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## A note on eigenvalues of ordinary di-erential operators

## Alan Ho

In this follow-up on the work of FS an improved condition for the discrete eigenvalues of the operator  $-a^{-}/ax^{-} + V(x)$  is established for  $V(x)$  satisfying certain hypotheses. The eigenvalue condition in [FS] establishes eigenvalues of this operator to within a small error. The order of a service due to C Fe Ferrican due to C Fe accuracy can be accuracy can be accuracy can be accura be improved if a certain condition is true. This paper improves on the result obtained in [FS] by showing that this condition does indeed hold.

The theorem proven here relies on a version of WKB theory developed in [FS] and applies to operators with large slowly varying potentials. For example, it applies to potentials of the form  $V(x) = \lambda^2 V_1(x)$ for fixed, smooth  $v_1$ , with  $v \to 0$ , v having a local minimum, and  $\lambda \gg 1$ . The theorem applies to more general potentials as well.

Standard WKB theory yields the statement that all eigenvalues E of the differential operator  $-a^2/ax^2 + V(x)$  satisfy

(1) 
$$
\int_{x_{\text{left}}}^{x_{\text{right}}} (E - V(x))^{1/2} dx = \pi \left( k + \frac{1}{2} \right) + O\left( \lambda^{-1} \right), \text{ for some } k \in \mathbb{Z},
$$

where  $x_{\text{left}}$  and  $x_{\text{right}}$  are the two solutions of  $E = V(x) = 0$ .

[FS] shows that this condition for eigenvalues can be improved so that given  $N > 0$ , there exists  $N > 0$  and complex functions  $n_l(E)$ defined in  $[FS]$  so that  $(1)$  becomes

$$
\int_{x_{\rm left}}^{x_{\rm right}} (E - V(x))^{1/2} \, dx + \text{Im} \log \left( 1 + \sum_{l=1}^{N'} h_l(E) \right)
$$

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(2) 
$$
= \pi \left( k + \frac{1}{2} \right) + O \left( \Lambda^{-N} \right),
$$

where  $\Lambda$ , which will be defined precisely in the theorem, plays a role analogous to h- is explicitly given in FS and is purely imaginary  $\Gamma$  for  $\ell$  is the critical property of higher property of higher property of  $\ell$  $O(\Lambda^{-l})$ , and the quantity  $\sum h_l(E)$  is  $O(\Lambda^{-1})$  in absolute value, and hence the Taylor series of log gives

(3) 
$$
\int_{x_{\text{left}}}^{x_{\text{right}}} (E - V(x))^{1/2} dx + i h_1(E) = \pi \left( k + \frac{1}{2} \right) + O\left( \Lambda^{-2} \right).
$$

But if we were to carry out the same calculation to order  $\Lambda^{-1}$ , then since

(4) 
$$
\log\left(1+\sum_{l=1}^{N'}h_l(E)\right)=h_1(E)+h_2(E)-\frac{1}{2}h_1^2(E)+O(\Lambda^{-3}),
$$

we have

(5) 
$$
\int_{x_{\text{left}}}^{x_{\text{right}}} (E - V(x))^{1/2} dx + \text{Im}(h_1(E) + h_2(E))
$$

$$
= \pi \left(k + \frac{1}{2}\right) + O\left(\Lambda^{-3}\right).
$$

Note  $n_{\bar{1}}$  is real and therefore makes no contribution to the left-hand side  $N=1$  and the more shown that hence  $N=1$  is purely in the  $N=1$  in the  $N=1$  to the simpler left-hand side of This improves upon since (5) holds to  $O(\Lambda^{-3})$  instead of  $O(\Lambda^{-2})$ . Using the above fact we obtain an improved version of part of the William Eigenvalue Theorem cf-part of the William p. 239.). For the reader's convenience and for completeness we repeat the hypotheses here

— — — — — — — — Suppose we are given positive functions Square functions and Bx on Bx on Bx and Bx and Bx on B I and a potential  $V(x)$  supported on a possibly unbounded interval I<sub>BVP</sub> with  $I \subset I_{\text{BVP}}$ . Furthermore, suppose we are given two real numbers  $E_0 \leq E_{\infty}$ , positive numbers  $\varepsilon < 1/100$ ,  $K > 1$  and  $N > K \varepsilon^{-10}$ . Define  $N = \frac{\mathbb{E}[N]}{200}$  and  $N = \frac{\mathbb{E}[N]}{2} = N - \frac{\mathbb{E}[N]}{200}$ . And suppose we have the following hypotheses:

Hyp0) If  $x, y \in I$  and  $|x - y| < c B(x)$ , then

$$
c < \frac{B(y)}{B(x)} < C \qquad and \qquad c < \frac{S(y)}{S(x)} < C.
$$

Hyp1) For  $x \in I$  and  $\alpha \geq 0$  we have

$$
\left| \left( \frac{d}{dx} \right)^{\alpha} V(x) \right| \leq C_{\alpha} S(x) B^{\alpha}(x) .
$$

 $\tau = \nu$   $\tau$  . The equation  $\tau$   $\tau$   $\tau$   $\tau$   $\tau$   $\tau$  is the solution  $\tau$  in the equation in the equation  $\tau$  $I$ , and they satisfy

$$
dist(x_{\text{left}}, \partial I) > c B(x_{\text{left}}), \quad dist(x_{\text{right}}, \partial I) > c B(x_{\text{right}}).
$$

Hyp3)

$$
-V'(x) > c S(xleft)B^{-1}(xleft), \qquad for \ x \in [xleft, xleft + c1B(xleft)]
$$

and

$$
V'(x) > c S(x_{\text{right}}) B^{-1}(x_{\text{right}}) , \qquad \text{for } x \in [x_{\text{right}} - c_1 B(x_{\text{right}}), x_{\text{right}}].
$$

Hyp4)

$$
c\,S(x) < E_0 - V(x) < CS(x)
$$

for  $x \in |x_{\text{left}} + c_1 B(x_{\text{left}}), x_{\text{right}} - c_1 B(x_{\text{right}})|$ .

To state the remaining hypotheses, it is convenient to establish some notation. Set  $\lambda(x) = S^{1/2}(x)B(x)$  for  $x \in I$ , and set

$$
B_{\text{left}} = B(x_{\text{left}}), \qquad S_{\text{left}} = S(x_{\text{left}}), \qquad \lambda_{\text{left}} = \lambda(x_{\text{left}}).
$$
  

$$
B_{\text{right}} = B(x_{\text{right}}), \qquad S_{\text{right}} = S(x_{\text{right}}), \qquad \lambda_{\text{right}} = \lambda(x_{\text{right}}).
$$

For  $|E - E_0| < c S_{\text{left}}$ , let  $x_{\text{left}}(E)$  be the solution of  $V(x) = E$  nearest to  $x_{\text{left}}$ , and for  $|E - E_0| < cS_{\text{right}}$ , let  $x_{\text{right}}(E)$  be the solution of  $\mathcal{L}$  , and the state to  $\mathcal{L}$  and  $\mathcal{L}$  are stated to  $\mathcal{L}$  . The state of  $\mathcal{L}$ 

$$
S_{\min} = \int_{x_{\text{left}} < x < x_{\text{right}}} S(x) \, dx
$$

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and

$$
\Lambda = \int_{x_{\rm left}}^{x_{\rm right}} (S^{1/2}(x)B^2(x))^{-1} dx.
$$

Our remaining hypotheses are as follows.

 $\bullet$  Assumptions on  $V(x)$  in all of IBVP:

Hyp5) If  $|E - E_0| < c_2 S_{\text{min}}$  and  $E \le E_{\infty}$ , then  $V(x) > E$  for all  $x \in I_{\rm BVP} - |x_{\rm left}(E), x_{\rm right}(E)|.$ 

Hyp6) If  $x \in I_{\text{BVP}}$  satisfies  $x < x_{\text{left}} - \lambda_{\text{left}}^{\mathbf{A}}B_{\text{left}}/2$  then  $V(x) \geq$  $E_{\infty}+100/|x-x_{\rm left}|^2$ , and if  $x \in I_{\rm BVP}$  satisfies  $x>x_{\rm right}+\lambda_{\rm right}^K B_{\rm right}/2$ , then  $V(x) \ge E_{\infty} + 100/|x - x_{\text{right}}|^2$ .

 $\bullet$  lecnnical Assumptions:

Hyp7)  $\max_{x \in I} S(x) \leq \lambda_{\text{left}}^{\mathbf{K}} S_{\text{left}}$  and  $\max_{x \in I} S(x) \leq \lambda_{\text{right}}^{\mathbf{K}} S_{\text{right}}$ . Hyp8)

$$
\int_{x_{\text{left}}}^{x_{\text{right}}} \left( \frac{dx}{S^{1/2}(x)} \right) \leq \Lambda^K \min \left\{ S_{\text{left}}^{-1/2} B_{\text{left}}, S_{\text{right}}^{-1/2} B_{\text{right}} \right\}.
$$

Hyp9)

$$
\Big(\int_{x_{\text{left}}}^{x_{\text{right}}}\frac{dx}{S^{1/2}(x)\,B^4(x)}\Big)\Big(\int_{x_{\text{left}}}^{x_{\text{right}}}\frac{dx}{S^{1/2}(x)}\Big)\leq \Lambda^K\;.
$$

 $\bullet$  WAB Condition:

Hyp10)  $\Lambda$  is bounded below by a positive constant depending only on K and N and on c C c- c C in Hyp
-Hyp-

Then if E is an eigenvalue of  $-a^{-}/ax^{+} + V(x)$ , we have that

$$
\int_{x_{\rm left}}^{x_{\rm right}} (E - V(x))^{1/2} dx + i h_1(E) = \pi \left( k + \frac{1}{2} \right) + \phi_{\rm error}(E) ,
$$

with  $|\phi_{\text{error}}| \leq C\Lambda^{-3}$  and

$$
h_1(E) = \frac{i}{48} \lim_{\delta \to 0} \left( \int_{x_{\text{left}} + \delta}^{x_{\text{right}} - \delta} V''(x) (E - V(x))^{- (3/2)} dx - q(E) \delta^{-1/2} \right)
$$

with  $q(E)$  uniquely specified by demanding the finiteness of the limit.

 $\mathcal{P}$  . The accomplex function fluid fluid particle function fluid particle particle function fluid particle function fluid particle function fluid fluid parameters  $\mathcal{P}$ property on the index l if it is real-valued for l even and purely imaginary for l odd. It suffices to show that  $h_l$  has the alternating parity property on the index l. Recall that the  $h_l$ 's are inductively determined by

(6) 
$$
u_k^{\text{left}}(x,E) = \sum_{l=0}^k h_l(E) u_{k-l}^{\text{right}}(x,E),
$$

where  $u_k$  is the canonical solution of the transport equations

$$
u_0 \equiv 1,
$$
  
\n
$$
2 i u'_{k+1} + \left(\frac{5}{16} (p')^2 p^{-5/2} - \frac{1}{4} p'' p^{-3/2}\right) u_k
$$
  
\n
$$
- \frac{1}{2} p' p^{-3/2} u'_k + p^{-1/2} u''_k = 0, \qquad 0 \le k < N'.
$$

In particular, since  $u_0^{new} = u_0^{new} = 1$ ,  $\mathbf{U}$   $\mathbf{U}$ 

$$
h_2(E) = u_2^{\text{left}}(x, E) - h_1(E) u_1^{\text{right}}(x, E).
$$

Since  $n_1$  is known to be purely imaginary, it sumees to show  $u_k^{\ldots}$  and  $u_k^{\text{re}}$  each have the alternating parity property on the index k. Let us show  $u_k^{\text{new}}$  has the alternating parity property; the proof for  $u_k^{\text{new}}$  is totally analogous

Lemma 10 of [FS] relate the canonical solution to the elementary solution of the transport equations in the following manner: if  $\alpha = \sqrt{\omega_0 \sqrt{\omega_1 \omega_2 \sqrt{2}}}$  is the calculation solution of the trains $p \sim p$  equations, and if  $\alpha$  and  $\alpha$  is  $\alpha$  is the elementary solutions, then

$$
u_k(x) = \sum_{l=0}^k w_{k-l,0} \, \tilde{u}_l(x) \,,
$$

where  $w_{kl}$  will be investigated in more detail below. Since the construction of the elementary solutions in [FS] makes it clear  $\tilde{u}_l$  has the alternation particle probabilities on the index limit of the index limit of the problem to showing  $w_{kl}$  has the alternating parity property on the index k. Equivalently, letting  $w_k(x) = \sum_{-3k \leq l} w_{kl} x^{l/2}$ , it suffices to show  $w_k$ has the alternating parity property on the index  $k$ .

Now all that is needed is to take account of the real and purely in a simulation of which are the construction of which  $\mathbb{R}$  arises in the construction of which  $\mathbb{R}$ 

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process which is the wind in the wave when  $\mathbb{N}$  as follows which is the written in the written in the wave we write  $\mathbb{N}$  and  $\mathbb{N}$  a terms of  $h_{kl}^{\prime\prime}, q_{kl}^{\prime\prime}$  and  $h_{kl}$  via the equation

$$
\left(1 + \sum_{k=1}^{N} \lambda^{-k} w_k(x)\right) = \left(\left(1 + \sum_{k=1}^{2N} \sum_{l=2-k}^{3N} h_{kl}^{\#} x^{l/2} \lambda^{-k} + O\left(\lambda^{-\varepsilon N/4}\right)\right) \right. \\
\left. \left(1 + \sum_{k=1}^{N} \sum_{l=-k}^{N} q_{kl}^{\#} x^l \lambda^{-2k} + O\left(\lambda^{-\varepsilon N/4}\right)\right) \right. \\
\left. \left(1 + \sum_{k=1}^{N} \sum_{l=-3k}^{N} \hat{h}_{kl} x^{l/2} \lambda^{-k} + O(\lambda^{-\varepsilon N/4})\right)\right).
$$

To prove what the alternation alternation parally property on the index will be index to the index k we will be want to show both  $h_{kl}^n$  and  $h_{kl}$  have this property on the index k and  $q_{kl}^{\scriptscriptstyle\prime}$  is real. Let us first look at  $h_{kl}^{\scriptscriptstyle\prime}$ . [FS] shows

(8)  

$$
\exp\left(\sum_{k=1}^{N} \sum_{l=-k}^{N} h_{kl} x^{l+3/2} \lambda^{-(2k-1)}\right)
$$

$$
= \left(1 + \sum_{k=1}^{2N} \sum_{l=2-k}^{3N} h_{kl}^{\#} x^{l/2} \lambda^{-k} + O\left(\lambda^{-\varepsilon N/4}\right)\right),
$$

where the right-hand side is a high-order Taylor expansion with remainder. Let us consider more carefully how  $h_{kl}^v$  depends on  $h_{kl}.$  Note that

(9)  

$$
\frac{2 i}{3} \lambda (y_0(x))^{3/2} \sum_{k=1}^N \sum_{l=-k}^N f_{kl}^{\# \#} x^l \lambda^{-2k}
$$

$$
= \sum_{k=1}^N \sum_{l=-k}^N h_{kl} x^{l+3/2} \lambda^{-(2k-1)} + O\left(\lambda^{-\varepsilon N/4}\right).
$$

Since  $y_0(x)$  and  $f_{kl}^{\mu\mu}$  are real,  $h_{kl}$  is purely imaginary since it depends only on these quantities multiplied by  $i$ . Now set

$$
X = \sum_{k=1}^{N} \sum_{l=-k}^{N} h_{kl} x^{l+3/2} \lambda^{-(2k-1)}.
$$

A sufficiently high power of  $\Lambda$  will be  $O\left(\lambda^{-\frac{1}{1-\gamma}-1}\right)$ , so the left-hand side of (8) has a Taylor expansion with remainder. Note that  $X^+$  is purely

imaginary if and only if  $s$  is odd. Since  $X$  contains nothing but odd  $p \in \{1, \ldots, n\}$  and the Taylor expansion terms of the Taylor expansion of th with respect to  $\lambda$  that the coefficients are purely imaginary for all odd powers of  $\lambda$ , real for all even powers of  $\lambda$ . This says precisely that  $h''_{\mu\nu}$ has the alternating parity property on the index  $k$ .

Now let us consider  $q_{kl}^n$ . Quite simply,  $q_{kl}^n$  is real since all the other quantities in the following equation are real

$$
\left(\frac{\partial y_N(x,\lambda)}{\partial x}\right)^{-1/2} (y_N(x,\lambda))^{-1/4}
$$
  
(10) 
$$
= (p(x))^{-1/4} \left(1 + \sum_{k=1}^N \sum_{l=-k}^N q_{kl}^{\#} \lambda^{-2k} + O\left(\lambda^{-\varepsilon N/4}\right)\right).
$$

**Finally, let us consider**  $n_{kl}$ **.** We have that

$$
\left(1 + \sum_{s=1}^{M} c_s \lambda^{-s} x^{-3s/2} \left( \sum_{k=0}^{N} \sum_{l=-k}^{N} h_{kl}^s x^l \lambda^{-2k} + O\left(\lambda^{-\varepsilon N/5}\right) \right) \right)
$$
  
(11)  

$$
= \left(1 + \sum_{k=1}^{N} \sum_{l=-3k}^{N} \hat{h}_{kl} x^{l/2} \lambda^{-k} + O\left(\lambda^{-\varepsilon N/6}\right)\right),
$$

where  $n_{kl}$  is real, and  $c_s$  has the alternating parity property on the index s. This is a consequence of the recurrence relation one finds upon substituting the asymptotic form of the Airey function

$$
A(t) = \text{Re}\left(\frac{e^{\pm i\pi/4}e^{2it^{3/2}/3}}{t^{1/4}}\left(1+\sum_{s=1}^{\infty}c_s t^{-(3/2)s}\right)\right)
$$

into the Airey equation

$$
\frac{d^2}{dy^2}A(y,\lambda) + \lambda^2 y A(y,\lambda) = 0.
$$

Collecting the even and odd powers of on the left-hand side of shows that  $n_{kl}$  has the alternating parity property on the index  $\kappa$ .

Putting what we know about  $h_{kl}^v, q_{kl}^v$  and  $h_{kl}$  into (7) reveals that coecients of even powers of a must involve the coefficients with even to the coefficients with  $\sim$ numbers of  $h_{kl}^{\pi}$ 's and  $h_{kl}$ 's. Note also that the  $q_{kl}^{\pi}$  are always accompanied by even powers of  $\lambda$ . Therefore the coefficients of even powers of  $\lambda$ 

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on the left-hand side of are real On the other hand the coecients of odd powers of  $\lambda$  are purely imaginary. Hence  $w_k$  has the alternating parity property on the index  $k$ .

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## References-

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