

Comparing two statistical ensembles of quadrangulations: a continued fraction approach

Éric Fusy and Emmanuel Guitter

Abstract. We use a continued fraction approach to compare two statistical ensembles of quadrangulations with a boundary, both controlled by two parameters. In the first ensemble, the quadrangulations are bicolored and the parameters control their numbers of vertices of both colors. In the second ensemble, the parameters control instead the number of vertices which are local maxima for the distance to a given vertex, and the number of those which are not. Both ensembles may be described either by their (bivariate) *generating functions at fixed boundary length* or, after some standard slice decomposition, by their (bivariate) *slice generating functions*. We first show that the fixed boundary length generating functions are in fact equal for the two ensembles. We then show that the slice generating functions, although different for the two ensembles, simply correspond to two different ways of encoding the same quantity as a continued fraction. This property is used to obtain explicit expressions for the slice generating functions in a constructive way.

Mathematics Subject Classification (2010). 05.

Keywords. Planar maps, quadrangulations, two-point functions, continued fractions, hard dimers, heaps theory.

Contents

1	Introduction	126
2	An equality between two bivariate generating functions for quadrangulations with a boundary	128
3	Slice decomposition and continued fractions	132
4	Getting the slice generating functions by solving recursion relations	138
5	Getting the slice generating functions by extracting continued fraction coefficients: generalities	141
6	Recovering (13) from the continued fraction formalism	145
7	Conserved quantities	155
8	Conclusion	161

Appendices	162
A A proof of formulas (25) and (26)	162
B A proof of the formulas (35)–(38)	172
References	175

1. Introduction

The study of planar maps has given rise in the recent years to a lot of remarkable enumeration results. A particularly fruitful approach consists in taking advantage of bijections between maps and tree-like objects called mobiles. This technique, initiated by Schaeffer [20, 10] (reinterpreting a bijection by Cori and Vauquelin [11]) was extended in many different directions [5, 8, 10, 17, 1, 9, 2] to deal with various refined map enumeration problems. It has also been a key tool to show that random planar maps have an explicit scaling limit called the Brownian map [16, 18, 19, 3]. Besides mobiles, another, slightly different, view on the problem consists in decomposing the maps into so-called *slices*, which are particular pieces of maps with nice combinatorial properties [7]. In particular, the generating functions for these slices were shown to obey discrete integrable systems of equations and in most cases, a solution of these equations could be obtained explicitly. Moreover, the slice decomposition of a map is intimately linked to its geodesic paths and the knowledge of slice generating functions directly gives explicit answers to a number of questions regarding the statistics of distances between random points within maps [4, 5, 6, 15].

A particularly important discovery was made in [7] where it was shown that slice generating functions happen to be simple coefficients in a suitable *continued fraction expansion* of standard map generating functions, making *de facto* a connection between the distance statistics within maps and some more global properties. On a computational point of view, this discovery provided a constructive way to obtain explicit solutions for the integrable systems at hand, by taking advantage of known results on continued fractions.

Quite recently, the slice decomposition technique was used in [15] to describe the distance statistics of general families of bicolored maps, and, in particular, of *bicolored quadrangulations*, with some simultaneous control on the numbers of vertices of both colors. Explicit expressions for the corresponding bivariate slice generating functions were obtained in a constructive way, leading in particular to explicit formulas for the distance dependent two-point function within bicolored quadrangulations. Remarkably, the expressions found for slice generating

functions are very similar to those obtained (via a mobile formalism) in another problem of quadrangulations considered in [1]. There, the discrimination between vertices no longer relies on their color but rather on their status with respect to the graph distance from a fixed origin vertex. Vertices namely come in two types: those, called *local maxima* which are further from the origin than all their neighbors, and the others. Bivariate slice generating functions can be defined so as to keep some independent control on the numbers of both types of vertices after the slice decomposition. Explicit expressions for these new bivariate slice generating functions were then *guessed* in [1] and, as just mentioned, their structure is very similar to that of their bicolored counterparts.

The aim of this paper is twofold: first, we establish a strong connection between the problem of quadrangulations with a control on the vertex color, as discussed in [15], and that of quadrangulations with a control on local maxima, as discussed in [1]. Then, we use a continued fraction formalism to re-derive, now in a constructive way, the explicit expressions found in [1].

The paper is organized as follows: in Section 2, we introduce the two ensembles of quadrangulations that we want to compare and define their generating functions at fixed boundary length. We then derive our first fundamental result which states that the fixed boundary length generating functions are in fact equal for the two ensembles. Section 3 presents the slice decomposition of the quadrangulations at hand and shows that the corresponding slice generating functions may be obtained as coefficients of the same quantity, once expanded as a continued fraction in two different ways. In Section 4, we recall the integrable systems which determine the slice generating functions of both ensembles as well as the explicit solutions of these systems obtained in [15] and [1]. Section 5 deals with results on continued fractions, and shows in particular how to extract their coefficients from those obtained via a standard series expansion. In one of the ensembles that we consider, the knowledge of this series expansion is not sufficient to get all the slice generating functions, a process which requires *the knowledge of some additional quantity*. An explicit expression for this latter quantity is conjectured in Section 6, based on simplifications observed in the case of finite continued fractions, and we then show how it allows to recover the explicit formulas for the slice generating functions found in [1]. Section 7 deals with another aspect of our problem, the existence of invariants, the so-called *conserved quantities*, as expected for discrete integrable systems. We show how to derive these invariants combinatorially for both ensembles and again emphasize the deep similarity existing between the conserved quantities for the two ensembles. We gather our concluding remarks in Section 8. Some side results or technical derivations are presented in Appendices A and B.

2. An equality between two bivariate generating functions for quadrangulations with a boundary

The aim of this section is to compare the generating functions of planar *quadrangulations with a boundary* weighted in two different ways, each of these weighting being *bivariate*, i.e. involving two independent parameters. As we shall see below, these two weightings, although fundamentally different, are intimately linked and some of the associated generating functions turn out to be equal.

2.1. Two bivariate generating functions for quadrangulations with a boundary. Recall that a planar quadrangulation with a boundary denotes a connected graph embedded on the sphere which is *rooted*, i.e. has a marked oriented edge (the root-edge) and is such that all its inner faces, i.e. all the faces except that lying on the right of the root-edge, have degree 4. As for the external face, which is the face lying on the right of the root-edge, its degree is arbitrary (but necessarily even). As customary, the origin of the root-edge will be called the root-vertex. Let us now consider two particular different ways to assign weights to these maps.

✓ *First weighting: bicoloring the map.* Since all their faces have even degree, planar quadrangulations with a boundary may be naturally bicolored in black and white in a unique way, by assigning the black color to the root-vertex and demanding that no two adjacent vertices have the same color. We may then enumerate these quadrangulations by assigning a weight t_\bullet to each black vertex and a weight t_\circ to each white vertex. For convenience, the root-vertex receives a weight 1 instead of t_\bullet . We shall then denote by $F_n \equiv F_n(t_\bullet, t_\circ)$ the corresponding generating function for these maps with a *boundary length* $2n$, i.e. with an external face of degree $2n$.

✓ *Second weighting: distinguishing local maxima of the distance.* Our second weighting consists in giving a special role to the local maxima of the distance from the root-vertex. More precisely, we may label each vertex v of the quadrangulation by its graph distance $d(v)$ from the root-vertex and look for the *local maxima* of this labeling, i.e. those vertices v having only neighbors with label $d(v) - 1$ (note that in all generality, neighbors of a vertex v may only be at distance $d(v) - 1$ or $d(v) + 1$ from the root-vertex). We decide to give a weight t_\circ to local maxima and a weight t_\bullet to the other vertices (see Figure 1 – left). As before, the root-vertex receives a weight 1 instead of t_\bullet (note that the root-vertex can never be a local maximum). We shall call $J_n \equiv J_n(t_\bullet, t_\circ)$ the generating function for these maps with a boundary of length $2n$.

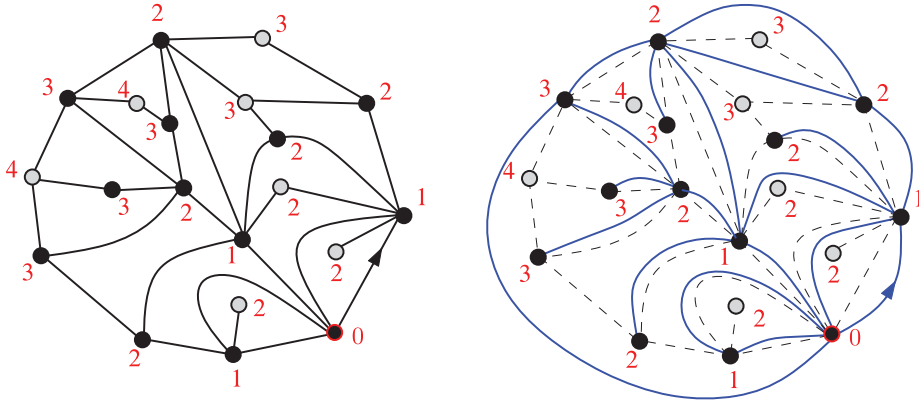


Figure 1. Left: an example of rooted quadrangulation with a boundary of length 10, where each vertex is labelled by its distance to the root-vertex. The local maxima for this distance are indicated in gray. Right: the associated rooted general map with a bridgeless boundary of length 5, obtained by applying the Ambjørn–Budd rule within each inner face (i.e. connecting the two corners within the face followed clockwise by a larger label) and, in the external face, connecting *cyclically* those corners followed by a larger label counterclockwise around the map.

The generating functions $F_n(t_\bullet, t_\circ)$ and $J_n(t_\bullet, t_\circ)$ may be understood as formal power series in t_\bullet and t_\circ , giving rise to convergent series for small enough t_\bullet, t_\circ . The first weighting is quite natural and was described in detail in [15]. We shall recall some of the corresponding results below. As for the second weighting, it may seem more artificial but, as explained in [1], it arises naturally in two contexts: first, letting $t_\circ \rightarrow 0$ (i.e. keeping the linear term in t_\circ) is a way to suppress local maxima of the distance, selecting quadrangulations arranged into layers between the root-vertex and a unique local maximum. These so-called Lorentzian or causal structures display a very different statistics from that of arbitrary quadrangulations [1]. As recalled below, the second weighting also arises naturally when enumerating *general planar maps* with a control on both their numbers of vertices and faces [1].

2.2. Equality of generating functions. Let us now prove a first fundamental equality, namely that

$$J_n(t_\bullet, t_\circ) = F_n(t_\bullet, t_\circ). \tag{1}$$

To this end, let us recall the so-called Ambjørn–Budd bijection of [1] between quadrangulations and general maps, slightly adapted to the case of quadrangulations with a boundary according to the rules of [6]. Starting with our quadran-

gulation with a boundary and labeling each vertex v by its distance $d(v)$ from the root-vertex, we associate to each inner face an edge as follows (see Figure 1 – right): looking at the corners¹ clockwise around the face, exactly two corners are followed by a corner with larger label. We connect these two corners by an edge lying inside the original face. As for the external face, looking again at the corner labels clockwise around the face, i.e. counterclockwise around the rest of the map, exactly n corners, including the root-corner (lying immediately to the right of the root-edge) are followed by a corner with larger label. We connect the n corners of this ensemble *cyclically* clockwise around the map, each edge connecting two successive corners in the ensemble (see Figure 1). Finally, we mark and orient away from the root-vertex the edge connecting the root-corner to its successor. As explained in [1] (and its extension [6]), the obtained edges form a rooted planar map together with those vertices of the original quadrangulation *which were not local maxima* for the distance from the root-vertex. Each inner face of this map surrounds exactly one of the original local maxima, which get disconnected in the construction. More precisely, from [1, 6], the above transformation provides a bijection between planar quadrangulations with a boundary of length $2n$ and rooted planar general (i.e. with faces of arbitrary degrees) maps with a *bridgeless* boundary² of length n (i.e. with external face – lying on the right of the root-edge – of degree n and without bridge). The vertices of the quadrangulation which are not local maxima for the distance from the root are in one-to-one correspondence with the vertices of the general map while the vertices of the quadrangulation which are local maxima are in one-to-one correspondence with the inner faces of the general map.

We may thus interpret $J_n(t_\bullet, t_\circ)$ as the generating function for rooted planar general maps with a bridgeless boundary of length n , weighted by t_\bullet per non-root-vertex and t_\circ per inner face.

As for the label $d(v)$ of a vertex v retained in this new map, it precisely corresponds to the *oriented* graph distance from the root-vertex to v on the new map, using paths oriented from the root-vertex to v which respect the following

¹ Recall that a corner is an angular sector between two successive half-edges around a given vertex. The label of a corner is that of the incident vertex.

² Strictly speaking, the extension [6] of [1] shows that corners followed by a smaller label should be connected *cyclically* within all faces, including the inner faces, so that the resulting object is a hypermap, made of alternating black and white faces with, in our case, all black faces of degree 2 but one, of degree n , which we choose as external face. The Ambjørn–Budd construction that we use here is recovered by squeezing all inner black faces, of degree 2, into simple edges while the external black face of degree n becomes the external face of the map. As for any face of a hypermap, its boundary is then necessarily without bridge.

edge orientation:³ all edges are oriented both ways except for the boundary-edges (i.e. the edges incident to the external face) which are oriented counterclockwise around the map.

Forgetting about distances and labels, we may now use a standard construction to rebuild a quadrangulation with a boundary out of our general map. Coloring the vertices of the map in black, we simply add a white vertex within each inner face and connect it to all the corners within the face.⁴ By doing so, we get a bicolored quadrangulation with a boundary twice larger as that of the general map we started from, which we root by picking the edge leaving the root-vertex within the corner immediately to the left of the root-edge of the general map, and orienting it from its black to its white extremity (see Figure 2– right). Again this construction provides a bijection between rooted planar general maps with a bridgeless boundary of length n and planar quadrangulations with a boundary of length $2n$, equipped with their (unique) bicolouration as defined in the previous section. The vertices of the general map are in one-to-one correspondence with the black vertices of the quadrangulation while the inner faces of the map are in one-to-one correspondence with the white vertices of the quadrangulation.

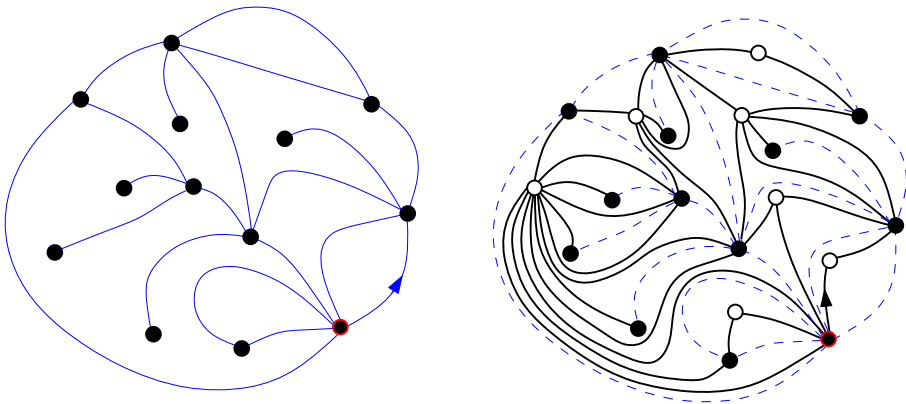


Figure 2. Left: the rooted map on the right of Figure 1, with a bridgeless boundary of length 5. Right: The associated rooted bicolored quadrangulation with a boundary of length 10 obtained by inserting a white vertex at the center of each inner face of the map and connecting it to all the incident vertices around the face.

³ In the underlying hypermap structure, the labels correspond to the distance using oriented paths going clockwise around the black faces. Squeezing the inner black faces of degree 2 results in simple edges oriented both ways, while the boundary-edges remain oriented oneway only.

⁴ Note that it is crucial that the boundary of the map be bridgeless for the obtained object to be connected.

We may thus interpret $F_n(t_\bullet, t_\circ)$ as the generating function for rooted planar general maps with a bridgeless boundary of length n , weighted by t_\bullet per non-root-vertex and t_\circ per face. eq. (1) follows.

3. Slice decomposition and continued fractions

3.1. Slice decomposition for maps enumerated by F_n . As explained in [15], the quadrangulations (with a boundary) enumerated by F_n may be decomposed into *slices* by some appropriate cutting procedure.

Labeling each boundary-vertex v by $d(v)$, the sequence of corner labels, read counterclockwise around the map starting from the root-corner, forms a directed path of length $2n$, made of elementary steps with height difference ± 1 , starting and ending at height 0 and remaining (weakly) above height 0 (see Figure 3). Drawing, for each boundary-vertex v , its *leftmost* geodesic (shortest) path to the root-vertex and cutting along these geodesics results into a decomposition of the map into pieces, called slices. More precisely, to each descending step $i \rightarrow i - 1$ of the path corresponds an i -slice, defined as follows (see [15] for details): it is a rooted map whose boundary is made of three parts (see Figure 4):

- (i) its base consisting of a single root-edge,
- (ii) a left boundary of length ℓ with $1 \leq \ell \leq i$ connecting the origin of the root-edge to another vertex, the *apex* and which is a geodesic path within the slice, and
- (iii) a right boundary of length $\ell - 1$ connecting the endpoint of the root-edge to the apex, and which is the unique geodesic path within the slice between these vertices.

The left and right boundaries do not meet before reaching the apex (which by convention is considered as part of the right boundary only). As a degenerate case when $\ell = 1$, the left boundary may stick to the base, in which case the slice is reduced to a single root-edge.

At this level, it is interesting to note that the distance $d(v)$ from the root-vertex to any vertex v in the quadrangulation is directly related to its distance $d_s(v)$, within the i -slice it lies in, from the apex of this slice via

$$d(v) = d_s(v) + i - \ell \tag{2}$$

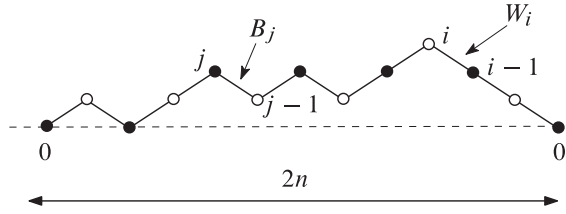


Figure 3. An example of directed path of length $2n = 12$, made of elementary steps with height difference ± 1 , starting and ending at height 0 and remaining (weakly) above height 0. The path is naturally colored in black and white. To compute F_n , we must sum over such bicolored paths with a weight B_i (resp. W_i) assigned to each descending step $i \rightarrow i - 1$ starting from a black height (resp. white height).

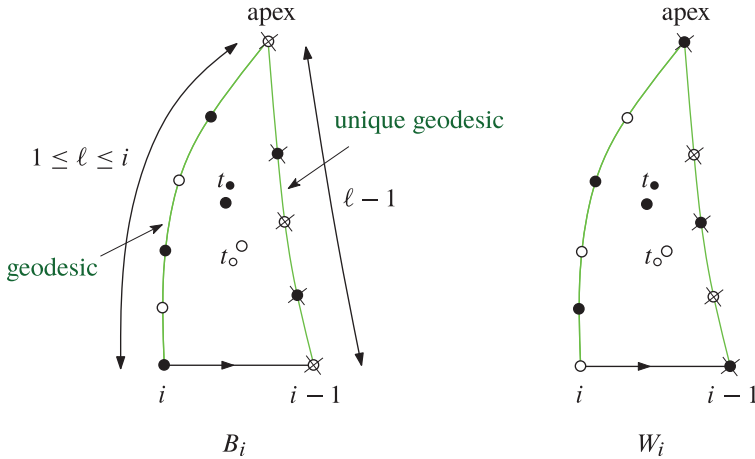


Figure 4. A schematic picture of an i -slice contributing to B_i (left) and to W_i (right).

if ℓ is the length of the left boundary of the slice. Indeed, it is clear by construction of the slices that either the root-vertex is the apex of the i -slice at hand or it does not belong to the slice at all and any path from v to this root-vertex must first reach one of the boundaries of the slice (possibly at the apex). In the first case, we have $d(v) = d_s(v)$ and $i = \ell$ so that (2) holds. In the second case, since the slice boundaries are part of geodesic paths to the root-vertex, $d(v)$ is equal to $d_s(v)$ plus the distance from the apex of the slice to the root-vertex. In other words, $d(v) - d_s(v)$ has a constant value within the i -slice, which is obtained by taking for v the origin of the root-edge of the slice, namely $d(v) - d_s(v) = i - \ell$, and (2) follows. Note that, in an i -slice, i only acts as an upper bound on the length ℓ of the left boundary. The vertices v of an i -slice may then be labelled by non-negative

integers in two natural ways: either by their distance $d_s(v)$ to the apex or by this distance plus $i - \ell$.

Let us call $B_i \equiv B_i(t_\bullet, t_\circ)$ (resp. $W_i \equiv W_i(t_\bullet, t_\circ)$) the generating function for i -slices whose root-vertex is black (resp. white), with a weight t_\bullet per black vertex and t_\circ per white vertex *except for the vertices of the right boundary* (including the endpoint of the root-edge and the apex) which receive a weight⁵ 1. The slice decomposition implies that ([15])

$$F_n = Z_{0,0}^+(2n; \{B_i\}_{i \geq 1}, \{W_i\}_{i \geq 1}), \tag{3}$$

where $Z_{0,0}^+(2n; \{B_i\}_{i \geq 1}, \{W_i\}_{i \geq 1})$ denotes the generating function of paths of length $2n$, made of elementary steps with height difference ± 1 , colored alternatively in black and white, starting and ending at black height 0 and remaining (weakly) above height 0, with each descending step from a black height i to a white height $i - 1$ weighted by B_i and each descending step from a white height i to a black height $i - 1$ weighted by W_i (see Figure 3).

The set of identities (3) for all $n > 0$ can then be translated [14] into the continued fraction expression

$$F(z) \equiv \sum_{n \geq 0} F_n z^n = \frac{1}{1 - z \frac{1}{1 - z \frac{W_1}{1 - z \frac{B_2}{1 - z \frac{W_3}{1 - z \frac{B_4}{1 - \dots}}}}} \tag{4}$$

with the convention that $F_0 = 0$ and where $F(z) = F(z; t_\bullet, t_\circ)$ implicitly depends on t_\bullet and t_\circ .

3.2. Slice decomposition for maps enumerated by J_n . Let us now play the same game with maps enumerated by J_n , which are the same maps as those enumerated by F_n , but now with the second weighting. We may again apply the same slice decomposition, resulting in the same i -slices as before. More precisely, labeling each boundary-vertex v by $d(v)$ gives rise to a path of length $2n$ (from height 0 to height 0, remaining above height 0) and each descending step $i \rightarrow i - 1$ gives rise to an i -slice. To assign the second weighting to the quadrangulation, we must label each vertex v of the i -slice by its distance $d(v)$ (in the quadrangulation) from the root-vertex of the quadrangulation. As explained above, if the i -slice has

⁵The reason why these vertices receive a weight 1 is to avoid double weighting upon regluing the slices into a quadrangulation. Indeed, all these vertices are already part of a left boundary, except for the root-vertex. In the end, only the root-vertex gets a weight 1, as wanted.

a left boundary length ℓ ($1 \leq \ell \leq i$), this amounts to label v by $d_s(v) + i - \ell$ where $d_s(v)$ is its distance (within the slice) from the apex of the slice. To recover the correct weights, we must first give weight 1 to all the vertices of the right boundary (including the endpoint of their root-edge and the apex) in order to avoid double weightings after regluing the slices. As for the vertices lying on the left boundary of the slice and *different from the root-vertex of the slice*, they cannot, as part of a geodesic of the original quadrangulation, be local maxima as they have a neighbor with larger label along the geodesic path. They receive a weight t_\bullet accordingly. Considering now vertices v lying strictly within the slice, they have all their original neighbors lying in the slice and, from (2), are local maxima for the distance $d(v)$ if and only if they are local maxima for the distance $d_s(v)$ within the slice. For all these vertices, we may thus use the distance $d_s(v)$ within the slice to detect the local maxima, and give them the weight t_\circ , while vertices that are not local maxima for $d_s(v)$ get the weight t_\bullet . The last vertex to consider is the origin w of the root-edge of the slice: for this vertex to be a local maximum of the original distance $d(v)$, it must both be a local maximum of the distance $d_s(v)$ within the slice and have no neighbor with larger label after regluing. Now two situations may occur: either the boundary-vertex preceding w along the boundary (oriented counterclockwise around the quadrangulation) has label $d(w) - 1$ and then the slice at hand sticks to the boundary so that w has all its neighbors within the slice. Then w is a local maximum for $d(v)$ if and only if it is a local maximum for $d_s(v)$. Or the boundary-vertex preceding w has label $d(w) + 1$ in which case w is not a local maximum for $d(v)$, irrespectively of whether or not it is one for $d_s(v)$.

To summarize, we are led to consider two different generating functions for i -slices. In the first generating function $Q_i \equiv Q_i(t_\bullet, t_\circ)$, all the vertices of the i -slice receive a weight t_\circ or t_\bullet according to whether or not they are a local maximum *for the distance $d_s(v)$ from the apex within the slice* (in particular the vertices of the left boundary different from the root-vertex of the slice get a weight t_\bullet as wanted), except for the vertices of the right boundary (including the endpoint of their root-edge and the apex) which receive a weight 1. In the second generating function $P_i \equiv P_i(t_\bullet, t_\circ)$, we assign exactly the same weights as in Q_i , except for the root-vertex of the slice, which gets the weight t_\bullet irrespectively of whether or not it is a local maximum for $d_s(v)$ (see Figure 5).

Returning to the slice decomposition, it now implies, for any positive n , that

$$J_n = \widehat{Z}_{0,0}^+(2n; \{P_i\}_{i \geq 1}, \{Q_i\}_{i \geq 1}) ,$$

where $\widehat{Z}_{0,0}^+(2n; \{P_i\}_{i \geq 1}, \{Q_i\}_{i \geq 1})$ denotes the generating function of paths of length $2n$, made of elementary steps of height difference ± 1 , starting and ending

at height 0 and remaining above height 0, with each descending step from height i to height $i - 1$ weighted by P_i if it follows a descending step $i + 1 \rightarrow i$ and by Q_i if it follows an ascending step $i - 1 \rightarrow i$ (see Figure 6).

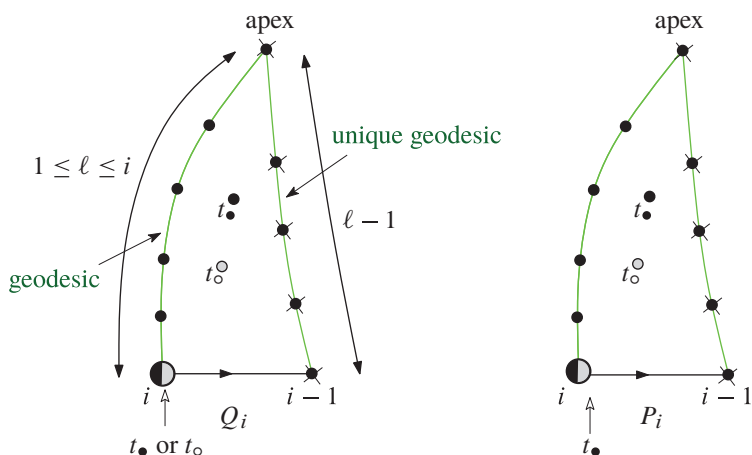


Figure 5. A schematic picture of an i -slice contributing to Q_i (left) and to P_i (right). Vertices that are (resp. are not) local maxima of the distance from the apex are indicated in gray (resp. in black). The root vertex may be a local maximum or not, and receives the weight $t_◦$ or $t_•$ accordingly in Q_i , while it always gets the weight $t_•$ in P_i .

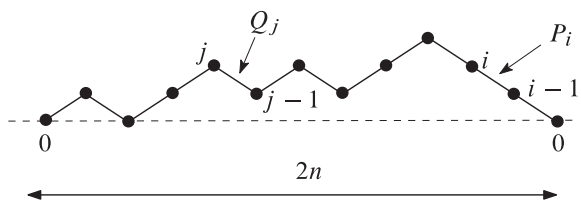


Figure 6. An example of directed path of length $2n = 12$, made of elementary steps with height difference ± 1 , starting and ending at height 0 and remaining (weakly) above height 0. To compute J_n , we must sum over such paths with a weight Q_i (resp. P_i) assigned to each descending step $i \rightarrow i - 1$ following an ascent (resp. following a descent).

Setting $J_0 = 1$ and using the shorthand notation $J(z) = J(z; t_\bullet, t_\circ)$, this is summarized into the new continued fraction expansion

$$\begin{aligned}
 J(z) &\equiv \sum_{n \geq 0} J_n z^n \\
 &= \frac{1}{1 - z(Q_1 - P_1) - z \frac{P_1}{1 - z(Q_2 - P_2) - z \frac{P_2}{1 - z(Q_3 - P_3) - z \frac{P_3}{1 - \dots}}}}
 \end{aligned} \tag{5}$$

which we may write as

$$J(z) = \frac{1}{1 - zY_1 - z \frac{Y_2}{1 - zY_3 - z \frac{Y_4}{1 - zY_5 - z \frac{Y_6}{1 - \dots}}}} \tag{6}$$

upon defining

$$Y_{2i-1} \equiv Q_i - P_i, \quad Y_{2i} = P_i \tag{7}$$

for $i \geq 1$.

To understand (5), or equivalently (6), we note that, expanding the right hand side of this latter equation, the term of order z^n builds the generating function $\check{Z}_{0,0}^+(2n; \{Y_i\}_{i \geq 1})$ of paths of length $2n$ starting and ending at height 0 and remaining above height 0, made of elementary (i.e. of horizontal length 1) steps of height difference ± 1 together with “elongated steps” of horizontal length 2 and height difference 0. Each elementary descending step from height i to height $i - 1$ ($i \geq 1$) receives a weight Y_{2i} while each elongated step at height $i - 1$ ($i \geq 1$) receives the weight Y_{2i-1} (see Figure 7). Deforming each elongated step at height $i - 1$ into a sequence of elementary steps $i - 1 \rightarrow i \rightarrow i - 1$, we recover paths made only of elementary steps of height difference ± 1 , and (after regrouping all paths with the same deformation) receiving a weight $Y_{2i-1} + Y_{2i} = Q_i$ for each sequence $i - 1 \rightarrow i \rightarrow i - 1$ or equivalently for each descending step $i \rightarrow i - 1$ following an ascent, and a weight $Y_{2i} = P_i$ for those elementary steps $i \rightarrow i - 1$ which are not part of a sequence $i - 1 \rightarrow i \rightarrow i - 1$, i.e. follow a descent. In other words, $\check{Z}_{0,0}^+(2n; \{Y_i\}_{i \geq 1}) = \hat{Z}_{0,0}^+(2n; \{P_i\}_{i \geq 1}, \{Q_i\}_{i \geq 1})$, which explains the identity (5).

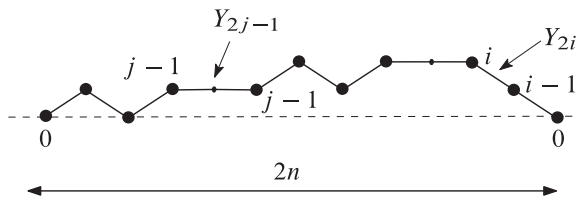


Figure 7. An example of directed path of length $2n = 12$, made of elementary steps with height difference ± 1 and elongated steps (of horizontal length 2) with height difference 0, starting and ending at height 0 and remaining (weakly) above height 0. To compute J_n , we must sum over such paths with a weight Y_{2i-1} (resp. Y_{2i}) assigned to each elongated step $i - 1 \rightarrow i - 1$ (resp. each descending step $i \rightarrow i - 1$).

To conclude this section, let us rewrite our fundamental equality (1) in the more compact form

$$J(z; t_\bullet, t_o) = F(z; t_\bullet, t_o).$$

4. Getting the slice generating functions by solving recursion relations

The slice generating functions B_i , W_i , P_i and Q_i satisfy systems of non linear recursive equations which may be derived by performing a slice decomposition of the slices themselves. Indeed, when the i -slice, of left boundary length ℓ , is not reduced to a single edge, we may look at the sequence of vertices encountered clockwise around the face lying on the left of the root-edge of the slice and draw the leftmost geodesic paths from these vertices to the apex. Using the labeling $d_s(v) + i - \ell$, the sequence of encountered labels, starting from its root-vertex, is either $i \rightarrow i - 1 \rightarrow i - 2 \rightarrow i - 1$ (if $i \geq 2$), $i \rightarrow i - 1 \rightarrow i \rightarrow i - 1$ or $i \rightarrow i + 1 \rightarrow i \rightarrow i - 1$ and, upon cutting along the leftmost geodesic paths, a new slice arises for each descending step of this sequence (see Figure 8).

For the first weighting, this decomposition, applied to i -slices enumerated by B_i and W_i , leads to the system

$$B_i = t_\bullet + B_i(W_{i-1} + B_i + W_{i+1}), \tag{8a}$$

$$W_i = t_o + W_i(B_{i-1} + W_i + B_{i+1}), \tag{8b}$$

for $i \geq 1$, with $B_0 = W_0 = 0$. For the second weighting, this decomposition, applied to i -slices enumerated by P_i and Q_i , leads similarly to the system

$$P_i = t_\bullet + P_i(P_{i-1} + Q_i + Q_{i+1}), \tag{9a}$$

$$Q_i = t_o + Q_i(P_{i-1} + Q_i) + P_i Q_{i+1}, \tag{9b}$$

for $i \geq 1$, with $P_0 = Q_0 = 0$.

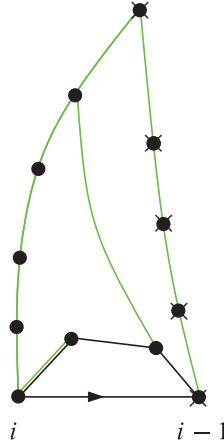


Figure 8. A schematic picture of the slice decomposition of a (non-trivial) i -slice into two slices, leading to the recursion relations (8) and (9).

The solution of (8) was derived in [15]. Parametrizing t_\bullet and t_o by x and γ via

$$t_\bullet = \frac{x(\gamma - x)^3(1 - \gamma x^3)}{(x + x^3 + \gamma - 6x^2\gamma + x^4\gamma + x\gamma^2 + x^3\gamma^2)^2}, \tag{10a}$$

$$t_o = \frac{x(\gamma - x^3)(1 - \gamma x^3)}{(x + x^3 + \gamma - 6x^2\gamma + x^4\gamma + x\gamma^2 + x^3\gamma^2)^2}, \tag{10b}$$

with⁶ $|x| \leq 1$, it was shown that

$$B_{2i} = B \frac{(1 - x^{2i})(1 - \gamma x^{2i+3})}{(1 - \gamma x^{2i+1})(1 - x^{2i+2})}, \tag{11a}$$

$$W_{2i+1} = W \frac{(1 - \gamma x^{2i+1})(1 - x^{2i+4})}{(1 - x^{2i+2})(1 - \gamma x^{2i+3})}, \tag{11b}$$

$$B_{2i+1} = B \frac{(1 - x^{2i+1}/\gamma)(1 - x^{2i+4})}{(1 - x^{2i+2})(1 - x^{2i+3}/\gamma)}, \tag{11c}$$

$$W_{2i} = W \frac{(1 - x^{2i})(1 - x^{2i+3}/\gamma)}{(1 - x^{2i+1}/\gamma)(1 - x^{2i+2})}, \tag{11d}$$

⁶The parametrization is invariant under $(x, \gamma) \rightarrow (1/x, 1/\gamma)$ so we may always choose $|x| \leq 1$.

for $i \geq 0$, where

$$B = \frac{x(\gamma - x)^2}{x + x^3 + \gamma - 6x^2\gamma + x^4\gamma + x\gamma^2 + x^3\gamma^2},$$

$$W = \frac{x(1 - \gamma x)^2}{x + x^3 + \gamma - 6x^2\gamma + x^4\gamma + x\gamma^2 + x^3\gamma^2}.$$

As for the solution of (9), it was guessed in [1]. Parametrizing now t_\bullet and t_\circ by y and α via

$$t_\bullet = \frac{y(1 - \alpha y)^3(1 - \alpha y^3)}{(1 + y + \alpha y - 6\alpha y^2 + \alpha y^3 + \alpha^2 y^3 + \alpha^2 y^4)^2}, \quad (12a)$$

$$t_\circ = \frac{\alpha y(1 - y)^3(1 - \alpha^2 y^3)}{(1 + y + \alpha y - 6\alpha y^2 + \alpha y^3 + \alpha^2 y^3 + \alpha^2 y^4)^2}, \quad (12b)$$

with $|y| \leq 1$, it was found that

$$P_i = P \frac{(1 - y^i)(1 - \alpha y^{i+3})}{(1 - y^{i+1})(1 - \alpha y^{i+2})}, \quad Q_i = Q \frac{(1 - y^i)(1 - \alpha^2 y^{i+3})}{(1 - \alpha y^{i+1})(1 - \alpha y^{i+2})}, \quad (13)$$

for $i \geq 0$, where

$$P = \frac{y(1 - \alpha y)^2}{1 + y + \alpha y - 6\alpha y^2 + \alpha y^3 + \alpha^2 y^3 + \alpha^2 y^4}, \quad (14a)$$

$$Q = \frac{\alpha y(1 - y)^2}{1 + y + \alpha y - 6\alpha y^2 + \alpha y^3 + \alpha^2 y^3 + \alpha^2 y^4}. \quad (14b)$$

Note that the two parametrizations (10) and (12) are actually equivalent providing we relate y and α to x and γ via

$$\alpha = \frac{1}{\gamma^2}, \quad y = \gamma x.$$

With this correspondence, we immediately deduce that

$$P = B, \quad Q = W.$$

This should not come as a surprise since, from (11) and (13), B , W , P and Q are the $i \rightarrow \infty$ limits of B_i , W_i , P_i and Q_i , enumerating slices *with no bound on their boundary lengths*. From (8) and (9), both pairs (B, W) , and (P, Q) are determined by the *same closed system*, namely:

$$B = t_\bullet + B(B + 2W), \quad W = t_\circ + W(W + 2B), \quad (15a)$$

$$P = t_\bullet + P(P + 2Q), \quad Q = t_\circ + Q(Q + 2P). \quad (15b)$$

Let us end this section by rewriting the results for P_i and Q_i in terms of Y_i , as defined in (7). First, eq. (9) may be rewritten as

$$Y_{2i} = t_{\bullet} + Y_{2i}(Y_{2i-2} + Y_{2i-1} + Y_{2i} + Y_{2i+1} + Y_{2i+2}), \quad (16a)$$

$$Y_{2i-1} = (t_{\circ} - t_{\bullet}) + Y_{2i-1}(Y_{2i-2} + Y_{2i-1} + Y_{2i}), \quad (16b)$$

for $i \geq 1$, with $Y_0 = 0$. From (13), we immediately deduce the solution

$$Y_{2i} = P \frac{(1-y^i)(1-\alpha y^{i+3})}{(1-y^{i+1})(1-\alpha y^{i+2})}, \quad Y_{2i+1} = Y \frac{(1-y^{i+1})(1-\alpha y^{i+3})}{(1-y^{i+2})(1-\alpha y^{i+2})}, \quad (17)$$

for $i \geq 0$, with

$$Y = Q - P.$$

The aim of this paper is to go beyond the guessing approach of [1] and to provide a constructive way to obtain this latter formula (17), and consequently (13), upon using general results for continued fractions of the type (6). This is indeed the constructive approach used in [15] to obtain the expressions (11) from general results for continued fractions of the type (4).

5. Getting the slice generating functions by extracting continued fraction coefficients: generalities

5.1. The Stieltjes type. Eq. (4) is a continued fraction of the so-called ‘‘Stieltjes type.’’ Its coefficients B_{2i} and W_{2i-1} for $i \geq 1$ are known to be related to the coefficients F_n via the relations

$$B_{2i} = \frac{h_i^{(0)}}{h_{i-1}^{(0)}} \Big/ \frac{h_{i-1}^{(1)}}{h_{i-2}^{(1)}}, \quad W_{2i-1} = \frac{h_{i-1}^{(1)}}{h_{i-2}^{(1)}} \Big/ \frac{h_{i-1}^{(0)}}{h_{i-2}^{(0)}}, \quad (18)$$

for $i \geq 1$, in terms of the Hankel determinants

$$h_i^{(0)} = \det(F_{n+m})_{0 \leq n, m \leq i}, \quad h_i^{(1)} = \det(F_{n+m+1})_{0 \leq n, m \leq i},$$

for $i \geq 0$, with the convention $h_{-1}^{(0)} = h_{-1}^{(1)} = 1$. These expressions were used in [15] to obtain the expressions (11) for B_{2i} and W_{2i+1} ($i \geq 0$). As for the expressions of B_{2i+1} and W_{2i} , it is clear from (8) that B_i and W_i play symmetric roles upon exchanging t_{\bullet} and t_{\circ} . The expressions (11) for B_{2i+1} and W_{2i} are simply deduced upon this transformation, which amounts to a change $\gamma \leftrightarrow 1/\gamma$, $B \leftrightarrow W$ in the formulas (see [15]). At this stage, it is important to note that the knowledge of the generating functions F_n is not sufficient to determine all

the B_i 's and W_i 's as the associated continued fraction involves only one parity of the index i (B_i 's with even index i and W_i 's with odd index i) and that we have to rely on a symmetry principle to get the other parity. Otherwise stated, the derivation of all the B_i 's and W_i 's requires in principle the knowledge of a second family of generating functions. In the present case, these generating functions are nothing but those of rooted quadrangulations with a boundary of length n , bicolored in such a way that their root-vertex is white instead of black. Of course, by symmetry, those are nothing but the $F_n(t_\bullet, t_\circ)$, $n \geq 1$ and a simple symmetry argument is sufficient to conclude.

5.2. The type of eq. (6). When dealing with a continued fraction of the type of eq. (6), a first remark should be emphasized: the knowledge of J_n is not sufficient to determine the coefficients Y_i . Indeed, expanding in z gives rise to the first equations:

$$J_1 = (Y_1 + Y_2), \quad J_2 = (Y_1 + Y_2)^2 + Y_2(Y_3 + Y_4), \quad \dots \quad (19)$$

and it is easily seen that, at each step, two new Y_i 's appear on the right hand side, so that the system is clearly underdetermined.

As shown in [12, 13], a full determination of the coefficients Y_i requires, in addition to the set of J_n for $n \geq 1$, the knowledge of Y_1 and of a second family of quantities $\tilde{J}_n \equiv \tilde{J}_n(t_\bullet, t_\circ)$, $n \geq 0$, satisfying

$$\tilde{J}(z) \equiv \sum_{n \geq 0} \tilde{J}_n z^n = \frac{1}{1 - z\tilde{Y}_1 - z \frac{\tilde{Y}_2}{1 - z\tilde{Y}_3 - z \frac{\tilde{Y}_4}{1 - z\tilde{Y}_5 - z \frac{\tilde{Y}_6}{1 - \dots}}}} \quad (20)$$

where we have defined (assuming $Y_i \neq 0$ for all $i \geq 1$)

$$\tilde{Y}_{2i-1} \equiv \frac{1}{Y_{2i-1}}, \quad \tilde{Y}_{2i} \equiv \frac{Y_{2i}}{Y_{2i-1}Y_{2i+1}}, \quad (21)$$

for $i \geq 1$. Expanding in z now gives rise to the first equations:

$$\tilde{J}_1 = \frac{(Y_2 + Y_3)}{Y_1 Y_3}, \quad \tilde{J}_2 = \frac{(Y_2 + Y_3)^2}{(Y_1 Y_3)^2} + \frac{Y_2(Y_4 + Y_5)}{Y_1 Y_3^2 Y_5}, \quad \dots \quad (22)$$

Knowing Y_1 , the first equation of (19) yields Y_2 , then the first equation of (22) yields Y_3 , the second equation of (19) yields Y_4 , and so on. The Y_i 's are now fully determined and a compact formula may be written as follows: define

$$j_n \equiv \begin{cases} 1 & \text{if } n = 0, \\ Y_1 J_{n-1} & \text{if } n \geq 1, \\ \tilde{J}_{-n} & \text{if } n \leq -1, \end{cases} \quad (23)$$

and the Hankel-type determinants

$$H_i^{(0)} = \det(j_{n+m-i-1})_{1 \leq n, m \leq i}, \quad (24a)$$

$$H_i^{(1)} = \det(j_{n+m-i})_{1 \leq n, m \leq i}. \quad (24b)$$

Then we have, for $i \geq 1$, the following formulas, reminiscent of (18),

$$Y_{2i} = \frac{H_{i-1}^{(0)}}{H_i^{(0)}} \bigg/ \frac{H_i^{(1)}}{H_{i+1}^{(1)}} \quad (25a)$$

$$Y_{2i-1} = \frac{H_i^{(1)}}{H_{i-1}^{(1)}} \bigg/ \frac{H_i^{(0)}}{H_{i-1}^{(0)}}, \quad (25b)$$

with the convention $H_0^{(0)} = H_0^{(1)} = 0$. A proof of these formulas can be found in [12, 13]. We present a slightly simpler proof in the Appendix A below. To summarize, when dealing with a continued fraction of the type of eq. (6), we may extract the coefficients Y_i if, in addition to $J(z)$, we also know Y_1 and $\tilde{J}(z)$. As we shall see in Section 7 below, getting a simple expression for Y_1 combinatorially may be achieved upon using a so-called conserved quantity. As for $\tilde{J}(z)$, we have not been able to obtain it via combinatorial arguments (as opposed to the previous section, we cannot rely here on any symmetry principle to get $\tilde{J}_n(t_\bullet, t_\circ)$ from $J_n(t_\bullet, t_\circ)$). Without the knowledge of $\tilde{J}(z)$, eq. (6) yields a much weaker system than the recursion equations (16). In fact, *any arbitrary choice of \tilde{J}_n will lead, through (25), to a set of Y_i 's satisfying eq. (6)*, while the actual Y_i 's, solution of eqs. (16), correspond to a unique value of the \tilde{J}_n 's, to be determined.

As we shall now explain, we may however *conjecture* a simple expression for $\tilde{J}(z)$, based on an explicit solution of the problem in the case of *finite* continued fractions. With this conjectured form of $\tilde{J}(z)$, we may then verify that the obtained Y_i 's precisely match their actual expressions (17) guessed in [1].

5.3. The case of finite continued fractions. In this section, let us briefly digress from our combinatorial problem and discuss the case of a *finite* continued fraction. More precisely, let

$$J(z) \equiv \frac{1}{1 - zY_1 - z \frac{Y_2}{1 - zY_3 - z \frac{Y_4}{\ddots \frac{Y_{2\alpha-2}}{1 - zY_{2\alpha-3} - z \frac{Y_{2\alpha-1}}{1 - zY_{2\alpha-1}}}}}}$$

where $Y_1, Y_2, \dots, Y_{2\alpha-1}$ denote *independent indeterminates*. We also define

$$\tilde{J}(z) \equiv \frac{1}{1 - z\tilde{Y}_1 - z \frac{\tilde{Y}_2}{1 - z\tilde{Y}_3 - z \frac{\tilde{Y}_4}{\ddots \frac{\tilde{Y}_{2\alpha-2}}{1 - z\tilde{Y}_{2\alpha-3} - z \frac{\tilde{Y}_{2\alpha-1}}{1 - z\tilde{Y}_{2\alpha-1}}}}}}$$

with

$$\tilde{Y}_{2i-1} = \frac{1}{Y_{2i-1}} \quad \text{for } 1 \leq i \leq \alpha$$

and

$$\tilde{Y}_{2i} = \frac{Y_{2i}}{Y_{2i-1}Y_{2i+1}} \quad \text{for } 1 \leq i < \alpha.$$

The rational function $J(z)$ is easily seen to be the ratio of a polynomial of degree $\alpha - 1$ in z by a polynomial of degree α in z , hence characterized by

$$(\alpha - 1 + 1) + (\alpha + 1) - 1 - 1 = 2\alpha - 1$$

coefficients (the last two -1 's correspond to removing a global factor in both the numerator and the denominator, and ensuring that $J(0) = 1$), depending on the $2\alpha - 1$ indeterminates $Y_1, Y_2, \dots, Y_{2\alpha-1}$. In this case, the knowledge of the function $J(z)$ alone therefore entirely determines all the coefficients of the continued fraction. This property may be reconciled with the apparently contrary statement of the previous section by noting that, in the present case of a finite continued fraction, both Y_1 and $\tilde{J}(z)$ (defined via (20) and (21)) can be deduced from $J(z)$. More precisely, we have the following relations, derived in Appendix A

below:

$$\tilde{J}(z) = -\frac{Y_1}{z} J\left(\frac{1}{z}\right), \quad Y_1 = -\frac{1}{\lim_{z \rightarrow \infty} z J(z)}. \tag{26}$$

Note that $\tilde{J}(z)$ is also a rational function of z and that the expression for Y_1 simply rephrases the desired property that $\tilde{J}(z) = 1 + O(z)$. Knowing $J(z)$, Y_1 and $\tilde{J}(z)$, we can then deduce the coefficients j_n for all integer n via their definition (23) and get $Y_2, Y_3, \dots, Y_{2\alpha-1}$ from eqs. (24) and (25) (which are also valid in the case of a finite continued fraction – see Appendix A). The relations (26) are proved in the Appendix A below.

6. Recovering (13) from the continued fraction formalism

6.1. A conjectured expression for Y_1 and $\tilde{J}(z)$. Returning now to our enumeration problem of i -slices with the second weighting, let us *conjecture* that, although our continued fraction is now infinite, *the relations (26) still hold* for the particular choice of Y_i we are interested in, namely the solution of (16). More precisely, $J(z)$, originally defined as a power series in z , is convergent for small enough real z (namely $0 \leq z < 1/(\sqrt{Q} + \sqrt{P})^2$, see explicit expressions below – here we assume that t_\bullet and t_\circ are small enough positive reals so that P and Q are positive reals) but may be analytically continued to large enough real z ($z > 1/(\sqrt{Q} - \sqrt{P})^2$). This allows us to define $J(1/z)$ for small real z (namely $0 \leq z < (\sqrt{Q} - \sqrt{P})^2$) and our conjecture is that, in this range, $\tilde{J}(z)$ is obtained via the relation $\tilde{J}(z) = -(Y_1/z)J(1/z)$ with a value of Y_1 adjusted so that $\tilde{J}(0) = 1$. Assuming this property, let us now see if we can then recover the desired expression (13), or equivalently (17).

Let us start by recalling the expression of $F(z)$, hence $J(z)$. From [15], we know that

$$F_n = \frac{B}{t_\bullet} (1 - B - W) Z_{0,0}^+(2n; B, W) - \frac{B}{t_\circ} Z_{0,0}^+(2n + 2; B, W) \tag{27}$$

where $Z_{0,0}^+(2n; B, W)$ denotes the generating function of paths of length $2n$, made of elementary steps with height difference ± 1 , colored alternatively in black and white, starting and ending at black height 0 and remaining (weakly) above height 0, with each descending step from a black height to a white height weighted by B and each descending step from a white height to a black height weighted by W . A derivation of this expression via slices is recalled in Section 7 below.

Equivalently, since $J_n = F_n$, $P = B$ and $Q = W$, we have

$$J_n = A_0 Z_{0,0}^+(2n; P, Q) + A_1 Z_{0,0}^+(2n + 2; P, Q), \tag{28a}$$

$$A_0 = \frac{P}{t_\bullet}(1 - P - Q), \quad A_1 = -\frac{P}{t_\bullet}. \tag{28b}$$

Let us introduce

$$Z(z; P, Q) \equiv \sum_{n \geq 0} Z_{0,0}^+(2n; P, Q) z^n,$$

which, by definition, is a solution of

$$Z(z; P, Q) = \frac{1}{1 - z \frac{Q}{1 - z P Z(z; P, Q)}}. \tag{29}$$

Note that this (quadratic) equation in Z is equivalent to the equation

$$Z(z; P, Q) = \frac{1}{1 - z(Q - P) - z P Z(z; P, Q)}$$

so that $Z_{0,0}^+(2n; P, Q)$ is also the generating function of paths of length $2n$, made of elementary steps with height difference ± 1 , with each descending step weighted by Q if it follows an ascending step and by P otherwise. Alternatively, $Z_{0,0}^+(2n; P, Q)$ enumerates paths of length $2n$, made of elementary steps of horizontal length 1 and height difference ± 1 , and elongated steps of horizontal length 2 and height difference 0, each elongated step receiving the weight $Y = (Q - P)$ and each elementary descending step the weight P (see Figure 9). As a continued fraction, we thus have

$$Z(z; P, Q) = \frac{1}{1 - z Y - z \frac{P}{1 - z Y - z \frac{P}{1 - z Y - z \frac{P}{1 - \dots}}}}, \quad Y = Q - P. \tag{30}$$

In terms of Z , we may write

$$J(z) = A_0 Z(z; P, Q) + A_1 \frac{Z(z; P, Q) - 1}{z}$$

and, in components

$$J_n = A_0 [z^n]Z(z; P, Q) + A_1 [z^{n+1}]Z(z; P, Q) \tag{31}$$

for $n \geq 0$.

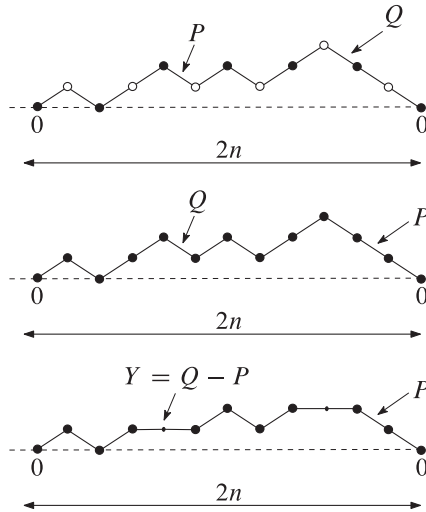


Figure 9. A sketch of the three interpretations of the generating function $Z_{0,0}^+(2n; P, Q)$.

At this stage, it is important to note that eq. (29) yields two branches for Z for real z , namely

$$Z_-(z; P, Q) = \frac{1 - Yz - \sqrt{(1 - Yz)^2 - 4Pz}}{2Pz},$$

$$Z_+(z; P, Q) = \frac{1 - Yz + \sqrt{(1 - Yz)^2 - 4Pz}}{2Pz}, \quad Y = Q - P,$$

for $|z| \leq 1/(\sqrt{Q} + \sqrt{P})^2$. To recover the coefficients J_n , we must expand $J(z)$, hence $Z(z; P, Q)$ at small z , which requires to choose

$$Z(z; P, Q) = Z_-(z; P, Q).$$

From $Z_-(z; P, Q) = 1 + Qz + O(z^2)$ we get $J(z) = 1 + O(z)$ (since from (15), $A_0 + A_1 Q = 1$), as wanted. Using (26), we find that

$$\tilde{J}(z) = -Y_1(A_0 \bar{Z}(z; P, Q) + A_1 (z \bar{Z}(z; P, Q) - 1)),$$

where

$$\bar{Z}(z; P, Q) \equiv \frac{1}{z} Z\left(\frac{1}{z}; P, Q\right).$$

Again, we have two possible branches for real z :

$$\begin{aligned}\bar{Z}_-(z; P, Q) &= \frac{z - Y - \sqrt{(z - Y)^2 - 4Pz}}{2Pz}, \\ \bar{Z}_+(z; P, Q) &= \frac{z - Y + \sqrt{(z - Y)^2 - 4Pz}}{2Pz}, \quad Y = Q - P,\end{aligned}$$

and, to get the \tilde{J}_n 's, we must, depending on whether $P > Q$ or $Q > P$, choose the first or second branch respectively to get rid of the $1/z$ term when $z \rightarrow 0$. Both situations yield actually the same expression for \tilde{J}_n . Assuming for instance $P > Q$, we get

$$\tilde{J}_n = -Y_1 A_0 [z^n] \bar{Z}_-(z; P, Q) - Y_1 A_1 [z^n] (z \bar{Z}_-(z; P, Q) - 1).$$

The value of Y_1 is obtained by ensuring that $\tilde{J}_0 = 1$. Using

$$\bar{Z}_-(z; P, Q) = -\frac{1}{Y} + O(z),$$

we deduce

$$Y_1 \left(\frac{A_0}{Y} + A_1 \right) = 1, \quad (32a)$$

hence

$$Y_1 = \frac{Y}{A_0 + A_1 Y} = \frac{(Q - P)(1 - P - 2Q)}{1 - 2Q}. \quad (32b)$$

This value matches that obtained directly via conserved quantities in Section 7.

Using the identity

$$\bar{Z}_-(z; P, Q) = \frac{1}{P - Q} Z_- \left(\frac{z}{(P - Q)^2}; P, Q \right),$$

we eventually arrive at

$$\begin{aligned}\tilde{J}_n &= -\frac{Y_1}{(P - Q)^{2n+1}} (A_0 [z^n] Z_-(z; P, Q) \\ &\quad + A_1 (P - Q)^2 [z^n] (z Z_-(z; P, Q) - 1)) \\ &= \frac{Y_1}{(Q - P)^{2n+1}} (A_0 Z_{0,0}^+(2n; P, Q) + A_1 (Q - P)^2 Z_{0,0}^+(2n - 2; P, Q))\end{aligned} \quad (33)$$

for $n > 0$. For $Q > P$, we must use instead $\bar{Z}_+(z; P, Q)$ but again in this case,

$$\bar{Z}_+(z; P, Q) = -\frac{1}{Y} + O(z)$$

and

$$\bar{Z}_+(z; P, Q) = \frac{1}{P - Q} Z_- \left(\frac{z}{(P - Q)^2}; P, Q \right)$$

so that the expressions (32) and (33) remain unchanged. The first line of (33) may be rewritten as

$$\begin{aligned} \tilde{J}_n &= \frac{Y_1}{Y} \left(A_0 [z^n] Z_- \left(\frac{z}{Y^2}; P, Q \right) + A_1 [z^n] \left(z Z_- \left(\frac{z}{Y^2}; P, Q \right) - 1 \right) \right) \\ &= \frac{Y_1}{Y} \left(A_0 [z^n] \tilde{Z}(z; P, Q) + A_1 [z^{n-1}] \tilde{Z}(z; P, Q) \right) \end{aligned} \tag{34}$$

for $n \geq 1$, with

$$\tilde{Z}(z; P, Q) \equiv Z_- \left(\frac{z}{Y^2}; P, Q \right) = \frac{1}{1 - z \tilde{Y} - z \frac{\tilde{P}}{1 - z \tilde{Y} - z \frac{\tilde{P}}{1 - z \tilde{Y} - z \frac{\tilde{P}}{1 - \dots}}}}$$

where

$$\tilde{Y} = \frac{1}{Y} \quad \text{and} \quad \tilde{P} = \frac{P}{Y^2}.$$

Upon defining

$$k_n \equiv \begin{cases} 1 & \text{if } n = 0, \\ Y Z_{n-1} & \text{if } n \geq 1, \\ \tilde{Z}_{-n} & \text{if } n \leq -1, \end{cases}$$

with the notations

$$Z_n \equiv [z^n] Z(z; P, Q) = Z_{0,0}^+(2n; P, Q)$$

and

$$\tilde{Z}_n \equiv [z^n] \tilde{Z}(z; P, Q) = Z_{0,0}^+(2n; P, Q) / Y^{2n}$$

($n \geq 0$), we may summarize (31), (32), and (34) into

$$j_n = A_0 \frac{Y_1}{Y} k_n + A_1 \frac{Y_1}{Y} k_{n+1}$$

for all integer n .

6.2. Computation of $H_i^{(0)}$ and $H_i^{(1)}$. The above expression for j_n for all integer n opens the way to compute $H_i^{(0)}$ and $H_i^{(1)}$ via (24). Indeed, as shown in Appendix B, the coefficients k_n satisfy a set of linear relations of the form

$$\sum_{m=0}^{i-1} k_{n-i+m} (-1)^m \frac{x_{i-1-m}^{(i-1)}}{x_{i-1}^{(i-1)}} = 0, \quad 2 \leq n \leq i, \quad (35)$$

while, for $n = 1$ and $n = i + 1$, we have

$$\sum_{m=0}^{i-1} k_{1-i+m} (-1)^m \frac{x_{i-1-m}^{(i-1)}}{x_{i-1}^{(i-1)}} = \frac{P^{i-1}}{Y^{2(i-1)}} \quad (36a)$$

$$\sum_{m=0}^{i-1} k_{1+m} (-1)^m \frac{x_{i-1-m}^{(i-1)}}{x_{i-1}^{(i-1)}} = (-1)^{i-1} \frac{P^{i-1}}{Y^{i-2}}. \quad (36b)$$

From these relations, replacing the first column $(j_{n-i})_{1 \leq n \leq i}$ of $H_i^{(0)}$ by the linear combination

$$\left(\sum_{m=0}^{i-1} j_{n-i+m} (-1)^m \frac{x_{i-1-m}^{(i-1)}}{x_{i-1}^{(i-1)}} \right)_{1 \leq n \leq i}$$

of this first column with the $i - 1$ last ones allows us to write

$$H_i^{(0)} = \begin{vmatrix} A_0 \frac{Y_1}{Y} \frac{P^{i-1}}{Y^{2(i-1)}} & j_{-i+2} & \cdots & \cdots & j_0 \\ 0 & j_{-i+3} & \cdots & \cdots & j_1 \\ \vdots & \vdots & \ddots & \ddots & \vdots \\ 0 & j_0 & \ddots & \ddots & j_{i-2} \\ (-1)^{i-1} A_1 \frac{Y_1}{Y} \frac{P^{i-1}}{Y^{i-2}} & j_1 & \cdots & \cdots & j_{i-1} \end{vmatrix}$$

$$= A_0 Y_1 \frac{P^{i-1}}{Y^{2i-1}} H_{i-1}^{(1)} + A_1 Y_1 \frac{P^{i-1}}{Y^{i-1}} H_{i-1}^{(0)}.$$

Alternatively, the coefficients k_n satisfy another set of linear relations of the form (see Appendix B)

$$\sum_{m=0}^{i-1} k_{n-m} (-1)^m x_m^{(i-1)} = 0, \quad 2 \leq n \leq i, \quad (37)$$

while, for $n = 1$ and $n = i + 1$, we have

$$\sum_{m=0}^{i-1} k_{1-m} (-1)^m x_m^{(i-1)} = (-1)^{i-1} \frac{P^{i-1}}{Y^{i-2}}, \quad (38a)$$

$$\sum_{m=0}^{i-1} k_{i+1-m} (-1)^m x_m^{(i-1)} = Y P^{i-1} (Y + P). \quad (38b)$$

From these relations, replacing the last column $(j_n)_{1 \leq n \leq i}$ of $H_i^{(1)}$ by the linear combination

$$\left(\sum_{m=0}^{i-1} j_{n-m} (-1)^m x_m^{(i-1)} \right)_{1 \leq n \leq i}$$

of this last column with the $i - 1$ first ones, and using $x_0^{(i-1)} = 1$ (see Appendix B), allows us to write

$$\begin{aligned} H_i^{(1)} &= \begin{vmatrix} j_{-i+2} & \cdots & \cdots & j_0 & (-1)^{i-1} A_0 \frac{Y_1}{Y} \frac{P^{i-1}}{Y^{i-2}} \\ j_{-i+3} & \cdots & \cdots & j_1 & 0 \\ \vdots & \vdots & \ddots & \ddots & \vdots \\ j_0 & \ddots & \ddots & j_{i-2} & 0 \\ j_1 & \cdots & \cdots & j_{i-1} & A_1 \frac{Y_1}{Y} Y P^{i-1} (Y + P) \end{vmatrix} \\ &= A_0 Y_1 \frac{P^{i-1}}{Y^{i-1}} H_{i-1}^{(1)} + A_1 Y_1 P^{i-1} (Y + P) H_{i-1}^{(0)}. \end{aligned}$$

To summarize, $H_i^{(0)}$ and $H_i^{(1)}$ are fully determined by the system

$$\begin{aligned} H_i^{(0)} &= A_0 Y_1 \frac{P^{i-1}}{Y^{2i-1}} H_{i-1}^{(1)} + A_1 Y_1 \frac{P^{i-1}}{Y^{i-1}} H_{i-1}^{(0)}, \\ H_i^{(1)} &= A_0 Y_1 \frac{P^{i-1}}{Y^{i-1}} H_{i-1}^{(1)} + A_1 Y_1 P^{i-1} (Y + P) H_{i-1}^{(0)}, \end{aligned}$$

for $i \geq 1$ with $H_0^{(0)} = H_0^{(1)} = 1$. Upon setting

$$L_i^{(0)} \equiv \left(\frac{Y}{P} \right)^{\frac{i(i-1)}{2}} H_i^{(0)}, \quad L_i^{(1)} \equiv \frac{1}{Y^{i-1}} \left(\frac{Y}{P} \right)^{\frac{i(i-1)}{2}} H_i^{(1)}, \quad (39)$$

these equations read

$$L_i^{(0)} = A_0 \frac{Y_1}{Y^2} L_{i-1}^{(1)} + A_1 Y_1 L_{i-1}^{(0)} \quad (40a)$$

$$L_i^{(1)} = A_0 \frac{Y_1}{Y} L_{i-1}^{(1)} + A_1 Y_1 (Y + P) L_{i-1}^{(0)}. \quad (40b)$$

Using the first line to express $L_i^{(1)}$ in terms of $L_i^{(0)}$ and re-injecting the result in the second line yields an equation for $L_i^{(0)}$ only, namely

$$L_{i+1}^{(0)} = Y_1 \left(\frac{A_0}{Y} + A_1 \right) L_i^{(0)} + A_0 A_1 \frac{Y_1^2}{Y^2} P L_{i-1}^{(0)}$$

for $i \geq 1$ with

$$L_0^{(0)} = 1$$

and

$$L_1^{(0)} = H_1^{(0)} = Y_1 \left(\frac{A_0}{Y} + A_1 \right).$$

Using (32) and setting

$$w \equiv A_0 A_1 \frac{Y_1^2}{Y^2} P,$$

we recover the well-known equation

$$L_{i+1}^{(0)} = L_i^{(0)} + w L_{i-1}^{(0)} \quad (41)$$

for $i \geq 1$ with $L_0^{(0)} = L_1^{(0)} = 1$, which allows us to interpret $L_i^{(0)}$ as the generating function of hard pieces on a linear graph with $i - 1$ vertices (see Figure 10), with a weight w per piece. The solution of this equation is known to be (see for instance [7] eq. (6.11))

$$L_i^{(0)} = \frac{1}{(1+y)^i} \frac{1-y^{i+1}}{1-y} \quad \text{where } w = -\frac{1}{y+y^{-1}+2}. \quad (42)$$

If we instead eliminate $L_i^{(0)}$ from the system (40), we obtain for $L_i^{(1)}$ the very same equation

$$L_{i+1}^{(1)} = L_i^{(1)} + w L_{i-1}^{(1)}$$

for $i \geq 1$, now with the initial conditions

$$L_0^{(1)} = Y$$

and

$$L_1^{(1)} = H_1^{(1)} = j_1 = Y_1$$

(this value can also be read from (40) as it yields $L_1^{(1)} = Y_1(A_0 + A_1(Y + P)) = Y_1(A_0 + A_1Q) = Y_1$). We immediately deduce

$$L_i^{(1)} = Y L_i^{(0)} + (Y_1 - Y)L_{i-1}^{(0)}$$

with the convention $L_{-1}^{(0)} = 0$. Using (42), we obtain

$$\begin{aligned} L_i^{(1)} &= \frac{1}{(1+y)^i} \frac{1}{1-y} \left(Y(1-y^{i+1}) + (Y_1 - Y)(1-y^i)(1+y) \right) \\ &= \frac{1}{(1+y)^i} Y_1(1+dy) \left(\frac{1-\alpha y^{i+2}}{1-y} \right), \end{aligned} \tag{43}$$

where

$$d \equiv \frac{Y_1 - Y}{Y_1} \quad \text{and} \quad \alpha \equiv \frac{1}{y^2} \frac{d+y}{1+dy}.$$

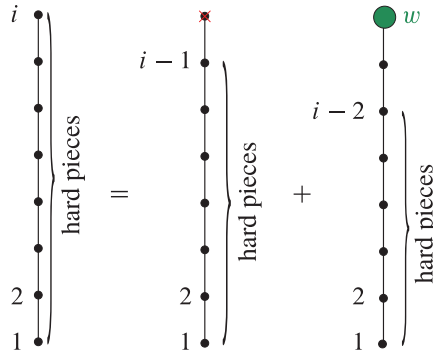


Figure 10. A schematic picture of eq. (41), identifying $L_i^{(0)}$ as the generating function of hard pieces on a linear graph with $i - 1$ vertices.

6.3. Comparison with formulas (13). Combining (39) and the explicit values (42) and (43), we obtain from (25) the desired expressions (17). It simply remains to show that our definitions for y and α of Section 6.2, given by (42) and (43) just above, match their definitions of Section 4 given by (12), or equivalently (14) and (15). If so, (17) is equivalent to (13) and we are done.

Using for y and α their definitions of Section 4 (through (14) and (15)), we obtain for $Y = Q - P$ and Y_1 (whose value in terms of P and Q is given by (32)) the parametrizations

$$\begin{aligned} Y &= \frac{(\alpha - 1)y(1 - \alpha y^2)}{1 + y + \alpha y - 6\alpha y^2 + \alpha y^3 + \alpha^2 y^3 + \alpha^2 y^4}, \\ Y_1 &= \frac{(\alpha - 1)y(1 - \alpha y^3)}{(1 + y)(1 + y + \alpha y - 6\alpha y^2 + \alpha y^3 + \alpha^2 y^3 + \alpha^2 y^4)}. \end{aligned}$$

so that

$$d = -\frac{y(1-\alpha y)}{1-\alpha y^3} \quad \text{and} \quad \frac{1}{y^2} \frac{d+y}{1+dy} = \alpha$$

as wanted to match the definition (43) of α of Section 6.2.

As for y , we use the expressions (28) of A_0 and A_1 to get the parametrizations

$$A_0 = \frac{(1-\alpha y^2)^2}{(1-\alpha y)(1-\alpha y^3)},$$

$$A_1 = -\frac{1+y+\alpha y-6\alpha y^2+\alpha y^3+\alpha^2 y^3+\alpha^2 y^4}{(1-\alpha y)(1-\alpha y^3)}.$$

so that

$$w = A_0 A_1 \frac{Y_1^2}{Y^2} P = -\frac{1}{y+y^{-1}+2}$$

as wanted to match the definition (42) of y of Section 6.2.

A more constructive approach consists in starting instead from the definitions of y and α of Section 6.2 (through (42) and (43)) and recovering the parametrization (14) of Section 4. From (28) (and $t_\bullet = P(1-P-2Q)$), A_0 can be expressed in terms of P and Q , as well as Y ($= Q - P$) and Y_1 via (32). This leads to

$$w \equiv A_0 A_1 \frac{Y_1^2}{Y^2} P = -\frac{P(1-Q-P)}{(1-2Q)^2} = -\frac{1}{y+y^{-1}+2}$$

hence we deduce the (so-called characteristic) equation

$$(1-2Q)^2 - (2+y+y^{-1})P(1-P-Q) = 0.$$

This in turn leads to

$$d \equiv \frac{Y_1 - Y}{Y_1} = -\frac{P}{1-P-2Q}, \quad \alpha \equiv \frac{1}{y^2} \frac{d+y}{1+dy} = -\frac{P-y(1-P-2Q)}{y^2(1-P-2Q-yP)}.$$

Using this latter equation to express Q in terms of P , y and α , namely

$$Q = -\frac{P(1+y)(1-\alpha y^2)}{2y(1-\alpha y)} + \frac{1}{2},$$

and plugging this value in the characteristic equation above, we find that P is determined by

$$-y(1-\alpha y)^2 + P(1+y+\alpha y-6\alpha y^2+\alpha y^3+\alpha^2 y^3+\alpha^2 y^4) = 0,$$

from which (14) follows.

7. Conserved quantities

The explicit formulas (11) (resp. (13) or equivalently (17)) are typical expressions for the solutions of *discrete integrable systems*. A deeper characterization of the integrability of the system (8) (resp. (9) or (16)) is the existence of a number of *discrete conserved quantities*, i.e. quantities whose expression depends explicitly on some positive integer (called d below) but whose value turns out to be independent of this integer. In the case of bicolored quadrangulations, it has already been recognized that these conserved quantities may be obtained by looking for a direct combinatorial derivation of F_n in the slice formalism. Let us first recall this construction and then see how it extends to the case of our second weighting governed by local maxima.

7.1. Conserved quantities for the first weighting. The slice decomposition described in [7, 15] applies more generally to *pointed rooted* quadrangulations with boundaries, that is, quadrangulations with a boundary, having a root-edge on the boundary oriented with the boundary-face on its right, and having a pointed vertex (which might not be incident to the boundary-face); the *label* $d(v)$ of each vertex v is now the distance from v to the pointed vertex v_0 . For Q such a map, the *canonical bicolouration* of Q is the vertex bicolouration in black and white where the root-vertex (origin of the root) is black and any two adjacent vertices have different colors. As before, a *local maximum* for the distance is a vertex v such that $d(v) = d(v') + 1$ for every neighbor v' of v . For $n \geq 1$ and $d \geq 0$, let $\mathcal{B}_n^{(d)}$ be the family of admissible pointed rooted quadrangulations with a boundary of length $2n$, where the root-vertex is at distance *at most* d from v_0 , and is *one* (possibly not unique) *of the boundary-vertices that reach the smallest distance from v_0* . Let $F_n^{(d)}$ be the generating function of $\mathcal{B}_n^{(d)}$ where each black vertex (resp. white vertex) receives weight t_\bullet (resp. t_\circ) except for the pointed vertex that receives weight 1. And let $Z_{d,d}^+(2n; \{B_i\}_{i \geq 1}, \{W_i\}_{i \geq 1})$ be the generating functions of paths of length $2n$ starting and ending at height d and staying at height at least d all along, made of elementary steps with height difference ± 1 , with each descending step from height i to height $i - 1$ weighted by B_i if $i \equiv d \pmod 2$ and weighted by W_i if $i \equiv d + 1 \pmod 2$.

Note that $\mathcal{B}_n^{(0)}$ is nothing but the set of rooted quadrangulations with a boundary of length $2n$ so that $F_n^{(0)} = F_n$. Then, as explained in [7, 15], the slice decomposition described in Section 3.1 for maps in $\mathcal{B}_n^{(0)}$ applies more generally for maps in $\mathcal{B}_n^{(d)}$ and yields

$$F_n^{(d)} = Z_{d,d}^+(2n; \{B_i\}_{i \geq 1}, \{W_i\}_{i \geq 1}).$$

Now let $\widehat{\mathcal{B}}_n^{(d)} = \mathcal{B}_n^{(d)} \setminus \mathcal{B}_n^{(0)}$ be the subfamily of $\mathcal{B}_n^{(d)}$ where the pointed vertex is *different from the root-vertex*, and let $\widehat{F}_n^{(d)}$ be the generating function for the subfamily $\widehat{\mathcal{B}}_n^{(d)}$ where the weights are specified as in $F_n^{(d)}$. For $Q \in \widehat{\mathcal{B}}_n^{(d)}$, with v_0 the pointed vertex and \vec{v} the root-vertex, let e be the first edge of the *leftmost* geodesic path from \vec{v} to v_0 . This edge cannot be a boundary-edge as otherwise, \vec{v} would not reach the smallest distance from v_0 among boundary-vertices. We may cut along e (starting from \vec{v}) so as to duplicate e into two edges e_1, e_2 (with e_2 before e_1 in ccw order around the new map) and duplicate \vec{v} into two vertices v_1, v_2 (see Figure 11).

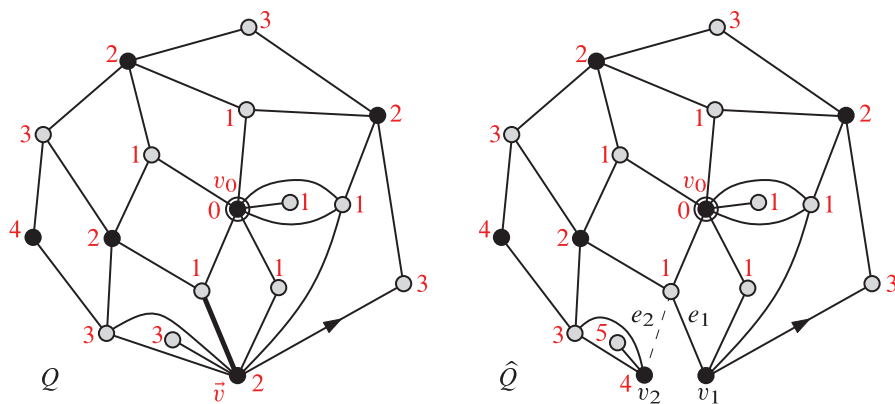


Figure 11. Left: a pointed rooted quadrangulation Q with a boundary. Right: the pointed rooted quadrangulation \widehat{Q} obtained by cutting along the first edge of the leftmost geodesic path from the root-vertex to the pointed vertex.

Let \widehat{Q} be the pointed rooted quadrangulation with a boundary of length $2n + 4$ that is obtained by erasing e_2 , taking v_1 as the new root-vertex, and keeping v_0 as the pointed vertex. Denoting by $\mathbf{d} = (d_1, \dots, d_{2n+4})$ the distances from v_0 of the successive boundary-vertices (starting from v_1) in ccw order around \widehat{Q} , we have the conditions that $d_{i+1} = d_i \pm 1$ for $i \in \{1, \dots, 2n + 3\}$, d_1 equals the distance of \vec{v} from v_0 in Q so that $d_1 \leq d$, $d_i \geq d_1$ for all $i \in \{1, \dots, 2n + 1\}$, $d_{2n+1} > d_1$ (indeed, by the effect of cutting along the first edge of the leftmost geodesic path, the distance of v_2 from v_0 is strictly larger than the distance of v_1 from v_0), and $d_{2n+4} = d_1 - 1$. The bipartiteness of \widehat{Q} implies that $d_{2n+1} \equiv d_1 \pmod{2}$, so that the last entries of \mathbf{d} must be $(d_1 + 2, d_1 + 1, d_1, d_1 - 1)$. Hence, if for $k \geq 1$ we denote by $Z_{d,d}^{+,k} \searrow (2n; \{B_i\}_{i \geq 1}, \{W_i\}_{i \geq 1})$ the generating function defined as $Z_{d,d}^+(2n; \{B_i\}_{i \geq 1}, \{W_i\}_{i \geq 1})$, but with the restriction that the k last steps of the

path are descending, then the slice decomposition applied to \widehat{Q} gives

$$\widehat{F}_n^{(d)} = \frac{1}{t_\bullet} Z_{d,d}^{+,2\searrow}(2n + 2; \{B_i\}_{i \geq 1}, \{W_i\}_{i \geq 1}) \cdot B_d,$$

where the factor $\frac{1}{t_\bullet}$ accounts for the (black) root-vertex being duplicated and the factor B_d accounts for the last descent $d_1, d_1 - 1$. Hence for each $n \geq 1$ we have the conserved quantity

$$F_n = Z_{d,d}^+(2n; \{B_i\}_{i \geq 1}, \{W_i\}_{i \geq 1}) - \frac{1}{t_\bullet} Z_{d,d}^{+,2\searrow}(2n + 2; \{B_i\}_{i \geq 1}, \{W_i\}_{i \geq 1}) \cdot B_d. \tag{44}$$

The first two conserved quantities, $n \in \{1, 2\}$, are (with $r = d + 1$): for all $r \geq 1$ (with $B_0 = 0$)

$$F_1 = W_r - \frac{1}{t_\bullet} B_{r+1} W_r B_{r-1},$$

$$F_2 = W_r^2 + B_{r+1} W_r - \frac{1}{t_\bullet} (W_r + B_{r+1} + W_{r+2}) B_{r+1} W_r B_{r-1}.$$

Shifting in (44) all path heights by $-d$ and replacing B_i and W_i by B_{i+d} and W_{i+d} so as to compensate this shift, we get, upon sending $d \rightarrow \infty$ the identity

$$F_n = Z_{0,0}^+(2n; B, W) - \frac{1}{t_\bullet} Z_{0,0}^{+,2\searrow}(2n + 2; B, W) \cdot B \tag{45}$$

(with some obvious notations) which, using the identities

$$Z_{0,0}^{+,2\searrow}(2n + 2; B, W) = Z_{0,0}^+(2n + 2; B, W) - Z_{0,0}^+(2n; B, W) \cdot W$$

and

$$B = t_\bullet + B(B + 2W),$$

is easily transformed into (27).

7.2. Conserved quantities for the second weighting. We may now play a similar game for the quantities J_n to obtain conserved quantities involving the generating functions $\{P_r, Q_r\}_{r \geq 1}$. Let $J_n^{(d)}$ be the generating function of $\mathcal{B}_n^{(d)}$ where vertices that are (resp. are not) local maxima receives weight t_\bullet (resp. t_\circ) except for the pointed vertex that receives weight 1. And let $\widehat{Z}_{d,d}^+(2n; \{P_i\}_{i \geq 1}, \{Q_i\}_{i \geq 1})$ be the generating function of paths of length $2n$ starting and ending at height d and staying at height at least d all along, made of elementary steps with height difference ± 1 , with each descending step from height i to height $i - 1$ weighted by P_i if just after a descent and weighted by Q_i if just after an ascent.

Again, the slice decomposition described in Section 3.2 for maps in $\mathcal{B}_n^{(0)}$ applies more generally for maps in $\mathcal{B}_n^{(d)}$ and yields

$$J_n^{(d)} = \widehat{Z}_{d,d}^+(2n; \{P_i\}_{i \geq 1}, \{Q_i\}_{i \geq 1}).$$

Let $\mathcal{M}_n^{(d)}$ be the family of rooted pointed general maps with a bridgeless boundary of length n , where the root-vertex is at distance at most d from the pointed vertex v_0 , and is at least as close from v_0 as any other boundary-vertex (here boundary-edges are directed ccw around the map while inner edges are bi-directed; the distance-label $d(v)$ is the length of a shortest directed path starting from v_0 and ending at v). The Ambjørn–Budd bijection described in Section 2.2 between $\mathcal{M}_n^{(0)}$ and $\mathcal{B}_n^{(0)}$ extends verbatim (using the same local rules, and having the same pointed vertex and the same root-vertex in corresponding maps, see [1, 6]) to a bijection between $\mathcal{B}_n^{(d)}$ and $\mathcal{M}_n^{(d)}$, so that $J_n^{(d)}$ is also the generating function of maps in $\mathcal{M}_n^{(d)}$ with a weight t_\bullet for each non-pointed vertex and a weight t_\circ for each inner face.

Let $\widehat{\mathcal{M}}_n^{(d)}$ be the subfamily of $\mathcal{M}_n^{(d)}$ where the pointed vertex is different from the root-vertex, and let $\widehat{J}_n^{(d)}$ be the generating function of the subfamily $\widehat{\mathcal{B}}_n^{(d)}$ where the weights are as in $J_n^{(d)}$. Then the Ambjørn–Budd bijection ensures that $\widehat{J}_n^{(d)}$ is also the generating function of $\widehat{\mathcal{M}}_n^{(d)}$ with a weight t_\bullet for each non-pointed vertex and a weight t_\circ for each inner face. For a map $M \in \widehat{\mathcal{M}}_n^{(d)}$, let e be the first edge on the leftmost geodesic path from the root-vertex \vec{v} to the pointed vertex v_0 (note that all the edges on this path are inner edges, and all vertices on this path are inner vertices except for \vec{v} , in particular the tip of e is an inner vertex). Again we can cut along e (starting from \vec{v}) so as to duplicate e into two edges e_1, e_2 (with e_2 before e_1 in ccw order around the map) and duplicate \vec{v} into two vertices v_1, v_2 , and take v_1 as the new root-vertex (see Figure 12).

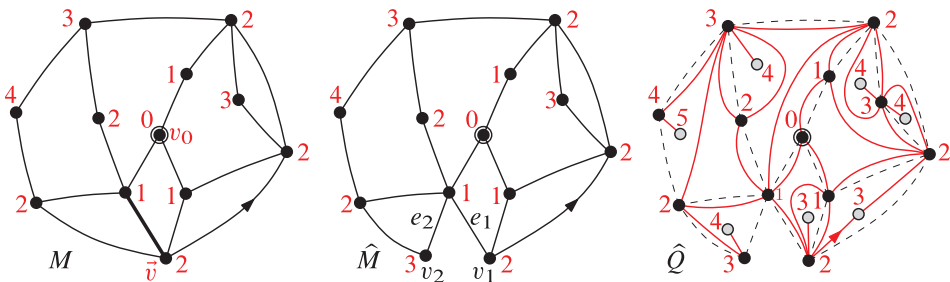


Figure 12. Left: a pointed rooted map M with a bridgeless boundary. Middle: the pointed rooted map \widehat{M} obtained by cutting along the first edge of the leftmost geodesic path from the root-vertex to the pointed vertex. Right: the associated quadrangulation \widehat{Q} with a boundary (see text).

The map \widehat{M} thus obtained (as opposed to the quadrangulated case we do not delete e_2) is a general map with a bridgeless boundary of length $n + 2$. If we denote by $\delta_1, \dots, \delta_{n+2}$ the distances from the pointed vertex v_0 of the successive boundary-vertices (starting with v_1) in ccw order around \widehat{M} , then δ_1 equals the distance of \bar{v} from v_0 in M so that $\delta_1 \leq d$, $\delta_i \geq \delta_1$ for $i \in \{1, \dots, n + 1\}$, $\delta_{i+1} \leq \delta_i + 1$ for $i \in \{1, \dots, n + 1\}$, $\delta_{n+1} > \delta_1$ (by the effect of cutting along the first edge of the leftmost geodesic path) and $\delta_{n+2} = \delta_1 - 1$. In particular if we reroot the map at the vertex between e_1 and e_2 , we get a map $M' \in \mathcal{M}_{n+2}^{(d-1)}$. We may then take the image $Q' \in \mathcal{B}_{n+2}^{(d-1)}$ of M' by the Ambjørn–Budd bijection, and denote by \widehat{Q} the quadrangulation with boundary obtained from Q' by shifting the root position by one in ccw order around Q' ; note also that the number of vertices that are (resp. are not) local maxima of \widehat{Q} equals the number of inner faces (resp. the number of vertices) of M' , which is also the number of inner faces (resp. the number of vertices plus 1) of M . Let again $\mathbf{d} = (d_1, \dots, d_{2n+4})$ be the distances from the pointed vertex v_0 of the successive boundary-vertices (starting with v_1) in ccw order around \widehat{Q} . By the local rules of the Ambjørn–Budd bijection, \mathbf{d} is obtained from the sequence $\delta_1, \dots, \delta_{n+2}$ where for each $i \in \{1, \dots, n + 1\}$, we insert between δ_i and δ_{i+1} the subsequence (of length $\delta_i - \delta_{i+1} + 1$) $\delta_i + 1, \delta_i, \dots, \delta_{i+1} + 1$. It is then easy to check that \mathbf{d} satisfies the following conditions: $d_1 = \delta_1$, $d_{i+1} = d_i \pm 1$ and $d_i \geq \delta_1$ for $i \in \{1, \dots, 2n + 3\}$, $d_{2n+4} = \delta_1 - 1$, and \mathbf{d} ends with $\delta_1 + 2, \delta_1 + 1, \delta_1, \delta_1 - 1$ (since $\delta_{n+1} > \delta_1$). Hence, if for $k \geq 1$ we denote by $\widehat{Z}_{d,d}^{+,k \searrow}(2n; \{P_i\}_{i \geq 1}, \{Q_i\}_{i \geq 1})$ the generating function defined as $\widehat{Z}_{d,d}^+(2n; \{P_i\}_{i \geq 1}, \{Q_i\}_{i \geq 1})$, but with the restriction that the k last steps of the path are descending, then the slice decomposition applied to \widehat{Q} gives

$$\widehat{J}_n^{(d)} = \frac{1}{t_\bullet} \widehat{Z}_{d,d}^{+,2 \searrow}(2n + 2; \{P_i\}_{i \geq 1}, \{Q_i\}_{i \geq 1}) \cdot P_d,$$

where the factor $\frac{1}{t_\bullet}$ accounts for the root-vertex of M being duplicated, and the factor P_d accounts for the last descent $\delta_1, \delta_1 - 1$. Hence for each $n \geq 1$ we have the conserved quantity

$$J_n = \widehat{Z}_{d,d}^+(2n; \{P_i\}_{i \geq 1}, \{Q_i\}_{i \geq 1}) - \frac{1}{t_\bullet} \widehat{Z}_{d,d}^{+,2 \searrow}(2n + 2; \{P_i\}_{i \geq 1}, \{Q_i\}_{i \geq 1}) \cdot P_d. \tag{46}$$

Remarkably this has exactly the same form as the bicolored conserved quantities (44), up to changing $\{P_i, Q_i\}_{i \geq 1}$ for $\{B_i, W_i\}_{i \geq 1}$ and taking the “hat” variants of the path generating functions. The first two invariants, $n \in \{1, 2\}$, are (with $r = d + 1$): for all $r \geq 1$ (with $P_0 = 0$)

$$J_1 = Q_r - \frac{1}{t_\bullet} Q_{r+1} P_r P_{r-1},$$

$$J_2 = Q_r^2 + Q_{r+1} P_r - \frac{1}{t_\bullet} ((Q_r + Q_{r+1}) Q_{r+1} + Q_{r+2} P_{r+1}) P_r P_{r-1}.$$

As before, upon sending $d \rightarrow \infty$ in (46), we get the expression

$$J_n = \widehat{Z}_{0,0}^+(2n; P, Q) - \frac{1}{t_\bullet} \widehat{Z}_{0,0}^{+,2\neg}(2n+2; P, Q) \cdot P$$

(with straightforward notations). Upon using $P = B$, $Q = W$ and comparing with (45), this provides another (computational) proof of the identity $J_n = F_n$ by noting that

$$\widehat{Z}_{0,0}^+(2n; P, Q) = Z_{0,0}^+(2n; P, Q)$$

and⁷

$$\widehat{Z}_{0,0}^{+,2\neg}(2n+2; P, Q) = Z_{0,0}^{+,2\neg}(2n+2; P, Q).$$

Finally, from (9), we get $Y_1 = (Q_1 - P_1) = t_\circ - t_\bullet + (Q_1 - P_1)Q_1$, hence $Y_1 = (t_\circ - t_\bullet)/(1 - Q_1)$. Using the first conserved quantity above, we deduce $Q_1 = J_1 = Q - QP^2/t_\bullet$ so that $Y_1 = t_\bullet(t_\circ - t_\bullet)/(t_\bullet - t_\bullet Q + QP^2)$ which upon expressing t_\bullet and t_\circ in terms of P and Q via (15), reproduces the expression (32) for Y_1 .

⁷ The identity

$$\widehat{Z}_{0,0}^+(2n; P, Q) = Z_{0,0}^+(2n; P, Q)$$

is easily proved by noting that the equation

$$\widehat{Z} = \frac{1}{(1 - z(Q - P) - zP \widehat{Z})}$$

that determines the generating function

$$\widehat{Z} \equiv \sum_{n \geq 0} \widehat{Z}_{0,0}^+(2n; P, Q) z^n$$

is identical to the equation

$$Z = \frac{1}{(1 - zQ/(1 - zP Z))}$$

that determines the generating function

$$Z \equiv \sum_{n \geq 0} Z_{0,0}^+(2n; P, Q) z^n,$$

hence $\widehat{Z} = Z$.

The identity

$$\widehat{Z}_{0,0}^{+,2\neg}(2n+2; P, Q) = Z_{0,0}^{+,2\neg}(2n+2; P, Q)$$

follows by noting that

$$\widehat{Z}_{0,0}^{+,2\neg}(2n+2; P, Q) = \widehat{Z}_{0,0}^+(2n+2; P, Q) - \widehat{Z}_{0,0}^+(2n; P, Q) \cdot Q$$

and similarly

$$Z_{0,0}^{+,2\neg}(2n+2; P, Q) = Z_{0,0}^+(2n+2; P, Q) - Z_{0,0}^+(2n; P, Q) \cdot Q.$$

8. Conclusion

In this paper, we presented a comparative study of two statistical ensembles of quadrangulations. We first showed how the corresponding slice generating functions (B_i, W_i for the first ensemble and P_i, Q_i for the second) appear as coefficients of the same quantity $F(z) = J(z)$, expanded as a continued fraction in two different ways. The slice generating functions may then be written as biratios of Hankel-type determinants and explicit formulas may be obtained, at the price of some conjectured expression for some intermediate quantity in the second ensemble.

To conclude, we would like to emphasize that our two ensembles may be viewed, in some sense, as the two extremal elements of a very general family of statistical ensembles as follows: by definition, the second ensemble gives a particular weight to those vertices which are local maxima for the distance to the root-vertex. Similarly, the first ensemble may be viewed as the ensemble which gives a particular weight to those vertices which are local maxima for the *distance to the root-vertex modulo 2*. Indeed, this distance modulo 2 is 0 for black vertices (recall that the root-vertex is black) and 1 for white vertices so that all white vertices are local maxima. In this respect, note also that performing the passage from the quadrangulation to the general map in the bijection of Figure 2 may be viewed as applying the Ambjørn–Budd rules, taking as labeling the distance modulo 2.

Denoting by $d(v)$ the distance from a vertex v to the root-vertex in a rooted quadrangulation with a boundary, we may more generally consider statistical ensembles which give a particular weight to those vertices which are local maxima for some labeling $\ell(d(v))$, with $d \mapsto \ell(d)$ some given function. Without loss of generality, we may set $\ell(0) = 0$ and, if we wish to apply the Ambjørn–Budd rules to transform our quadrangulation into a general map, we need that $|\ell(d) - \ell(d - 1)| = 1$ (it also seems natural to impose that $\ell(d)$ remains non negative so that the root-vertex cannot be a local maximum). It is likely that slice generating functions in this ensemble may appear as coefficient of $F(z) = J(z)$, once expanded as a continued fraction with some appropriate structure, being a mixture of the Stieljes-type and of our new encountered type. At this stage, it is interesting to notice that, in their study of *finite* continued fractions [12, 13], Di Francesco and Kedem introduced precisely a whole family of such “mixed” fractions as well as some passage rules on their coefficients to go from one to the other without changing the actual value of the fraction. It is very tempting to speculate that their study may be extended to infinite continued fractions to describe our more general ensembles.

Appendices

A. A proof of formulas (25) and (26)

As in [12, 13], our proof of formulas (25) and (26) is based on the theory of *heaps of pieces*. The reader is invited to consult [21] for the basics of this theory.

Let us simply recall what we mean by a heap of pieces on a graph \mathcal{G} , supposedly connected, planar, and drawn in a *horizontal* plane for simplicity. Imagine to complete the graph by a set of *vertical* half-lines, with a half-line starting from each vertex of the graph. Informally speaking, a heap is a collection of pieces threaded along these half-lines. Each piece therefore sits on top of a given vertex and may move freely along the corresponding vertical half-line as long as it does not meet another piece. More precisely, the pieces are supposed to be designed so that two pieces may not pass each other if and only if they sit on top of the same vertex or on top of adjacent vertices.

Given a subset \mathcal{B} of the set of vertices of \mathcal{G} , a heap of pieces is said to be of *base* \mathcal{B} if, moving its pieces as far as possible to the bottom of the half-lines, the set of those vertices hit by a piece forms a subset of \mathcal{B} (see Figure 13).

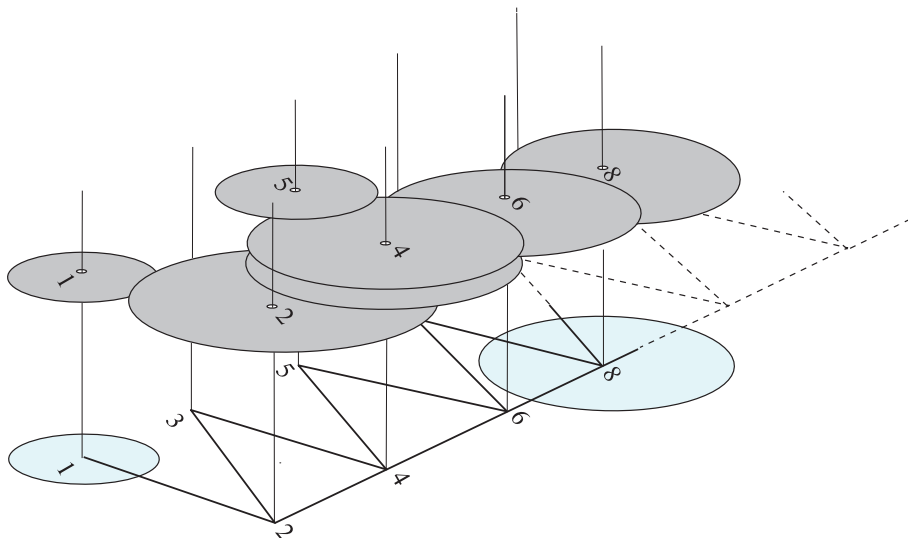


Figure 13. An example of heap of 7 pieces sitting on top of the graph \mathcal{G} of Figure 14 with base $\{1, 8\}$ (we indicated in light blue the “shadow” of those pieces which can move freely and hit the vertices of the graph). The diameter of the pieces is adjusted so that pieces sitting on top of vertices which are adjacent in \mathcal{G} cannot pass through each other.

A fundamental remark is that, from the relation (6), $J(z)$ may be viewed as the generating function for heaps of pieces on the *semi-infinite* graph \mathcal{G} of Figure 14, with a weight zY_i per piece sitting at position i along the graph, and whose base is $\{1, 2\}$. Similarly, we may interpret $\tilde{J}(z)$ as the generating function for the very same heaps, but now with a weight $z\tilde{Y}_i$ per piece sitting at position i . Let us finally introduce the quantity

$$K(z) \equiv 1 + z Y_1 J(z) = \frac{1}{1 - z \frac{Y_1}{1 - z \frac{Y_2}{1 - z \frac{Y_4}{1 - z Y_3 - z \frac{Y_6}{1 - z Y_5 - z \frac{Y_6}{1 - \dots}}}}}}$$

which is the generating function for heaps of pieces on the graph \mathcal{G} again with a weight zY_i per piece sitting at position i along the graph, but now with base $\{1\}$.

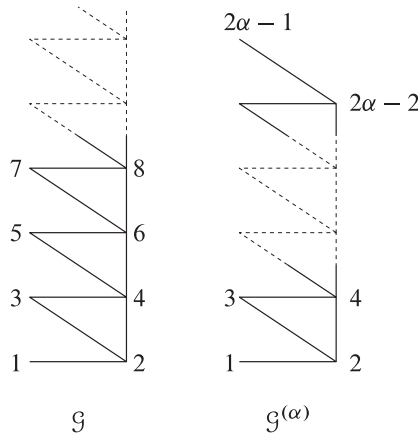


Figure 14. The semi-infinite graph \mathcal{G} and the finite graph $\mathcal{G}^{(\alpha)}$.

From the definition (23) of the j_n 's, we have

$$K(z) = \sum_{n \geq 0} j_n z^n, \quad \tilde{J}(z) = \sum_{n \geq 0} j_{-n} z^n,$$

so that all the j_n 's have a direct interpretation as enumerating heap configurations made of $|n|$ pieces.

Let us now consider the analogs $J^{(\alpha)}(z)$, $K^{(\alpha)}(z)$ and $\tilde{J}^{(\alpha)}(z)$ of $J(z)$, $K(z)$ and $\tilde{J}(z)$ respectively, viewed as heap generating functions, now defined on the *finite* graph $\mathcal{G}^{(\alpha)}$ of Figure 14. In other words, we set

$$J^{(\alpha)}(z) \equiv \frac{1}{1 - zY_1 - z \frac{Y_2}{1 - zY_3 - z \frac{Y_4}{\ddots \frac{Y_{2\alpha-2}}{1 - zY_{2\alpha-3} - z \frac{Y_{2\alpha-2}}{1 - zY_{2\alpha-1}}}}}}$$

$$\tilde{J}^{(\alpha)}(z) \equiv \frac{1}{1 - z\tilde{Y}_1 - z \frac{\tilde{Y}_2}{1 - z\tilde{Y}_3 - z \frac{\tilde{Y}_4}{\ddots \frac{\tilde{Y}_{2\alpha-2}}{1 - z\tilde{Y}_{2\alpha-3} - z \frac{\tilde{Y}_{2\alpha-2}}{1 - z\tilde{Y}_{2\alpha-1}}}}}}$$

$$\begin{aligned} K^{(\alpha)}(z) &\equiv 1 + zY_1 J^{(\alpha)}(z) \\ &= \frac{1}{1 - z \frac{Y_1}{1 - z \frac{Y_2}{1 - zY_3 - z \frac{Y_4}{\ddots \frac{Y_{2\alpha-2}}{1 - zY_{2\alpha-3} - z \frac{Y_{2\alpha-2}}{1 - zY_{2\alpha-1}}}}}}}. \end{aligned}$$

We finally define the analogs $j_n^{(\alpha)}$ of j_n via

$$\begin{aligned} K^{(\alpha)}(z) &= \sum_{n \geq 0} j_n^{(\alpha)} z^n, \\ \tilde{J}^{(\alpha)}(z) &= \sum_{n \geq 0} j_{-n}^{(\alpha)} z^n, \end{aligned}$$

so that $j_n^{(\alpha)}$ ($n \geq 0$) enumerates heap configurations of n pieces on $\mathcal{G}^{(\alpha)}$ with weights Y_i and base $\{1\}$, and $j_{-n}^{(\alpha)}$ ($n \geq 0$) enumerates heap configurations of n pieces on $\mathcal{G}^{(\alpha)}$ with weights \tilde{Y}_i and base $\{1, 2\}$.

It is now a standard result of the theory of heaps of pieces [21] that⁸

$$K^{(\alpha)}(z) = \frac{X^{(\alpha)}(0, -zY_2, -zY_3, \dots, -zY_{2\alpha-1})}{X^{(\alpha)}(-zY_1, -zY_2, -zY_3, \dots, -zY_{2\alpha-1})} \tag{47a}$$

$$\tilde{J}^{(\alpha)}(z) = \frac{X^{(\alpha)}(0, 0, -z\tilde{Y}_3, \dots, -z\tilde{Y}_{2\alpha-1})}{X^{(\alpha)}(-z\tilde{Y}_1, -z\tilde{Y}_2, -z\tilde{Y}_3, \dots, -z\tilde{Y}_{2\alpha-1})} \tag{47b}$$

where $X^{(\alpha)}(y_1, y_2, y_3, \dots, y_{2\alpha-1})$ denotes the generating function of *hard pieces* on the graph $\mathcal{G}^{(\alpha)}$, each piece sitting at position i receiving the weight y_i . Recall that, by definition, in a configuration of hard pieces, each vertex of the graph is occupied by at most one piece, with *no two adjacent vertices occupied simultaneously*. Note that the positions of the 0's in the numerators correspond to the location of the vertices of the corresponding base of the heaps ($\{1\}$ and $\{1, 2\}$ respectively). Clearly, on the graph $\mathcal{G}^{(\alpha)}$, we can put at most α hard pieces. Moreover, this maximal situation is achieved by a single configuration with all sites with odd index occupied (see Figure (15)).

The quantity $X^{(\alpha)}(-zY_1, -zY_2, -zY_3, \dots, -zY_{2\alpha-1})$ is therefore a polynomial of degree α in z that we write

$$X^{(\alpha)}(-zY_1, -zY_2, -zY_3, \dots, -zY_{2\alpha-1}) = \sum_{m=0}^{\alpha} (-z)^m X_m^{(\alpha)}(Y_1, Y_2, Y_3, \dots, Y_{2\alpha-1})$$

where $X_m^{(\alpha)}(y_1, y_2, y_3, \dots, y_{2\alpha-1})$ denotes the generating function of exactly m *hard pieces* on the graph $\mathcal{G}^{(\alpha)}$, each piece sitting at position i receiving the weight y_i . Clearly, both

$$X^{(\alpha)}(0, -zY_2, -zY_3, \dots, -zY_{2\alpha-1}) \quad \text{and} \quad X^{(\alpha)}(0, 0, -z\tilde{Y}_3, \dots, -z\tilde{Y}_{2\alpha-1})$$

are polynomials of degree $\alpha - 1$ in z .

⁸ A sketch of the proof is as follows: given \mathcal{B} , consider *pairs* $(\mathcal{H}, \mathcal{H}\mathcal{P})$ made of a heap configuration \mathcal{H} of base \mathcal{B} together with a configuration $\mathcal{H}\mathcal{P}$ of hard pieces, drawn on top of the heap. For such a pair, let \mathcal{E} be the set of pieces that can be moved up freely to infinity, and when pushed downward either are blocked by a piece (that has to be in \mathcal{H}) or hit a vertex of the base \mathcal{B} . Consider the following transformation: if \mathcal{E} is not empty, pick the piece $p \in \mathcal{E}$ of smallest index and change its status (from \mathcal{H} to $\mathcal{H}\mathcal{P}$ if $p \in \mathcal{H}$, from $\mathcal{H}\mathcal{P}$ to \mathcal{H} if $p \in \mathcal{H}\mathcal{P}$); if \mathcal{E} is empty do nothing. This transformation is easily seen to be an involution (which leaves \mathcal{E} invariant), and, if we assign a weight z per piece in the heap and $-z$ per piece in the configuration of hard pieces, the weight is multiplied by -1 for each configuration which changes under the involution. The generating function for the pairs, which is the product of the generating function for heaps with a weight z per piece times that of configurations of hard pieces with a weight $-z$ per piece, therefore reduces to those pairs for which \mathcal{E} is empty. It is easily seen that this situation corresponds to an empty heap and a configuration of hard pieces made of pieces which do not belong to \mathcal{B} . The corresponding generating function is nothing but that of configurations of hard pieces with a weight $-z$ per piece not in \mathcal{B} and 0 per piece in \mathcal{B} .

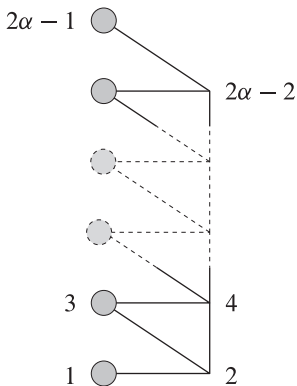


Figure 15. The maximally occupied configuration of hard pieces on the graph $\mathcal{G}^{(\alpha)}$, made of α pieces (here represented by gray circles) sitting on all the sites with an odd index.

Let us now come to our fundamental identities. We have

$$X_m^{(\alpha)} = X_\alpha^{(\alpha)} \tilde{X}_{\alpha-m}^{(\alpha)}, \tag{48a}$$

$$X_m^{(\alpha)}(0) = X_\alpha^{(\alpha)}(\tilde{X}_{\alpha-m}^{(\alpha)} - \tilde{X}_{\alpha-m}^{(\alpha)}(0, 0)), \tag{48b}$$

with the short-hand notations

$$X_m^{(\alpha)} \equiv X_m^{(\alpha)}(Y_1, Y_2, Y_3, \dots, Y_{2\alpha-1}),$$

$$X_m^{(\alpha)}(0) \equiv X_m^{(\alpha)}(0, Y_2, Y_3, \dots, Y_{2\alpha-1}),$$

$$\tilde{X}_m^{(\alpha)} \equiv X_m^{(\alpha)}(\tilde{Y}_1, \tilde{Y}_2, \tilde{Y}_3, \dots, \tilde{Y}_{2\alpha-1}),$$

$$\tilde{X}_m^{(\alpha)}(0, 0) \equiv X_m^{(\alpha)}(0, 0, \tilde{Y}_3, \dots, \tilde{Y}_{2\alpha-1}).$$

To explain these identities, let us analyze the structure of a configuration C of hard pieces on $\mathcal{G}^{(\alpha)}$. In C , a number k of pieces occupy even sites $2j_1, 2j_2, \dots, 2j_k$ with $1 \leq j_1 < j_2 < \dots < j_k \leq \alpha - 1$ and $j_{\ell+1} - j_\ell > 1$ for $\ell = 1, \dots, k - 1$. The set of available odd sites is

$$\text{Odd} = \{1, 3, 5, \dots, 2\alpha - 1\} \setminus \{2j_1 - 1, 2j_1 + 1, 2j_2 - 1, 2j_2 + 1, \dots, 2j_k - 1, 2j_k + 1\}$$

and satisfies

$$|\text{Odd}| = \alpha - 2k.$$

A number k' of pieces occupy a subset $\{2j'_1 - 1, 2j'_2 - 1, \dots, 2j'_{k'} - 1\}$ of this set. In $X_m^{(\alpha)}$ (corresponding to a situation where $k + k' = m$), any occupied site i receives the weight Y_i , so the weight of the configuration C is $Y_{2j_1} Y_{2j_2} \dots Y_{2j_k} \times Y_{2j'_1-1} Y_{2j'_2-1} \dots Y_{2j'_{k'}-1}$. Let us now consider instead the configuration \tilde{C} where, again, the sites $2j_1, 2j_2, \dots, 2j_k$ are occupied but now the complementary of $\{2j'_1 - 1, 2j'_2 - 1, \dots, 2j'_{k'} - 1\}$ in Odd (namely $\text{Odd} \setminus \{2j'_1 - 1, 2j'_2 - 1, \dots, 2j'_{k'} - 1\}$) is covered by pieces. Clearly, going in from C to \tilde{C} provides a bijection between configurations C with

$$k + k' = m$$

pieces and configurations \tilde{C} with

$$k + (\alpha - 2k) - k' = \alpha - m$$

pieces. In $\tilde{X}_{\alpha-m}$, the configuration \tilde{C} receives the weight

$$\begin{aligned} & \tilde{Y}_{2j_1} \tilde{Y}_{2j_2} \dots \tilde{Y}_{2j_k} \times \frac{\prod_{i \in \text{Odd}} \tilde{Y}_i}{\tilde{Y}_{2j'_1-1} \tilde{Y}_{2j'_2-1} \dots \tilde{Y}_{2j'_{k'}-1}} \\ &= \frac{Y_{2j_1} Y_{2j_2} \dots Y_{2j_k}}{Y_{2j_1-1} Y_{2j_1+1} Y_{2j_2-1} Y_{2j_2+1} \dots Y_{2j_k-1} Y_{2j_k+1}} \times \frac{Y_{2j'_1-1} Y_{2j'_2-1} \dots Y_{2j'_{k'}-1}}{\prod_{i \in \text{Odd}} Y_i} \\ &= \frac{Y_{2j_1} Y_{2j_2} \dots Y_{2j_k} \times Y_{2j'_1-1} Y_{2j'_2-1} \dots Y_{2j'_{k'}-1}}{Y_1 Y_3 Y_5 \dots Y_{2\alpha-1}} \\ &= \frac{Y_{2j_1} Y_{2j_2} \dots Y_{2j_k} \times Y_{2j'_1-1} Y_{2j'_2-1} \dots Y_{2j'_{k'}-1}}{X_\alpha^{(\alpha)}} \end{aligned}$$

since $X_\alpha^{(\alpha)} = Y_1 Y_3 Y_5 \dots Y_{2\alpha-1}$. From the bijection $C \mapsto \tilde{C}$, we therefore deduce immediately the first equality in (48). To get the second equality, we note that $X_m^{(\alpha)}(0)$ enumerates configurations C with m pieces such that *the site 1 is not occupied by a piece*. Two situations may then occur: either site 2 is occupied or not. In the first case, the bijection $C \mapsto \tilde{C}$ will generate a configuration \tilde{C} where site 2 is occupied (and site 1 does not belong to Odd) while in the second case, it will generate a configuration where site 2 is empty and site 1 (which belongs to Odd) is necessarily occupied (since it was empty in C and the empty and occupied sites get exchanged in the bijection for those odd sites belonging to Odd). To summarize, in the configuration \tilde{C} , either site 1 or site 2 must be occupied. The restriction of $\tilde{X}_{\alpha-m}^{(\alpha)}$ to these configurations yields $\tilde{X}_{\alpha-m}^{(\alpha)} - \tilde{X}_{\alpha-m}^{(\alpha)}(0, 0)$, hence the second equality.

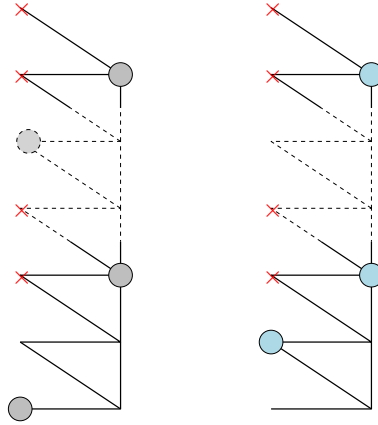


Figure 16. Left: an example of configuration C of hard pieces (represented in gray) on $\mathcal{G}^{(\alpha)}$. Right: the associated configuration \tilde{C} of hard pieces (represented in light blue), obtained by keeping the particles sitting on even vertices and, in the ensemble of odd vertices which are not adjacent to the occupied even vertices, exchanging the occupied and un-occupied sites.

From (48), we deduce

$$\begin{aligned}
 & X^{(\alpha)}\left(-\frac{Y_1}{z}, -\frac{Y_2}{z}, -\frac{Y_3}{z}, \dots, -\frac{Y_{2\alpha-1}}{z}\right) \\
 &= \left(-\frac{1}{z}\right)^\alpha X_\alpha^{(\alpha)} \cdot X^{(\alpha)}(-z\tilde{Y}_1, -z\tilde{Y}_2, -z\tilde{Y}_3, \dots, -z\tilde{Y}_{2\alpha-1}), \\
 & X^{(\alpha)}\left(0, -\frac{Y_2}{z}, -\frac{Y_3}{z}, \dots, -\frac{Y_{2\alpha-1}}{z}\right) \\
 &= \left(-\frac{1}{z}\right)^\alpha X_\alpha^{(\alpha)} \cdot (X^{(\alpha)}(-z\tilde{Y}_1, -z\tilde{Y}_2, -z\tilde{Y}_3, \dots, -z\tilde{Y}_{2\alpha-1}) \\
 &\quad - X^{(\alpha)}(0, 0, -z\tilde{Y}_3, \dots, -z\tilde{Y}_{2\alpha-1})),
 \end{aligned}$$

and therefore, taking the ratio of the two lines and using (47),

$$1 + \frac{Y_1}{z} J^{(\alpha)}\left(\frac{1}{z}\right) = K^{(\alpha)}\left(\frac{1}{z}\right) = 1 - \tilde{J}^{(\alpha)}(z).$$

The finite continued fraction case of Section 5.3 corresponds precisely to a situation where $J(z) = J^{(\alpha)}(z)$ and $\tilde{J}(z) = \tilde{J}^{(\alpha)}(z)$. The above formula explains the first identity in (26) while the second identity is guaranteed by the relation $\tilde{J}^{(\alpha)}(z) \rightarrow 1$ when $z \rightarrow 0$. This concludes the proof of (26).

We now prove (25) by computing explicitly the determinants $H_i^{(0)}$ and $H_i^{(1)}$ in terms of the Y_i 's. More precisely, let us show that, for $i \geq 1$,

$$H_i^{(0)} \equiv \begin{vmatrix} j_{-(i-1)} & \cdots & \cdots & j_0 \\ \vdots & \ddots & \ddots & j_1 \\ \vdots & \ddots & \ddots & \vdots \\ j_0 & j_1 & \cdots & j_{i-1} \end{vmatrix} = \left(\frac{Y_2}{Y_3}\right)^{i-1} \left(\frac{Y_4}{Y_5}\right)^{i-2} \cdots \left(\frac{Y_{2i-4}}{Y_{2i-3}}\right)^2 \left(\frac{Y_{2i-2}}{Y_{2i-1}}\right), \tag{49a}$$

$$H_i^{(1)} \equiv \begin{vmatrix} j_{-(i-2)} & \cdots & \cdots & j_1 \\ \vdots & \ddots & \ddots & j_2 \\ \vdots & \ddots & \ddots & \vdots \\ j_1 & j_2 & \cdots & j_i \end{vmatrix} = Y_1 Y_3 Y_5 \cdots Y_{2i-1} H_i^{(0)}. \tag{49b}$$

Once these formulas are proved, the relations (25) indeed follow immediately.

A first crucial point is the existence of a linear relation between the $j_n^{(\alpha)}$'s, namely

$$\sum_{m=0}^{\alpha} (-1)^m X_m^{(\alpha)} j_{n-m}^{(\alpha)} = 0 \text{ for all integers } n. \tag{50}$$

Indeed, writing the first identity in (47) as

$$\begin{aligned} K^{(\alpha)}(z) X^{(\alpha)}(-zY_1, -zY_2, -zY_3, \dots, -zY_{2\alpha-1}) \\ = X^{(\alpha)}(0, -zY_2, -zY_3, \dots, -zY_{2\alpha-1}), \end{aligned}$$

and extracting the term of order z^n , we immediately see that (50) holds for any positive integer $n \geq \alpha$ since $X^{(\alpha)}(0, -zY_2, -zY_3, \dots, -zY_{2\alpha-1})$ is a polynomial of degree $\alpha - 1$. Similarly, writing the second identity in (47) as

$$\tilde{J}^{(\alpha)}(z) \times X^{(\alpha)}(-z\tilde{Y}_1, -z\tilde{Y}_2, \dots, -z\tilde{Y}_{2\alpha-1}) = X^{(\alpha)}(0, 0, -z\tilde{Y}_3, \dots, -z\tilde{Y}_{2\alpha-1}),$$

a polynomial of degree $\alpha - 1$, we find that

$$\sum_{m=0}^{\alpha} (-1)^m \tilde{X}_m^{(\alpha)} j_{-(n'-m)}^{(\alpha)} = 0 = \frac{(-1)^\alpha}{X_\alpha^{(\alpha)}} \sum_{m'=0}^{\alpha} (-1)^{m'} X_{m'}^{(\alpha)} j_{(\alpha-n')-m'}^{(\alpha)}$$

for $n' \geq \alpha$. Here we have set $m' = \alpha - m$ and used $\tilde{X}_{\alpha-m'}^{(\alpha)} = X_{m'}^{(\alpha)} / X_\alpha^{(\alpha)}$. Setting $n = \alpha - n' \leq 0$, we deduce that (50) also holds for any non-positive integer.

It remains to show that it is valid in the range $1 \leq n \leq \alpha - 1$. For n in this range, we have

$$\begin{aligned} \sum_{m=0}^{\alpha} (-1)^m X_m^{(\alpha)} j_{n-m}^{(\alpha)} &= \sum_{m=0}^n (-1)^m X_m^{(\alpha)} j_{n-m}^{(\alpha)} + \sum_{m=n+1}^{\alpha} (-1)^m X_m^{(\alpha)} j_{n-m}^{(\alpha)} \\ &= (-1)^n X_n^{\alpha}(0) + (-1)^{\alpha} \sum_{p=0}^{\alpha-n-1} (-1)^p X_{\alpha-p}^{(\alpha)} j_{p-(\alpha-n)}^{(\alpha)} \\ &= (-1)^n X_n^{\alpha}(0) + (-1)^{\alpha} X_{\alpha}^{(\alpha)} \sum_{p=0}^{\alpha-n-1} (-1)^p \tilde{X}_p^{(\alpha)} j_{-((\alpha-n)-p)}^{(\alpha)}. \end{aligned}$$

Now since $1 \leq \alpha - n \leq \alpha - 1$, we also have

$$\sum_{p=0}^{\alpha-n} (-1)^p \tilde{X}_p^{(\alpha)} j_{-((\alpha-n)-p)}^{(\alpha)} = (-1)^{\alpha-n} \tilde{X}_{\alpha-n}^{(\alpha)}(0, 0)$$

so that we eventually get

$$\begin{aligned} \sum_{m=0}^{\alpha} (-1)^m X_m^{(\alpha)} j_{n-m}^{(\alpha)} &= (-1)^n X_n^{\alpha}(0) + (-1)^{\alpha} X_{\alpha}^{(\alpha)} ((-1)^{\alpha-n} \tilde{X}_{\alpha-n}^{(\alpha)}(0, 0) - (-1)^{\alpha-n} \tilde{X}_{\alpha-n}^{(\alpha)}) \\ &= (-1)^n (X_n^{\alpha}(0) - X_{\alpha}^{(\alpha)} (\tilde{X}_{\alpha-n}^{(\alpha)} - \tilde{X}_{\alpha-n}^{(\alpha)}(0, 0))) \\ &= 0. \end{aligned}$$

The linear relation (50) therefore holds for all integers n , as stated.

Let us now come to the computation of $H_i^{(0)}$ and $H_i^{(1)}$. Since j_n , $n \geq 1$, enumerates heaps of n pieces on the graph \mathcal{G} with base $\{1\}$, the pieces cannot reach sites with index more than 1 for $n = 1$ and $2n - 2$ for $n \geq 2$. In other words, j_n enumerates heaps which “live” on $\mathcal{G}^{(n)}$, therefore on $\mathcal{G}^{(i-1)}$ for all $n \leq i - 1$. As for j_n , $n \leq -1$, it enumerates heaps of $|n|$ pieces on the graph \mathcal{G} with base $\{1, 2\}$ so that the pieces cannot reach sites with index more than $2|n|$, therefore “live” on $\mathcal{G}^{(|n|+1)}$, therefore on $\mathcal{G}^{(i-1)}$ for all $|n| \leq i - 2$. In other words, we have

$$\begin{aligned} j_n &= j_n^{(i-1)} \quad \text{for } n = 0, 1, 2, \dots, i - 1, \\ j_{-n} &= j_{-n}^{(i-1)} \quad \text{for } n = 0, 1, 2, \dots, i - 2. \end{aligned}$$

In the determinant $H_i^{(0)}$, the only term which does not “live” on $\mathcal{G}^{(i-1)}$ is $j_{-(i-1)}$ and it is easily seen that

$$j_{-(i-1)} = j_{-(i-1)}^{(i-1)} + \tilde{Y}_2 \tilde{Y}_4 \dots \tilde{Y}_{2i-2}$$

with an additional term corresponding to the unique heap that hits position $2i - 2$.

Using the linear relation (50) for $\alpha = i - 1$, we may thus rewrite $H_i^{(0)}$ as

$$H_i^{(0)} = \begin{vmatrix} \tilde{Y}_2 \tilde{Y}_4 \dots \tilde{Y}_{2i-2} & j_{-(i-2)} & \cdots & j_0 \\ 0 & j_{-(i-3)} & \ddots & j_1 \\ \vdots & \ddots & \ddots & \vdots \\ 0 & j_1 & \cdots & j_{i-1} \end{vmatrix} = \tilde{Y}_2 \tilde{Y}_4 \dots \tilde{Y}_{2i-2} H_{i-1}^{(1)}.$$

Similarly, in the determinant $H_i^{(1)}$, the only term which does not “live” on $\mathcal{G}^{(i-1)}$ is j_i and it is easily seen that

$$j_i = j_i^{(i-1)} + Y_1(Y_2 Y_4 \dots Y_{2i-2})$$

with again an additional term corresponding to the unique heap that hits position $2i - 2$. We may thus rewrite $H_i^{(1)}$ as

$$H_i^{(1)} \equiv \begin{vmatrix} j_{-(i-2)} & \cdots & j_0 & 0 \\ \vdots & \ddots & \ddots & \vdots \\ j_0 & \ddots & \ddots & 0 \\ j_1 & \cdots & j_{i-1} & Y_1(Y_2 Y_4 \dots Y_{2i-2}) \end{vmatrix} = Y_1(Y_2 Y_4 \dots Y_{2i-2}) H_{i-1}^{(0)}.$$

Combining the two above formulas and replacing the \tilde{Y}_i 's by their value in terms of the Y_i 's, we deduce the recursion relation

$$H_i^{(0)} = \left(\frac{Y_2 Y_4}{Y_3 Y_5} \cdots \frac{Y_{2i-4}}{Y_{2i-3}} \right)^2 \frac{Y_{2i-2}}{Y_{2i-1}} H_{i-2}^{(0)}$$

for $i \geq 3$ with initial conditions $H_1^{(0)} = 1$ and

$$H_2^{(0)} = j_{-1} j_1 - j_0^2 = (\tilde{Y}_1 + \tilde{Y}_2) Y_1 - 1 = \frac{Y_2}{Y_3}.$$

The first line of eq (49) follows immediately. As for the second line, it follows from

$$\begin{aligned} \frac{H_i^{(1)}}{H_i^{(0)}} &= Y_1(Y_2 Y_4 \dots Y_{2i-2}) \frac{H_{i-1}^{(0)}}{H_i^{(0)}} \\ &= Y_1(Y_2 Y_4 \dots Y_{2i-2}) \left(\frac{Y_3 Y_5}{Y_2 Y_4} \cdots \frac{Y_{2i-1}}{Y_{2i-2}} \right) \\ &= Y_1 Y_3 Y_5 \dots Y_{2i-1}. \end{aligned}$$

The above derivation of eq. (25) extends verbatim to the case of the finite continued fraction of Section 5.3 by limiting to $i \leq \alpha - 1$ the range of allowed values for the index i in $H_i^{(0)}$ and $H_i^{(1)}$. This range is precisely what is needed to compute $Y_1, Y_2, \dots, Y_{2\alpha-1}$.

B. A proof of the formulas (35)–(38)

The quantities $Z(z; P, Q)$, $\tilde{Z}(z; P, Q)$ are specializations of $J(z)$ and $\tilde{J}(z)$ (viewed as defined from the Y_i 's through their continued fraction expansions) to the case where

$$Y_{2i-1} = Y, \quad Y_{2i} = P, \quad \tilde{Y}_{2i-1} = \frac{1}{Y} = \tilde{Y}, \quad \tilde{Y}_{2i} = \frac{P}{Y^2} = \tilde{P}$$

for all $i \geq 1$. Consequently, Z_n , \tilde{Z}_n , and k_n are the corresponding specializations of J_n , \tilde{J}_n and j_n . The analysis of Appendix A applies to arbitrary Y_i 's. In particular, k_n has a direct interpretation in terms of heaps of pieces on the graph \mathcal{G} for all n . For $n > 0$, k_n enumerates heaps of n pieces of base $\{1\}$, with weights Y and P for pieces on odd or even sites respectively. For $n < 0$, k_n enumerates heaps of $|n|$ pieces of base $\{1, 2\}$, with weights \tilde{Y} and \tilde{P} for pieces on odd or even sites respectively. Let us thus introduce the quantities $k_n^{(\alpha)}$ (analogs of $j_n^{(\alpha)}$) corresponding to a restriction of the heaps to the graph $\mathcal{G}^{(\alpha)}$. From (50), we deduce immediately

$$\sum_{m=0}^{\alpha} (-1)^m x_m^{(\alpha)} k_{n-m}^{(\alpha)} = 0 \quad \text{for all integers } n,$$

where $x_m^{(\alpha)} = X_m^{(\alpha)}(Y, P, Y, P, \dots, Y)$ is the generating function of configurations of exactly m hard pieces on the graph $\mathcal{G}^{(\alpha)}$. Setting $\alpha = i - 1$, this equation reads equivalently

$$\sum_{m=0}^{i-1} k_{n-i+m}^{(i-1)} (-1)^m \frac{x_{i-1-m}^{(i-1)}}{x_{i-1}^{(i-1)}} = 0 \quad \text{for all integers } n.$$

Now, from their heap interpretation, it is clear that $k_n = k_n^{(i-1)}$ for $n = 0, 1, \dots, i - 1$ as well as for $n = -1, -2, \dots, -(i - 2)$. For $2 \leq n \leq i$, all the $k_p^{(i-1)}$'s appearing in the above formula may thus be replaced by k_p 's and (35) follows. For $n = 1$, the only term which gets out of the graph $\mathcal{G}^{(i-1)}$ is for $m = 0$, since $k_{-(i-1)}^{(i-1)} = k_{-(i-1)} - \tilde{P}^{i-1}$ (the two indeed differ by the contribution of the heap made of one piece on each even site from 2 to $2(i - 1)$). This explains the right hand side $\tilde{P}^{i-1} = P^{i-1}/Y^{2(i-1)}$ in the first line of (36). For $n = i + 1$, the only term which gets out of the graph $\mathcal{G}^{(i-1)}$ is for $m = i - 1$, since $k_i^{(i-1)} = k_i - YP^{i-1}$ (the two indeed differ by the contribution of the heap made of one piece on site 1 and one piece on each even site from 2 to $2(i - 1)$). This explains the right hand side $(-1)^{i-1} YP^{i-1} (x_0^{(i-1)}/x_{i-1}^{(i-1)}) = (-1)^{i-1} P^{i-1}/Y^{i-2}$ in the second line of (36).

Eqs. (37) and (38) can be proved in the same way but their proof now relies on a restriction of the heaps to the graph $\mathcal{G}'^{(\alpha)}$ of Figure 17.

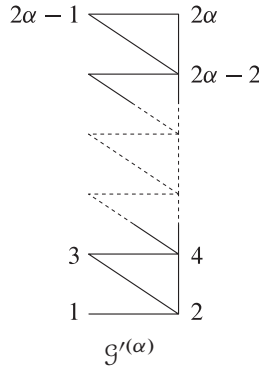


Figure 17. The graph $\mathcal{G}'^{(\alpha)}$.

Let us first analyze the heap generating functions on this graph: denoting, for $n > 0$, the generating function $j_n^{(\alpha)}$ of heaps of n pieces on $\mathcal{G}'^{(\alpha)}$, of base $\{1\}$ and with weight Y_i per piece sitting on site i . It is easily seen that $j_n^{(\alpha)}$ also corresponds to enumerating heaps of n pieces on the graph $\mathcal{G}^{(\alpha)}$ provided we assign, instead of $Y_{2\alpha-1}$, a weight $Y'_{2\alpha-1} \equiv Y_{2\alpha-1} + Y_{2\alpha}$ to pieces sitting at position $2\alpha - 1$. The same remark holds for configurations of m hard pieces on $\mathcal{G}'^{(\alpha)}$ which are enumerated by $X_m^{(\alpha)}(Y_1, Y_2, \dots, Y_{2\alpha-2}, Y'_{2\alpha-1})$. In order to use directly our previous results (obtained for $\mathcal{G}^{(\alpha)}$), we are thus led to define $j_n^{(\alpha)}$ for $n < 0$ as enumerating heaps of n pieces with base $\{1, 2\}$ on the graph $\mathcal{G}^{(\alpha)}$, with weights \tilde{Y}_i built via the same expression (21) as before with $Y_{2\alpha-1}$ replaced by $Y'_{2\alpha-1}$. Getting back to $\mathcal{G}'^{(\alpha)}$, $j_n^{(\alpha)}$ for $n < 0$ therefore enumerates heaps of n pieces with base $\{1, 2\}$ on this graph, with weights \tilde{Y}_i as in (21) for pieces on sites $1, 2, 3, \dots, 2\alpha - 3$ and the special weights $Y_{2\alpha-2}/(Y_{2\alpha-3}(Y_{2\alpha-1} + Y_{2\alpha}))$ for pieces on the site $2\alpha - 2$, $1/(Y_{2\alpha-1} + Y_{2\alpha})$ for pieces on the site $2\alpha - 1$ and 0 for pieces on the site 2α . With this definition, we have the analog of (50), namely

$$\sum_{m=0}^{\alpha} (-1)^m X_m^{(\alpha)} j_{n-m}^{(\alpha)} = 0 \quad \text{for all integers } n.$$

where $X_m^{(\alpha)}$ enumerates configurations of m hard pieces on $\mathcal{G}'^{(\alpha)}$.

Let us now specialize this result upon introducing, for $n > 0$, the generating function $k_n^{(\alpha)}$ of heaps of n pieces on $\mathcal{G}^{(\alpha)}$, of base $\{1\}$ and with weights Y and P for pieces on odd or even sites respectively. From the above discussion, for $n < 0$, $k_{-n}^{(\alpha)}$ must be defined as enumerating heaps of n pieces on $\mathcal{G}^{(\alpha)}$ with weight \tilde{Y} for pieces on odd sites $1, 3, 5, \dots, 2\alpha - 3$, \tilde{P} for pieces on even sites $2, 4, \dots, 2\alpha - 4$ and the special weights $P/(Y(Y + P))$ for pieces on the site $2\alpha - 2$, $1/(Y + P)$ for pieces on the site $2\alpha - 1$ and 0 for pieces on the site 2α . With these definitions (and $k_0^{(\alpha)} \equiv 1$), we have

$$\sum_{m=0}^{\alpha} (-1)^m x_m^{(\alpha)} k_{n-m}^{(\alpha)} = 0 \quad \text{for all integers } n$$

where $x_m^{(\alpha)}$ enumerates configurations of m hard pieces on $\mathcal{G}^{(\alpha)}$ with weight Y (resp. P) per piece sitting on an odd (resp. even) site (in particular $x_0^{(\alpha)} = 1$). Setting $\alpha = i - 1$, the above equation becomes

$$\sum_{m=0}^{i-1} (-1)^m x_m^{(i-1)} k_{n-m}^{(i-1)} = 0 \quad \text{for all integers } n.$$

Again, from their heap interpretation, it is clear that $k_n = k_n^{(i-1)}$ for $n = 0, 1, \dots, i$ as well as for $n = -1, -2, \dots, -(i - 3)$. For $2 \leq n \leq i$, all the $k_p^{(i-1)}$'s appearing in the above formula may thus be replaced by k_p 's and (37) follows. For $n = 1$, the only term for which this substitution fails is for $m = i - 1$, since $k_{-(i-2)}^{(i-1)} = k_{-(i-2)} - \tilde{P}^{i-3}(\tilde{P} - P/(Y(Y + P)))$ (the two indeed differ by the contribution of the last piece, at position $2(i - 2)$, in the heap made of one piece on each even site from 2 to $2(i - 2)$). This explains the right hand side $(-1)^{i-1} \tilde{P}^{i-3}(\tilde{P} - P/(Y(Y + P)))x_{i-1}^{(i-1)} = (-1)^{i-1} P^{i-1}/Y^{i-2}$ in the first line of (38) (note that $x_{i-1}^{(i-1)} = Y^{i-2}(Y + P)$). For $n = i + 1$, the only term for which this substitution fails is for $m = 0$, since $k_{i+1}^{(i-1)} = k_{i+1} - YP^{i-1}(Y + P)$ (the two indeed differ by the contribution of the two heaps made of one piece on site 1, one piece on each even site from 2 to $2(i - 1)$ and a last piece at position $2i - 1$ or $2i$). This explains the right hand side $YP^{i-1}(Y + P)$ in the second line of (38).

Acknowledgements. We warmly thank P. Di Francesco and R. Kedem for useful discussions. The work of É. Fusy was partly supported by the ANR grant ‘‘Carta-plus’’ 12-JS02-001-01 and the ANR grant ‘‘EGOS’’ 12-JS02-002-01.

References

- [1] J. Ambjørn and T. G. Budd, Trees and spatial topology change in causal dynamical triangulations. *J. Phys. A* **46** (2013), no. 31, article id. 315201, 33 pp. [MR 3090757](#) [Zbl 1273.81224](#)
- [2] J. Bettinelli, A bijection for nonorientable general maps. Preprint 2015. [arXiv:1512.02208](#) [math.CO]
- [3] J. Bettinelli and G. Miermont, Compact Brownian surfaces I. Brownian disks. *Probab. Theory Related Fields* **167** (2017), no. 3-4, 555–614. [MR 3627425](#)
- [4] J. Bouttier, P. Di Francesco, and E. Guitter, Geodesic distance in planar graphs. *Nuclear Phys. B* **663** (2003), no. 3, 535–567. [MR 1987861](#) [Zbl 1022.05022](#)
- [5] J. Bouttier, P. Di Francesco, and E. Guitter, Planar maps as labeled mobiles. *Electron. J. Combin.* **11** (2004), no. 1, Research Paper 69, 27 pp. [MR 2097335](#) [Zbl 1060.05045](#)
- [6] J. Bouttier, É. Fusy, and E. Guitter, On the two-point function of general planar maps and hypermaps. *Ann. Inst. Henri Poincaré D* **1** (2014), no. 3, 265–306. [MR 3239273](#) [Zbl 1301.05079](#)
- [7] J. Bouttier and E. Guitter, Planar maps and continued fractions. *Comm. Math. Phys.* **309** (2012), no. 3, 623–662. [MR 2885603](#) [Zbl 1235.05067](#)
- [8] G. Chapuy, Asymptotic enumeration of constellations and related families of maps on orientable surfaces. *Combin. Probab. Comput.* **18** (2009), no. 4, 477–516. [MR 2507734](#) [Zbl 1221.05204](#)
- [9] G. Chapuy and M. Dołęga, A bijection for rooted maps on general surfaces. *J. Combin. Theory Ser. A* **145** (2017), 252–307. [MR 3551653](#) [Zbl 1355.05094](#)
- [10] G. Chapuy, M. Marcus, and G. Schaeffer, A bijection for rooted maps on orientable surfaces. *SIAM J. Discrete Math.* **23** (2009), no. 3, 1587–1611. [MR 2563085](#) [Zbl 1207.05087](#)
- [11] R. Cori and B. Vauquelin, Planar maps are well labeled trees. *Canad. J. Math.* **33** (1981), no. 5, 1023–1042. [MR 0638363](#) [Zbl 0415.05020](#)
- [12] P. Di Francesco and R. Kedem, Q -systems, heaps, paths and cluster positivity. *Comm. Math. Phys.* **293** (2010), no. 3, 727–802. [MR 2566162](#) [Zbl 1194.05165](#)
- [13] P. Di Francesco and R. Kedem, Non-commutative integrability, paths and quasi-determinants. *Adv. Math.* **228** (2011), no. 1, 97–152. [MR 2822229](#) [Zbl 1252.37069](#)
- [14] P. Flajolet, Combinatorial aspects of continued fractions. *Discrete Math.* **32** (1980), no. 2, 125–161. [MR 0592851](#) [Zbl 0445.05014](#)
- [15] É. Fusy and E. Guitter, The two-point function of bicolored planar maps. *Ann. Inst. Henri Poincaré D* **2** (2015), no. 4, 335–412. [MR 3428405](#) [Zbl 1327.05077](#)
- [16] J.-F. Le Gall, Uniqueness and universality of the Brownian map. *Ann. Probab.* **41** (2013), no. 4, 2880–2960. [MR 3112934](#) [Zbl 1282.60014](#)

- [17] G. Miermont, Tessellations of random maps of arbitrary genus. *Ann. Sci. Éc. Norm. Supér. (4)* **42** (2009), no. 5, 725–781. [MR 2571957](#) [Zbl 1228.05118](#)
- [18] G. Miermont, The Brownian map is the scaling limit of uniform random plane quadrangulations. *Acta Math.* **210** (2013), no. 2, 319–401. [MR 3070569](#) [Zbl 1278.60124](#)
- [19] J. Miller and S. Sheffield, An axiomatic characterization of the Brownian map. Preprint 2015. [arXiv:1506.03806](#) [math.PR]
- [20] G. Schaeffer, *Conjugaison d'arbres et cartes combinatoires aléatoires*. Ph.D. thesis. Université Bordeaux I, Bordeaux, 1998.
- [21] G. X. Viennot, Heaps of pieces. I. Basic definitions and combinatorial lemmas. In M. F. Capobianco, M. G. Guan, D. F. Hsu and F. Tian. *Graph theory and its applications: East and West*. Proceedings of the First China–USA International Conference held in Jinan, June 9–20, 1986. Annals of the New York Academy of Sciences, 576. New York Academy of Sciences, New York, 1989, 542–570. [MR 1110852](#) [Zbl 0792.05012](#)

© European Mathematical Society

Communicated by Adrian Tanasă

Received July 20, 2015; accepted March 8, 2016

Éric Fusy, LIX, École Polytechnique, 91120 Palaiseau, France

ACADEMIC YEARS 2014-2016: PIMS-CNRS, University of British Columbia, Vancouver, BC, Canada

e-mail: fusy@lix.polytechnique.fr

Emmanuel Guitter, Institut de Physique Théorique, Université Paris Saclay, CEA, CNRS, 91191 Gif-sur-Yvette, France

e-mail: emmanuel.gutter@cea.fr