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Harmonic measure and approximation of uniformly rectifiable sets

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Abstract. Let $E \subset \mathbb{R}^{n+1}$, $n \geq 1$, be a uniformly rectifiable set of dimension n. We show E that has big pieces of boundaries of a class of domains which satisfy a 2-sided corkscrew condition, and whose connected components are all chord-arc domains (with uniform control of the various constants). As a consequence, we deduce that E has big pieces of sets for which harmonic measure belongs to weak- A_{∞} .

1. Introduction

The results in this paper grew out of a project to prove higher dimensional, quantitative versions of the classical Frigyes and Marcel Riesz theorem [29]. The latter states that for a simply connected domain in the complex plane, with a rectifiable boundary, harmonic measure is absolutely continuous with respect to arclength measure. A quantitative version of this theorem (again in the plane) was obtained by Lavrentiev [24]. We note that some connectivity hypothesis is essential to these results: indeed, Bishop and Jones [5] have presented a counter-example to show that the result of [29] may fail in the absence of sufficient connectivity. Thus, roughly speaking, rectifiability plus connectivity implies absolute continuity, but rectifiability alone does not.

In higher dimensions, quantitative (scale-invariant) versions of the F. and M. Riesz theorem were obtained by Dahlberg [10] in Lipschitz domains, and by David and Jerison [11], and independently, by Semmes [30], in NTA domains with ADR (Ahlfors-David regular) boundaries (all terminology and notation to be defined below). In these quantitative results, the conclusion is that harmonic measure satisfies a scale invariant version of absolute continuity with respect to surface measure, namely that it belongs to the Muckenhoupt class A_{∞} . To draw the analogy with the result of [29] more precisely, we note that recently, in [3] it has been shown that for a domain Ω satisfying a scale invariant connectivity hypothesis (the

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so called "uniform" condition, which is a unilateral version of the NTA property), whose boundary is UR (Uniformly Rectifiable, a scale invariant version of rectifiability which entails, in particular, the ADR property), then in fact Ω is NTA, so that the result of [11] and [30] applies. An earlier, direct proof of the scale invariant absolute continuity of harmonic measure with respect to surface measure, in a uniform domain with a UR boundary, appears in [19]. The converse is also true, see [21].

As noted above, by the counter-example of [5], such results cannot hold in the absence of some connectivity hypothesis. Nonetheless, in this paper, we obtain a structure theorem for uniformly rectifiable sets of co-dimension 1, which yields in particular that the F. and M. Riesz theorem holds for every such set E (viewed as the boundary of an open set $\Omega = \mathbb{R}^{n+1} \setminus E$), in a "big pieces" sense. Our main result (the structure theorem) is the following (our terminology and notation will be defined in the sequel; in particular, however, we let $\mathbb{D}(E)$ denote the collection of "dyadic cubes" on the set E, as per David and Semmes [12] and M. Christ [9]; see Lemma 1.16 below).

Theorem 1.1. Let $E \subset \mathbb{R}^{n+1}$ be a UR (uniformly rectifiable) set of dimension n. Then for each $Q \in \mathbb{D}(E)$ there exists an open set $\widetilde{\Omega} = \widetilde{\Omega}_Q \subset \Omega := \mathbb{R}^{n+1} \setminus E$, with $\operatorname{diam}(\widetilde{\Omega}) \approx \operatorname{diam}(Q)$, such that $\widetilde{\Omega}$ has an ADR (Ahlfors-David regular) boundary, satisfies a 2-sided corkscrew condition, and

(1.1)
$$\sigma(\partial \Omega \cap Q) \gtrsim \sigma(Q).$$

Moreover, each connected component of $\widetilde{\Omega}$ is an NTA domain with ADR boundary. The various NTA, ADR, and implicit constants are uniformly controlled, and depend only on dimension and on the UR character of E.

We remark that, in particular, Theorem 1.1 says that E has big pieces of sets satisfying a 2-sided corkscrew condition, and thus, by a result of David and Jerison [11] (see also [14]), has "big pieces of big pieces of Lipschitz graphs" (BP²(LG); see Definition 1.6 below). Theorem 1.1 therefore yields as an immediate corollary the co-dimension 1 case of a result of Azzam and Schul [4].

Corollary 1.2. Let $E \subset \mathbb{R}^{n+1}$ be a UR (uniformly rectifiable) set of co-dimension 1. Then $E \in BP^2(LG)$.

We should note that, in fact, the result of [4] establishes $BP^2(LG)$ for UR sets in all co-dimensions, whereas our arguments do not address the case of co-dimension greater than 1. On the other hand, in the co-dimension 1 case, our Theorem 1.1 yields extra structure which allows us to obtain estimates for harmonic measure. More precisely, we have the following.

Theorem 1.3. Let $E \subset \mathbb{R}^{n+1}$ be an n-dimensional UR set. Let $\Omega := \mathbb{R}^{n+1} \setminus E$. Then $E \in \text{IBP}(\text{GHME})$, that is, E has "interior big pieces of good harmonic measure estimates", in the following sense: for each $Q \in \mathbb{D}(E)$ there exists an open set $\tilde{\Omega} = \tilde{\Omega}_Q \subset \Omega$, with diam $(\tilde{\Omega}) \approx \text{diam}(Q)$, such that $\tilde{\Omega}$ satisfies a strong 2-sided corkscrew condition along with estimate (1.1), and for each surface ball $\Delta = \Delta(x, r) := B(x, r) \cap \partial \widetilde{\Omega}$, with $x \in \partial \widetilde{\Omega}$ and $r \in (0, \operatorname{diam}(\widetilde{\Omega}))$, and with interior corkscrew point X_{Δ} , it holds that $\omega^{X_{\Delta}} := \omega_{\widetilde{\Omega}}^{X_{\Delta}}$, the harmonic measure for $\widetilde{\Omega}$ with pole at X_{Δ} , belongs to weak- $A_{\infty}(\Delta)$.

Thus, every UR set of co-dimension 1 has big pieces of sets satisfying a quantitative, scale invariant F. and M. Riesz theorem. We remark that this fact actually characterizes uniformly rectifiable sets of co-dimension 1, as the second named author will show in a forthcoming joint paper with J. M. Martell [16].

For the sake of clarity, we note that in the statements of Theorems 1.1 and 1.3, we use the notation

$$\Omega := \mathbb{R}^{n+1} \setminus E \,,$$

where $E \subset \mathbb{R}^{n+1}$ is in particular an *n*-dimensional ADR set (hence closed); thus Ω is open, but need not be a connected domain. The open set $\widetilde{\Omega} \subset \Omega$ in Theorem 1.3 is the one that we construct in Theorem 1.1.

We remark that the weak- A_{∞} conclusion of Theorem 1.3 is in the nature of best possible. Indeed, fix $X \in \Omega$, let $\hat{x} \in E$ be such that $|X - \hat{x}| = \operatorname{dist}(X, E) =:$ $\delta(X)$, and consider the ball $B_X := B(\hat{x}, 10\delta(X))$, and corresponding surface ball $\Delta_X := B_X \cap E$. Choose $Q \in \mathbb{D}(E)$ such that $\operatorname{diam}(Q) \approx \delta(X)$, with $Q \subset \Delta_X$. Our construction in the proof of Theorem 1.1 will yield that $X \in \widetilde{\Omega}_Q$, and in fact is a corkscrew point for a surface ball $\widetilde{\Delta} \subset \partial \widetilde{\Omega}_Q$, of radius $r \approx \delta(X)$, which contains $Q \cap \partial \widetilde{\Omega}_Q$. Consequently, if $\omega_{\widetilde{\Omega}}^X$ were in $A_{\infty}(\widetilde{\Delta})$ (rather than merely weak- A_{∞}), then by the maximum principle, letting ω^X denote harmonic measure for Ω , we would have

(1.2)
$$A \subset \Delta_X, \ H^n(A) \ge (1-\eta)H^n(\Delta_X) \implies \omega^X(A) \ge \omega_{\widetilde{\Omega}}^X(A \cap \widetilde{\Delta}) \ge c > 0,$$

for some uniform positive constant c, provided that $\eta \in (0, 1)$ was sufficiently small depending only on the ADR constants for E and for $\partial \tilde{\Omega}$, and the implicit constant in (1.1). In turn, by the result of [6], (1.2) for every $X \in \Omega$ would then imply that $\omega^X \in \text{weak-}A_{\infty}(\Delta_X)$, which cannot hold in general, by the example of [5].

Let us now define the terms used in the statements of our theorems. Most of the following notions have meaning in co-dimensions greater than 1, but here we shall discuss only the co-dimension 1 case that is of interest to us in the present work.

Definition 1.4 (ADR, aka *Ahlfors–David regular*). We say that a set $E \subset \mathbb{R}^{n+1}$, of Hausdorff dimension n, is ADR if it is closed, and if there is some uniform constant C such that

(1.3)
$$\frac{1}{C}r^n \le \sigma(\Delta(x,r)) \le Cr^n, \quad \forall r \in (0, \operatorname{diam}(E)), \ x \in E,$$

where diam(E) may be infinite. Here, $\Delta(x,r) := E \cap B(x,r)$ is the "surface ball" of radius r, and $\sigma := H^n|_E$ is the "surface measure" on E, where H^n denotes *n*-dimensional Hausdorff measure.

Definition 1.5 (UR, aka uniformly rectifiable). An *n*-dimensional ADR (hence closed) set $E \subset \mathbb{R}^{n+1}$ is UR if and only if it contains "big pieces of Lipschitz images" of \mathbb{R}^n (BPLI). This means that there are positive constants θ and M_0 , such that for each $x \in E$ and each $r \in (0, \operatorname{diam}(E))$, there is a Lipschitz mapping $\rho = \rho_{x,r} : \mathbb{R}^n \to \mathbb{R}^{n+1}$, with Lipschitz constant no larger than M_0 , such that

$$H^n(E \cap B(x,r) \cap \rho\left(\{z \in \mathbb{R}^n : |z| < r\}\right)) \ge \theta r^n$$

We recall that *n*-dimensional rectifiable sets are characterized by the property that they can be covered, up to a set of H^n measure 0, by a countable union of Lipschitz images of \mathbb{R}^n ; thus, BPLI is a quantitative version of this fact.

There are numerous other characterizations of UR sets (many of which remain valid in higher co-dimensions), see [12], [13]. In particular, at least among the class of ADR sets, the UR sets are precisely those for which all "sufficiently nice" singular integrals are L^2 -bounded [12]. In fact, for *n*-dimensional ADR sets in \mathbb{R}^{n+1} , the L^2 boundedness of certain special singular integral operators (the "Riesz transforms"), suffices to characterize uniform rectifiability (see [27] for the case n = 1, and [28] in general). We remark that there exist sets that are ADR (and that even form the boundary of a domain satisfying interior corkscrew and Harnack chain conditions), but that are totally non-rectifiable (e.g., see the construction of Garnett's "4-corners Cantor set" in Chapter 1 of [13]).

Definition 1.6 (BP(S) and BP²(LG)). Let S be a collection of subsets of \mathbb{R}^{n+1} . We say an *n*-dimensional ADR set $E \subset \mathbb{R}^{n+1}$ has big pieces of S (" $E \in BP(S)$ ") if there exists a positive constant θ such that for each $x \in E$ and $r \in (0, \operatorname{diam}(E))$, there is a set $S \in S$ with

$$H^n(B(x,r) \cap E \cap S) \ge \theta r^n.$$

A Lipschitz graph in \mathbb{R}^{n+1} is a set of the form

$$\{y + \rho(y) : y \in P\}$$

where P is an *n*-plane and ρ is a Lipschitz mapping onto a line perpendicular to P. We say that E has big pieces of Lipschitz graphs ("BP(LG)") if there exists a positive constant M_0 such that $E \in BP(S)$, where S is the collection of all Lipschitz graphs with Lipschitz constant no greater than M_0 .

Finally, if E has BP(S), where S is a collection of sets satisfying BP(LG), with uniform bounds on the various constants, then we say that $E \in BP^2(LG)$.

Definition 1.7 ("UR character"). Given a UR set $E \subset \mathbb{R}^{n+1}$, its "UR character" is just the pair of constants (θ, M_0) involved in the definition of uniform rectifiability, along with the ADR constant; or equivalently, the quantitative bounds involved in any particular characterization of uniform rectifiability.

1.1. Further notation and definitions

• We use the letters c, C to denote harmless positive constants, not necessarily the same at each occurrence, which depend only on dimension and the constants

appearing in the hypotheses of the theorems (which we refer to as the "allowable parameters"). We shall also sometimes write $a \leq b$ and $a \approx b$ to mean, respectively, that $a \leq Cb$ and $0 < c \leq a/b \leq C$, where the constants c and C are as above, unless explicitly noted to the contrary. At times, we shall designate by M a particular constant whose value will remain unchanged throughout the proof of a given lemma or proposition, but which may have a different value during the proof of a different lemma or proposition.

- Given a closed set $E \subset \mathbb{R}^{n+1}$, we shall use lower case letters x, y, z, etc., to denote points on E, and capital letters X, Y, Z, etc., to denote generic points in \mathbb{R}^{n+1} (especially those in $\mathbb{R}^{n+1} \setminus E$).
- The open (n+1)-dimensional Euclidean ball of radius r will be denoted B(x,r) when the center x lies on E, or B(X,r) when the center $X \in \mathbb{R}^{n+1} \setminus E$. A "surface ball" is denoted $\Delta(x,r) := B(x,r) \cap E$.
- Given a Euclidean ball B or surface ball Δ , its radius will be denoted r_B or r_{Δ} , respectively.
- Given a Euclidean or surface ball B = B(X, r) or $\Delta = \Delta(x, r)$, its concentric dilate by a factor of $\kappa > 0$ will be denoted $\kappa B := B(X, \kappa r)$ or $\kappa \Delta := \Delta(x, \kappa r)$.
- Given a (fixed) closed set $E \subset \mathbb{R}^{n+1}$, for $X \in \mathbb{R}^{n+1}$, we set $\delta(X) := \operatorname{dist}(X, E)$.
- We let H^n denote *n*-dimensional Hausdorff measure, and let $\sigma := H^n |_E$ denote the "surface measure" on a closed set *E* of co-dimension 1.
- For a Borel set $A \subset \mathbb{R}^{n+1}$, we let 1_A denote the usual indicator function of A, i.e. $1_A(x) = 1$ if $x \in A$, and $1_A(x) = 0$ if $x \notin A$.
- For a Borel set $A \subset \mathbb{R}^{n+1}$, we let int(A) denote the interior of A.
- We shall use the letter I (and sometimes J) to denote a closed (n+1)-dimensional Euclidean dyadic cube with sides parallel to the co-ordinate axes, and we let $\ell(I)$ denote the side length of I. If $\ell(I) = 2^{-k}$, then we set $k_I := k$. Given an ADR set $E \subset \mathbb{R}^{n+1}$, we use Q to denote a dyadic "cube" on E. The latter exist (cf. [12], [9]), and enjoy certain properties which we enumerate in Lemma 1.16 below.

Definition 1.8 (Corkscrew point). Following [23], given an open set $\Omega \subset \mathbb{R}^{n+1}$, and a ball B = B(x, r), with $x \in \partial \Omega$ and $0 < r < \operatorname{diam}(\partial \Omega)$, we say that a point $X = X_B \in \Omega$ is a corkscrew point relative to B with constant c > 0, if there is a ball $B(X, cr) \subset B(x, r) \cap \Omega$.

Definition 1.9 (2-sided corkscrew condition). We say that an open set Ω satisfies the 2-sided corkscrew condition if for some uniform constant c > 0 (the "corkscrew constant"), and for every $x \in \partial\Omega$ and $0 < r < \operatorname{diam}(\partial\Omega)$, there are two corkscrew points X_1 and X_2 relative to B(x, r), with constant c (as in Definition 1.8), and two distinct connected components of $\mathbb{R}^n \setminus \partial\Omega$, \mathcal{O}_1 and \mathcal{O}_2 , with $B_1 = B(X_1, cr) \subset \mathcal{O}_1$ and $B_2 = B(X_2, cr) \subset \mathcal{O}_2$. We recall that this property is called "Condition B" in the work of David and Semmes [14]. We refer to the balls B_1 and B_2 as corkscrew balls. **Definition 1.10** (Strong 2-sided corkscrew condition). We say that an open set Ω satisfies the strong 2-sided corkscrew condition if Ω satisfies the 2-sided corkscrew condition, and one of the balls $B_1 \subset \Omega$ or $B_2 \subset \Omega$.

Definition 1.11 (Harnack chain condition). Following [23], we say that Ω satisfies the Harnack chain condition if there is a uniform constant C such that for every $\rho > 0, \Lambda \ge 1$, and every pair of points $X, X' \in \Omega$ with $\delta(X), \delta(X') \ge \rho$ and $|X - X'| < \Lambda \rho$, there is a chain of open balls $B_1, \ldots, B_N \subset \Omega, N \le C(\Lambda)$, with $X \in B_1, X' \in B_N, B_k \cap B_{k+1} \ne \emptyset$ and $C^{-1} \operatorname{diam}(B_k) \le \operatorname{dist}(B_k, \partial\Omega) \le C \operatorname{diam}(B_k)$. The chain of balls is called a "Harnack chain".

Definition 1.12 (NTA). Again following [23], we say that a domain $\Omega \subset \mathbb{R}^{n+1}$ is NTA ("non-tangentially accessible") if it satisfies the Harnack chain condition, and the strong 2-sided corkscrew condition.

Definition 1.13 (Chord-arc domain). An NTA domain with an ADR boundary is said to be a chord-arc domain.

Definition 1.14 $(A_{\infty} \text{ and weak-}A_{\infty})$. Given an ADR set $E \subset \mathbb{R}^{n+1}$, and a surface ball $\Delta_0 := B_0 \cap E$, we say that a Borel measure μ defined on E belongs to $A_{\infty}(\Delta_0)$ if there are positive constants C and θ such that for each surface ball $\Delta = B \cap E$, with $B \subseteq B_0$, we have

(1.4)
$$\mu(F) \le C\left(\frac{\sigma(F)}{\sigma(\Delta)}\right)^{\theta} \mu(\Delta), \text{ for every Borel set } F \subset \Delta.$$

Similarly, $\mu \in \text{weak-}A_{\infty}(\Delta_0)$, with $\Delta_0 = B_0 \cap \partial\Omega$, if for every $\Delta = B \cap \partial\Omega$ with $2B \subseteq B_0$ we have

(1.5)
$$\mu(F) \le C\left(\frac{\sigma(F)}{\sigma(\Delta)}\right)^{\theta} \mu(2\Delta), \text{ for every Borel set } F \subset \Delta.$$

In the case that $\mu = \omega$ is harmonic measure for an open set Ω satisfying an interior corkscrew condition, setting $E = \partial \Omega$, we shall say that ω belongs to A_{∞} (resp., weak- A_{∞}), if for every surface ball Δ_0 , and for any corkscrew point $X_{\Delta_0} \in \Omega$ relative to Δ_0 , harmonic measure $\omega^{X_{\Delta_0}}$, with pole at X_{Δ_0} , belongs to $A_{\infty}(\Delta_0)$ (resp., weak- $A_{\infty}(\Delta_0)$), in the sense above.

Definition 1.15 (IBP(GHME)). When the collection S in Definition 1.6 consists of boundaries of domains $\widetilde{\Omega} \subset \mathbb{R}^{n+1} \setminus E$, for which the associated harmonic measures belong to weak- A_{∞} , and if the various boundaries $\{\partial \widetilde{\Omega}\}$ are ADR, with uniform control of the ADR and weak- A_{∞} constants, then we say that E has "interior big pieces of good harmonic measure estimates", and we write $E \in \text{IBP}(\text{GHME})$.

Lemma 1.16 (Existence and properties of the "dyadic grid", [12], [13], [9]). Suppose that $E \subset \mathbb{R}^{n+1}$ is a closed n-dimensional ADR set. Then there exist constants $a_0 > 0, \gamma > 0$ and $C_1 < \infty$, depending only on dimension and the ADR constant, such that for each $k \in \mathbb{Z}$, there is a collection of Borel sets ("cubes")

$$\mathbb{D}_k := \{ Q_j^k \subset E : j \in \mathfrak{I}_k \},\$$

where \mathfrak{I}_k denotes some (possibly finite) index set depending on k, satisfying

- (i) $E = \bigcup_j Q_j^k$ for each $k \in \mathbb{Z}$.
- (ii) If $m \ge k$ then either $Q_i^m \subset Q_j^k$ or $Q_i^m \cap Q_j^k = \emptyset$.
- (iii) For each (j,k) and each m < k, there is a unique i such that $Q_j^k \subset Q_i^m$.
- (iv) diam $\left(Q_j^k\right) \leq C_1 2^{-k}$.
- (v) Each Q_i^k contains some "surface ball" $\Delta(x_i^k, a_0 2^{-k}) := B(x_i^k, a_0 2^{-k}) \cap E$.
- (vi) $H^n(\{x \in Q_j^k : \operatorname{dist}(x, E \setminus Q_j^k) \le \varrho \, 2^{-k}\}) \le C_1 \, \varrho^{\gamma} \, H^n(Q_j^k)$, for all k, j and for all $\varrho \in (0, a_0)$.

A few remarks are in order concerning this lemma.

- In the setting of a general space of homogeneous type, this lemma has been proved by Christ [9], with the dyadic parameter 1/2 replaced by some constant $\delta \in (0, 1)$. In fact, one may always take $\delta = 1/2$ (see the proof of Proposition 2.12 in [22]). In the presence of the Ahlfors–David property (1.3), the result already appears in [12], [13].
- For our purposes, we may ignore those $k \in \mathbb{Z}$ such that $2^{-k} \gtrsim \text{diam}(E)$, in the case that the latter is finite.
- We shall denote by $\mathbb{D} = \mathbb{D}(E)$ the collection of all relevant Q_j^k , i.e.,

$$\mathbb{D} := \cup_k \mathbb{D}_k,$$

where, if diam(E) is finite, the union runs over those k such that $2^{-k} \leq \text{diam}(E)$.

• Properties (iv) and (v) imply that for each cube $Q \in \mathbb{D}_k$, there is a point $x_Q \in E$, a Euclidean ball $B(x_Q, r)$ and a surface ball $\Delta(x_Q, r) := B(x_Q, r) \cap E$ such that $r \approx 2^{-k} \approx \operatorname{diam}(Q)$ and

(1.6)
$$\Delta(x_Q, r) \subset Q \subset \Delta(x_Q, Cr),$$

for some uniform constant C. We shall denote this ball and surface ball by

(1.7)
$$B_Q := B(x_Q, r), \qquad \Delta_Q := \Delta(x_Q, r),$$

and we shall refer to the point x_Q as the "center" of Q.

- For each cube $Q \in \mathbb{D}$, we let X_Q be a corkscrew point relative to B_Q , and refer to this as a corkscrew point relative to Q. Such a corkscrew point exists, since Eis *n*-dimensional ADR (with the constant *c* in Definition 1.8 depending only on dimension and the ADR constants).
- For a dyadic cube $Q \in \mathbb{D}_k$, we shall set $\ell(Q) = 2^{-k}$, and we shall refer to this quantity as the "length" of Q. Evidently, $\ell(Q) \approx \operatorname{diam}(Q)$.
- For a dyadic cube $Q \in \mathbb{D}$, we let k(Q) denote the "dyadic generation" to which Q belongs, i.e., we set k = k(Q) if $Q \in \mathbb{D}_k$; thus, $\ell(Q) = 2^{-k(Q)}$.

2. A bilateral corona decomposition and corona type approximation by chord-arc domains

In this section, we state a bilateral variant of the "corona decomposition" of David and Semmes [12], [13]. The bilateral version was proved in [20], Lemma 2.2. We first recall the notions of "coherency" and "semi-coherency":

Definition 2.1 ([13]). Let $\mathbf{S} \subset \mathbb{D}(E)$. We say that \mathbf{S} is "coherent" if the following conditions hold:

- (a) **S** contains a unique maximal element $Q(\mathbf{S})$ which contains all other elements of **S** as subsets.
- (b) If Q belongs to **S**, and if $Q \subset \widetilde{Q} \subset Q(\mathbf{S})$, then $\widetilde{Q} \in \mathbf{S}$.
- (c) Given a cube $Q \in \mathbf{S}$, either all of its children belong to \mathbf{S} , or none of them do.

We say that \mathbf{S} is "semi-coherent" if only conditions (a) and (b) hold.

The bilateral "corona decomposition" is as follows.

Lemma 2.2 (Lemma 2.2 in [20]). Suppose that $E \subset \mathbb{R}^{n+1}$ is n-dimensional UR. Then given any positive constants $\eta \ll 1$ and $K \gg 1$, there is a disjoint decomposition $\mathbb{D}(E) = \mathcal{G} \cup \mathcal{B}$, satisfying the following properties.

- The "good" collection G is further subdivided into disjoint stopping time regimes, such that each such regime S is coherent (cf. Definition 2.1).
- (2) The "bad" cubes, as well as the maximal cubes $Q(\mathbf{S})$, satisfy a Carleson packing condition:

$$\sum_{Q' \subset Q, \, Q' \in \mathcal{B}} \sigma(Q') + \sum_{\mathbf{S}: Q(\mathbf{S}) \subset Q} \sigma(Q(\mathbf{S})) \leq C_{\eta, K} \, \sigma(Q) \,, \quad \forall Q \in \mathbb{D}(E) \,.$$

(3) For each **S**, there is a Lipschitz graph $\Gamma_{\mathbf{S}}$, with Lipschitz constant at most η , such that, for every $Q \in \mathbf{S}$,

(2.1)
$$\sup_{x \in \Delta_Q^*} \operatorname{dist}(x, \Gamma_{\mathbf{S}}) + \sup_{y \in B_Q^* \cap \Gamma_{\mathbf{S}}} \operatorname{dist}(y, E) < \eta \,\ell(Q) \,,$$

where
$$B_Q^* := B(x_Q, K\ell(Q))$$
 and $\Delta_Q^* := B_Q^* \cap E$.

In this section, we construct the same domains as in [20], for each stopping time regime **S** in Lemma 2.2, a pair of NTA domains $\Omega_{\mathbf{S}}^{\pm}$, with ADR boundaries, which provide a good approximation to E, at the scales within **S**, in some appropriate sense. To be a bit more precise, $\Omega_{\mathbf{S}} := \Omega_{\mathbf{S}}^+ \cup \Omega_{\mathbf{S}}^-$ will be constructed as a sawtooth region relative to some family of dyadic cubes, and the nature of this construction will be essential to the dyadic analysis that we will use below. In this section, we follow essentially verbatim the construction in [20], which we reproduce here for the reader's convenience.

We first discuss some preliminary matters. We shall utilize the notation and constructions of [20] (and essentially that of [19] and [21]).

Let $\mathcal{W} = \mathcal{W}(\mathbb{R}^{n+1} \setminus E)$ denote a collection of (closed) dyadic Whitney cubes of $\mathbb{R}^{n+1} \setminus E$, so that the cubes in \mathcal{W} have pairwise disjoint interiors and cover $\mathbb{R}^{n+1} \setminus E$, and furthermore satisfy

(2.2)
$$4 \operatorname{diam}(I) \le \operatorname{dist}(4I, E) \le \operatorname{dist}(I, E) \le 40 \operatorname{diam}(I), \quad \forall I \in \mathcal{W}$$

(just dyadically divide the standard Whitney cubes, as constructed in [31], Chapter VI, into cubes with side length 1/8 as large) and also

$$(1/4) \operatorname{diam}(I_1) \leq \operatorname{diam}(I_2) \leq 4 \operatorname{diam}(I_1),$$

whenever I_1 and I_2 touch.

Let E be an n -dimensional ADR set and pick two parameters $\eta \ll 1$ and $K \gg 1.$ Define

(2.3)
$$\mathcal{W}_Q^0 := \{ I \in \mathcal{W} : \eta^{1/4} \ell(Q) \le \ell(I) \le K^{1/2} \ell(Q), \operatorname{dist}(I,Q) \le K^{1/2} \ell(Q) \}.$$

Remark 2.3. We note that \mathcal{W}_Q^0 is non-empty, provided that we choose η small enough, and K large enough, depending only on dimension and the ADR constant of E.

Assume now that E is UR and make the corresponding bilateral corona decomposition of Lemma 2.2 with $\eta \ll 1$ and $K \gg 1$. Given $Q \in \mathbb{D}(E)$, for this choice of η and K, we set (as above) $B_Q^* := B(x_Q, K\ell(Q))$, where we recall that x_Q is the "center" of Q (see (1.6)–(1.7)). For a fixed stopping time regime **S**, we choose a co-ordinate system so that $\Gamma_{\mathbf{S}} = \{(z, \varphi_{\mathbf{S}}(z)) : z \in \mathbb{R}^n\}$, where $\varphi_{\mathbf{S}} : \mathbb{R}^n \mapsto \mathbb{R}$ is a Lipschitz function with $\|\varphi\|_{\text{Lip}} \leq \eta$.

Claim 2.4. If $Q \in \mathbf{S}$, and $I \in \mathcal{W}_Q^0$, then I lies either above or below $\Gamma_{\mathbf{S}}$. Moreover, $\operatorname{dist}(I, \Gamma_{\mathbf{S}}) \geq \eta^{1/2} \ell(Q)$ (and therefore, by (2.1), $\operatorname{dist}(I, \Gamma_{\mathbf{S}}) \approx \operatorname{dist}(I, E)$, with implicit constants that may depend on η and K).

Proof of Claim 2.4. Suppose by way of contradiction that $\operatorname{dist}(I, \Gamma_{\mathbf{S}}) \leq \eta^{1/2} \ell(Q)$. Then we may choose $y \in \Gamma_{\mathbf{S}}$ such that

$$\operatorname{dist}(I, y) \le \eta^{1/2} \ell(Q) \,.$$

By construction of \mathcal{W}_Q^0 , it follows that for all $Z \in I$, $|Z-y| \lesssim K^{1/2}\ell(Q)$. Moreover, $|Z - x_Q| \lesssim K^{1/2}\ell(Q)$, and therefore $|y - x_Q| \lesssim K^{1/2}\ell(Q)$. In particular, $y \in B_Q^* \cap \Gamma_{\mathbf{S}}$, so by (2.1), dist $(y, E) \leq \eta \ell(Q)$. On the other hand, choosing $Z_0 \in I$ such that $|Z_0-y| = \text{dist}(I, y) \leq \eta^{1/2}\ell(Q)$, we obtain dist $(I, E) \leq 2\eta^{1/2}\ell(Q)$. For η small, this contradicts the Whitney construction, since dist $(I, E) \approx \ell(I) \geq \eta^{1/4}\ell(Q)$. \Box

Next, given $Q \in \mathbf{S}$, we augment \mathcal{W}_Q^0 . We split $\mathcal{W}_Q^0 = \mathcal{W}_Q^{0,+} \cup \mathcal{W}_Q^{0,-}$, where $I \in \mathcal{W}_Q^{0,+}$ if I lies above $\Gamma_{\mathbf{S}}$, and $I \in \mathcal{W}_Q^{0,-}$ if I lies below $\Gamma_{\mathbf{S}}$. Choosing K large and η small enough, by (2.1), we may assume that both $\mathcal{W}_Q^{0,\pm}$ are non-empty. We focus on $\mathcal{W}_Q^{0,+}$, as the construction for $\mathcal{W}_Q^{0,-}$ is the same. For each $I \in \mathcal{W}_Q^{0,+}$,

let X_I denote the center of I. Fix one particular $I_0 \in \mathcal{W}_Q^{0,+}$, with center $X_Q^+ := X_{I_0}$. Let \tilde{Q} denote the dyadic parent of Q, unless $Q = Q(\mathbf{S})$; in the latter case we simply set $\tilde{Q} = Q$. Note that $\tilde{Q} \in \mathbf{S}$, by the coherency of \mathbf{S} . By Claim 2.4, for each Iin $\mathcal{W}_Q^{0,+}$, or in $\mathcal{W}_{\tilde{Q}}^{0,+}$, we have

(2.4)
$$\operatorname{dist}(I, E) \approx \operatorname{dist}(I, Q) \approx \operatorname{dist}(I, \Gamma_{\mathbf{S}}),$$

where the implicit constants may depend on η and K. Thus, for each such I, we may fix a Harnack chain, call it \mathcal{H}_I , relative to the Lipschitz domain

$$\Omega_{\Gamma_{\mathbf{S}}}^{+} := \left\{ (x,t) \in \mathbb{R}^{n+1} : t > \varphi_{\mathbf{S}}(x) \right\} \,,$$

connecting X_I to X_Q^+ . By the bilateral approximation condition (2.1), the definition of \mathcal{W}_Q^0 , and the fact that $K^{1/2} \ll K$, we may construct this Harnack chain so that it consists of a bounded number of balls (depending on η and K), and stays a distance at least $c\eta^{1/2}\ell(Q)$ away from $\Gamma_{\mathbf{S}}$ and from E. We let $\mathcal{W}_Q^{*,+}$ denote the set of all $J \in \mathcal{W}$ which meet at least one of the Harnack chains \mathcal{H}_I , with $I \in \mathcal{W}_Q^{0,+} \cup \mathcal{W}_{\widetilde{Q}}^{0,+}$ (or simply $I \in \mathcal{W}_Q^{0,+}$, if $Q = Q(\mathbf{S})$), i.e.,

(2.5)
$$\mathcal{W}_Q^{*,+} := \left\{ J \in \mathcal{W} : \exists I \in \mathcal{W}_Q^{0,+} \cup \mathcal{W}_{\widetilde{Q}}^{0,+} \text{ for which } \mathcal{H}_I \cap J \neq \emptyset \right\},$$

where as above, \widetilde{Q} is the dyadic parent of Q, unless $Q = Q(\mathbf{S})$, in which case we simply set $\widetilde{Q} = Q$ (so the union is redundant). We observe that, in particular, each $I \in \mathcal{W}_Q^{0,+} \cup \mathcal{W}_{\widetilde{Q}}^{0,+}$ meets \mathcal{H}_I , by definition, and therefore

(2.6)
$$\mathcal{W}_Q^{0,+} \cup \mathcal{W}_{\widetilde{Q}}^{0,+} \subset \mathcal{W}_Q^{*,+}$$

Of course, we may construct $\mathcal{W}_Q^{*,-}$ analogously. We then set

$$\mathcal{W}_Q^* := \mathcal{W}_Q^{*,+} \cup \mathcal{W}_Q^{*,-}$$
 .

It follows from the construction of the augmented collections $\mathcal{W}_Q^{*,\pm}$ that there are uniform constants c and C such that

(2.7)
$$c\eta^{1/2}\ell(Q) \le \ell(I) \le CK^{1/2}\ell(Q), \quad \forall I \in \mathcal{W}_Q^*,$$
$$\operatorname{dist}(I,Q) \le CK^{1/2}\ell(Q), \quad \forall I \in \mathcal{W}_Q^*.$$

Observe that $\mathcal{W}_Q^{*,\pm}$ and hence also \mathcal{W}_Q^* have been defined for any Q that belongs to some stopping time regime **S**, that is, for any Q belonging to the "good" collection \mathcal{G} of Lemma 2.2. On the other hand, we have defined \mathcal{W}_Q^0 for arbitrary $Q \in \mathbb{D}(E)$.

We now set

(2.8)
$$\mathcal{W}_Q := \begin{cases} \mathcal{W}_Q^*, \ Q \in \mathcal{G}, \\ \mathcal{W}_Q^0, \ Q \in \mathcal{B}, \end{cases}$$

and for $Q \in \mathcal{G}$ we shall henceforth simply write \mathcal{W}_Q^{\pm} in place of $\mathcal{W}_Q^{*,\pm}$.

Next, we choose a small parameter $\tau_0 > 0$, so that for any $I \in \mathcal{W}$, and any $\tau \in (0, \tau_0]$, the concentric dilate $I^*(\tau) := (1 + \tau)I$ still satisfies the Whitney property

(2.9)
$$\operatorname{diam} I \approx \operatorname{diam} I^*(\tau) \approx \operatorname{dist} \left(I^*(\tau), E\right) \approx \operatorname{dist}(I, E), \quad 0 < \tau \le \tau_0.$$

Moreover, for $\tau \leq \tau_0$ small enough, and for any $I, J \in \mathcal{W}$, we have that $I^*(\tau)$ meets $J^*(\tau)$ if and only if I and J have a boundary point in common, and that, if $I \neq J$, then $I^*(\tau)$ misses (3/4)J. Given an arbitrary $Q \in \mathbb{D}(E)$, we may define an associated Whitney region U_Q (not necessarily connected), as follows:

(2.10)
$$U_Q = U_{Q,\tau} := \bigcup_{I \in \mathcal{W}_Q} I^*(\tau)$$

For later use, it is also convenient to introduce some fattened version of U_Q : if $0 < \tau \leq \tau_0/2$,

(2.11)
$$\widehat{U}_Q = U_{Q,2\tau} := \bigcup_{I \in \mathcal{W}_Q} I^*(2\tau).$$

If $Q \in \mathcal{G}$, then U_Q splits into exactly two connected components

(2.12)
$$U_Q^{\pm} = U_{Q,\tau}^{\pm} := \bigcup_{I \in \mathcal{W}_Q^{\pm}} I^*(\tau)$$

When the particular choice of $\tau \in (0, \tau_0]$ is not important, for the sake of notational convenience, we may simply write I^* , U_Q , and U_Q^{\pm} in place of $I^*(\tau)$, $U_{Q,\tau}$, and $U_{Q,\tau}^{\pm}$. We note that for $Q \in \mathcal{G}$, each U_Q^{\pm} is Harnack chain connected, by construction (with constants depending on the implicit parameters τ, η and K); moreover, for a fixed stopping time regime **S**, we claim that if Q' is a child of Q, with both $Q', Q \in \mathbf{S}$, then $U_{Q'}^+ \cup U_Q^+$ is Harnack chain connected, and similarly for $U_{Q'}^- \cup U_Q^-$. Indeed, by construction (see (2.4), the ensuing discussion, and (2.5), as well as (2.8) and (2.10)), each of $U_{Q'}^+$ and U_Q^+ is a union of fattened Whitney boxes, such that the center of each such fattened Whitney box I^* may be connected via a Harnack path to X_Q^+ ; moreover, any two such I_1^* and I_2^* either overlap, or are separated by a distance comparable to $\ell(Q)$. The claimed Harnack chain property follows readily.

We may also define "Carleson boxes" relative to any $Q \in \mathbb{D}(E)$, by

(2.13)
$$T_Q = T_{Q,\tau} := \operatorname{int} \left(\bigcup_{Q' \in \mathbb{D}_Q} U_{Q,\tau} \right),$$

where

$$(2.14) \qquad \qquad \mathbb{D}_Q := \{Q' \in \mathbb{D}(E) : Q' \subset Q\} \ .$$

Let us note that we may choose K large enough so that, for every Q,

(2.15)
$$T_Q \subset B_Q^* := B\left(x_Q, K\ell(Q)\right) \,.$$

For future reference, we also introduce dyadic sawtooth regions as follows. Given a family \mathcal{F} of disjoint cubes $\{Q_j\} \subset \mathbb{D}$, we define the "global discretized sawtooth" relative to \mathcal{F} by

$$(2.16) \qquad \qquad \mathbb{D}_{\mathcal{F}} := \mathbb{D} \setminus \bigcup_{\mathcal{F}} \mathbb{D}_{Q_i},$$

i.e., $\mathbb{D}_{\mathcal{F}}$ is the collection of all $Q \in \mathbb{D}$ that are not contained in any $Q_j \in \mathcal{F}$. Given some fixed cube Q, the "local discretized sawtooth" relative to \mathcal{F} by

(2.17)
$$\mathbb{D}_{\mathcal{F},Q} := \mathbb{D}_Q \setminus \bigcup_{\mathcal{F}} \mathbb{D}_{Q_j} = \mathbb{D}_{\mathcal{F}} \cap \mathbb{D}_Q.$$

Note that in this way $\mathbb{D}_Q = \mathbb{D}_{\emptyset,Q}$.

Similarly, we may define geometric sawtooth regions as follows. Given a family \mathcal{F} of disjoint cubes $\{Q_j\} \subset \mathbb{D}$, we define the "global sawtooth" and the "local sawtooth" relative to \mathcal{F} by respectively

(2.18)
$$\Omega_{\mathcal{F}} := \operatorname{int}\left(\bigcup_{Q' \in \mathbb{D}_{\mathcal{F}}} U_{Q'}\right), \quad \Omega_{\mathcal{F},Q} := \operatorname{int}\left(\bigcup_{Q' \in \mathbb{D}_{\mathcal{F},Q}} U_{Q'}\right).$$

Notice that $\Omega_{\emptyset,Q} = T_Q$. For the sake of notational convenience, given a pairwise disjoint family $\mathcal{F} \in \mathbb{D}$, and a cube $Q \in \mathbb{D}_{\mathcal{F}}$, we set

(2.19)
$$\mathcal{W}_{\mathcal{F}} := \bigcup_{Q' \in \mathbb{D}_{\mathcal{F}}} \mathcal{W}_{Q'}, \quad \mathcal{W}_{\mathcal{F},Q} := \bigcup_{Q' \in \mathbb{D}_{\mathcal{F},Q}} \mathcal{W}_{Q'},$$

so that in particular, we may write

(2.20)
$$\Omega_{\mathcal{F},Q} = \operatorname{int}\left(\bigcup_{I \in \mathcal{W}_{\mathcal{F},Q}} I^*\right).$$

Remark 2.5. Given a stopping time regime **S** as in Lemma 2.2, for any semicoherent subregime (cf. Definition 2.1) $\mathbf{S}' \subset \mathbf{S}$ (including, of course, **S** itself), we now set

(2.21)
$$\Omega_{\mathbf{S}'}^{\pm} = \operatorname{int}\left(\bigcup_{Q \in \mathbf{S}'} U_Q^{\pm}\right),$$

and let $\Omega_{\mathbf{S}'} := \Omega_{\mathbf{S}'}^+ \cup \Omega_{\mathbf{S}'}^-$. Note that implicitly, $\Omega_{\mathbf{S}'}$ depends upon τ (since U_Q^{\pm} has such dependence). When it is necessary to consider the value of τ explicitly, we shall write $\Omega_{\mathbf{S}'}(\tau)$.

It is helpful to introduce some terminology now whose utility will become clear later. Let $Q \in \mathbb{D}$ define the following

(2.22)
$$\mathcal{I}(Q) := \{I \in \mathcal{W} : I \cap T_Q \neq \emptyset\}$$

and also

(2.23)
$$V(Q) = \operatorname{int}\left(\bigcup_{I \in \mathcal{I}(Q)} I^*\right).$$

We note that, trivially, $T_Q \subset V(Q)$. Notice also that if $int(I^*) \subset V(Q)$ then

- (2.24) $\operatorname{dist}(I^*, Q) \lesssim \ell(Q)$ and
- (2.25) $\ell(I) \lesssim \ell(Q).$

Lemma 2.6. Let $Q_1, Q_2 \in \mathbb{D}(E)$. If U_{Q_1} meets U_{Q_2} then

(2.26)
$$\operatorname{dist}(Q_1, Q_2) \lesssim \min\{\ell(Q_1), \ell(Q_2)\}$$

with implicit constant depending only on K, η , and dimension. Moreover there exists a constant Υ depending only on K, η , and dimension such that if

$$\operatorname{dist}(Q_1, Q_2) > \Upsilon \max\{\ell(Q_1), \ell(Q_2)\}$$

then

(2.27)
$$\overline{V(Q_1)} \cap \overline{V(Q_2)} = \emptyset.$$

Proof. Suppose that U_{Q_1} meets U_{Q_2} then we have that there exists a cube $I_1^* \in U_{Q_1}$ and $I_2^* \in U_{Q_2}$ such that $I_1^* \cap I_2^* \neq \emptyset$. Since I_1^* and I_2^* are Whitney cubes that meet we have that

(2.28)
$$\ell(I_1) \approx \ell(I_2).$$

Then by construction of U_{Q_1} and $U_{Q'_2}$ we have

(2.29)
$$\ell(Q_1) \approx \ell(Q_2),$$

(2.30)
$$\operatorname{dist}(Q_1, I_1^*) \lesssim \ell(Q_1),$$

(2.31)
$$\operatorname{dist}(Q_2, I_2^*) \lesssim \ell(Q_2^*).$$

So that (2.28), (2.29), (2.30), and (2.31) yield

 $\operatorname{dist}(Q_1, Q_2) \lesssim \ell(Q_1).$

One sees (2.24) and (2.25) imply, if $\overline{V(Q_1)}$ meets $\overline{V(Q_2)}$,

(2.32) $\operatorname{dist}(Q_1, Q_2) \lesssim \max\{\ell(Q_1), \ell(Q_2)\};\$

this shows (2.27).

In addition we can even put a distance between $V(Q_1)$ and $V(Q_2)$ on the order of $\max\{\ell(Q_1), \ell(Q_2)\}$ by making Υ larger. \Box

3. Carleson measures: proof of the Theorem 1.1

The proof will utilize the method of "extrapolation of Carleson measures". This method was first used by J. L. Lewis [25], whose work was influenced by the corona construction of Carleson [7] and the work of Carleson and Garnett [8] (see also [17], [1], [2], [18], [19].) We will apply this method to the (discrete) packing measure from the bilateral corona decomposition. Let $E \subset \mathbb{R}^{n+1}$ be a UR set of co-dimension 1. We fix positive numbers $\eta \ll 1$, and $K \gg 1$, and for these values of η and K, we perform the bilateral corona decomposition of $\mathbb{D}(E)$ guaranteed by Lemma 2.2. Let $\mathcal{M} := \{Q(\mathbf{S})\}_{\mathbf{S}}$ denotes the collection of cubes which are the maximal elements of the stopping time regimes in \mathcal{G} .

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Given a cube $Q \in \mathbb{D}(E)$, we set

(3.1)
$$\alpha_Q := \begin{cases} \sigma(Q) \,, & \text{if } Q \in \mathcal{M} \cup \mathcal{B}, \\ 0 \,, & \text{otherwise.} \end{cases}$$

Given any collection $\mathbb{D}' \subset \mathbb{D}(E)$, we define

(3.2)
$$\mathfrak{m}(\mathbb{D}') := \sum_{Q \in \mathbb{D}'} \alpha_Q.$$

We recall that \mathbb{D}_Q is the "discrete Carleson region relative to Q", defined in (2.14). Then by Lemma 2.2(2), we have the discrete Carleson measure estimate

(3.3)
$$\mathfrak{m}(\mathbb{D}_Q) := \sum_{Q' \subset Q, \, Q' \in \mathcal{B}} \sigma(Q') + \sum_{\mathbf{S}: Q(\mathbf{S}) \subset Q} \sigma(Q(\mathbf{S})) \leq C_{\eta, K} \, \sigma(Q) \,, \quad \forall Q \in \mathbb{D}(E) \,.$$

Given a family $\mathcal{F} := \{Q_j\} \subset \mathbb{D}(E)$ of pairwise disjoint cubes, we recall that the "discrete sawtooth" $\mathbb{D}_{\mathcal{F}}$ is the collection of all cubes in $\mathbb{D}(E)$ that are not contained in any $Q_j \in \mathcal{F}$ (cf. (2.16)), and define the "restriction of \mathfrak{m} to the sawtooth $\mathbb{D}_{\mathcal{F}}$ " by

(3.4)
$$\mathfrak{m}_{\mathcal{F}}(\mathbb{D}') := \mathfrak{m}(\mathbb{D}' \cap \mathbb{D}_{\mathcal{F}}) = \sum_{Q \in \mathbb{D}' \setminus (\cup_{\mathcal{F}} \mathbb{D}_{Q_j})} \alpha_Q \, .$$

We define the usual Carleson norm, and the Carleson norm restricted to a cube Q, as follows:

$$\|\mathfrak{m}\|_{\mathcal{C}} := \sup_{Q \in \mathbb{D}(E)} \frac{\mathfrak{m}(\mathbb{D}_Q)}{\sigma(Q)}, \quad \|\mathfrak{m}\|_{\mathcal{C}(Q)} := \sup_{Q' \subset Q} \frac{\mathfrak{m}(\mathbb{D}_{Q'})}{\sigma(Q')}.$$

Note that by the way that we have defined \mathfrak{m} , we have

$$(3.5) $\|\mathfrak{m}\|_{\mathcal{C}} \le C_{\eta,K}$$$

The following lemma will be one of two crucial lemmas in proving Theorem 1.1.

Lemma 3.1 (Lemma 7.2 in [19]). Suppose that E is ADR. Fix $Q \in \mathbb{D}(E)$ and \mathfrak{m} as above. Let $a \ge 0$ and b > 0, and suppose that $\mathfrak{m}(\mathbb{D}_Q) \le (a+b) \sigma(Q)$. Then there is a family $\mathcal{F} = \{Q_j\} \subset \mathbb{D}_Q$ of pairwise disjoint cubes, and a constant C depending only on dimension and the ADR constant such that

(3.6)
$$\|\mathfrak{m}_{\mathcal{F}}\|_{\mathcal{C}(Q)} \le C b,$$

and

(3.7)
$$\sigma(B) \le \frac{a+b}{a+2b} \,\sigma(Q) \,,$$

where B is the union of those $Q_j \in \mathcal{F}$ such that $\mathfrak{m}(\mathbb{D}_{Q_j} \setminus \{Q_j\}) > a \sigma(Q_j)$.

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The other crucial lemma is the following.

Lemma 3.2 (Lemma 3.24 in [20]). Let **S** be a given stopping time regime as in Lemma 2.2, and let **S'** be any nonempty, semi-coherent subregime of **S**. Then for $0 < \tau \leq \tau_0$, with τ_0 small enough, each of $\Omega_{\mathbf{S'}}^{\pm}$ is an NTA domain, with ADR boundary. The constants in the NTA and ADR conditions depend only on n, τ , η , K, and the ADR/UR constants for E.

The following standard covering type lemma will be required.

Lemma 3.3. Fix $Q_0 \in \mathbb{D}(E)$ and let $\mathcal{F} = \{Q_j\} \subset \mathbb{D}_{Q_0}$ be any pairwise disjoint family of cubes. Then for any positive constant κ we may find a sub-collection $\mathcal{G} = \{\widetilde{Q}_i\} \subset \mathcal{F}$ with the following properties:

(3.8)
$$\sigma(\cup_{\mathcal{G}} Q_i) \ge C\sigma(\cup_{\mathcal{F}} Q_j),$$

(3.9)
$$\operatorname{dist}(\widetilde{Q}_i, \widetilde{Q}_k) \ge \kappa \max\{\ell(\widetilde{Q}_i), \ell(\widetilde{Q}_k)\},$$

where C depends on κ , dimension and ADR.

Proof of Theorem 1.1. First we fix η and K so that that Lemma 3.2 holds. The proof will follow by induction. For any $a \ge 0$ we have the induction hypothesis H(a), defined in the following way.

H(a): There exists $\eta_a > 0$ such that for all $Q_0 \in \mathbb{D}(E)$ satisfying $\mathfrak{m}(\mathbb{D}_{Q_0}) \leq a\sigma(Q_0)$, there is a collection $\mathbb{D}' \subseteq \mathbb{D}_{Q_0}$, and an open set $\widetilde{\Omega}$ of the form

(3.10)
$$\Omega := \operatorname{int} \left(\bigcup_{Q_j \in \mathbb{D}'} U_{Q_j} \right) \subset T_{Q_0} \,,$$

which has an ADR boundary and satisfies the strong 2-sided corkscrew condition for open sets with corkscrew balls lying in $V(Q_0)$ (see (2.23)), and in addition

(3.11)
$$\sigma(\partial \hat{\Omega} \cap Q_0) \ge \eta_a \sigma(Q_0).$$

Moreover, each connected component of $\widetilde{\Omega}$ is an NTA domain with ADR boundary with uniform constants possibly depending on a.

To prove the theorem it is enough to show that $H(M_0)$ holds where M_0 is the Carleson norm of \mathfrak{m} , that is, $M_0 = C_{\eta,K}$ in Lemma 2.2. We do this by showing first that H(0) holds and then that H(a) implies H(a+b) for some fixed constant b depending only on dimension and the ADR constants. This way we will only use finitely many steps to get to $H(M_0)$. We set $b = \gamma/C$, where C is the constant in (3.6) and γ is a small positive number to be chosen.

The fact that H(0) holds is somewhat trivial since this would imply that Q_0 and all of its descendants are in the same stopping time regime and therefore Q_0 coincides with a Lipschitz graph. Here we can also directly apply the results in [20]. We are then left with showing that H(a) implies H(a + b).

Proof that $H(a) \implies H(a+b)$. Suppose that H(a) holds. Let $Q_0 \in \mathbb{D}(E)$ be such that $\mathfrak{m}(\mathbb{D}_{Q_0}) \leq (a+b)\sigma(Q_0)$. First let C_2 be an integer so large that if $Q_1 \subseteq Q_2$ with $k(Q_2) + C_2 - 5 < k(Q_1)$ then we have that $U_{Q_2} \cap V(Q_1) = \emptyset$. We obtain via Lemma 3.1 a collection $\mathcal{F} = \{Q_j\}_{j=1}^{\infty}$ such that

(3.12)
$$\|\mathfrak{m}_{\mathcal{F}}\|_{C(Q_0)} \le C \, b = \gamma \quad \text{and}$$

(3.13)
$$\sigma(B) \le \frac{a+b}{a+2b} \,\sigma(Q_0) \,,$$

where B is the union of those Q_j in \mathcal{F} such that $\mathfrak{m}(\mathbb{D}_{Q_j} \setminus \{Q_j\}) > a\sigma(Q_j)$, call this collection \mathcal{F}_{bad} . Define $\mathcal{F}_{\text{good}} := \mathcal{F} \setminus \mathcal{F}_{\text{bad}}$. Then by pigeon-holing, for each $Q_j \in \mathcal{F}_{\text{good}}$ we may find a child of Q_j to which we can apply the induction hypothesis H(a).

Iterating the pigeon-holing argument, we may find a cube Q'_j that is C_2 generations down from Q_j (i.e., so that $\ell(Q'_j) = 2^{-C_2}\ell(Q_j)$), to which we may apply the induction hypothesis.

Remark 3.4. Choosing γ , small enough in (3.12) (in fact, $\gamma = 1/2$ will suffice), we obtain that $\mathbb{D}_{\mathcal{F},Q_0}$ does not contain any $Q' \in \mathcal{M} \cup \mathcal{B}$, where $\mathcal{M} := \{Q(\mathbf{S})\}_{\mathbf{S}}$. Fixing such a γ , we find therefore that every $Q' \in \mathbb{D}_{\mathcal{F},Q_0}$ is a good cube, and moreover all such Q' belong to the same \mathbf{S} , a stopping time regime as in Lemma 2.2. Thus, $\mathbb{D}_{\mathcal{F},Q_0}$ is a semi-coherent (see Definition 2.1) subregime of that \mathbf{S} , and therefore $\Omega_{\mathcal{F},Q_0}$ splits into two disjoint NTA domains with ADR boundary, by Lemma 3.2.

We may clearly assume that $a < M_0$. Set $\eta := 1 - \frac{M_0 + b}{M_0 + 2b}$ and $A := Q \setminus (\bigcup_{\mathcal{F}} Q_j)$ and $G := (\bigcup_{\mathcal{F}} Q_j) \setminus B$. Then by (3.13), we see immediately that

$$\sigma(A \cup G) \ge \eta \sigma(Q_0).$$

If $\sigma(A) > (\eta/2)\sigma(Q_0)$, we set $\widetilde{\Omega} = \Omega_{\mathcal{F},Q_0}$, and note that $A \subseteq \partial \widetilde{\Omega}$. Then H(a+b) holds in this case, by Remark 3.4.

Therefore it is enough to consider the case when $\sigma(G) \geq (\eta/2)\sigma(Q_0)$. Suppose first that $\mathcal{F} = \{Q_0\}$. Then necessarily, $Q_0 \in \mathcal{F}_{good}$, and in this case we may apply the induction hypothesis to a child of Q_0 to see that H(a + b) holds. Thus, we may assume that the collection $\mathcal{F} \neq \{Q_0\}$. We now apply Lemma 3.3 to \mathcal{F}_{good} with $\kappa \geq \Upsilon$, where κ is to be chosen momentarily, and where Υ is the constant in Lemma 2.6, to obtain a subcollection $\widetilde{\mathcal{F}} \subset \mathcal{F}_{good}$ with the following properties:

$$\sigma(\cup_{\widetilde{\mathcal{F}}} Q_i) \gtrsim \sigma(\cup_{\mathcal{F}_{\text{good}}} Q_j) \ge (\eta/2)\sigma(Q_0),$$

$$\operatorname{dist}(Q_i, Q_k) \ge \kappa \max\{\ell(Q_i), \ell(Q_k)\}, \quad \forall Q_i, Q_k \in \widetilde{\mathcal{F}}, \quad i \neq k.$$

Now for each $Q_j \in \widetilde{\mathcal{F}}$, we define two families as follows: let Q_j^* be the parent of Q_j , and let Q'_j be the cube C_2 generations down to which we can apply the induction hypothesis. Now set $\mathcal{F}' := \{Q'_j\}_{Q_j \in \widetilde{\mathcal{F}}}$, and $\mathcal{F}^* := \{Q^*_j\}_{Q_j \in \widetilde{\mathcal{F}}}$. Notice first that all of the cubes Q^*_j are in $\mathbb{D}_{\mathcal{F},Q_0}$ and that \mathcal{F}' has the same properties as $\widetilde{\mathcal{F}}$, namely,

- (3.14) $\sigma(\cup_{\mathcal{F}'}Q'_i) \gtrsim \sigma(\cup_{\mathcal{F}_{\text{good}}}Q_j) \gtrsim (\eta/2)\sigma(Q_0)$ and
- (3.15) $\operatorname{dist}(Q'_i, Q'_k) \ge \kappa \max\{\ell(Q'_i), \ell(Q'_k)\}, \quad \forall Q'_i, Q'_k \in \mathcal{F}', \quad i \neq k.$

For each $Q'_j \in \mathcal{F}'$ we apply the induction hypothesis to obtain an open set as in H(a) and call this set $\widetilde{\Omega}_i$.

Next, we construct two "large" NTA domains with ADR boundary with the help of Lemma 3.24 in [20]. For each $Q'_j \in \mathcal{F}'$, let \mathcal{I}_j be the collection of fattened Whitney cubes I^* such that $\operatorname{int}(I^*) \subset \Omega_{\mathcal{F},Q_0}$, and I^* meets $V(Q'_j)$. Let \mathcal{B}_j be the collection of $Q \in \mathbb{D}_{Q_0}$ such that there exists an $I^* \subset U_Q$ with $I^* \in \mathcal{I}_j$. Now we define \mathcal{F}_{∞} as the cubes in $\mathcal{F} \cup (\cup_j \mathcal{B}_j)$ which are maximal with respect to containment.

By construction (see Remark 3.4), $\mathbf{S}' := \mathbb{D}_{\mathcal{F}_{\infty},Q_0}$ is a semi-coherent subregime of some stopping time regime \mathbf{S} as in Lemma 2.2. Thus, setting $\widetilde{\Omega}_0 = \Omega_{F_{\infty},Q_0}$, by Lemma 3.2 we obtain that Ω_0 is the union of two disjoint NTA domains with ADR boundaries, whose diameters are comparable to $\ell(Q_0)$.

We will need to know that we did not remove too many cubes, namely, we do not want to remove any $Q_i^* \in \mathcal{F}^*$.

Claim 3.5. If κ is chosen large enough then for every $j \ge 1$ we have for each $Q_j^* \in \mathcal{F}^*$ that $Q_j^* \in \mathbb{D}_{\mathcal{F}_{\infty},Q_0}$.

Proof of claim: Note that by our choice of C_2 , U_Q does not meet $V(Q'_j)$, for any $Q \supseteq Q_j^*$. Moreover, by construction, we have that $Q_j^* \notin \mathcal{F}$. Suppose now for the purposes of contradiction, that there is a $j \ge 1$ such that $Q_j^* \subseteq Q$, with $Q \in \mathcal{F}_{\infty}$. Then there exists a $k \neq j$ such that $Q \in \mathcal{B}_k$ and U_Q meets $V(Q'_k)$ for some $Q'_k \in \mathcal{F}'$. Then we have immediately that $\ell(Q) \lesssim \ell(Q'_k) < \ell(Q_k)$ and $\operatorname{dist}(Q'_k, Q) \lesssim \ell(Q'_k) < \ell(Q_k)$ so that

$$\operatorname{dist}(Q'_k, Q'_i) \lesssim \operatorname{dist}(Q'_k, Q) + \ell(Q) \lesssim \ell(Q_k)$$

a contradiction for κ large enough. We remind the reader that we are only considering the cubes extracted with a covering lemma so that they separated (see Lemma 3.3 and (3.15)).

Remark 3.6. We note for future reference that $V(Q'_j) \cap (\tilde{\Omega}_0 \cup \partial \tilde{\Omega}_0) = \emptyset$, for every $j \geq 1$. Indeed, this follows by our observation, in the preceding paragraph, that by choice of C_2 large enough, U_Q does not meet $V(Q'_j)$, for any $Q \supseteq Q_j^*$. We further note that $\overline{V(Q'_j)} \cap \overline{V(Q'_k)} = \emptyset$, for all $j \neq k$, with $j, k \geq 1$, by (3.15), Lemma 2.6, and our choice of $\kappa \geq \Upsilon$.

Recall that for $j \geq 1$, $\widetilde{\Omega}_j$ is the open set associated to $Q'_j \in \mathcal{F}'$ via the induction hypothesis, and $\widetilde{\Omega}_0 = \Omega_{F_\infty,Q_0}$. We now set $\widetilde{\Omega} := \bigcup_{j=0}^{\infty} \widetilde{\Omega}_j$. We shall show that $\widetilde{\Omega}$ has all the desired properties. First we show that $\widetilde{\Omega}$ satisfies a 2-sided corkscrew condition. Note that by Remark 3.6, the distinct components of $\widetilde{\Omega}_j$ remain distinct components in $\widetilde{\Omega}$. We let $x \in \partial \widetilde{\Omega}$ and $0 < r \leq \ell(Q_0)$, let M be a large number to be chosen, and set $\delta(x) = \operatorname{dist}(x, E)$. We note that $\partial \widetilde{\Omega} \subseteq (\bigcup_{j=0}^{\infty} \partial \widetilde{\Omega}_j) \cup \Sigma$, where $\Sigma \subset \partial \Omega$, and for each $x \in \Sigma$, there is a subsequence of \mathcal{F}' , call it $\{Q'_{j_k}\}$, with $\operatorname{diam}(Q'_{j_k}) \to 0$, and a sequence of points $y_{j_k} \in Q'_{j_k}$ converging to x. We break into two cases. Case 1: $x \in \partial \widetilde{\Omega}_0$.

Recall that $\widetilde{\Omega}_0$ splits into two NTA domains $\widetilde{\Omega}_0^{\pm}$. Moreover, following the construction in [20], the interior and exterior corkscrew points for the domain $\widetilde{\Omega}_{0}^{+}$ are found as follows. Without loss of generality we may assume that $x \in \partial \widetilde{\Omega}_{0}^{+}$. For M sufficiently large (depending only on allowable constants) the argument distinguishes between two cases, when $r < M\delta(x)$ and when $r \geq M\delta(x)$. In the case that $r \geq M\delta(x)$ we find one corkscrew point in the domain $\widetilde{\Omega}_0^-$ and one corkscrew point in $\tilde{\Omega}_0^+$, and these serve as corkscrew points in separate components for $\hat{\Omega}$ as well. In the case that $r < M\delta(x)$, we have that $\delta(x) > 0$, and x lies on the face of a fattened Whitney cube I^* whose interior lies in $\widetilde{\Omega}_0^+$. Moreover, $x \in J$, for some Whitney cube $J \notin (\bigcup_{Q \in \mathbb{D}_{\mathcal{F}_{\infty},Q_0}} \mathcal{W}_Q)$; we then have one corkscrew point in I^* , and a second in $J \setminus \widetilde{\Omega}_0^+$. Clearly, the first of these is also a corkscrew point for $\widetilde{\Omega}$, in the component $\widetilde{\Omega}_0^+$. To see that the second is a corkscrew point relative to $\partial \widetilde{\Omega}$, it remains to show that J misses $\partial \widetilde{\Omega} \setminus \partial \widetilde{\Omega}_0^+$. If not, then J must intersect $\partial \widetilde{\Omega}_j$ for some $j \geq 1$, and therefore $T_{Q'_i}$ meets J, so that $J \subset V(Q'_i)$. On the other hand, Jalso meets I^* , so that $I^* \in \mathcal{I}_j$, hence there is a cube $Q \in \mathbb{D}_{\mathcal{F}_{\infty},Q_0}$ that belongs to \mathcal{B}_i , a contradiction.

Case 2: $x \in \partial \widetilde{\Omega}_j$ for some $j \ge 1$, or $x \in \Sigma$.

Here $\widetilde{\Omega}_j$ is associated to Q'_j . Let us note that by definition, if $x \in \Sigma$, then there is a "small cube near x", more precisely, there is a cube $Q'_j \in \mathcal{F}'$ with $r \geq M\ell(Q'_j)$, and a point $y_j \in Q'_j$ with $|x - y_j| < r/4$.

Case 2a: $x \in \partial \tilde{\Omega}_j$ for some $j \geq 1$, with $r \geq M\ell(Q'_j)$, or $x \in \Sigma$. In the latter case, we let Q'_j be "the small cube near x", whose existence was noted in the preceeding paragraph. In either of these scenarios, since $Q^*_j \in \mathbb{D}_{\mathcal{F}_{\infty},Q_0}$, we have that $\operatorname{int}(U^{\pm}_{Q^*_j}) \subset \tilde{\Omega}^{\pm}_0$, so for M large enough depending on C_2 , there exists points $y^{\pm} \in \partial \tilde{\Omega}^{\pm}_0$, with $|x - y^{\pm}| < r/2$. We may then apply Case 1 to each of the balls $B(y^{\pm}, r/2) \subset B(x, r)$.

Case 2b: $x \in \partial \widetilde{\Omega}_j$ for some $j \geq 1$, with $r < M\ell(Q'_j)$. From the induction hypothesis we have two corkscrew balls B_1 and B_2 that satisfy the strong 2-sided corkscrew condition, at scale r/M, relative to the open set $\widetilde{\Omega}_j$ (see Definition 1.10). Without loss of generality we may assume that $B_1 \subset \widetilde{\Omega}_j$, hence also $B_1 \subset \widetilde{\Omega}$. We must show that these corkscrew balls satisfy the strong 2-sided corkscrew condition for open sets, with $\widetilde{\Omega}$ the open set in question, thus, it remains to show that $B_1, B_2 \subset \mathbb{R}^{n+1} \setminus \partial \widetilde{\Omega}$. To this end, we simply observe that B_1 and B_2 do not meet $\partial \widetilde{\Omega}_j$ by hypothesis, nor they do not meet $\partial \widetilde{\Omega}_0$ or $\partial \widetilde{\Omega}_i$ for some $i \neq j$, by Remark 3.6, since by the induction hypothesis B_1, B_2 are in $V(Q'_j)$, and since $\partial \widetilde{\Omega}_i \subset \overline{V(Q'_i)}$, by construction.

Next, we observe that $\partial \Omega_0$ is ADR, (see Appendix A in [20]), and that each $\partial \Omega_j$ is ADR, with uniform control of the ADR constants, by the induction hypothesis. Thus, we are left with showing that $\tilde{\Omega}$ has ADR boundary and that condition (3.11) holds. We begin by verifying the upper ADR condition for $\partial \tilde{\Omega}$. Let $x \in \partial \tilde{\Omega}$ and $0 < r \leq \operatorname{diam}(Q_0)$. Recall that $\partial \widetilde{\Omega} \subseteq (\bigcup_{j=0}^{\infty} \partial \widetilde{\Omega}_j) \cup \Sigma$, where in particular $\Sigma \subset \partial \Omega$. Since $\partial \Omega$ is ADR by hypothesis, it suffices to check the upper ADR bound for $A := \bigcup_{j=0}^{\infty} \partial \widetilde{\Omega}_j$. To this end, we observe that, trivially,

(3.16)
$$H^n(B(x,r) \cap A) \le \sum_{j=0}^{\infty} H^n(B(x,r) \cap \partial \widetilde{\Omega}_j).$$

If B(x,r) meets $\partial \widetilde{\Omega}_0$, then there is an $x_0 \in B(x,r) \cap \partial \widetilde{\Omega}_0$, and $B(x_0,2r) \supset B(x,r)$. By the ADR property for $\widetilde{\Omega}_0$

(3.17)
$$H^{n}(B(x,r) \cap \partial \widetilde{\Omega}_{0}) \leq H^{n}(B(x_{0},2r) \cap \partial \widetilde{\Omega}_{0}) \lesssim r^{n}$$

Next, we consider the contributions of $\partial \widetilde{\Omega}_j$, $j \geq 1$, which we write as

$$\sum_{j=1}^{\infty} H^n(B(x,r) \cap \partial \widetilde{\Omega}_j) = \sum_{j: \, \ell(Q'_j) > r} + \sum_{j: \, \ell(Q'_j) \le r} =: \mathbf{I} + \mathbf{II} \,.$$

We recall that by construction $\partial \widetilde{\Omega}_j \subset \overline{V(Q'_j)}$, so by Remark 3.6, the boundaries $\partial \widetilde{\Omega}_j$ are pairwise disjoint. Thus, only a bounded number of terms can appear in the sum I, so the desired bound $I \leq r^n$ follows by the ADR property of each $\partial \widetilde{\Omega}_j$. Moreover, for each j, the diameter of $\widetilde{\Omega}_j$ is comparable to $\ell(Q'_j)$, and therefore the cubes Q'_j appearing in II are all contained in B(x, Cr), for some sufficiently large constant C depending only on allowable parameters. Consequently, by the ADR property of $\partial \widetilde{\Omega}_j$ and of E,

$$\mathrm{II} \lesssim \sum_{j:\,Q_j' \subset B(x,Cr)} H^n(\partial \widetilde{\Omega}_j) \, \approx \sum_{j:\,Q_j' \subset B(x,Cr)} \sigma(Q_j') \, \leq \, \sigma(E \cap B(x,Cr)) \, \approx r^n \, .$$

Thus, we obtain the upper ADR bound.

With the upper ADR bound in hand we now know that $\partial \Omega$ is a set of locally finite perimeter, so by the isoperimetric inequality (see Theorem 2 on page 190 of [15]), and the strong two sided corkscrew condition for open sets, the lower ADR bound follows. The last and easiest thing to show is that condition (3.11) holds for Ω , but this follows readily from (3.14), and the fact that by the induction hypothesis, (3.11) hold for each Q'_i .

4. Proof of Theorem 1.3

The proof of Theorem 1.3 will be immediate from the following lemma.

Lemma 4.1. Let Ω be a bounded open set in \mathbb{R}^{n+1} , with n-dimensional ADR boundary, such that $\Omega := \bigcup_j \Omega_j$ is the union of its connected components Ω_j . Suppose further that each component is a chord-arc domain with uniform bounds on the chord-arc constants. For $x_0 \in \partial\Omega$ and $r_0 < \frac{1}{2} \operatorname{diam} \Omega$, and for $Y \in \Omega \setminus B(x_0, 2r_0)$, let ω^Y denote the harmonic measure associated to Ω , and set $\Delta_0 := B(x_0, r_0) \cap \partial\Omega$. Then $\omega^Y \in \operatorname{weak} A_\infty(\Delta_0)$, with uniform control on the weak- A_∞ constants. Proof. Fix $B_0 := B(x_0, r_0)$, with $x_0 \in \partial\Omega$, and $r_0 < \operatorname{diam}(\partial\Omega)$. Set $\Delta = \Delta(x, r)$ and $2\Delta = \Delta(x, 2r)$, with $x \in \partial\Omega$, and suppose that $B(x, 2r) \subset B_0$ (thus, $r_0 \ge 2r$). Recalling the definition of weak- A_∞ , we need to show there exist uniform positive constants C and θ such that for each Borel set $A \subseteq \Delta$

(4.1)
$$\omega^{Y}(A) \le C \left(\frac{\sigma(A)}{\sigma(\Delta)}\right)^{\theta} \omega(2\Delta) \,,$$

whenever $Y \in \Omega \setminus B(x_0, 2r_0)$. Let us fix such a point Y. We note that $Y \in \Omega_j$ for some j and therefore ω^Y is just the harmonic measure associated to the domain Ω_j . Thus, if $A \cap \partial \Omega_j = \emptyset$, then 4.1 holds trivially. We may therefore assume that this is not the case. Let $z \in A \cap \partial \Omega_j$, set $A' = A \cap \partial \Omega_j$, and set $\sigma_j = H^n|_{\partial \Omega_j}$. Notice that dist $(z, Y) \ge r_0 \ge 2r$, since in particular, $z \in B_0$, while $Y \in \Omega \setminus 2B_0$. Thus, the diameter of Ω_j must be greater than r. Moreover, by the result obtained independently in [11] and in [30], $\omega^Y \in A_\infty(2\Delta_*)$, where $\Delta_* = B(z, r) \cap \Omega_j$, $2\Delta_* = B(z, 2r) \cap \Omega_j$, and in particular, we have the doubling estimate

(4.2)
$$\omega^Y(2\Delta_\star) \lesssim \omega^Y(\Delta_\star) \,.$$

Note further that $\Delta_{\star} \subset 2\Delta$, that $A' \subset 2\Delta_{\star}$, and also that by the uniform ADR property, $\sigma(\Delta) \approx \sigma_j(\Delta_{\star}) = \sigma(\Delta_{\star})$.

Since $\omega^Y \in A_{\infty}(2\Delta_{\star})$, we therefore have that

$$\omega^{Y}(A) = \omega^{Y}(A') \leq C \left(\frac{\sigma_{j}(A')}{\sigma_{j}(\Delta_{\star})}\right)^{\theta} \omega^{Y}(2\Delta_{\star})$$
$$\leq C \left(\frac{\sigma(A')}{\sigma(\Delta_{\star})}\right)^{\theta} \omega^{Y}(\Delta_{\star}) \lesssim \left(\frac{\sigma(A)}{\sigma(\Delta)}\right)^{\theta} \omega^{Y}(2\Delta),$$

where in the next to last inequality, we have used (4.2).

To prove Theorem 1.3, we apply the preceding lemma to each domain Ω constructed in Theorem 1.1. We need only verify that the point $Y \in \Omega \setminus 2B_0$ may be replaced by any corkscrew point $X_{\Delta_0} \in \widetilde{\Omega}$, relative to the surface ball Δ_0 . To this end, we fix such a corkscrew point X_{Δ_0} , and observe that $X_{\Delta_0} \in \Omega \setminus 2B_1$, for any ball B_1 meeting B_0 , centered on $\partial \widetilde{\Omega}$, with radius $r_1 = \eta r_0$, if $\eta > 0$ is chosen small enough depending only on the corkscrew constant for $\widetilde{\Omega}$. Then for each such B_1 , by Lemma 4.1, $\omega_0 := \omega^{X_{\Delta_0}}$ belongs to weak- $A_{\infty}(\Delta_1)$, where $\Delta_1 = B_1 \cap \partial \widetilde{\Omega}$. Now suppose that $B(x, 2r) \subset B_0$, with $x \in \partial \widetilde{\Omega}$, and observe that B := B(x, r) may be covered by a collection of balls $\mathcal{F} = \{B^i := B(x^i, r^i)\}$, of bounded cardinality depending only on dimension and η , such that $r^i \approx r$, $2B^i \subset 2B$, and each B^i is contained in some ball $B_1 = B_1^i$ as above. Set $\Delta = B \cap \partial \widetilde{\Omega}$, $\Delta^i := B^i \cap \partial \widetilde{\Omega}$, and for $A \subset \Delta$, let $A_i := A \cap \Delta^i$. We then apply (4.1) to each A_i, Δ^i , and use that for each $i, H^n(\Delta) \approx H^n(\Delta^i)$ (depending on η and the ADR constants). Since $\#\mathcal{F}$ is bounded, we may then sum in i to obtain Theorem 1.3.

We conclude by observing that Lemma 4.1, and hence Theorem 1.1, also apply to the Riesz measure ("*p*-harmonic measure") associated to the *p*-Laplace equation

$$\Delta_p u := \operatorname{div} \left(|\nabla u|^{p-2} \nabla u \right) = 0.$$

Definition 4.2. (*p*-harmonic measure). Let $\Omega \subset \mathbb{R}^{n+1}$ be open. For $x \in \partial\Omega$, $0 < r < \frac{1}{8} \operatorname{diam}(\Omega)$, suppose that $u \ge 0$, with $u \equiv 0$ on $\partial\Omega$, $\Delta_p u = 0$ in $4B \cap \Omega$. We define the *p*-harmonic measure μ associated to *u*, as the unique finite positive Borel measure such that

(4.3)
$$-\iint_{\Omega} |\nabla u|^{p-2} \nabla u \cdot \nabla \Phi \, dx = \int_{\partial \Omega} \Phi \, d\mu \quad \forall \Phi \in C_0^{\infty}(4B).$$

Proposition 4.3. Let Ω be a bounded open set in \mathbb{R}^{n+1} , with n-dimensional ADR boundary, such that $\Omega := \bigcup_j \Omega_j$ is the union of its connected components Ω_j , and satisfies a strong 2-sided corkscrew condition. Suppose further that each component Ω_j is a chord-arc domain with uniform bounds on the chord-arc constants. Then for $x \in \partial \Omega$ and $r < \frac{1}{8} \operatorname{diam}(\Omega)$, let u and μ be as above, set $\hat{\mu} := \mu|_{\partial \Omega_j}$, and let \hat{c} be the corkscrew constant for the set Ω . If $\operatorname{diam}(\Omega_j) > 2\hat{c}r$, then $\hat{\mu} \in$ weak- $A_{\infty}(\Delta)$. In particular, given $x \in \partial \Omega$ and $r \in (0, \frac{1}{8} \operatorname{diam}(\Omega))$, if we let Ω_{Δ} be some component which contains an interior corkscrew point relative to the ball B(x, r), then $\mu|_{\partial \Omega_{\Delta}} \in \operatorname{weak-} A_{\infty}(\Delta)$.

Proof. The proof is the same as that of Lemma 4.1. Lewis and Nyström [26] showed that for an NTA domain D with ADR boundary, if μ is p-harmonic measure associated to a non-trivial positive p-harmonic function u in $B(x, 4r) \cap D$, with $x \in \partial D, r < \operatorname{diam}(D)/4$, then $\mu \in A_{\infty}(\partial D \cap B(x, r))$, with respect to $\sigma_D := H^n|_{\partial D}$; we use this fact in place of the result of [11], [30].

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