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Free boundary regularity for a problem with right hand side

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We consider a one-phase free boundary problem with variable coefficients and nonzero right hand side. We prove that flat free boundaries are $C^{1,\alpha}$ using a different approach than the classical supconvolution method of Caffarelli. We use this result to deduce that Lipschitz free boundaries are $C^{1,\alpha}$.

1. Introduction

Consider the following one-phase free boundary problem with variable coefficients and nonzero right hand side:

$$\begin{cases} \sum_{i,j} a_{ij}(x)u_{ij} = f & \text{in } \Omega^+(u) := \{x \in \Omega : u(x) > 0\}, \\ |\nabla u| = g & \text{on } F(u) := \partial \Omega^+(u) \cap \Omega, \end{cases}$$
(1.1)

with Ω a bounded domain in \mathbb{R}^n , the coefficients $a_{ij} \in C^{0,\beta}(\Omega)$, $f \in C(\Omega) \cap L^{\infty}(\Omega)$, and $g \in C^{0,\beta}(\Omega)$, $g \ge 0$.

In this paper we are concerned with the regularity of the set F(u), called the free boundary of u. There is an extensive literature on the regularity of the free boundary for this type of problem when $f \equiv 0$. In the case of the Laplace operator, Caffarelli proved in his pioneer work [C1] that Lipschitz free boundaries are $C^{1,\alpha}$, while in [C2] he showed that "flat" free boundaries are Lipschitz. The key step of the method in [C1, C2] consists in finding a family of comparison subsolutions using supconvolutions on balls of variable radii.

Higher regularity of the free boundary follows from the classical work of Kinderlehrer and Nirenberg [KN].

Regularity results in the spirit of [C1, C2] have been subsequently proved for more general operators. In [W1, W2] Wang considered concave fully nonlinear uniformly elliptic operators of the form $\mathcal{F}(D^2u)$. The work [C1] was extended by Feldman [F1, F2] to a class of nonconcave fully nonlinear uniformly elliptic operators of the type $\mathcal{F}(D^2u, Du)$ and to certain nonisotropic problems. For operators with variable coefficients regularity results are proved in the work of Cerruti, Ferrari and Salsa [CFS], and Ferrari and Salsa [FS1, FS2]. Also, Ferrari and then Argiolas and Ferrari in [Fe, AF] considered a class of fully nonlinear operators of the form $\mathcal{F}(D^2u, x)$ with Hölder dependence on x.

The results cited above follow the guidelines of [C1, C2]. One purpose of this paper is to provide a different method of proving that flat free boundaries are $C^{1,\alpha}$. The approach we use is quite flexible since it easily applies to more general nonlinear operators, even degenerate ones, and it also applies to two-phase problems.

In particular, when dealing with operators with variable coefficients we easily deduce that Lipschitz free boundaries are $C^{1,\alpha}$. In fact our flatness result allows us to use a blow-up argument

and reduce the problem to the case of constant coefficients operators. Our strategy is largely inspired by the work of Savin [S].

We now state our main results (for the precise definition of viscosity solutions we refer the reader to Section 2). We assume that the matrix $(a_{ij}(x))$ is positive definite.

THEOREM 1.1 (Flatness implies $C^{1,\alpha}$) Let u be a viscosity solution to (1.1) in B_1 . Assume that $0 \in F(u), g(0) = 1$ and $a_{ij}(0) = \delta_{ij}$. There exists a universal constant $\overline{\varepsilon} > 0$ such that, if the graph of u is $\overline{\varepsilon}$ -flat in B_1 , i.e.

$$(x_n - \bar{\varepsilon})^+ \leqslant u(x) \leqslant (x_n + \bar{\varepsilon})^+, \quad x \in B_1,$$
(1.2)

and

$$[a_{ij}]_{C^{0,\beta}(B_1)} \leqslant \bar{\varepsilon}, \quad \|f\|_{L^{\infty}(B_1)} \leqslant \bar{\varepsilon}, \quad [g]_{C^{0,\beta}(B_1)} \leqslant \bar{\varepsilon}, \tag{1.3}$$

then F(u) is $C^{1,\alpha}$ in $B_{1/2}$.

THEOREM 1.2 (Lipschitz implies $C^{1,\alpha}$) Let u be a viscosity solution to (1.1). Assume that $0 \in F(u)$ and g(0) > 0. If F(u) is a Lipschitz graph in a neighborhood of 0, then F(u) is $C^{1,\alpha}$ in a (smaller) neighborhood of 0.

In the theorem above, the size of the neighborhood where F(u) is $C^{1,\alpha}$ depends on the radius ρ of the ball B_{ρ} where F(u) is Lipschitz, on the Lipschitz norm of F(u), on $[a_{ij}]_{C^{0,\beta}(B_{\rho})}$, $||g||_{C^{0,\beta}(B_{\rho})}$, and $||f||_{L^{\infty}(B_{\rho})}$.

We remark that the assumptions on the coefficients $a_{ij}(x)$ in Theorem 1.1 can be weakened to a Cordes–Nirenberg type condition

$$||a_{ij} - \delta_{ij}||_{L^{\infty}(B_1)} \leq \delta(n).$$

As already pointed out, our strategy of the proof of Theorem 1.1 is inspired by [S]. The main idea is to show that the graph of u enjoys an "improvement of flatness" property, that is, if the graph of u oscillates ε away from a hyperplane in B_1 , then in B_{r_0} it oscillates $\varepsilon r_0/2$ away from possibly a different hyperplane. The key tool in proving this property will be a Harnack type inequality for solutions to a one-phase free boundary problem.

Theorem 1.2 will follow via a blow-up argument from Theorem 1.1 and the classical theory in [C1].

The problem (1.1), in which a right hand side appears, is not specifically dealt with in any of the previous cited works. Our interest in this problem arises in connection with the question of the regularity of the free surface which occurs in the classical hydrodynamical problem for traveling two-dimensional gravity water-waves with vorticity. There has been considerable interest in this problem in recent years, starting with the systematic study of Constantin and Strauss [CS].

The physical situation is the following: a traveling wave of an incompressible, inviscid, heavy fluid moves with constant speed over a horizontal surface. Since the fluid is incompressible, the flow can be described by a stream function u which solves the following free boundary problem (in 2D):

$$\begin{aligned} \Delta u &= -\gamma(u) \quad \text{in } \Omega := \{ (x, y) \in \mathbb{R}^2 : 0 < u(x, y) < B \}, \\ u &= B \quad \text{on } y = 0, \\ |\nabla u|^2 + 2gy = Q \quad \text{on } S := \{ u = 0 \}, \end{aligned}$$

with *B*, *g* fixed constants, γ a given vorticity function and *Q* a parameter. Of special interest are those free boundaries which are given by the graph of a function $y = \psi(x)$. In the regions where ψ

is decreasing (resp. increasing) the free boundary is Lipschitz with respect to the direction $e_1 + e_2$ (resp. $e_2 - e_1$) and moreover Q - 2gy > 0. As a consequence of Theorem 1.2, the free boundary is smooth in these regions.

The free boundary is not expected to be smooth at the so-called stagnation points where Q = 2gy. At such points, the profile of an irrotational wave ($\gamma \equiv 0$) has a corner with included angle of 120°. This was conjectured by Stokes and it was proved by Amick, Fraenkel, and Toland [AFT], and by Plotnikov [P]. The case $\gamma \neq 0$ was investigated by Varvaruca in [V] and recently by Varvaruca and Weiss in [VW].

The paper is organized as follows. In Section 2 we introduce notation and definitions and we prove a regularity result for viscosity solutions to a Neumann problem which we will use in the proof of Theorem 1.1. Next, in Section 3, we present the statement of our Harnack inequality and we exhibit its proof. In Section 4, we state and prove the "improvement of flatness" lemma. Finally, in Section 5, we provide the proofs of Theorems 1.1 and 1.2. We conclude the paper with an Appendix in which we prove the standard Lipschitz continuity and nondegeneracy of solutions to a one-phase free boundary problem.

2. Preliminaries

In this section we provide notation and definitions used throughout the paper. We also present an auxiliary result which will be used in the proof of our main Theorem 1.1.

Notation. For any continuous function $u : \Omega \subset \mathbb{R}^n \to \mathbb{R}$ we denote

$$\Omega^+(u) := \{ x \in \Omega : u(x) > 0 \}, \quad F(u) := \partial \Omega^+(u) \cap \Omega.$$

We refer to the set F(u) as the free boundary of u, while $\Omega^+(u)$ is its positive phase (or side).

We now state the definition of viscosity solution to the problem under consideration, that is,

$$\begin{cases} \sum_{i,j} a_{ij}(x)u_{ij} = f & \text{in } \Omega^+(u), \\ |\nabla u| = g & \text{on } F(u). \end{cases}$$
(2.1)

Here Ω is a bounded domain in \mathbb{R}^n , $a_{ij} \in C^{0,\beta}(\Omega)$, $f \in C(\Omega) \cap L^{\infty}(\Omega)$, $g \in C^{0,\beta}(\Omega)$, and $g \ge 0$.

First we need the following standard notion.

DEFINITION 2.1 Given $u, \varphi \in C(\Omega)$, we say that φ touches u from below (resp. *above*) at $x_0 \in \Omega$ if $u(x_0) = \varphi(x_0)$ and

 $u(x) \ge \varphi(x)$ (resp. $u(x) \le \varphi(x)$) in a neighborhood *O* of x_0 .

If this inequality is strict in $O \setminus \{x_0\}$, we say that φ touches *u strictly* from below (resp. above).

DEFINITION 2.2 Let *u* be a nonnegative continuous function in Ω . We say that *u* is a *viscosity solution* to (2.1) in Ω if the following conditions are satisfied:

(i) $\sum_{i,j} a_{ij}(x)u_{ij} = f$ in $\Omega^+(u)$ in the viscosity sense, i.e. if $\varphi \in C^2(\Omega^+(u))$ touches u from below (resp. above) at $x_0 \in \Omega^+(u)$ then

$$\sum_{i,j} a_{ij}(x_0)\varphi_{ij}(x_0) \leqslant f(x_0) \quad (\text{resp.} \sum_{i,j} a_{ij}(x_0)\varphi_{ij}(x_0) \geqslant f(x_0)).$$

(ii) If $\varphi \in C^2(\Omega)$ and φ^+ touches *u* from below (resp. above) at $x_0 \in F(u)$ and $|\nabla \varphi|(x_0) \neq 0$ then

$$|\nabla \varphi|(x_0) \leq g(x_0)$$
 (resp. $|\nabla \varphi|(x_0) \geq g(x_0)$).

Viscosity solutions are introduced to be able to use comparison techniques. To this end, we will need the following notion of comparison subsolution/supersolution.

DEFINITION 2.3 Let $v \in C^2(\Omega)$. We say that v is a *strict* (comparison) subsolution (resp. supersolution) to (2.1) in Ω if the following conditions are satisfied:

(i) $\sum_{i,j} a_{ij}(x)v_{ij} > f(x)$ (resp. < f(x)) in $\Omega^+(v)$. (ii) If $x_0 \in F(v)$, then

$$|\nabla v|(x_0) > g(x_0)$$
 (resp. $0 < |\nabla v|(x_0) < g(x_0)$).

Notice that, by the implicit function theorem, if v is a strict subsolution/supersolution then F(v) is a C^2 hypersurface.

The following lemma is an immediate consequence of the definitions above.

LEMMA 2.4 Let u, v be respectively a solution and a strict subsolution to (2.1) in Ω . If $u \ge v^+$ in Ω then $u > v^+$ in $\Omega^+(v) \cup F(v)$.

Notation. Hereafter $B_{\rho}(x_0) \subset \mathbb{R}^n$ denotes the open ball of radius ρ centered at x_0 , and $B_{\rho} = B_{\rho}(0)$. A positive constant depending only on the dimension *n* is called a *universal constant*. We often use *c*, *c_i* to denote small universal constants, and *C*, *C_i* to denote large universal constants.

Our main Theorem 1.1 will follow from the regularity properties of solutions to the classical Neumann problem for the Laplace operator. Precisely, we consider the following boundary value problem:

$$\begin{cases} \Delta \tilde{u} = 0 & \text{in } B_{\rho} \cap \{x_n > 0\}, \\ \tilde{u}_n = 0 & \text{on } B_{\rho} \cap \{x_n = 0\}. \end{cases}$$

$$(2.2)$$

We use the notion of viscosity solution to (2.2). For completeness (and for lack of references), we recall standard notions and we prove regularity of viscosity solutions.

DEFINITION 2.5 Let \tilde{u} be a continuous function on $B_{\rho} \cap \{x_n \ge 0\}$. We say that \tilde{u} is a viscosity solution to (2.2) if given a quadratic polynomial P(x) touching \tilde{u} from below (resp. above) at $\bar{x} \in B_{\rho} \cap \{x_n \ge 0\}$,

(i) if $\bar{x} \in B_{\rho} \cap \{x_n > 0\}$ then $\Delta P \leq 0$ (resp. $\Delta P \geq 0$), i.e. \tilde{u} is harmonic in the viscosity sense; (ii) if $\bar{x} \in B_{\rho} \cap \{x_n = 0\}$ then $P_n(\bar{x}) \leq 0$ (resp. $P_n(\bar{x}) \geq 0$).

REMARK Notice that in the definition above we can choose polynomials *P* that touch \tilde{u} strictly from below/above (replace *P* by $P_{\eta}(x) = P(x) - \eta(x_n - \bar{x}_n)^2$ and then let η go to 0).

Also, it suffices to verify that (ii) holds for polynomials \tilde{P} with $\Delta \tilde{P} > 0$. Indeed, let P touch \tilde{u} from below at \bar{x} . Then

$$\tilde{P} = P - \eta (x_n - \bar{x}_n) + C(\eta) (x_n - \bar{x}_n)^2$$

touches \tilde{u} from below at \bar{x} (for a sufficiently small constant $\eta > 0$ and a large constant C > 0 depending on η) and satisfies

$$\Delta \tilde{P} > 0, \quad \tilde{P}_n(\bar{x}) = P_n(\bar{x}) - \eta.$$

If (ii) holds for strictly subharmonic polynomials, we get $\tilde{P}_n(\bar{x}) \leq \eta$, which by letting η go to 0 implies $P_n(\bar{x}) \leq 0$.

LEMMA 2.6 Let \tilde{u} be a viscosity solution to (2.2). Then \tilde{u} is a classical solution to (2.2). In particular, $\tilde{u} \in C^{\infty}(B_{\rho} \cap \{x_n \ge 0\})$.

Proof. Let

$$u^*(x) = \begin{cases} \tilde{u}(x) & \text{if } x \in B_\rho \cap \{x_n \ge 0\}, \\ \tilde{u}(x', -x_n) & \text{if } x \in B_\rho \cap \{x_n < 0\}, \end{cases}$$

where $x' = (x_1, ..., x_{n-1})$.

We claim that u^* is harmonic (in the viscosity sense), and hence smooth, in B_ρ . Indeed, let P be a polynomial touching u^* at $\bar{x} \in B_\rho$ strictly from below. We need to show that $\Delta P \leq 0$. Clearly, we only need to consider the case when $\bar{x} \in \{x_n = 0\}$.

Consider the polynomial

$$S(x) = \frac{P(x) + P(x', -x_n)}{2}.$$

Then

$$\Delta S = \Delta P, \quad S_n(x', 0) = 0. \tag{2.3}$$

Also, S still touches u^* strictly from below at \bar{x} . Now, consider the family of polynomials

$$S_{\varepsilon} = S + \varepsilon x_n, \quad \varepsilon > 0.$$

For ε small S_{ε} will touch u^* from below at some point x_{ε} .

If x_{ε} belongs to $\{x_n = 0\}$, since S_{ε} touches \tilde{u} from below at x_{ε} and $\tilde{u}_n(x', 0) = 0$ in the viscosity sense, we obtain

$$(S_{\varepsilon})_n(x'_{\varepsilon}, 0) \leq 0, \quad \text{i.e.} \quad S_n(x'_{\varepsilon}, 0) + \varepsilon \leq 0,$$

contradicting (2.3).

Thus $x_{\varepsilon} \in B_{\rho} \setminus \{x_n = 0\}$ and hence $\Delta S = \Delta P \leq 0$. In conclusion, u^* is harmonic in B_{ρ} and our statement immediately follows.

3. A Harnack inequality

In this section we will prove a Harnack type inequality for a solution *u* to our problem

$$\begin{cases} \sum_{i,j} a_{ij}(x)u_{ij} = f & \text{in } \Omega^+(u), \\ |\nabla u| = g & \text{on } F(u), \end{cases}$$
(3.1)

under the assumption $(0 < \varepsilon < 1)$

$$\|f\|_{L^{\infty}(\Omega)} \leqslant \varepsilon^{2}, \quad \|g-1\|_{L^{\infty}(\Omega)} \leqslant \varepsilon^{2}, \quad \|a_{ij}-\delta_{ij}\|_{L^{\infty}(\Omega)} \leqslant \varepsilon.$$
(3.2)

This theorem roughly says that if the graph of u oscillates εr away from x_n^+ in B_r , then it oscillates $(1 - c)\varepsilon r$ in $B_{r/20}$. A corollary of this theorem will be a key tool in the proof of Theorem 1.1.

THEOREM 3.1 (Harnack inequality) There exists a universal constant $\bar{\varepsilon}$ such that if u solves (3.1)–(3.2), and for some point $x_0 \in \Omega^+(u) \cup F(u)$,

$$(x_n + a_0)^+ \leqslant u(x) \leqslant (x_n + b_0)^+ \quad \text{in } B_r(x_0) \subset \Omega$$
(3.3)

with

$$b_0 - a_0 \leqslant \varepsilon r, \qquad \varepsilon \leqslant \overline{\varepsilon},$$

then

$$(x_n + a_1)^+ \le u(x) \le (x_n + b_1)^+$$
 in $B_{r/20}(x_0)$

with

$$a_0 \leqslant a_1 \leqslant b_1 \leqslant b_0$$
, $b_1 - a_1 \leqslant (1 - c)\varepsilon r$,

and 0 < c < 1 universal.

From this statement we immediately get the desired corollary to be used in the proof of our main result. Precisely, if u satisfies (3.3) with r = 1, then we can apply the Harnack inequality repeatedly to obtain

$$(x_n + a_m)^+ \leq u(x) \leq (x_n + b_m)^+$$
 in $B_{20^{-m}}(x_0)$

with

$$b_m - a_m \leqslant (1 - c)^m \varepsilon$$

for all m's such that

$$(1-c)^m 20^m \varepsilon \leqslant \bar{\varepsilon}.$$

This implies that for all such m's, the oscillation of the function

$$\tilde{u}_{\varepsilon}(x) = \frac{u(x) - x_n}{\varepsilon}$$

in $(\Omega^+(u) \cup F(u)) \cap B_r(x_0), r = 20^{-m}$, is less than $(1-c)^m = 20^{-\gamma m} = r^{\gamma}$. Thus, the following corollary holds.

COROLLARY 3.2 Let *u* be a solution to (3.1)–(3.2) satisfying (3.3) for r = 1. Then in $B_1(x_0)$, \tilde{u}_{ε} has a Hölder modulus of continuity at x_0 , outside the ball of radius $\varepsilon/\bar{\varepsilon}$, i.e. for all $x \in (\Omega^+(u) \cup F(u)) \cap B_1(x_0)$ with $|x - x_0| \ge \varepsilon/\bar{\varepsilon}$,

$$|\tilde{u}_{\varepsilon}(x) - \tilde{u}_{\varepsilon}(x_0)| \leq C|x - x_0|^{\gamma}.$$

The proof of the Harnack inequality relies on the following lemma.

LEMMA 3.3 There exists a universal constant $\bar{\varepsilon} > 0$ such that if *u* is a solution to (3.1)–(3.2) in *B*₁ with $0 < \varepsilon \leq \bar{\varepsilon}$ and *u* satisfies

$$p(x)^{+} \leq u(x) \leq (p(x) + \varepsilon)^{+}, \quad x \in B_{1}, \ p(x) = x_{n} + \sigma, \ |\sigma| < 1/10,$$
 (3.4)

then if at $\bar{x} = \frac{1}{5}e_n$,

$$u(\bar{x}) \ge (p(\bar{x}) + \varepsilon/2)^+, \tag{3.5}$$

then

$$u \ge (p + c\varepsilon)^+ \quad \text{in } \overline{B}_{1/2} \tag{3.6}$$

for some 0 < c < 1. Analogously, if

$$u(\bar{x}) \leqslant (p(\bar{x}) + \varepsilon/2)^+,$$

then

$$u \leq (p + (1 - c)\varepsilon)^+$$
 in $\overline{B}_{1/2}$.

Proof. We prove the first statement. Clearly, from (3.4),

$$u \ge p \quad \text{in } B_1. \tag{3.7}$$

Let

$$w = c(|x - \bar{x}|^{-\gamma} - (3/4)^{-\gamma})$$

be defined in the closure of the annulus

$$A := B_{3/4}(\bar{x}) \setminus \overline{B}_{1/20}(\bar{x}).$$

The constant c is such that w satisfies the boundary conditions

$$\begin{cases} w = 0 & \text{on } \partial B_{3/4}(\bar{x}), \\ w = 1 & \text{on } \partial B_{1/20}(\bar{x}) \end{cases}$$

Also, since $||a_{ij} - \delta_{ij}||_{L^{\infty}(B_1)} \leq \varepsilon$ the matrix (a_{ij}) is uniformly elliptic and we can choose the constant γ universal so that

$$\sum_{i,j} a_{ij}(x) w_{ij} \ge \delta > 0 \quad \text{ in } A,$$

with δ universal. Extend w to be equal to 1 on $B_{1/20}(\bar{x})$. Notice that since $|\sigma| < 1/10$ using (3.7) we get

$$B_{1/10}(\bar{x}) \subset B_1^+(u). \tag{3.8}$$

Also,

$$B_{1/2} \subset B_{3/4}(\bar{x}) \subset B_1$$

Since in view of (3.7)–(3.8), $u - p \ge 0$ and solves a uniformly elliptic equation in $B_{1/10}(\bar{x})$ with right-hand side f, we can apply the Harnack inequality to obtain

$$u(x) - p(x) \ge c(u(\bar{x}) - p(\bar{x})) - C ||f||_{L^{\infty}} \quad \text{in } \overline{B}_{1/20}(\bar{x}).$$
(3.9)

From (3.5) and the first inequality in (3.2) we conclude that (for ε small enough)

$$u - p \ge c\varepsilon - C\varepsilon^2 \ge c_0 \varepsilon \quad \text{in } \overline{B}_{1/20}(\bar{x}).$$
 (3.10)

Now set

$$v(x) = p(x) + c_0 \varepsilon(w(x) - 1), \quad x \in \overline{B}_{3/4}(\bar{x}),$$
 (3.11)

and for $t \ge 0$,

$$v_t(x) = v(x) + t, \quad x \in B_{3/4}(\bar{x})$$

Notice that

$$\sum_{ij} a_{ij}(x)(v_t)_{ij} \ge c_0 \delta \varepsilon > \varepsilon^2 \quad \text{in } A.$$

According to (3.7) and the definition of v_t we have

$$v_0(x) = v(x) \leqslant p(x) \leqslant u(x), \quad x \in \overline{B}_{3/4}(\overline{x}).$$

Let \bar{t} be the largest $t \ge 0$ such that

$$v_t(x) \leq u(x)$$
 in $B_{3/4}(\bar{x})$.

We want to show that $\overline{t} \ge c_0 \varepsilon$. Then, using the definition (3.11) of v(x) we get

$$u(x) \ge v(x) + \overline{t} \ge p(x) + c_0 \varepsilon w(x)$$

and hence, since on $\overline{B}_{1/2} \subset B_{3/4}(\bar{x})$ one has $w(x) \ge c_2$ for some universal constant c_2 , we obtain

$$u(x) - p(x) \ge c\varepsilon$$
 on $B_{1/2}$

as desired.

Suppose $\overline{t} < c_0 \varepsilon$. Then at some $\widetilde{x} \in \overline{B}_{3/4}(\overline{x})$ we have

$$v_{\bar{t}}(\tilde{x}) = u(\tilde{x}).$$

We show that such a touching point can only occur on $\overline{B}_{1/20}(\bar{x})$. Indeed, since $w \equiv 0$ on $\partial B_{3/4}(\bar{x})$, from the definition of v_t we get

$$v_{\overline{t}}(x) = p(x) - c_0 \varepsilon + \overline{t} \quad \text{on } \partial B_{3/4}(\overline{x}).$$

Using that $\overline{t} < c_0 \varepsilon$ together with the fact that $u \ge p$ we then obtain

$$v_{\bar{t}} < u$$
 on $\partial B_{3/4}(\bar{x})$.

We now show that \tilde{x} cannot belong to the annulus A. As already observed,

$$\sum_{ij} a_{ij}(x)(v_{\bar{t}})_{ij} > \varepsilon^2 \quad \text{in } A$$

and also

$$|\nabla v_{\bar{t}}| \geqslant |v_n| = |1 + c_0 \varepsilon w_n| \quad \text{in } A.$$
(3.12)

We claim that

 $w_n(x) \ge c_1$ on $\{v_{\bar{t}} \le 0\} \cap A$,

for a universal constant c_1 .

Indeed, since w is radially symmetric,

$$w_n(x) = |\nabla w(x)| v_x \cdot e_n, \quad x \in A$$

where v_x is the unit direction of $x - \bar{x}$. Clearly from the formula for w we get $|\nabla w| > c$ on A. Also, $v_x \cdot e_n$ is bounded below in the region $\{v_{\bar{t}} \leq 0\} \cap A$, since for ε small enough,

$$\{v_{\bar{t}} \leq 0\} \cap A \subset \{p \leq c_0 \varepsilon\} = \{x_n \leq -\sigma + c_0 \varepsilon\} \subset \{x_n < 3/20\},\$$

and $\bar{x} = \frac{1}{5}e_n$.

Hence, from (3.12) we deduce that

$$|\nabla v_{\bar{t}}| \ge 1 + c_2 \varepsilon$$
 on $\{v_{\bar{t}} \le 0\} \cap A$.

In particular, for ε small enough and in view of the second inequality in (3.2),

$$|\nabla v_{\bar{t}}|(x) > 1 + \varepsilon^2 \ge g(x) \quad \text{for } x \in A \cap F(v_{\bar{t}}).$$

Thus, $v_{\bar{t}}$ is a strict subsolution to (3.1) in A and according to Lemma 2.4 since u solves (3.1) in B_1 , \tilde{x} cannot belong to A. Therefore, $\tilde{x} \in \overline{B}_{1/20}(\bar{x})$ and

$$u(\tilde{x}) = v_{\bar{t}}(\tilde{x}) \leqslant p(\tilde{x}) + \bar{t} < p(\tilde{x}) + c_0 \varepsilon_1$$

 $u(\tilde{x}) - p(\tilde{x}) < c_0 \varepsilon,$

which implies

contradicting (3.10).

The second statement follows by a similar argument.

We are now ready to give the proof of the Harnack inequality.

Proof of Theorem 3.1. Assume without loss of generality

 $x_0 = 0, \quad r = 1.$

According to (3.3),

$$p(x)^+ \leq u(x) \leq (p(x) + \varepsilon)^+$$
 in B_1

with $p(x) = x_n + a_0$. If $|a_0| < 1/10$ then we can apply the previous Lemma 3.3, and the desired statement immediately follows.

Now suppose otherwise. If $a_0 < -1/10$, then (for ε small) 0 belongs to the zero phase of $(p(x) + \varepsilon)^+$, which implies that 0 also belongs to the zero phase of u, a contradiction.

If $a_0 > 1/10$ then $B_{1/10} \subset B_1^+(u)$, and the conclusion follows by the classical Harnack inequality in $B_{1/10}$ as long as ε is small enough.

4. Improvement of flatness

In this section we present the main "improvement of flatness" lemma, from which Theorem 1.1 will easily follow via an iterative argument.

LEMMA 4.1 (Improvement of flatness) Let u be a solution to (3.1)–(3.2) in B_1 satisfying

$$(x_n - \varepsilon)^+ \leqslant u(x) \leqslant (x_n + \varepsilon)^+ \quad \text{for } x \in B_1, \tag{4.1}$$

and with $0 \in F(u)$. If $0 < r \leq r_0$ for r_0 a universal constant and $0 < \varepsilon \leq \varepsilon_0$ for some ε_0 depending on *r*, then

$$x \cdot \nu - r\varepsilon/2)^+ \leqslant u(x) \leqslant (x \cdot \nu + r\varepsilon/2)^+ \quad \text{for } x \in B_r,$$
(4.2)

with $|\nu| = 1$ and $|\nu - e_n| \leq C\varepsilon^2$ for a universal constant *C*.

Proof. We divide the proof into three steps. We use the following notation:

$$\Omega_{\rho}(u) := (B_1^+(u) \cup F(u)) \cap B_{\rho}$$

Step 1: Compactness. Fix $r \leq r_0$ with r_0 universal (the precise r_0 will be given in Step 3). Assume for contradiction that we can find a sequence $\varepsilon_k \to 0$ and a sequence u_k of solutions to (3.1) in B_1 with coefficients a_{ij}^k , right hand side f_k and free boundary condition g_k satisfying (3.2), such that u_k satisfies (4.1), i.e.

$$(x_n - \varepsilon_k)^+ \leqslant u_k(x) \leqslant (x_n + \varepsilon_k)^+ \quad \text{for } x \in B_1, 0 \in F(u_k),$$
(4.3)

but it does not satisfy the conclusion (4.2) of the lemma.

Set

$$\widetilde{u}_k(x) = \frac{u_k(x) - x_n}{\varepsilon_k}, \quad x \in \Omega_1(u_k).$$

Then (4.3) gives

$$-1 \leqslant \tilde{u}_k(x) \leqslant 1 \quad \text{for } x \in \Omega_1(u_k).$$

$$(4.4)$$

From Corollary 3.2, it follows that the function \tilde{u}_k satisfies

$$|\tilde{u}_k(x) - \tilde{u}_k(y)| \leqslant C|x - y|^{\gamma} \tag{4.5}$$

for C universal and

$$|x-y| \ge \varepsilon_k/\bar{\varepsilon}, \quad x, y \in \Omega_{1/2}(u_k).$$

From (4.3) it clearly follows that $F(u_k)$ converges to $B_1 \cap \{x_n = 0\}$ in the Hausdorff distance. This fact and (4.5) together with Ascoli–Arzelà give that as $\varepsilon_k \to 0$ the graphs of the \tilde{u}_k over $\Omega_{1/2}(u_k)$ converge (up to a subsequence) in the Hausdorff distance to the graph of a Hölder continuous function \tilde{u} over $B_{1/2} \cap \{x_n \ge 0\}$.

Step 2: Limiting solution. We now show that \tilde{u} solves

$$\begin{cases} \Delta \tilde{u} = 0 & \text{in } B_{1/2} \cap \{x_n > 0\}, \\ \tilde{u}_n = 0 & \text{on } B_{1/2} \cap \{x_n = 0\}, \end{cases}$$
(4.6)

in the sense of Definition 2.5.

Let P(x) be a quadratic polynomial touching \tilde{u} at $\bar{x} \in B_{1/2} \cap \{x_n \ge 0\}$ strictly from below. We need to show that

- (i) if $\bar{x} \in B_{1/2} \cap \{x_n > 0\}$ then $\Delta P(\bar{x}) \leq 0$;
- (ii) if $\bar{x} \in B_{1/2} \cap \{x_n = 0\}$ then $P_n(\bar{x}) \leq 0$.

Since $\tilde{u}_k \to \tilde{u}$ in the sense specified above, there exist points $x_k \in \Omega_{1/2}(u_k)$, $x_k \to \bar{x}$, and constants $c_k \to 0$ such that

$$P(x_k) + c_k = \tilde{u}_k(x_k) \tag{4.7}$$

and

$$\tilde{u}_k \ge P + c_k$$
 in a neighborhood of x_k . (4.8)

From the definition of \tilde{u}_k , (4.7) and (4.8) read

$$u_k(x_k) = Q(x_k)$$

and

$$u_k(x) \ge Q(x)$$
 in a neighborhood of x_k

where

$$Q(x) = \varepsilon_k(P(x) + c_k) + x_n.$$

We now distinguish two cases.

(i) If $\bar{x} \in B_{1/2} \cap \{x_n > 0\}$ then $x_k \in B_{1/2}^+(u_k)$ (for k large) and hence, since Q touches u_k from below at x_k , we get

$$\sum_{i,j} a_{ij}^k(x_k) Q_{ij} = \varepsilon_k \sum_{i,j} a_{ij}^k(x_k) P_{ij} \leqslant f_k(x_k) \leqslant \varepsilon_k^2.$$

Thus, in view of the last inequality in (3.2),

$$\Delta P = \sum_{i,j} (\delta_{ij} - a_{ij}^k(x_k)) P_{ij} + \sum_{i,j} a_{ij}^k(x_k) P_{ij} \leq C \varepsilon_k.$$

Passing to the limit as $k \to +\infty$ we conclude that $\Delta P \leq 0$ as desired.

(ii) If $\bar{x} \in B_{1/2} \cap \{x_n = 0\}$, as observed in the Remark following Definition 2.5, we can assume that $\Delta P > 0$. We claim that for k large enough, $x_k \in F(u_k)$. Otherwise $x_{k_n} \in B_1^+(u_{k_n})$ for a subsequence $k_n \to \infty$ and as in case (i),

$$\Delta P \leqslant C \varepsilon_{k_n}$$
.

Letting $k_n \to \infty$ we contradict the fact that *P* is strictly subharmonic. Thus $x_k \in F(u_k)$ for *k* large. Now notice that

$$\nabla Q = \varepsilon_k \nabla P + e_n$$

thus, for k large, $|\nabla Q| > 0$. Since Q^+ touches u_k from below,

$$|\nabla Q|(x_k) \leqslant g_k(x_k) \leqslant 1 + \varepsilon_k^2,$$

which gives

$$|\nabla Q|^2(x_k) = \varepsilon_k^2 |\nabla P|^2(x_k) + 1 + 2\varepsilon_k P_n(x_k) \leq 1 + 3\varepsilon_k^2,$$

and thus (after division by ε_k)

$$\varepsilon_k |\nabla P|^2(x_k) - 3\varepsilon_k + 2P_n(x_k) \leq 0.$$

Passing to the limit as $k \to +\infty$ we obtain $P_n(\bar{x}) \leq 0$ as desired.

Step 3: Improvement of flatness. From the previous step, \tilde{u} solves (4.6) and from (4.4),

$$-1 \leq \tilde{u} \leq 1$$
 in $B_{1/2} \cap \{x_n \geq 0\}$.

From Lemma 2.6 and the bound above we find that, for the given r,

$$|\tilde{u}(x) - \tilde{u}(0) - \nabla \tilde{u}(0) \cdot x| \leq C_0 r^2 \quad \text{in } B_r \cap \{x_n \ge 0\},\$$

for a universal constant C_0 . In particular, since $0 \in F(\tilde{u})$ and also $\tilde{u}_n(0) = 0$, we obtain

$$x' \cdot \tilde{\nu} - C_0 r^2 \leqslant \tilde{u}(x) \leqslant x' \cdot \tilde{\nu} + C_0 r^2 \quad \text{in } B_r \cap \{x_n \ge 0\},$$

with $\tilde{\nu}_i = \tilde{u}_i(0), i = 1, ..., n - 1, |\tilde{\nu}| \leq \tilde{C}, \tilde{C}$ a universal constant. Therefore, for k large enough we get

$$x' \cdot \tilde{\nu} - C_1 r^2 \leq \tilde{u}_k(x) \leq x' \cdot \tilde{\nu} + C_1 r^2$$
 in $\Omega_r(u_k)$.

From the definition of \tilde{u}_k the inequality above reads

$$\varepsilon_k x' \cdot \tilde{\nu} + x_n - \varepsilon_k C_1 r^2 \leqslant u_k \leqslant \varepsilon_k x' \cdot \tilde{\nu} + x_n + \varepsilon_k C_1 r^2 \quad \text{in } \Omega_r(u_k).$$
(4.9)

Set

$$\nu = \frac{(\varepsilon_k \tilde{\nu}, 1)}{\sqrt{\varepsilon_k^2 + 1}}.$$

Since, for *k* large, $1 \leq \sqrt{\varepsilon_k^2 + 1} \leq 1 + \varepsilon_k^2/2$, we deduce from (4.9) that

$$x \cdot v - \varepsilon_k^2 r/2 - C_1 r^2 \varepsilon_k \leqslant u_k \leqslant x \cdot v + \varepsilon_k^2 r/2 + C_1 r^2 \varepsilon_k$$
 in $\Omega_r(u_k)$.

In particular, if r_0 is such that $C_1 r_0 \leq 1/4$ and also k is large enough so that $\varepsilon_k \leq 1/2$ we obtain

$$x \cdot v - \varepsilon_k r/2 \leq u_k \leq x \cdot v + \varepsilon_k r/2$$
 in $\Omega_r(u_k)$,

which together with (4.3) implies that

$$(x \cdot v - \varepsilon_k r/2)^+ \leq u_k \leq (x \cdot v + \varepsilon_k r/2)^+$$
 in B_r .

Thus the u_k satisfy the conclusion of the lemma, and we have reached a contradiction.

5. The proofs of Theorems 1.1 and 1.2

In this section we finally present the proofs of our main theorems.

Proof of Theorem 1.1. Let *u* be a viscosity solution to (1.1) in B_1 with $0 \in F(u)$, g(0) = 1 and $a_{ij}(0) = \delta_{ij}$. Consider the sequence of rescalings

$$u_k(x) := \frac{u(\rho_k x)}{\rho_k}, \quad x \in B_1,$$

with $\rho_k = \bar{r}^k$, k = 0, 1, ..., for a fixed \bar{r} such that

$$\bar{r}^{\beta} \leqslant 1/4, \quad \bar{r} \leqslant r_0,$$

with r_0 the universal constant of Lemma 4.1.

Each u_k solves (1.1) in B_1 with coefficients $a_{ij}^k(x) = a_{ij}(\rho_k x)$, right hand side $f_k(x) := \rho_k f(\rho_k x)$, and free boundary condition $g_k(x) := g(\rho_k x)$. For the chosen \bar{r} , by taking $\bar{\varepsilon} = \varepsilon_0(\bar{r})^2$ the assumption (3.2) holds for $\varepsilon = \varepsilon_k := 2^{-k} \varepsilon_0(\bar{r})$. Indeed, in B_1 , in view of (1.3),

$$|f_k(x)| \leq ||f||_{L^{\infty}} \rho_k \leq \bar{\varepsilon} \bar{r}^k \leq \varepsilon_k^2,$$

$$|g_k(x) - 1| = |g(\rho_k x) - g(0)| \leq [g]_{0,\beta} \rho_k^\beta \leq \bar{\varepsilon} \bar{r}^{k\beta} \leq \varepsilon_k^2,$$

$$|a_{ij}^k(x) - \delta_{ij}| = |a_{ij}(\rho_k x) - a_{ij}(0)| \leq [a_{ij}]_{0,\beta} \rho_k^\beta \leq \bar{\varepsilon} \bar{r}^{k\beta} \leq \varepsilon_k.$$

The hypothesis (1.2) guarantees that for k = 0 also the flatness assumption (4.1) in Lemma 4.1 is satisfied by u_0 . Then it easily follows by induction on k and Lemma 4.1 that each u_k is ε_k -flat in B_1 in the sense of (4.1). Now, a standard iteration argument gives the desired statement.

Proof of Theorem 1.2. Let *u* be a viscosity solution to (1.1) with $0 \in F(u)$ and g(0) > 0. Without loss of generality, assume g(0) = 1. Also, for simplicity we take $a_{ij}(0) = \delta_{ij}$.

Consider the blow-up sequence

$$u_k := u_{\delta_k}(x) = u(\delta_k x)/\delta_k,$$

with $\delta_k \to 0$ as $k \to \infty$. As in the previous theorem, each u_k solves (1.1) with coefficients $a_{ij}^k(x) = a_{ij}(\delta_k x)$, right hand side $f_k(x) := \delta_k f(\delta_k x)$, and free boundary condition $g_k(x) := g(\delta_k x)$. For k large, the assumption (1.3) is satisfied for the universal constant $\bar{\varepsilon}$. In fact, in B_1 ,

$$|f_k(x)| = \delta_k |f(\delta_k x)| \leq \delta_k ||f||_{L^{\infty}} \leq \bar{\varepsilon},$$

$$|g_k(x) - 1| = |g_k(x) - g(0)| \leq [g_k]_{0,\beta} = \delta_k^{\beta} [g]_{0,\beta} \leq \bar{\varepsilon},$$

$$|a_{ij}^k(x) - \delta_{ij}| = |a_{ij}(\delta_k x) - a_{ij}(0)| \leq [a_{ij}(\delta_k x)]_{0,\beta} = \delta_k^{\beta} [a_{ij}]_{0,\beta} \leq \bar{\varepsilon}.$$

Thus, using nondegeneracy and uniform Lipschitz continuity of the u_k 's (see Appendix for a proof of these properties), standard arguments (see for example [AC]) imply that (up to extracting a subsequence):

- (i) $u_k \to u_0 < 1$,
- (ii) $\partial \{u_k > 0\} \rightarrow \partial \{u_0 > 0\}$ locally in the Hausdorff distance,

for a globally defined function $u_0 : \mathbb{R}^n \to \mathbb{R}$. The blow-up limit u_0 is a global solution to the free boundary problem

$$\begin{cases} \Delta u_0 = 0 & \text{in } \{u_0 > 0\}, \\ |\nabla u_0| = 1 & \text{on } F(u_0), \end{cases}$$
(5.1)

and since F(u) is a Lipschitz graph in a neighborhood of 0 we also see from (i)–(ii) that $F(u_0)$ is Lipschitz continuous. Thus, it follows from [C1] that u_0 is a so-called one-plane solution, i.e. (up to rotations) $u_0 = x_n^+$. Combining the facts above, one concludes that for all k large enough, u_k is $\bar{\varepsilon}$ -flat say in B_1 , i.e.

$$(x_n - \overline{\varepsilon})^+ \leq u_k(x) \leq (x_n + \overline{\varepsilon})^+, \quad x \in B_1.$$

Thus u_k satisfies the assumptions of Theorem 1.1, and our conclusion follows.

6. Appendix

We sketch here the proof of a standard result, that is, Lipschitz continuity and nondegeneracy of a solution u to (3.1)–(3.2).

LEMMA 6.1 Let *u* be a solution to (3.1)–(3.2) with $\varepsilon \leq \tilde{\varepsilon}$ a universal constant. If $F(u) \cap B_1 \neq \emptyset$ and F(u) is a Lipschitz graph in B_2 , then *u* is Lipschitz and nondegenerate in $B_1^+(u)$, i.e.

$$c_0 d(z) \leq u(z) \leq C_0 d(z)$$
 for all $z \in B_1^+(u)$,

with d(z) = dist(z, F(u)) and c_0, C_0 universal constants.

Proof. Assume without loss of generality that $0 \in B_1^+(u)$ and set d := d(0).

Consider the rescaled function

$$\tilde{u}(x) = u(dx)/d, \quad x \in B_1.$$

Clearly \tilde{u} still satisfies (1.1) in B_1 with coefficients $\tilde{a}_{ij}(x) = a_{ij}(dx)$, right hand side $\tilde{f}(x) = df(dx)$ and free boundary condition $\tilde{g}(x) = g(dx)$. Since $d \leq 1$, the assumption (3.2) holds. We wish to show that

$$c_0 \leq \tilde{u}(0) \leq C_0.$$

Assume for contradiction that $\tilde{u}(0) > C_0$, with C_0 to be made precise later.

To construct a subsolution, we use the same function as in Lemma 3.3. Precisely, let

$$G(x) = C(|x|^{-\gamma} - 1)$$

be defined on the closure of the annulus $B_1 \setminus \overline{B}_{1/2}$. In view of the uniform ellipticity of the coefficients, we can choose γ large universal so that (for ε small)

$$\sum_{ij} \tilde{a}_{ij} G_{ij} > \varepsilon^2 \quad \text{ on } B_1 \setminus \overline{B}_{1/2}$$

and we can choose C so that

$$G = 1$$
 on $\partial B_{1/2}$.

By the Harnack inequality (see (3.9)), using the contradiction hypothesis we get (for ε small)

$$\tilde{u} \ge c\tilde{u}(0)$$
 on $\overline{B}_{1/2}$.

Thus, by the maximum principle,

$$\tilde{u}(x) \ge v(x) = c\tilde{u}(0)G(x)$$
 on $B_1 \setminus B_{1/2}$.

Hence at the point z where d(0) is achieved we have

$$|\nabla v|(z) \leqslant g(z) \leqslant 1 + \varepsilon^2 \leqslant 2,$$

which contradicts $\tilde{u}(0) > C_0$ if C_0 is large enough.

To prove the lower bound, let

$$G(x) = \eta(1 - G(x))$$

with η (depending on γ) such that

$$|\nabla \tilde{G}| < 1 - \varepsilon^2$$
 on $\partial B_{1/2}$.

Assume without loss of generality that F(u) is a Lipschitz graph in the x_n direction with Lipschitz constant equal to 1. We translate the graph of \tilde{G} by $-4e_n$. Notice that it is above the graph of \tilde{u} since $\tilde{u} \equiv 0$ in $B_1(-4e_n)$. We slide the graph of \tilde{G} in the e_n direction till we touch the graph of \tilde{u} . Since \tilde{G} is a strict supersolution to our free boundary problem, the touching point \tilde{z} can occur only on the η level set with $\tilde{d} := d(\tilde{z}, F(u)) \leq 1$. From the first part, \tilde{u} is Lipschitz continuous and hence $\tilde{u}(\tilde{z}) = \eta \leq C\tilde{d}$. Thus

$$C^{-1}\eta \leqslant \tilde{d} \leqslant 1,$$

that is, \tilde{d} is comparable to 1. Since F(u) is Lipschitz we can connect 0 and \tilde{z} with a chain of intersecting balls included in the positive side of \tilde{u} with radii comparable to 1. The number of balls is bounded by a universal constant. Then we can apply the Harnack inequality to obtain (for ε small)

$$\tilde{u}(0) \geqslant c\tilde{u}(\tilde{z}) = c_0,$$

as desired.

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