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# Multiplicity Results for Classes of Infinite Positone Problems

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Abstract. We study positive solutions to the singular boundary value problem

$$\begin{cases} -\Delta_p u = \lambda \frac{f(u)}{u^{\beta}} & \text{in } \Omega\\ u = 0 & \text{on } \partial\Omega, \end{cases}$$

where  $\Delta_p u = \text{div}(|\nabla u|^{p-2}\nabla u), p > 1, \lambda > 0, \beta \in (0, 1)$  and  $\Omega$  is a bounded domain in  $\mathbb{R}^N, N \ge 1$ . Here  $f: [0, \infty) \to (0, \infty)$  is a continuous nondecreasing function such that  $\lim_{u\to\infty} \frac{f(u)}{u^{\beta+p-1}} = 0$ . We establish the existence of multiple positive solutions for certain range of  $\lambda$  when f satisfies certain additional assumptions. A simple model that will satisfy our hypotheses is  $f(u) = e^{\frac{\alpha u}{\alpha+u}}$  for  $\alpha \gg 1$ . We also extend our results to classes of systems when the nonlinearities satisfy a combined sublinear condition at infinity. We prove our results by the method of sub-supersolutions.

**Keywords.** Singular boundary value problems, infinite positone problems, multiplicity of positive solutions, sub-supersolutions

Mathematics Subject Classification (2000). 35J25, 35J55

## 1. Introduction

We first consider the boundary value problem

$$\begin{cases} -\Delta_p u = \lambda \frac{f(u)}{u^{\beta}} & \text{in } \Omega\\ u = 0 & \text{on } \partial\Omega, \end{cases}$$
(1)

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where  $\Delta_p u = \operatorname{div}(|\nabla u|^{p-2}\nabla u)$  is the *p*-Laplacian of  $u, p > 1, \beta \in (0, 1), \lambda$  is a positive parameter and  $\Omega$  is a bounded domain with a smooth boundary in  $\mathbb{R}^N, N \geq 1$ . We assume that f is a  $C([0, \infty))$ -function satisfying the following assumptions:

- (H1) f(u) > 0 for all  $u \ge 0$ ,
- (H2)  $\lim_{u\to\infty} \frac{f(u)}{u^{\beta+p-1}} = 0.$

We note that  $\lim_{u\to 0} \frac{f(u)}{u^{\beta}} = \infty$ , and hence (1) is a singular boundary value problem which we call here as an infinite positone problem. Our results in this paper are motivated by the problem:

$$\begin{cases} -\Delta_p u = \lambda \frac{\exp[\frac{\alpha u}{\alpha + u}]}{u^{\beta}} & \text{in } \Omega\\ u = 0 & \text{on } \partial\Omega. \end{cases}$$
(2)

When  $\beta = 0$  for every  $\alpha > 0$  and  $\lambda > 0$  it is known that there exist a positive solution and when  $\alpha \gg 1$  there exists a range of  $\lambda$  for which there exist at least three positive solutions [8]. In this paper, we extend this study to the singular case when  $0 < \beta < 1$ . In particular, we establish the existence of a positive solution for all  $\alpha > 0$  and for all  $\lambda > 0$  and a multiplicity result for certain range of  $\lambda$  when  $\alpha \gg 1$ . However, our multiplicity result is restricted two positive solutions. In [1], the author studied this singular problem (2) when p = 2 by treating it as a limit problem of the class of non-singular problems defined by  $-\Delta u_{\epsilon} = \lambda \frac{e^{\frac{\alpha}{\alpha+u}}}{(u+\epsilon)^{\beta}}$  in  $\Omega$  and  $u_{\epsilon} = 0$  on  $\partial\Omega$ . Here we establish the our results for all p > 1 directly by method of sub- and supersolutions associated with such singular problems. Also our proofs easily extend to classes of system where the nonlinearities satisfy a combined sublinear condition at infinity.

By a subsolution of (1) we mean a function  $\psi : \overline{\Omega} \to \mathbb{R}$  such that  $\psi \in W^{1,p}(\Omega) \bigcap C(\overline{\Omega})$  and satisfies:

$$-\Delta_p \psi \le \lambda \frac{f(\psi)}{\psi^{\beta}} \quad \text{in } \Omega$$
$$\psi > 0 \qquad \text{in } \Omega$$
$$\psi = 0 \qquad \text{on } \partial\Omega$$

and by a supersolution of (1) we mean a function  $\phi : \overline{\Omega} \to \mathbb{R}$  such that  $\phi \in W^{1,p}(\Omega) \bigcap C(\overline{\Omega})$  and satisfies:

$$-\Delta_p \phi \ge \lambda \frac{f(\phi)}{\phi^\beta} \quad \text{in } \Omega$$
$$\phi > 0 \qquad \text{in } \Omega$$
$$\phi = 0 \qquad \text{on } \partial\Omega$$

Then we have the following lemma:

**Lemma 1.1** (See [2, 5, 9]). If there exist a subsolution  $\psi$  and a supersolution  $\phi$  of (1) such that  $\psi \leq \phi$  on  $\overline{\Omega}$ , then (1) has at least one solution  $u \in W^{1,p}(\Omega) \bigcap C(\overline{\Omega})$  satisfying  $\psi \leq u \leq \phi$  on  $\overline{\Omega}$ .

We first establish:

**Theorem 1.2.** Assume (H1) – (H2). Then (1) has a positive solution for all  $\lambda > 0$ .

We refer to [7] for a more general existence result for (1). However, for certain classes of f we can get at least two positive solutions for certain range of  $\lambda$ . To state this multiplicity result, for any 0 < a < d we define

$$Q(a,d) := \frac{a^{\beta+p-1}}{f(a)} \frac{f(d)}{d^{\beta+p-1}}.$$

Further, let

$$A := \left(\frac{(N+p-1)^{N+p-1}}{N^N}\right)^{\frac{1}{p-1}}$$

Throughout this paper,  $w \in W^{1,p}(\Omega) \bigcap C(\overline{\Omega})$  (see [4, Lemma 3.1]) is the unique solution of

$$\begin{cases} -\Delta_p w = \frac{1}{w^{\beta}} & \text{in } \Omega\\ w = 0 & \text{on } \partial\Omega. \end{cases}$$
(3)

We now assume that f further satisfies:

- (H3) f(u) is nondecreasing for all  $u \ge 0$
- (H4) There exist a and b such that  $0 < a < \frac{p}{A}b$  and  $\frac{f(u)}{u^{\beta}}$  is nondecreasing on (a, b).

We establish:

**Theorem 1.3.** Assume (H1) - (H4). Further assume that there exists d such that

$$a < d < \frac{p}{A}b$$
 and  $Q(a,d) > \frac{A^{p-1}N \|w\|_{\infty}^{\beta+p-1}}{(p-1)^{p-1}R^p} := C(\beta, N, \Omega),$ 

where R is the radius of the largest inscribed ball  $B_R$  in  $\Omega$ . Then (1) has at least two positive solutions for  $\lambda_* < \lambda < \lambda^*$ , where

$$\lambda_* = \frac{d^{\beta+p-1}}{f(d)} \frac{A^{p-1}N}{(p-1)^{p-1}R^p} \quad and$$
$$\lambda^* = \min\left\{\frac{d^{\beta}}{f(d)} \frac{N}{R^p} \left(\frac{p}{p-1}\right)^{p-1} b^{p-1}, \ \frac{a^{\beta+p-1}}{f(a)} \frac{1}{\|w\|_{\infty}^{\beta+p-1}}\right\}.$$



Figure 1: Graph of the function  $\frac{f(u)}{u^{\beta}}$ 

**Remark 1.4.** Since  $d < \frac{p}{A}b$ , we have  $\frac{d^{\beta+p-1}}{f(d)} \frac{A^{p-1}N}{(p-1)^{p-1}R^p} < \frac{d^{\beta}}{f(d)} \frac{N}{R^p} (\frac{p}{p-1})^{p-1} b^{p-1}$ and since  $Q(a,d) > \frac{A^{p-1}N ||w||_{\infty}^{\beta+p-1}}{(p-1)^{p-1}R^p}$ , we obtain  $\frac{d^{\beta+p-1}}{f(d)} \frac{A^{p-1}N}{(p-1)^{p-1}R^p} < \frac{a^{\beta+p-1}}{f(a)} \frac{1}{||w||_{\infty}^{\beta+p-1}}$ . Therefore,  $(\lambda_*, \lambda^*)$  is not empty.

**Remark 1.5.** A simple example satisfying the hypotheses of Theorem 1.2 and Theorem 1.3 is

$$\begin{cases} -\Delta_p u = \lambda \frac{e^{\frac{\alpha u}{\alpha + u}}}{u^{\beta}} & \text{in } \Omega\\ u = 0 & \text{on } \partial\Omega. \end{cases}$$

Clearly,  $f(u) := e^{\frac{\alpha u}{\alpha + u}}$  satisfies hypotheses (H1) – (H3). Choosing  $a = 1, d = \alpha$ and  $b = \frac{\alpha^2}{2}$ , we can easily show that  $\frac{f(u)}{u^{\beta}}$  is nondecreasing on (a, b) for  $\alpha \gg 1$ . Further  $Q(a, d) = \frac{a^{\beta + p - 1}}{f(a)} \frac{f(d)}{d^{\beta + p - 1}} = [\frac{1}{\alpha}]^{\beta + p - 1} \exp[\frac{\alpha}{2} - \frac{\alpha}{\alpha + 1}]$  and hence, for any given  $\Omega$ , we have  $a < d < \frac{p}{A}b$  and  $Q(1, \alpha) > C(\beta, N, \Omega)$  for  $\alpha$  large.

Next we note that the method of sub- and supersolutions discussed in Lemma 1.1 extends to the system:

$$\begin{cases} -\Delta_p u = \lambda \frac{f(v)}{u^{\beta}} & \text{in } \Omega \\ -\Delta_p v = \lambda \frac{g(u)}{v^{\beta}} & \text{in } \Omega \\ u = v = 0 & \text{on } \partial\Omega. \end{cases}$$
(4)

This follows by using the result in [5]. For the system (4) by a subsolution we mean a pair of functions  $(\psi, \bar{\psi}) : \overline{\Omega} \to \mathbb{R} \times \mathbb{R}$  such that  $(\psi, \bar{\psi}) \in$   $(W^{1,p}(\Omega) \cap C(\overline{\Omega})) \times (W^{1,p}(\Omega) \cap C(\overline{\Omega}))$  and satisfy

$$-\Delta_p \psi \le \lambda \frac{f(\bar{\psi})}{\psi^{\beta}} \quad \text{in } \Omega$$
$$-\Delta_p \bar{\psi} \le \lambda \frac{g(\psi)}{\bar{\psi}^{\beta}} \quad \text{in } \Omega$$
$$\psi > 0, \ \bar{\psi} > 0 \quad \text{in } \Omega$$
$$\psi = \bar{\psi} = 0 \quad \text{on } \partial\Omega$$

By a supersolution we mean a pair of functions  $(\phi, \overline{\phi}) : \overline{\Omega} \to \mathbb{R} \times \mathbb{R}$  such that  $(\phi, \overline{\phi}) \in (W^{1,p}(\Omega) \cap C(\overline{\Omega})) \times (W^{1,p}(\Omega) \cap C(\overline{\Omega}))$  and satisfy

$$\begin{aligned} -\Delta_p \phi &\geq \lambda \frac{f(\bar{\phi})}{\phi^\beta} & \text{in } \Omega \\ -\Delta_p \bar{\phi} &\geq \lambda \frac{g(\phi)}{\bar{\phi}^\beta} & \text{in } \Omega \\ \phi &> 0, \ \bar{\phi} &> 0 & \text{in } \Omega \\ \phi &= \bar{\phi} &= 0 & \text{on } \partial\Omega. \end{aligned}$$

We now assume that f and g are  $C([0, \infty))$  functions satisfying the following assumptions:

- (G1) f and g are nondecreasing and f(0) > 0 and g(0) > 0
- (G2)  $\lim_{x\to\infty} \frac{f(Mg(x))}{x^{\beta+p-1}} = 0$  for all M > 0 (a combined sublinear condition at infinity).

We establish:

**Theorem 1.6.** Assume (G1) – (G2). Then (4) has a positive solution for all  $\lambda > 0$ .

Next, under certain combined nonlinear effects of  $\frac{x^{\beta+p-1}}{f(x)}$  and  $\frac{x^{\beta+p-1}}{g(x)}$  we study the existence of multiple positive solutions to (4). To state the multiplicity result, for any 0 < a < d we define

$$Q_1(a,d) := \frac{a^{\beta+p-1}}{g(a)} \frac{f(d)}{d^{\beta+p-1}}.$$

We also assume:

(G3)  $f(u) \le g(u)$  for all  $u \ge 0$ 

(G4) There exist a and b with 0 < a < b such that  $a < \frac{p}{A}b$  and  $\frac{f(u)}{u^{\beta}}$  is nondecreasing on (a, b).

We establish:

**Theorem 1.7.** Assume (G1) – (G4). Further assume there exists d such that  $a < d < \frac{p}{A}b$  and  $Q_1(a, d) > C(\beta, N, \Omega)$ , where  $C(\beta, N, \Omega)$  is as defined in Theorem 1.3. Then (4) has at least two positive solutions for  $\lambda_* < \lambda < \lambda^*$ , where

$$\lambda_* = \frac{d^{\beta+p-1}}{f(d)} \frac{A^{p-1}N}{(p-1)^{p-1}R^p} \quad and$$
$$\lambda^* = \min\left\{\frac{d^{\beta}}{f(d)} \frac{N}{R^p} \left(\frac{p}{p-1}\right)^{p-1} b^{p-1}, \ \frac{a^{\beta+p-1}}{g(a)} \frac{1}{\|w\|_{\infty}^{\beta+p-1}}\right\}.$$

**Remark 1.8.** Since  $d < \frac{p}{A}b$ , we have  $\frac{d^{\beta+p-1}}{f(d)} \frac{A^{p-1}N}{(p-1)^{p-1}R^p} < \frac{d^{\beta}}{f(d)} \frac{N}{R^p} (\frac{p}{p-1})^{p-1} b^{p-1}$ and since  $Q_1(a,d) > \frac{A^{p-1}N ||w||_{\infty}^{\beta+p-1}}{(p-1)^{p-1}R^p}$ , we obtain  $\frac{d^{\beta+p-1}}{f(d)} \frac{A^{p-1}N}{(p-1)^{p-1}R^p} < \frac{a^{\beta+p-1}}{g(a)} \frac{1}{||w||_{\infty}^{\beta+p-1}}$ . Therefore,  $(\lambda_*, \lambda^*)$  is not empty.

**Remark 1.9.** A simple example satisfying the hypotheses of Theorem 1.6 and Theorem 1.7 is  $\alpha v$ 

$$-\Delta_p u = \lambda \frac{e^{\overline{\alpha + v}}}{u^{\beta}} \quad \text{in } \Omega$$
$$-\Delta_p v = \lambda \frac{u^q + M}{v^{\beta}} \quad \text{in } \Omega$$
$$u = 0 = v \quad \text{on } \partial\Omega$$

where q > 0 and  $M \gg 1$  so that (G3) is satisfied. Clearly,  $f(u) := e^{\frac{\alpha u}{\alpha + u}}$  and  $g(u) := u^q + M$  satisfy hypotheses (G1) – (G3). Choosing  $a = 1, d = \alpha$  and  $b = \frac{\alpha^2}{2}$ , we can easily show that  $\frac{f(u)}{u^\beta}$  is nondecreasing on (a, b) for  $\alpha \gg 1$ . Further  $Q_1(a, d) = \frac{a^{\beta+p-1}}{g(a)} \frac{f(d)}{d^{\beta+p-1}} = \left(\frac{1}{1+M}\right) \left(\frac{1}{\alpha}\right)^{\beta+p-1} \exp[\frac{\alpha}{2}]$  and hence, for any given  $\Omega$  we have  $a < d < \frac{p}{A}b$  and  $Q_1(1, \alpha) > C(\beta, N, \Omega)$  for  $\alpha$  large.

We will prove Theorems 1.2 and 1.3 in the Section 2 and the Theorem 1.6 and 1.7 in Section 3.

## 2. Proof of Theorem 1.2 and Theorem 1.3

*Proof of Theorem* 1.2. We construct a positive supersolution  $\phi_1$  of (1).

Let  $f^*(u) = \max_{0 \le x \le u} f(x)$ . Then  $f^*(u)$  is nondecreasing and  $\frac{f^*(u)}{u^{\beta+p-1}} \to 0$ as  $u \to \infty$ , since  $\frac{f(u)}{u^{\beta+p-1}} \to 0$  as  $u \to \infty$ . So there exists  $M_{\lambda} \gg 1$  such that

$$\frac{f^*(M_\lambda \|w\|_\infty)}{(M_\lambda \|w\|_\infty)^{\beta+p-1}} \le \frac{1}{\lambda \|w\|_\infty^{\beta+p-1}}.$$

Let  $\phi_1 = M_{\lambda} w$ , where w is defined in (3). We have

$$-\Delta_p \phi_1 = \frac{M_{\lambda}^{p-1}}{w^{\beta}} \ge \lambda \frac{f^*(M_{\lambda} \|w\|_{\infty})}{(M_{\lambda} w)^{\beta}} \ge \lambda \frac{f^*(M_{\lambda} w)}{(M_{\lambda} w)^{\beta}} \ge \lambda \frac{f(M_{\lambda} w)}{(M_{\lambda} w)^{\beta}} = \lambda \frac{f(\phi_1)}{\phi_1^{\beta}},$$

showing that  $\phi_1$  is a positive supersolution of (1).

Now we construct a positive subsolution  $\psi_1$ . Let  $\lambda_1$  be the first eigenvalue of  $-\Delta_p$  with Dirichlet boundary condition and e > 0 be a corresponding eigenfunction. Hence e and  $\lambda_1$  satisfy:

$$\begin{cases} -\Delta_p e = \lambda_1 e^{p-1} & \text{in } \Omega\\ e = 0 & \text{on } \partial \Omega \end{cases}$$

Since  $\frac{f(u)}{u^{\beta}} \to \infty$  as  $u \to 0$ , there exists a sufficiently small  $m_{\lambda}$  such that

$$\lambda_1(m_\lambda e)^{p-1} \le \lambda \frac{f(m_\lambda e)}{(m_\lambda e)^{\beta}}$$
 for all  $\lambda > 0$ .

Let  $\psi_1 = m_{\lambda} e$ . We have  $-\Delta_p \psi_1 = \lambda_1 (m_{\lambda} e)^{p-1} \leq \lambda \frac{f(m_{\lambda} e)}{(m_{\lambda} e)^{\beta}} = \lambda \frac{f(\psi_1)}{\psi_1^{\beta}}$ . Thus  $\psi_1$  is subsolution of (1), and if  $m_{\lambda}$  is chosen sufficiently small, then  $\psi_1 \leq \phi_1$ . Hence, Theorem 1.2 is proven.

Proof of Theorem 1.3. Here we construct a second positive supersolution  $\phi_2$  of (1) with  $\|\phi_2\|_{\infty} = a$  when  $\lambda \leq \frac{a^{\beta+p-1}}{f(a)} \frac{1}{\|w\|_{\infty}^{\beta+p-1}}$ . Let  $\phi_2 = a \frac{w}{\|w\|_{\infty}}$ , where w is defined in (3). Since  $\lambda \leq \frac{a^{\beta+p-1}}{f(a)} \frac{1}{\|w\|_{\infty}^{\beta+p-1}}$ ,

$$-\Delta_p \phi_2 = \frac{a^{p-1}}{\|w\|_{\infty}} \frac{1}{w^{\beta}} = \frac{\|w\|_{\infty}^{\beta}}{a^{\beta} w^{\beta}} \frac{a^{\beta+p-1}}{\|w\|_{\infty}^{\beta+p-1}} \ge \lambda \frac{f(a)}{\phi_2^{\beta}} \ge \lambda \frac{f\left(a\frac{w}{\|w\|_{\infty}}\right)}{\phi_2^{\beta}} = \lambda \frac{f(\phi_2)}{\phi_2^{\beta}}.$$

Next we construct a second positive subsolution  $\psi_2$  of (1) when

$$\frac{d^{\beta+p-1}}{f(d)} \frac{A^{p-1}N}{(p-1)^{p-1}R^p} < \lambda < \frac{d^{\beta}}{f(d)} \frac{N}{R^p} \left(\frac{p}{p-1}\right)^{p-1} b^{p-1}$$

Let  $a^* \in (0, a]$  be such that  $f(a^*) = \min_{0 < x \le a} f(x)$  and define  $h \in C([0, \infty))$  such that

$$h(u) = \begin{cases} \frac{f(a^*)}{(a^*)^{\beta}}, & u \le a^* \\ \frac{f(u)}{u^{\beta}}, & u \ge a, \end{cases}$$

so that h is nondecreasing on (0, a] and  $h(u) \leq \frac{f(u)}{u^{\beta}}$  for all  $u \geq 0$  (See Figure 2).

Consider the following nonsingular problem:

$$\begin{cases} -\Delta_p u = \lambda h(u) & \text{in } \Omega\\ u = 0 & \text{on } \partial\Omega. \end{cases}$$
(5)



Figure 2: Graph of the function h(u) below  $\frac{f(u)}{u^{\beta}}$ 

Let R be the radius of the largest inscribed ball  $B_R$  of  $\Omega$ . For  $0 < \epsilon < R$ , and  $\delta, \mu > 1$ , define  $\rho(r) : [0, R] \to [0, 1]$  by

$$\rho(r) = \begin{cases} 1, & 0 \le r \le \epsilon \\ 1 - \left(1 - \left(\frac{R-r}{R-\epsilon}\right)^{\mu}\right)^{\delta}, & \epsilon < r \le R. \end{cases}$$

Then

$$\rho'(r) = \begin{cases} 0, & 0 \le r \le \epsilon \\ -\frac{\delta\mu}{R-\epsilon} \left(1 - \left(\frac{R-r}{R-\epsilon}\right)^{\mu}\right)^{\delta-1} \left(\frac{R-r}{R-\epsilon}\right)^{\mu-1}, & \epsilon < r \le R. \end{cases}$$

Let  $v(r) = d\rho(r)$ . Here note that  $|v'(r)| \leq \frac{d\delta\mu}{R-\epsilon}$  since  $|\rho'(r)| \leq \frac{\delta\mu}{R-\epsilon}$ . Define  $\psi$  as the radially symmetric solution of

$$\begin{cases} -\Delta_p \psi(x) = \lambda h(v(|x|)) & \text{in } B(0, R) \\ \psi = 0 & \text{on } \partial B(0, R). \end{cases}$$

Then  $\psi$  satisfies

$$\begin{cases} -(r^{N-1}G(\psi'(r)))' = \lambda r^{N-1}h(v(r))\\ \psi'(0) = 0, \quad \psi(R) = 0, \end{cases}$$

where  $G(t) = |t|^{p-2}t$  for all  $t \in \mathbb{R}$ . Integrating once, for 0 < r < R, we get

$$-G(\psi'(r)) = \frac{\lambda}{r^{N-1}} \int_0^r s^{N-1} h(v(s)) \, ds.$$
(6)

Since G is monotone,  $G^{-1}$  is also continuous and monotone. Hence, we have

$$-\psi'(r) = G^{-1}\left(\frac{\lambda}{r^{N-1}} \int_0^r s^{N-1} h(v(s)) \, ds\right). \tag{7}$$

We claim that

$$\psi(r) \ge v(r), \quad \forall \ 0 \le r \le R \tag{8}$$

and

$$\|\psi\|_{\infty} \le b,\tag{9}$$

when  $\frac{d^{\beta+p-1}}{f(d)} \frac{A^{p-1}N}{(p-1)^{p-1}R^p} < \lambda < \frac{d^{\beta}}{f(d)} \frac{N}{R^p} (\frac{p}{p-1})^{p-1} b^{p-1}$ . If our claim is true,  $\psi$  is a positive subsolution of the nonsingular problem (5) since  $-\Delta_p \psi = \lambda h(v) \leq \lambda h(\psi)$ . In order to show (8), since  $\psi(R) = v(R) = 0$ , it is enough to show that

$$\psi'(r) \le v'(r), \quad \forall \ 0 \le r \le R.$$
(10)

Note that for  $0 \le r \le \epsilon$ , clearly  $\psi'(r) \le 0 = v'(r)$ . Now for  $r > \epsilon$ , from (6)

$$\begin{split} -G(\psi'(r)) &= \frac{\lambda}{r^{N-1}} \int_0^r s^{N-1} h(v(s)) \, ds \\ &> \frac{\lambda}{R^{N-1}} \int_0^\epsilon s^{N-1} h(v(s)) \, ds \\ &= \frac{\lambda}{R^{N-1}} h(d) \frac{\epsilon^N}{N} \\ &= \frac{\lambda}{R^{N-1}} \frac{f(d)}{d^\beta} \frac{\epsilon^N}{N}. \end{split}$$

So, we have  $-\psi'(r) > G^{-1}\left(\frac{\lambda}{R^{N-1}}\frac{f(d)}{d^{\beta}}\frac{\epsilon^{N}}{N}\right)$ . Thus, (10) will hold for all  $\epsilon \ge r \ge R$ , if  $G^{-1}\left(\frac{\lambda}{R^{N-1}}\frac{f(d)}{d^{\beta}}\frac{\epsilon^{N}}{N}\right) \ge \frac{\delta\mu}{R-\epsilon}d$ , which is same as

$$\frac{\lambda}{R^{N-1}} \frac{f(d)}{d^{\beta}} \frac{\epsilon^N}{N} \ge G\left(\frac{\delta\mu}{R-\epsilon}d\right) = \left(\frac{\delta\mu}{R-\epsilon}d\right)^{p-1}.$$

Thus, if  $\lambda \geq \frac{d^{\beta+p-1}}{f(d)} \frac{NR^{N-1}(\delta\mu)^{p-1}}{\epsilon^{N}(R-\epsilon)^{p-1}}$ , inequality (10) will hold for all  $\epsilon \leq r \leq R$ . Note that

$$\inf \frac{d^{\beta+p-1}}{f(d)} \frac{NR^{N-1}(\delta\mu)^{p-1}}{\epsilon^N (R-\epsilon)^{p-1}} = \frac{d^{\beta+p-1}}{f(d)} \frac{A^{p-1}N}{(p-1)^{p-1}R^p} (\delta\mu)^{p-1}$$

and is achieved at  $\epsilon = \frac{NR}{N+p-1}$ . Hence, if  $\lambda > \frac{d^{\beta+p-1}}{f(d)} \frac{A^{p-1}N}{(p-1)^{p-1}R^p}$ , then in the definition of the function  $\rho$  we can choose  $\epsilon = \frac{NR}{N+p-1}$  and values for  $\delta(>1)$  and  $\mu(>1)$  so that  $\lambda \geq \frac{d^{\beta+p-1}}{f(d)} \frac{NR^{N-1}(\delta\mu)^{p-1}}{\epsilon^N(R-\epsilon)^{p-1}}$  and hence (10) will hold for all  $\epsilon \leq r \leq R$ .

In order to obtain (9), integrating (7) from t to R, we have

$$\int_t^R -\psi'(r)dr = \int_t^R G^{-1}\left(\frac{\lambda}{r^{N-1}}\left(\int_0^r s^{N-1}h(v(s))\,ds\right)\right)dr$$

for  $0 \leq t \leq R$ . Hence

$$\begin{split} \psi(t) &= \int_t^R G^{-1} \left( \frac{\lambda}{r^{N-1}} \left( \int_0^r s^{N-1} h(v(s)) \, ds \right) \right) dr \\ &\leq \int_t^R G^{-1} \left( \frac{\lambda}{r^{N-1}} h(d) \frac{r^N}{N} \right) dr \\ &= \int_t^R \left( \frac{\lambda}{N} h(d) \right)^{\frac{1}{p-1}} r^{\frac{1}{p-1}} dr \\ &\leq \left( \frac{\lambda}{N} h(d) \right)^{\frac{1}{p-1}} \int_0^R r^{\frac{1}{p-1}} dr \\ &= \frac{p-1}{p} \left( \frac{\lambda R^p}{N} \frac{f(d)}{d^\beta} \right)^{\frac{1}{p-1}}. \end{split}$$

from which we have  $\|\psi\|_{\infty} \leq \frac{p-1}{p} \left(\frac{\lambda R^p}{N} \frac{f(d)}{d^{\beta}}\right)^{\frac{1}{p-1}}$ . Since  $\lambda < \frac{d^{\beta}}{f(d)} \frac{N}{R^p} \left(\frac{p}{p-1}\right)^{p-1} b^{p-1}$ , we obtain  $\|\psi\|_{\infty} \leq b$ . Thus  $\psi$  satisfies

$$\begin{cases} -\Delta_p \psi \le \lambda h(\psi) & \text{in } B(0,R) \\ \psi = 0 & \text{on } \partial B(0,R) \end{cases}$$

and  $d \leq \|\psi\|_{\infty} \leq b$ .

Now, let  $z(x) = \psi(x)$ , if  $x \in B_R$  and z(x) = 0, if  $x \in \Omega - B_R$ . Then  $z \in W^{1,p}(\Omega) \bigcap C(\overline{\Omega})$  and z = 0 on  $\partial\Omega$ , which is subsolution of the nonsigular problem (5) in  $\Omega$ . However, z is not strictly positive in  $\Omega$ . To obtain a strictly positive subsolution of (5) in  $\Omega$  we iterate this subsolution z once in a suitable manner. By the properties of h, there exists  $\sigma_{\lambda} > 0$  such that  $\lambda h(z) + \sigma_{\lambda} G(z)$  is increasing for all  $z \geq 0$ . Define  $\psi_2$  to be the solution of

$$\begin{cases} -\Delta_p \psi_2 + \sigma_\lambda G(\psi_2) = \tilde{h}(z) & \text{in } \Omega\\ \psi_2 = 0 & \text{on } \partial\Omega \end{cases}$$

with  $\tilde{h}(z) = \lambda h(z) + \sigma_{\lambda} G(z)$ . Then since the operator  $-\Delta_p + \sigma_{\lambda} G$  satisfies the weak comparison principle (see [3]), we can have  $z \leq \psi_2$  (see [6]). Further we get  $\psi_2(x) > 0$  for all  $x \in \Omega$  since  $\tilde{h}(0) > 0$ . Hence by the monotonicity of  $\tilde{h}$  we have

$$-\Delta_p \psi_2 + \sigma_\lambda G(\psi_2) = \tilde{h}(z) \le \tilde{h}(\psi_2) = \lambda h(\psi_2) + \sigma_\lambda G(\psi_2),$$

which implies that  $\psi_2$  is a subsolution of the nonsingular problem (5) such that  $\psi_2 > 0$  in  $\Omega$ . Since  $h(u) \leq \frac{f(u)}{u^{\beta}}$  for all  $u \geq 0$ , we have  $-\Delta_p \psi_2 \leq \lambda h(\psi_2) \leq \lambda \frac{f(\psi_2)}{\psi_2^{\beta}}$ , showing that  $\psi_2$  is a positive subsolution of our singular problem (1). Therefore, for

$$\frac{d^{1+\beta}}{f(d)} \frac{A^{p-1}N}{(p-1)^{p-1}R^p} < \lambda < \min\left\{\frac{d^{\beta}}{f(d)} \frac{N}{R^p} \left(\frac{p}{p-1}\right)^{p-1} b^{p-1}, \ \frac{a^{\beta+p-1}}{f(a)} \frac{1}{\|w\|_{\infty}^{\beta+p-1}}\right\}$$

we obtain a positive subsolution  $\psi_2$  and a positive supersolution  $\phi_2$  be such that  $\psi_2 \nleq \phi_2$ .

From the proof of Theorem 1.2 we note that we have a sufficiently small positive subsolution  $\psi_1$  such that  $\psi_1 \leq \phi_2$  and a sufficiently large positive supersolution  $\phi_1$  such that  $\psi_2 \leq \phi_1$ . Hence, there exist a positive solution  $u_1$  of (1) such that  $\psi_1 \leq u_1 \leq \phi_2$  and a positive solution  $u_2$  of (1) such that  $\psi_2 \leq u_2 \leq \phi_1$ . Since  $\psi_2 \nleq \phi_2$ , we have  $u_1 \neq u_2$ . Therefore, there exist at least two positive solutions of (1) for  $\lambda \in (\lambda_*, \lambda^*)$  and Theorem 1.3 is proven.

## 3. Proof of Theorem 1.6 and Theorem 1.7

*Proof of Theorem* 1.6. We construct a positive supersolution  $(\phi_1, \phi_1)$  of (4).

If both f and g are bounded, let  $(\phi_1, \overline{\phi_1}) = (\lambda M_\lambda w, \lambda M_\lambda w)$  and choose  $M_\lambda$ so large that  $M_\lambda^{p-1} \ge \frac{1}{\lambda^{p-2}} \max\{\|f\|_\infty, \|g\|_\infty\}$ . Then for  $M_\lambda \gg 1$  we have

$$-\Delta_p \phi_1 = \lambda^{p-1} M_{\lambda}^{p-1} \frac{1}{w^{\beta}} \ge \lambda \frac{\|f\|_{\infty}}{w^{\beta}} \ge \lambda \frac{f(\lambda M_{\lambda} w)}{(\lambda M_{\lambda} w)^{\beta}} = \lambda \frac{f(\bar{\phi_1})}{\phi_1^{\beta}}$$

and

$$-\Delta_p \bar{\phi_1} = \lambda^{p-1} M_{\lambda}^{p-1} \frac{1}{w^{\beta}} \ge \lambda \frac{\|g\|_{\infty}}{w^{\beta}} \ge \lambda \frac{g(\lambda M_{\lambda} w)}{(\lambda M_{\lambda} w)^{\beta}} = \lambda \frac{g(\phi_1)}{\bar{\phi_1}^{\beta}},$$

showing that  $(\phi_1, \bar{\phi_1})$  is a positive supersolution of (4). Suppose that  $g(x) \to \infty$ as  $x \to \infty$ , let  $(\phi_1, \bar{\phi_1}) = \left(M_{\lambda}w, \lambda^{\frac{1}{\beta+p-1}}g(M_{\lambda}\|w\|_{\infty})^{\frac{1}{\beta+p-1}}w\right)$ . Then by (G2), we can choose  $M_{\lambda}$  large so that

$$\frac{f\left(\lambda^{\frac{1}{\beta+p-1}} \|w\|_{\infty} g(M_{\lambda} \|w\|_{\infty})^{\frac{1}{\beta+p-1}}\right)}{(M_{\lambda} \|w\|_{\infty})^{\beta+p-1}} \le \frac{1}{\lambda \|w\|_{\infty}^{\beta+p-1}}.$$

Then we have

$$-\Delta_p \phi_1 = \frac{M_{\lambda}^{p-1}}{w^{\beta}}$$

$$\geq \lambda \frac{f\left(\lambda^{\frac{1}{\beta+p-1}} \|w\|_{\infty} g(M_{\lambda} \|w\|_{\infty})^{\frac{1}{\beta+p-1}}\right)}{(M_{\lambda}w)^{\beta}}$$

$$\geq \lambda \frac{f\left(\lambda^{\frac{1}{\beta+p-1}} g(M_{\lambda} \|w\|_{\infty})^{\frac{1}{\beta+p-1}}w\right)}{(M_{\lambda}w)^{\beta}}$$

$$= \lambda \frac{f(\bar{\phi_1})}{\phi_1^{\beta}}.$$

We also have

$$-\Delta_{p}\bar{\phi_{1}} = \lambda^{\frac{p-1}{\beta+p-1}}g(M_{\lambda}\|w\|_{\infty})^{\frac{p-1}{\beta+p-1}}\frac{1}{w^{\beta}}$$
$$= \lambda \frac{g(M_{\lambda}\|w\|_{\infty})}{\lambda^{\frac{\beta}{\beta+p-1}}g(M_{\lambda}\|w\|_{\infty})^{\frac{\beta}{\beta+p-1}}w^{\beta}}$$
$$\geq \lambda \frac{g(M_{\lambda}w)}{\left(\lambda^{\frac{1}{\beta+p-1}}g(M_{\lambda}\|w\|_{\infty})^{\frac{1}{\beta+p-1}}w\right)^{\beta}}$$
$$= \lambda \frac{g(\phi_{1})}{\bar{\phi_{1}}^{\beta}},$$

showing that  $(\phi_1, \bar{\phi_1})$  is a supersolution of (4). (If g is bounded and  $f(x) \to \infty$ as  $x \to \infty$ , then  $\lim_{x\to\infty} \frac{g(Mf(x))}{x^{\beta+p-1}} = 0$  for all M > 0 and we can prove that  $(\phi_1, \bar{\phi_1}) = \left(\lambda^{\frac{1}{\beta+p-1}} f(M_\lambda ||w||_\infty)^{\frac{1}{\beta+p-1}} w, M_\lambda w\right)$  is a supersolution of (4)).

Now, we construct a positive subsolution  $(\psi_1, \bar{\psi}_1)$  of (4). Let e and  $\lambda_1$  be as in the proof of Theorem 1.2. Since  $\lim_{x\to 0} \frac{f(0)}{x^{\beta}} = \infty = \lim_{x\to 0} \frac{g(0)}{x^{\beta}}$ , there exist sufficiently small  $m_{\lambda}$  and  $m'_{\lambda}$  such that

$$\lambda_1(m_\lambda e)^{p-1} \le \lambda \frac{f(0)}{(m_\lambda e)^{\beta}}$$
 and  $\lambda_1(m'_\lambda e)^{p-1} \le \lambda \frac{g(0)}{(m'_\lambda e)^{\beta}}$ 

Let  $(\psi_1, \bar{\psi_1}) = (m_\lambda e, m'_\lambda e)$ . Since f and g are nondecreasing, we have

$$-\Delta_p \psi_1 = \lambda_1 (m_\lambda e)^{p-1} \le \lambda \frac{f(0)}{(m_\lambda e)^\beta} \le \lambda \frac{f(m'_\lambda e)}{(m_\lambda e)^\beta} = \lambda \frac{f(\bar{\psi_1})}{\psi_1^\beta}$$

and

$$-\Delta_p \bar{\psi_1} = \lambda_1 (m'_{\lambda} e)^{p-1} \le \lambda \frac{g(0)}{(m'_{\lambda} e)^{\beta}} \le \lambda \frac{g(m_{\lambda} e)}{(m'_{\lambda} e)^{\beta}} = \lambda \frac{g(\psi_1)}{\bar{\psi_1}^{\beta}}$$

Thus  $(\psi_1, \bar{\psi}_1)$  is a positive subsolution of (4), and if  $m_{\lambda}$  and  $m'_{\lambda}$  are sufficiently small then  $(\psi_1, \bar{\psi}_1) \leq (\phi_1, \bar{\phi}_1)$ . Hence Theorem 1.6 is proven.  $\Box$ 

Proof of Theorem 1.7. We construct a second positive supersolution  $(\phi_2, \bar{\phi_2})$ of (4) when  $\lambda \leq \frac{a^{\beta+p-1}}{g(a)} \frac{1}{\|w\|_{\infty}^{\beta+p-1}}$ . Let  $(\phi_2, \bar{\phi_2}) = \left(a \frac{w}{\|w\|_{\infty}}, a \frac{w}{\|w\|_{\infty}}\right)$ . Since  $\lambda \leq \frac{1}{\|w\|_{\infty}^{\beta+p-1}} \frac{a^{\beta+p-1}}{g(a)}$  and  $g(x) \geq f(x)$  for all  $x \geq 0$ , we have

$$-\Delta_p \phi_2 = \frac{a^{p-1}}{\|w\|_{\infty}^{p-1}} \frac{1}{w^{\beta}} \ge \lambda \frac{g(a)}{\left(a\frac{w}{\|w\|_{\infty}}\right)^{\beta}} \ge \lambda \frac{f\left(a\frac{w}{\|w\|_{\infty}}\right)}{\left(a\frac{w}{\|w\|_{\infty}}\right)^{\beta}} = \lambda \frac{f(\bar{\phi_2})}{\phi_2^{\beta}}$$

and

$$-\Delta_p \bar{\phi_2} = \frac{a^{p-1}}{\|w\|_{\infty}^{p-1}} \frac{1}{w^{\beta}} \ge \lambda \frac{g(a)}{\left(a\frac{w}{\|w\|_{\infty}}\right)^{\beta}} \ge \lambda \frac{g\left(a\frac{w}{\|w\|_{\infty}}\right)}{\left(a\frac{w}{\|w\|_{\infty}}\right)^{\beta}} = \lambda \frac{g(\phi_2)}{\bar{\phi_2}^{\beta}}.$$

Hence,  $(\phi_2, \bar{\phi}_2)$  is a positive supersolution of (4) with  $\|\phi_2\|_{\infty} = a$  and  $\|\bar{\phi}_2\|_{\infty} = a$ when  $\lambda \leq \frac{a^{\beta+p-1}}{g(a)} \frac{1}{\|w\|_{\infty}^{\beta+p-1}}$ .

Now, when  $\frac{d^{\beta+p-1}}{f(d)} \frac{A^{p-1}N}{(p-1)^{p-1}R^p} < \lambda < \frac{d^{\beta}}{f(d)} \frac{N}{R^p} \left(\frac{p}{p-1}\right)^{p-1} b^{p-1}$ , we construct a second positive subsolution  $(\psi_2, \bar{\psi}_2)$  of (4). Let  $h, \rho, v, \psi, z$  and consequently  $\psi_2$  be as defined in the proof of Theorem 1.3. We note that  $\psi_2 > 0$  in  $\Omega$  and for this range of  $\lambda$  it satisfies

$$\begin{cases} -\Delta_p \psi_2 \le \lambda \frac{f(\psi_2)}{\psi_2^\beta} & \text{in } \Omega\\ \psi_2 = 0 & \text{on } \partial\Omega \end{cases}$$

Now choosing  $\bar{\psi}_2 = \psi_2$ , we have

$$-\Delta_p \psi_2 \le \lambda \frac{f(\psi_2)}{\psi_2^\beta} = \lambda \frac{f(\bar{\psi}_2)}{\psi_2^\beta}$$

and

$$-\Delta_p \bar{\psi}_2 \le \lambda \frac{f(\bar{\psi}_2)}{\bar{\psi}_2^{\beta}} \le \lambda \frac{g(\psi_2)}{\bar{\psi}_2^{\beta}}$$

since  $f(u) \leq g(u)$  for all  $u \geq 0$ . Hence,  $(\psi_2, \bar{\psi}_2)$  is a positive subsolution of (4), when  $\frac{d^{\beta+p-1}}{f(d)} \frac{A^{p-1}N}{(p-1)^{p-1}R^p} < \lambda < \frac{d^{\beta}}{f(d)} \frac{N}{R^p} \left(\frac{p}{p-1}\right)^{p-1} b^{p-1}$ . Therefore, we obtain a positive supersolution  $(\phi_2, \bar{\phi}_2)$  and a positive subsolution  $(\psi_2, \bar{\psi}_2)$  such that for

$$\frac{d^{\beta+p-1}}{f(d)} \frac{A^{p-1}N}{(p-1)^{p-1}R^p} < \lambda < \min\left\{\frac{d^{\beta}}{f(d)} \frac{N}{R^p} (\frac{p}{p-1})^{p-1} b^{p-1}, \ \frac{1}{\|w\|_{\infty}^{\beta+p-1}} \frac{a^{\beta+p-1}}{g(a)}\right\},$$

 $(\psi_2, \bar{\psi}_2) \nleq (\phi_2, \bar{\phi}_2).$ 

From the proof of Theorem 1.6 we note that we have a sufficiently small positive subsolution  $(\psi_1, \bar{\psi}_1)$  such that  $(\psi_1, \bar{\psi}_1) \leq (\phi_2, \bar{\phi}_2)$  and a sufficiently large positive supersolution  $(\phi_1, \bar{\phi}_1)$  such that  $(\psi_2, \bar{\psi}_2) \leq (\phi_1, \bar{\phi}_1)$ . Hence, there exist a positive solution  $(u_1, \bar{u}_1)$  of (4) such that  $(\psi_1, \bar{\psi}_1) \leq (u_1, \bar{u}_1) \leq (\phi_2, \bar{\phi}_2)$  and a positive solution  $(u_2, \bar{u}_2)$  of (4) such that  $(\psi_2, \bar{\psi}_2) \leq (u_2, \bar{u}_2) \leq (\phi_1, \bar{\phi}_1)$ . Since  $(\psi_2, \bar{\psi}_2) \nleq (\phi_2, \bar{\phi}_2)$ , we have  $(u_1, \bar{u}_1) \neq (u_2, \bar{u}_2)$ . Therefore, there exist at least two positive solutions of (4) for  $\lambda \in (\lambda_*, \lambda^*)$  and Theorem 1.7 is proven.

Acknowledgement. Eun Kyoung Lee was supported by the National Research Foundation of Korea Grant funded by the Korean Government [NRF-2009-353-C00042] .

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Received January 16, 2010; revised June 20, 2010