\Omega$  whenever (13) holds ■

We invite the reader to formulate appropriate generalizations of other known non-oscillation criteria.

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