

# Blow-Up Solutions and Global Existence for a Kind of Quasilinear Reaction-Diffusion Equations

*Lingling Zhang, Na Zhang and Lixiang Li*

**Abstract.** In this paper, we study the blow-up solutions and global existence for a quasilinear reaction-diffusion equation including a gradient term and nonlinear boundary condition:

$$\begin{cases} (g(u))_t = \nabla \cdot (a(u)\nabla u) + f(x, u, |\nabla u|^2, t) & \text{in } D \times (0, T) \\ \frac{\partial u}{\partial n} = r(u) & \text{on } \partial D \times (0, T) \\ u(x, 0) = u_0(x) > 0 & \text{in } \bar{D}, \end{cases}$$

where  $D \subset R^N$  is a bounded domain with smooth boundary  $\partial D$ . The sufficient conditions are obtained for the existence of a global solution and a blow-up solution. An upper bound for the “blow-up time”, an upper estimate of the “blow-up rate”, and an upper estimate of the global solution are specified under some appropriate assumptions for the nonlinear system functions  $f, g, r, a$ , and initial value  $u_0$  by constructing suitable auxiliary functions and using maximum principles.

**Keywords.** Reaction-diffusion equation, blow-up solution, global solution

**Mathematics Subject Classification (2010).** Primary 35K57, secondary 35B44

## 1. Introduction

Everyone knows that blow-up solutions and global existence for reaction diffusion equations had played an important role in many fields. So many authors always focus on the study of this field ([4–6, 9, 12]). In this paper, we study the blow-up solution and global existence for the following initial-boundary-value problem of quasilinear reaction-diffusion equation with a gradient term

---

L. Zhang, N. Zhang, L. Li: Department of Mathematics, Taiyuan University of Technology, Taiyuan, P. R. China; zllww@126.com; znmath@126.com; 704761476@qq.com

and nonlinear boundary condition:

$$\begin{cases} (g(u))_t = \nabla \cdot (a(u)\nabla u) + f(x, u, q, t) & \text{in } D \times (0, T) \\ \frac{\partial u}{\partial n} = r(u) & \text{on } \partial D \times (0, T) \\ u(x, 0) = u_0(x) > 0 & \text{in } \bar{D}, \end{cases} \quad (1)$$

where  $q = |\nabla u|^2$ ,  $D \subset R^N$  is a bounded domain with smooth boundary  $\partial D$ ,  $\frac{\partial}{\partial n}$  represents the outward normal derivative on  $\partial D$ ,  $u_0$  is the initial value,  $T$  is the maximal existence time of  $u$ . Set  $R^+ = (0, +\infty)$ , we assume that  $f(x, s, d, t)$  is a nonnegative  $C^1(\bar{D} \times R^+ \times R^+ \times R^+)$  function,  $g(s)$  is a  $C^2(R^+)$  function,  $g'(s) > 0$  for any  $s > 0$ ,  $r(s)$  is a positive  $C^2(R^+)$  function,  $a(s)$  is a  $C^2(R^+)$  function. Under the assumptions above, the problem (1) has a unique classical solution  $u(x, t)$  with some  $T > 0$ , and the solution is positive over  $\bar{D} \times [0, T)$ .

The problem (1) describes many physical phenomena in mechanics, physics and biology, etc. We refer to [5, 10] and the reference therein for many other applications. Souplet et al. [13] deal with the blow-up and global solutions of initial value problems for the reaction-diffusion equations with a gradient term. [1, 3, 14] study the existence of blow-up and global solutions for the reaction-diffusion equations with a gradient term and initial-Dirichlet-boundary-value. Zhang [16] and Ding and Guo [7] investigate the blow-up and global solutions for the reaction-diffusion equations with gradient terms and initial-Neumann-boundary-values.

Some special cases of (1) are also treated. Walter [15] studies the following problem:

$$\begin{cases} u_t = \Delta u & \text{in } D \times (0, T) \\ \frac{\partial u}{\partial n} = r(u) & \text{on } \partial D \times (0, T) \\ u(x, 0) = u_0(x) > 0 & \text{in } \bar{D}, \end{cases}$$

where  $D \subset R^N$  is a bounded domain with smooth boundary. The sufficient conditions characterized by function  $r$  are given for the existence of blow-up and global solutions. Zhang [17] considers the following problem:

$$\begin{cases} (g(u))_t = \Delta u + f(u) & \text{in } D \times (0, T) \\ \frac{\partial u}{\partial n} = r(u) & \text{on } \partial D \times (0, T) \\ u(x, 0) = u_0(x) > 0 & \text{in } \bar{D}, \end{cases}$$

where  $D \subset R^N$  is a bounded domain with smooth boundary. The sufficient conditions are obtained there for the existence of a global solution and a blow-up solution. Meanwhile, the upper estimate of the global solution, the upper

bound of the “blow-up time”, and the upper estimate of the “blow-up rate” are also given. Ding and Guo [8] consider the following problem:

$$\begin{cases} (g(u))_t = \Delta u + f(x, u, q, t) & \text{in } D \times (0, T) \\ \frac{\partial u}{\partial n} = r(u) & \text{on } \partial D \times (0, T) \\ u(x, 0) = u_0(x) > 0 & \text{in } \bar{D}, \end{cases}$$

where  $q = |\nabla u|^2$ ,  $D \subset R^N$  is a bounded domain with smooth boundary. The sufficient conditions for the existence of a blow-up solution, an upper bound for the “blow-up time”, an upper estimate of the “blow-up rate”, the sufficient conditions for the existence of the global solution, and an upper estimate of the global solution are specified under some appropriate assumptions on the nonlinear functions  $f, g, r$ , and initial value  $u_0$ .

In this paper, we consider blow-up solutions and global solutions of (1). We obtain some existence theorems for blow-up solutions, upper bounds of the blow-up time, upper estimates of the blow-up rate, existence theorems for global solutions, and upper estimates of global solutions. The results generalize and deepen ones from corresponding work in [8, 15, 17].

The plan of this paper is as follows. In Section 2 we give the proofs for the main results. A few examples are presented in Section 3 to illustrate the applications of the abstract results.

## 2. Main results

Our first result Theorem 2.1 is about the existence of a blow-up solution.

**Theorem 2.1.** *Let  $u$  be a solution of (1). Assume that the following conditions (i)–(iii) are satisfied:*

(i) *the initial value condition:*

$$\beta = \min_{\bar{D}} \frac{a(u_0)[a'(u_0)q_0 + a(u_0)\Delta u_0 + f(x, u_0, q_0, 0)]}{r(u_0)g'(u_0)} > 0, \tag{2}$$

$$q_0 = |\nabla u_0|^2;$$

(ii) *further restrictions for functions involved:*

*for any  $(x, s, d, t) \in D \times R^+ \times \bar{R}^+ \times R^+$ ,*

$$\begin{aligned} 2f_d(x, s, d, t) \left(\frac{r(s)}{a(s)}\right)' - \frac{r'^2(s)}{a(s)} \left(\frac{a(s)}{r'(s)}\right)' &\geq 0, \\ \left(\frac{a(s)}{g'(s)}\right)' &\geq 0, \quad \left(\frac{f(x, s, d, t)}{r(s)}\right)'_s &\geq 0, \\ f_t(x, s, d, t) &\geq 0, \quad a'(s) &\geq 0; \end{aligned} \tag{3}$$

(iii) the integration condition:

$$\int_{M_0}^{+\infty} \frac{a(s)}{r(s)} ds < \infty, \quad M_0 = \max_{\bar{D}} u_0(x). \tag{4}$$

Then the solution  $u$  of (1) must blow up in a finite time  $T$ , and

$$T \leq \frac{1}{\beta} \int_{M_0}^{+\infty} \frac{a(s)}{r(s)} ds, \quad u(x, t) \leq H^{-1}(\beta(T - t)),$$

where  $H(z) = \int_z^{+\infty} \frac{a(s)}{r(s)} ds$ ,  $z > 0$ , and  $H^{-1}$  is the inverse function of  $H$ .

*Proof.* Consider the auxiliary function

$$\Psi(x, t) = -\frac{1}{r(u)}u_t + \beta\frac{1}{a(u)}. \tag{5}$$

We find that

$$\nabla\Psi = \frac{r'}{r^2}u_t\nabla u - \frac{\nabla u_t}{r} - \frac{\beta a'\nabla u}{a^2}, \tag{6}$$

$$\begin{aligned} \Delta\Psi &= \nabla \cdot (\nabla\Psi) \\ &= \frac{\nabla \cdot (u_t r' \nabla u) r^2 - 2r r' u_t |\nabla u|^2}{r^4} - \frac{r \nabla \cdot (\nabla u_t) - r' \nabla u \cdot \nabla u_t}{r^2} \\ &\quad - \frac{a^2 \beta \nabla \cdot (a' \nabla u) - 2a \cdot a'^2 \beta |\nabla u|^2}{a^4} \end{aligned} \tag{7}$$

$$\begin{aligned} &= \frac{r'' |\nabla u|^2 u_t + r' u_t \Delta u + 2r' \nabla u \cdot \nabla u_t}{r^2} - \frac{2r'^2 |\nabla u|^2 u_t}{r^3} - \frac{\Delta u_t}{r} \\ &\quad - \frac{\beta(a'' |\nabla u|^2 + a' \Delta u)}{a^2} + \frac{2a'^2 |\nabla u|^2 \beta}{a^3}, \\ \Psi_t &= \frac{r'}{r^2}(u_t)^2 - \frac{(u_t)_t}{r} - \frac{\beta a' u_t}{a^2}. \end{aligned}$$

By (1) we have  $u_t = \frac{a'|\nabla u|^2 + a\Delta u}{g'} + \frac{f}{g'}$ ,

$$\begin{aligned} \Psi_t &= \frac{r'}{r^2}(u_t)^2 - \frac{(u_t)_t}{r} - \frac{\beta a' u_t}{a^2} \\ &= \frac{r'}{r^2}(u_t)^2 - \frac{1}{r} \left( \frac{a'|\nabla u|^2 + a\Delta u}{g'} + \frac{f}{g'} \right)_t - \frac{\beta a'}{a^2} \left( \frac{a'|\nabla u|^2 + a\Delta u}{g'} + \frac{f}{g'} \right) \\ &= \frac{r'}{r^2}(u_t)^2 - \frac{a\Delta u_t}{r g'} + \left( \frac{a' g''}{r g'^2} - \frac{a''}{r g'} \right) u_t |\nabla u|^2 + \left( \frac{a g''}{r g'^2} - \frac{a'}{r g'} \right) u_t \Delta u \\ &\quad + \left( \frac{f g''}{r g'^2} - \frac{f_u}{r g'} \right) u_t - \frac{2a' + 2f'_q}{r g'} \nabla u_t \nabla u - \frac{f_t}{r g'} - \frac{\beta a'^2}{a^2 g'} |\nabla u|^2 - \frac{\beta a'}{a g'} \Delta u - \frac{a' \beta f}{a^2 g'}. \end{aligned} \tag{8}$$

It follows from (7) and (8) that

$$\begin{aligned} \frac{a}{g'}\Delta\Psi &= \left(\frac{ar''}{g'r^2} - \frac{2ar'^2}{g'r^3}\right)|\nabla u|^2u_t + \frac{ar'}{g'r^2}u_t\Delta u + \frac{2ar'}{g'r^2}\nabla u \cdot \nabla u_t \\ &\quad - \frac{a}{g'r}\Delta u_t + \left(\frac{2\beta a'^2}{g'a^2} - \frac{\beta a''}{g'a}\right)|\nabla u|^2 - \frac{\beta a'}{g'a}\Delta u, \end{aligned}$$

$$\begin{aligned} \frac{a}{g'}\Delta\Psi - \Psi_t &= \left(\frac{ar''}{g'r^2} - \frac{2ar'^2}{g'r^3}\right)|\nabla u|^2u_t + \frac{ar'}{g'r^2}u_t\Delta u + \frac{2ar'}{g'r^2}\nabla u \cdot \nabla u_t - \frac{a}{g'r}\Delta u_t \\ &\quad + \left(\frac{2\beta a'^2}{g'a^2} - \frac{\beta a''}{g'a}\right)|\nabla u|^2 - \frac{\beta a'}{g'a}\Delta u - \frac{r'}{r^2}(u_t)^2 + \frac{a\Delta u_t}{rg'} \\ &\quad - \left(\frac{a'g''}{rg'^2} - \frac{a''}{rg'}\right)u_t|\nabla u|^2 - \left(\frac{ag''}{rg'^2} - \frac{a'}{rg'}\right)u_t\Delta u - \left(\frac{fg''}{rg'^2} - \frac{f_u}{rg'}\right)u_t \\ &\quad + \frac{2a' + 2f_q}{rg'}\nabla u_t \cdot \nabla u + \frac{f_t}{rg'} + \frac{\beta a'^2}{a^2g'}|\nabla u|^2 + \frac{\beta a'}{ag'}\Delta u + \frac{a'\beta f}{a^2g} \\ &= \left(\frac{ar''}{g'r^2} - \frac{2ar'^2}{g'r^3} - \frac{a'g''}{rg'^2} + \frac{a''}{rg'}\right)u_t|\nabla u|^2 + \left(\frac{ar'}{g'r^2} - \frac{ag''}{rg'^2} + \frac{a'}{rg'}\right)u_t\Delta u \\ &\quad + \left(\frac{2ar'}{g'r^2} + \frac{2a' + 2f_q}{rg'}\right)\nabla u \cdot \nabla u_t + \left(\frac{2\beta a'^2}{g'a^2} - \frac{\beta a''}{g'a} + \frac{\beta a'^2}{a^2g'}\right)|\nabla u|^2 \\ &\quad - \frac{r'}{r^2}(u_t)^2 - \left(\frac{fg''}{rg'^2} - \frac{f_u}{rg'}\right)u_t + \frac{f_t}{rg'} + \frac{a'\beta f}{a^2g'}. \end{aligned} \tag{9}$$

In view of (6), we have

$$\nabla u_t = \frac{r'}{r}u_t\nabla u - \frac{r\beta a'\nabla u}{a^2} - r\nabla\Psi. \tag{10}$$

Substitute (10) into (9) to obtain

$$\begin{aligned} \frac{a}{g'}\Delta\Psi - \Psi_t &= \left(\frac{ar'' + (2a + 2f_q)r'}{r^2g'} - \frac{a'g''}{rg'^2} + \frac{a''}{rg'}\right)u_t|\nabla u|^2 + \left(\frac{ar'}{g'r^2} - \frac{ag''}{rg'^2} + \frac{a'}{rg'}\right)u_t\Delta u \\ &\quad + \left(\frac{\beta a'^2}{g'a^2} - \frac{2f_q\beta a'}{g'a^2} - \frac{2\beta a'r'}{g'ra} - \frac{\beta a''}{g'a}\right)|\nabla u|^2 - \left(\frac{2ar'}{g'r} + \frac{2a' + 2f_q}{g'}\right)\nabla u \cdot \nabla\Psi \\ &\quad - \frac{r'}{r^2}(u_t)^2 - \left(\frac{fg''}{rg'^2} - \frac{f_u}{rg'}\right)u_t + \frac{f_t}{rg'} + \frac{a'\beta f}{a^2g'}, \end{aligned}$$

$$\begin{aligned}
 & \frac{a}{g'} \Delta \Psi + \left( \frac{2ar'}{g'r} + \frac{2a' + 2f_q}{g'} \right) \nabla u \cdot \nabla \Psi - \Psi_t \\
 &= \left( \frac{ar'' + (2a' + 2f_q)r'}{r^2g'} - \frac{a'g''}{rg'^2} + \frac{a''}{rg'} \right) u_t |\nabla u|^2 \\
 &+ \left( \frac{ar'}{g'r^2} - \frac{ag''}{rg'^2} + \frac{a'}{rg'} \right) u_t \Delta u + \left( \frac{\beta a'^2}{g'a^2} - \frac{2f_q \beta a'}{g'a^2} - \frac{2\beta a'r'}{g'ra} - \frac{\beta a''}{g'a} \right) |\nabla u|^2 \\
 &- \frac{r'}{r^2} (u_t)^2 - \left( \frac{fg''}{rg'^2} - \frac{f_u}{rg'} \right) u_t + \frac{f_t}{rg'} + \frac{a'\beta f}{a^2g'}.
 \end{aligned} \tag{11}$$

By (1) we have

$$\Delta u = \frac{g'u_t - f}{a} - \frac{a'}{a} |\nabla u|^2. \tag{12}$$

Substitute (12) into (11), to get

$$\begin{aligned}
 & \frac{a}{g'} \Delta \Psi + \left( \frac{2ar'}{g'r} + \frac{2a' + 2f_q}{g'} \right) \nabla u \cdot \nabla \Psi - \Psi_t \\
 &= \left( \frac{ar'' + (a' + 2f_q)r'}{r^2g'} - \frac{a'^2}{rag'} + \frac{a''}{rg'} \right) u_t |\nabla u|^2 + \left( \frac{a'}{ra} - \frac{g''}{rg'} \right) u_t^2 \\
 &+ \left( \frac{f_u}{rg'} - \frac{fr'}{g'r^2} - \frac{fa'}{arg'} \right) u_t + \left( \frac{\beta a'^2}{g'a^2} - \frac{\beta a''}{g'a} - \frac{2\beta a'r'}{g'ra} - \frac{2\beta a'f_q}{g'a^2} \right) |\nabla u|^2 \\
 &+ \frac{f_t}{rg'} + \frac{a'\beta f}{a^2g'}.
 \end{aligned} \tag{13}$$

With (5), we have

$$u_t = -r\Psi + \frac{r\beta}{a}. \tag{14}$$

Substitution of (14) into (13) gives

$$\begin{aligned}
 & \frac{a}{g'} \Delta \Psi + \left( \frac{2ar'}{g'r} + \frac{2a' + 2f_q}{g'} \right) \nabla u \cdot \nabla \Psi \\
 &- \left[ \left( -\frac{ar'' + (a' + 2f_q)r'}{rg'} - \frac{a''}{g'} + \frac{a'^2}{ag'} \right) |\nabla u|^2 \right. \\
 &+ \left. \frac{2\beta rg''}{ag'} - \frac{2a'r\beta}{a^2} + \frac{fr'}{g'r} + \frac{fa'}{ag'} - \frac{f_u}{g'} + \left( \frac{ra'}{a} - \frac{rg''}{g'} \right) \Psi \right] \Psi - \Psi_t \\
 &= \frac{\beta}{rg'} \left( 2f_q \left( \frac{r}{a} \right)' - \frac{r'^2}{a} \left( \frac{a}{r'} \right)' \right) |\nabla u|^2 + \frac{\beta^2 rg'}{a^3} \left( \frac{a}{g'} \right)' + \frac{\beta r}{ag'} \left( \frac{f}{r} \right)'_u + \frac{f_t}{rg'}.
 \end{aligned} \tag{15}$$

From assumptions (2) and (3), the right-hand side of (15) is nonnegative, i.e.

$$\begin{aligned}
 & \frac{a}{g'} \Delta \Psi + \left( \frac{2ar'}{g'r} + \frac{2a' + 2f_q}{g'} \right) \nabla u \cdot \nabla \Psi - \left[ \left( -\frac{ar'' + (a' + 2f_q)r'}{rg'} - \frac{a''}{g'} + \frac{a'^2}{ag'} \right) |\nabla u|^2 \right. \\
 &+ \left. \frac{2\beta rg''}{ag'} - \frac{2a'r\beta}{a^2} - \frac{f_u}{g'} + \frac{fr'}{g'r} + \frac{fa'}{ag'} + \left( \frac{ra'}{a} - \frac{rg''}{g'} \right) \Psi \right] \Psi - \Psi_t \\
 &\geq 0.
 \end{aligned} \tag{16}$$

Now by (2), we have

$$\begin{aligned} \max_{\bar{D}} \Psi(x, 0) &= \max_{\bar{D}} \left( -\frac{1}{r(u_0)} u_t + \beta \frac{1}{a(u_0)} \right) \\ &= \max_{\bar{D}} \left( -\frac{a'(u_0)q_0 + a(u_0)\Delta u_0 + f(x, u_0, q_0, 0)}{r(u_0)g'(u_0)} + \frac{\beta}{a(u_0)} \right) \\ &= 0. \end{aligned} \tag{17}$$

It follows from (1) that, on  $\partial D \times (0, T)$ ,

$$\begin{aligned} \frac{\partial \Psi}{\partial n} &= \frac{r'}{r^2} u_t \frac{\partial u}{\partial n} - \frac{1}{r} \frac{\partial u_t}{\partial n} - \frac{\beta a'}{a^2} \frac{\partial u}{\partial n} \\ &= \frac{r'}{r} u_t - \frac{1}{r} r_t - \frac{\beta a' r}{a^2} \\ &= \frac{r'}{r} u_t - \frac{r'}{r} u_t - \frac{\beta a'}{a^2} r \\ &= -\frac{\beta a'}{a^2} r \\ &\leq 0. \end{aligned} \tag{18}$$

Combining (16)–(18), and applying the maximum principles [11], we know that the maximum of  $\Psi$  in  $D \times [0, T)$  is zero. Thus  $\Psi \leq 0$ , in  $\bar{D} \times [0, T)$ , and

$$\frac{a(u)}{\beta r(u)} u_t \geq 1. \tag{19}$$

At the point  $x_0 \in \bar{D}$  where  $u_0(x_0) = M_0$ , integrate (19) over  $[0, t]$  to produce

$$\frac{1}{\beta} \int_0^t \frac{a(u)}{r(u)} u_t dt = \frac{1}{\beta} \int_{M_0}^{u(x_0, t)} \frac{a(s)}{r(s)} ds \geq \int_0^t ds = t.$$

This together with assumption (4) shows that  $u$  must blow up in the finite time  $T$  and

$$T \leq \frac{1}{\beta} \int_{M_0}^{+\infty} \frac{a(s)}{r(s)} ds.$$

By integrating the inequality (19) over  $[t, s] (0 < t < s < T)$ , one has, for each fixed  $x$ , that

$$H(u(x, t)) \geq H(u(x, t)) - H(u(x, s)) = \int_{u(x, t)}^{u(x, s)} \frac{a(s)}{r(s)} ds \geq \beta(s - t),$$

passing to the limit as  $s \rightarrow T$  yields  $H(u(x, t)) \geq \beta(T - t)$ , which implies that

$$u(x, t) \leq H^{-1}(\beta(T - t)).$$

The proof is complete. □

The result on the global solution is stated as Theorem 2.2 below.

**Theorem 2.2.** *Let  $u$  be a solution of (1). Assume that the following conditions are satisfied:*

(i) *the initial value condition:*

$$\begin{aligned} \mu &= \max_{\overline{D}} \frac{a(u_0)[a'(u_0)q_0 + a(u_0)\Delta u_0 + f(x, u_0, q_0, 0)]}{r(u_0)g'(u_0)} > 0, \\ q_0 &= |\nabla u_0|^2 \end{aligned} \tag{20}$$

(ii) *further restrictions on functions involved:  
for any  $(x, s, d, t) \in D \times R^+ \times \overline{R^+} \times R^+$ ,*

$$\begin{aligned} 2f_d(x, s, d, t) \left(\frac{r(s)}{a(s)}\right)' - \frac{r'^2(s)}{a(s)} \left(\frac{a(s)}{r'(s)}\right)' &\leq 0, \\ \left(\frac{a(s)}{g'(s)}\right)' &\leq 0, \quad \left(\frac{f(x, s, d, t)}{r(s)}\right)'_s \leq 0, \\ f_t(x, s, d, t) &\leq 0, \quad a'(s) \leq 0; \end{aligned} \tag{21}$$

(iii) *the integration condition:*

$$\int_{m_0}^{+\infty} \frac{a(s)}{r(s)} ds = +\infty, \quad m_0 = \min_{\overline{D}} u_0(x). \tag{22}$$

Then the solution  $u$  of (1) must be a global solution and

$$u(x, t) \leq G^{-1}(\mu t + G(u_0(x))),$$

where  $G(z) = \int_{m_0}^z \frac{a(s)}{r(s)} ds$ ,  $z \geq m_0$ , and  $G^{-1}$  is the inverse function of  $G$ .

*Proof.* Construct an auxiliary function

$$\Phi(x, t) = -\frac{1}{r(u)}u_t + \mu \frac{1}{a(u)}. \tag{23}$$

Replacing  $\Psi$  and  $\beta$  with  $\Phi$  and  $\mu$  in (15), we have

$$\begin{aligned} &\frac{a}{g'}\Delta\Phi + \left(\frac{2ar'}{g'r} + \frac{2a' + 2f_q}{g'}\right)\nabla u \cdot \nabla\Phi \\ &- \left[ \left(-\frac{ar'' + (a' + 2f_q)r'}{rg'} - \frac{a''}{g'} + \frac{a'^2}{ag'}\right)|\nabla u|^2 \right. \\ &\left. + \frac{2\mu rg''}{ag'} - \frac{2a'r\mu}{a^2} - \frac{f_u}{g'} + \frac{fr'}{g'r} + \frac{fa'}{ag'} + \left(\frac{ra'}{a} - \frac{rg''}{g'}\right)\Phi \right] \Phi - \Phi_t \\ &= \frac{\mu}{rg'} \left( 2f_q \left(\frac{r}{a}\right)' - \frac{r'^2}{a} \left(\frac{a}{r'}\right)' \right) |\nabla u|^2 + \frac{\mu^2 rg'}{a^3} \left(\frac{a}{g'}\right)' + \frac{\mu r}{ag'} \left(\frac{f}{r}\right)'_u + \frac{f_t}{rg'}. \end{aligned} \tag{24}$$

It is seen from assumptions (20) and (21) that the right-hand side of (24) is nonpositive, i.e.

$$\begin{aligned} & \frac{a}{g'} \Delta \Phi + \left( \frac{2ar'}{g'r} + \frac{2a' + 2f_q}{g'} \right) \nabla u \cdot \nabla \Phi \\ & - \left[ \left( -\frac{ar'' + (a' + 2f_q)r'}{rg'} - \frac{a''}{g'} + \frac{a'^2}{ag'} \right) |\nabla u|^2 \right. \\ & \left. + \frac{2\mu rg''}{ag'} - \frac{2a'r\mu}{a^2} - \frac{f_u}{g'} + \frac{fr'}{g'r} + \frac{fa'}{ag'} + \left( \frac{ra'}{a} - \frac{rg''}{g'} \right) \Phi \right] \Phi - \Phi_t \\ & \leq 0. \end{aligned} \tag{25}$$

By (20), we have

$$\min_{\bar{D}} \Phi(x, 0) = \min_{\bar{D}} \left( -\frac{a'(u_0)q_0 + a(u_0)\Delta u_0 + f(x, u_0, q_0, 0)}{r(u_0)g'(u_0)} + \frac{\mu}{a(u_0)} \right) = 0. \tag{26}$$

From (1) it follows that

$$\frac{\partial \Phi}{\partial n} = \frac{r'}{r^2} u_t \frac{\partial u}{\partial n} - \frac{1}{r} \frac{\partial u_t}{\partial n} - \frac{\mu a'}{a^2} \frac{\partial u}{\partial n} = \frac{r'}{r} u_t - \frac{1}{r} r_t - \frac{\mu a' r}{a^2} = -\frac{\mu a'}{a^2} r \geq 0 \tag{27}$$

on  $\partial D \times (0, T)$ . Combining (25)–(27) and applying the maximum principles, we know that the minimum of  $\Phi$  in  $\bar{D} \times [0, T]$  is zero. Hence  $\Phi \geq 0$  in  $\bar{D} \times [0, T]$ , i.e.

$$\frac{a(u)}{\mu r(u)} u_t \leq 1. \tag{28}$$

For each fixed  $x \in \bar{D}$ , integrate (28) over  $[0, t]$  to get

$$\frac{1}{\mu} \int_0^t \frac{a(u)}{r(u)} u_t dt = \frac{1}{\mu} \int_{u_0(x)}^{u(x,t)} \frac{a(s)}{r(s)} ds \leq t.$$

This together with (22) shows that  $u$  must be a global solution. Moreover, (28) implies that

$$G(u(x, t)) - G(u_0(x)) = \int_{m_0}^{u(x,t)} \frac{a(s)}{r(s)} ds - \int_{m_0}^{u_0(x)} \frac{a(s)}{r(s)} ds = \int_0^t \frac{a(u)}{r(u)} u_t dt \leq \mu t.$$

Therefore

$$u(x, t) \leq G^{-1}(\mu t + G(u_0(x))).$$

The proof is complete. □

### 3. Applications

**Example 3.1.** Let  $u$  be a solution of the following problem:

$$\left\{ \begin{array}{ll} (\sqrt{u} + u)_t = \nabla \cdot (u \nabla u) + \left( tu^2 + q + \sum_{i=1}^3 x_i^2 \right) u^3 & \text{in } D \times (0, T) \\ \frac{\partial u}{\partial n} = \frac{u^3}{4} & \text{on } \partial D \times (0, T) \\ u(x, 0) = 1 + \sum_{i=1}^3 x_i^2 & \text{in } \bar{D}, \end{array} \right. \quad (29)$$

where  $q = |\nabla u|^2$ ,  $D = \{x = (x_1, x_2, x_3) \mid \sum_{i=1}^3 x_i^2 < 1\}$  is the unit ball of  $R^3$ . Here

$$\begin{aligned} g(u) &= \sqrt{u} + u, & f(x, u, q, t) &= \left( tu^2 + q + \sum_{i=1}^3 x_i^2 \right) u^3, \\ a(u) &= u, & r(u) &= \frac{u^3}{4}, \end{aligned}$$

and

$$\begin{aligned} \beta &= \min_D \frac{a(u_0)[a'(u_0)q_0 + a(u_0)\Delta u_0 + f(x, u_0, q_0, 0)]}{r(u_0)g'(u_0)} \\ &= 8 \min_{1 \leq u_0 \leq 2} \frac{5u_0^4 - 5u_0^3 + 10u_0 - 4}{u_0^{\frac{3}{2}} + 2u_0^2} \\ &= 16. \end{aligned}$$

It is easy to check that (3) and (4) hold. It follows from Theorem 2.1 that  $u$  must blow up in a finite time  $T$ , and

$$T \leq \frac{1}{\beta} \int_{M_0}^{+\infty} \frac{a(s)}{r(s)} ds = \frac{1}{8}, \quad u(x, t) \leq H^{-1}(\beta(T - t)) = \frac{1}{4(T - t)}.$$

**Example 3.2.** Let  $u$  be a solution of the following problem:

$$\left\{ \begin{array}{ll} (ue^u)_t = \nabla \cdot (e^{-u} \nabla u) + e^{-u} \left( e^{-t} + q + \sum_{i=1}^3 x_i^2 \right) & \text{in } D \times (0, T) \\ \frac{\partial u}{\partial n} = 2e^{2-u} & \text{on } \partial D \times (0, T) \\ u(x, 0) = 1 + \sum_{i=1}^3 x_i^2 & \text{in } \bar{D}, \end{array} \right. \quad (30)$$

where  $q = |\nabla u|^2$ ,  $D = \{x = (x_1, x_2, x_3) \mid \sum_{i=1}^3 x_i^2 < 1\}$  is the unit ball of  $R^3$ . Now we have

$$\begin{aligned} g(u) &= ue^u, & f(x, u, q, t) &= e^{-u} \left( e^{-t} + q + \sum_{i=1}^3 x_i^2 \right), \\ a(u) &= e^{-u}, & r(u) &= 2e^{2-u}, \end{aligned}$$

and

$$\begin{aligned} \mu &= \max_D \frac{a(u_0)[a'(u_0)q_0 + a(u_0)\Delta u_0 + f(x, u_0, q_0, 0)]}{r(u_0)g'(u_0)} \\ &= \max_{1 \leq u_0 \leq 2} \frac{u_0 e^{-u_0} + 6e^{-u_0}}{2e^2(e^{u_0} + u_0 e^{u_0})} \\ &= \frac{7}{4} e^{-4}. \end{aligned}$$

It is easy to check that (21) and (22) hold. It follows from Theorem 2.2 that  $u$  must be a global solution and

$$u(x, t) \leq G^{-1}(\mu t + G(u_0(x))) = \frac{7}{2} e^{-2t} + u_0(x).$$

**Acknowledgement.** This work was supported by the National Natural Science Foundation of China (No. 61250011), the Natural Science Foundation of Shanxi Province (No. 2012011004-4), China Postdoctoral Science Foundation (No. 2012M510786) and Research Project supported by Shanxi Scholarship Council of China (No. 2011-011).

## References

- [1] Chen, S., Global existence and blow-up of solutions for a parabolic equation with a gradient term. *Proc. Amer. Math. Soc.* 129 (2001), 975 – 981.
- [2] Chen, S., Boundedness and blow-up for nonlinear degenerate parabolic equations. *Nonlinear Anal.* 70 (2009), 1087 – 1095.
- [3] Chipot, M. and Weissler, F., Some blow-up results for a nonlinear parabolic equation with a gradient term. *SIAM J. Math. Anal.* 20 (1989), 886 – 907.
- [4] Deng, K. and Levine, H., The role of critical exponents in blow-up theorems: the sequel. *J. Math. Anal. Appl.* 243 (2000), 85 – 126.
- [5] Diaz, J. and Thélin, F., On a nonlinear parabolic problem arising in some model related to turbulent flows. *SIAM J. Math. Anal.* 25 (1994), 1085 – 1111.
- [6] Ding, J. and Guo, B., Blow-up and global existence for nonlinear parabolic equations with Neumann boundary conditions. *Comput. Math. Appl.* 60 (2010), 670 – 679.

- [7] Ding, J. and Guo, B., Global and blow-up solutions for nonlinear parabolic equations with a gradient term. *Houston J. Math.* 37 (2011), 1265 – 1277.
- [8] Ding, J. and Guo, B., Global existence and blow-up solutions for quasilinear reaction-diffusion equations with a gradient term. *Appl. Math. Lett.* 24 (2011), 936 – 942.
- [9] Hollis, S. and Morgan, J., On the blow-up of solutions to some semilinear and quasilinear reaction-diffusion systems. *Rocky Mountain J. Math.* 14 (1994), 1447 – 1465.
- [10] Jüngel, A., *Quasi-Hydrodynamic Semiconductor Equations*. Basel: Birkhäuser 2001.
- [11] Protter, M. and Weinberger, H., *Maximum Principles in Differential Equations*. Englewood Cliffs (NJ): Prentice-Hall 1967.
- [12] Qi, Y., Wang, M. and Wang, Z., Existence and non-existence of global solutions of diffusion systems with nonlinear boundary conditions. *Proc. Roy. Soc. Edinburgh Sect. A* 134 (2004), 1199 – 1217.
- [13] Souplet, P., Tayachi, S. and Weissler, F., Exact self-similar blow-up of solutions of a semilinear parabolic equation with a nonlinear gradient term. *Indiana Univ. Math. J.* 45 (1996), 655 – 682.
- [14] Souplet, P. and Weissler, F., Self-similar subsolutions and blow-up for nonlinear parabolic equations. *J. Math. Anal. Appl.* 212 (1997), 60 – 74.
- [15] Walter, W., On existence and nonexistence in the large of solutions of parabolic differential equation with nonlinear boundary condition. *SIAM J. Math. Anal.* 24 (1975), 85 – 90.
- [16] Zhang, L., Blow-up of solutions for a class of nonlinear parabolic equations. *Z. Anal. Anwend.* 25 (2006), 479 – 486.
- [17] Zhang, H., Blow-up solutions and global solutions for nonlinear parabolic problems. *Nonlinear Anal.* 69 (2008), 4567 – 4575.

Received April 23, 2013; revised February 3, 2014