



Partial Differential Equations — *An obstacle problem with gradient term and asymptotically linear reaction*, by B. ABDELLAOUI*, S. M. BOUGUIMA and I. PERAL, communicated on 10 December 2010.

In memoria di Giovanni Prodi matematico e gentiluomo.

ABSTRACT. — We will consider the following obstacle problem

$$\int_{\Omega} \nabla u \nabla T_k(v - u) \, dx + \int_{\Omega} h(u) |\nabla u|^q T_k(v - u) \, dx \geq \int_{\Omega} (g(x, u) + f) T_k(v - u) \, dx,$$

with the condition that $u \geq \psi$ a.e in Ω . Under suitable condition relating g , h and q , we show the existence of a solution for all $f \in L^1(\Omega)$.

The main feature is, assuming that $g(x, s)$ is asymptotically linear as $|s| \rightarrow \pm\infty$ and independently of the values of

$$\lim_{s \rightarrow \pm\infty} \frac{g(x, s)}{s},$$

to obtain a solution for all $\lambda > 0$ and $f \in L^1(\Omega)$. In this sense we could say that the first order term break down any resonant effect.

KEY WORDS: Nonlinear obstacle problems, existence and nonexistence, regularization, resonance.

2000 MATHEMATICS SUBJECT CLASSIFICATION: 35D05, 35D10, 35J20, 35J25, 35J70, 46E30, 46E35.

1. INTRODUCTION

In this paper we deal with a nonlinear elliptic obstacle problem of the form

$$(1.1) \quad \begin{cases} u \geq \psi \text{ a.e in } \Omega, & \text{for all } v \in \mathcal{K}(\psi), \text{ we have} \\ \int_{\Omega} \nabla u \nabla(v - u) \, dx + \int_{\Omega} h(u) |\nabla u|^q (v - u) \, dx \geq \int_{\Omega} (g(x, u) + f)(v - u) \, dx, \end{cases}$$

where $\Omega \subset \mathbb{R}^N$ is an open bounded domain, $1 < q \leq 2$, ψ is a bounded function such that $\psi \in W_0^{1,2}(\Omega)$ and

$$\mathcal{K}(\psi) = \{v \in W_0^{1,2} \cap L^\infty(\Omega) : v \geq \psi \text{ in } \Omega\}.$$

Work partially supported by project MTM2007-65018, MICINN, Spain and projet A/0016546/08, PCI, Spain.

*First author also is partially supported by a grant from the ICTP of Trieste, Italy.

We suppose that $f \in L^1(\Omega)$, g is a Caratheodory function asymptotically linear, that is, verifying

$$(1.2) \quad |g(x, s)| \leq \lambda g_0(x)|s| + g_1(x)$$

where $g_1 \in L^1(\Omega)$ and g_0 satisfies

$$(1.3) \quad \begin{cases} g_0 \not\equiv 0, \\ g_0 \in L^1(\Omega), \\ C(g_0, q) > 0, \text{ where } C(g_0, q) = \inf_{\phi \in W_0^{1,q}(\Omega) \setminus \{0\}} \frac{(\int_{\Omega} |\nabla \phi|^q dx)^{1/q}}{\int_{\Omega} g_0 |\phi| dx}. \end{cases}$$

It is easy to check that if g_0 satisfies (1.3), then $g_0 \in W^{-1,q'}(\Omega) \cap L^1(\Omega)$, $q' = \frac{q}{q-1}$.

We say that g_0 is an *admissible weight* if (1.3) holds.

If $\psi \equiv 0$ and we consider the equation

$$(1.4) \quad -\Delta u = g(u) + f, \quad \text{in } \Omega, \quad u \geq 0 \text{ and } u \in W_0^{1,2}(\Omega),$$

where g is a lipschitz function such that $g(0) = 0$ and verifying the condition

$$(1.5) \quad \lim_{s \rightarrow \pm\infty} \frac{g(s)}{s} = \lambda_{\pm},$$

for $\lambda_- < \lambda_1 < \lambda_+ < \lambda_2$, λ_1 and λ_2 are the first and second eigenvalue of the Laplacian. The problem (1.4) was solved in the famous work by Ambrosetti-Prodi [2]. The authors establish a sharp existence, nonexistence and multiplicity result related to the value of the projection of the datum $f \in L^2(\Omega)$ on the first positive normalized eigenfunction of the Laplacian,

$$\int_{\Omega} f(x) \phi_1(x) dx = t.$$

More precisely they prove that there exists a threshold \bar{t} such that, if $t > \bar{t}$ there is no solution, if $t = \bar{t}$ there exist a solution and if $t < \bar{t}$ there exist two solutions.

One of the goals of this paper is to prove that under some hypotheses on g , for all $f \in L^1(\Omega)$, g satisfying (1.2) (1.3) and h with some structural conditions, there exists a solution to the variational inequality (1.1). In particular in the Ambrosetti-Prodi context we prove that the gradient term give a solution without any condition on λ_{\pm} or the projection of f on ϕ_1 .

As a precedent we have the case of an equation with gradient term. It was proved in [1], for the case $g(x, u) = \lambda g_0(x)u$, under a suitable condition on q and g_0 , $h(u) \equiv 1$, that the absorption term $|\nabla u|^q$ is sufficient to break down any resonant effect of the linear zero order term and then the existence of a solution is

obtained for all $\lambda > 0$ and $f \in L^1(\Omega)$. In this sense this paper could be understood, in particular, as the extension of the result in [1] to variational inequalities with g verifying (1.2) (1.3) and h verifying (3.2) below.

Unilateral problems with gradient term has been largely studied in the literature, we refer, for instance, to [4], [8], [14] and the references therein. In [4] it is studied the existence of unbounded solutions for an obstacle problem with natural growth in the gradient.

To prove the existence of solutions for unilateral problems with L^1 datum, it is necessary to consider entropy solution in the sense that $u \geq \psi$ and

$$\int_{\Omega} \nabla u \nabla (T_k(v - u)) \, dx \geq \int_{\Omega} f(T_k(v - u)) \, dx$$

for all $v \in \mathcal{K}(\psi)$. See for instance [7].

We organize the contents as follows.

In Section 2 we consider a simple model where $\psi \geq 0$, $f \geq 0$, $h \equiv 1$ and $g(x, s) \equiv \lambda g_0(x)u$, with $g_0 \geq 0$. Then for all $\lambda > 0$, we prove the existence of a nonnegative solution. More precisely we show that if g_0 is a nonnegative admissible weight in the sense of condition (1.3), then we have a solution for all $\lambda > 0$ and all $f \in L^1(\Omega)$.

To prove the main result we use a convenient approximate problems and uniform estimates in order to pass to the limit. In Subsection 2.1 a partial uniqueness result is given for $q = 2$ and $\psi \equiv 0$.

Section 3, is devoted to obstacle problem (1.1) without any sign condition on f and ψ . The term $|\nabla u|^q$ will be substituted by the more general $h(u)|\nabla u|^q$ and we will consider the general nonlinearity $g(x, u)$ satisfying (1.5). Under suitable conditions on h we will prove the existence of entropy solution for all $f \in L^1$ and without any restriction on λ_{\pm} . In this sense the result can be seen as *breaking of resonance* for the Ambrosetti-Prodi obstacle problem.

It is worthy to point out that in the problem without constraint, condition (1.3) is optimal. It is sufficient to consider $g(x) = |x|^{-2}$, the Hardy potential, for which we have the classical inequality

$$\int_{\Omega} |\nabla u|^2 \, dx \geq \Lambda_N \int_{\Omega} \frac{u^2}{|x|^2} \, dx, \quad \text{for all } u \in \mathcal{C}_0^{\infty}(\Omega) \text{ where } \Lambda_N = \left(\frac{N-2}{2}\right)^2.$$

In this case condition (1.3) holds if and only if $q > \frac{N}{N-1}$. Then if $q < \frac{N}{N-1}$ and $\lambda > \Lambda_N = \left(\frac{N-2}{2}\right)^2$, there is no solution to the obstacle problem. (See Theorem 3.1 in [1] for details).

We will use the following notation. For a measurable function u we define the truncation $T_k(u)$ by

$$T_k(u) = \max\{-k, \min\{k, u\}\}.$$

We set $G_k(u) = u - T_k(u)$.

2. EXISTENCE OF NONNEGATIVE SOLUTIONS TO THE SIMPLE MODEL

In this section we deal with the simple case where $\psi \geq 0$, $f \geq 0$, $h \equiv 1$ and $g(x, s) \equiv \lambda g_0(x)u$, with $g_0 \geq 0$. Define the convex set

$$\mathcal{K}(\psi) = \{v \in W_0^{1,2} \cap L^\infty(\Omega) : v \geq \psi \text{ in } \Omega\}.$$

We find the following result.

THEOREM 2.1. *Assume that g_0 is an admissible weight in the sense of condition (1.3), then for all $\lambda > 0$ and for all $f \in L^1(\Omega)$, there exists a positive $u \geq \psi$ such that $|\nabla u|^q \in L^1(\Omega)$, $T_k(u) \in W_0^{1,2}(\Omega)$ for all $k > 0$ and for all $v \in \mathcal{K}(\psi)$ we have*

$$(2.1) \quad \begin{aligned} & \int_{\Omega} \nabla u \nabla (T_k(v - u)) \, dx + \int_{\Omega} |\nabla u|^q (T_k(v - u)) \, dx \\ & \geq \int_{\Omega} (\lambda g u + f)(T_k(v - u)) \, dx. \end{aligned}$$

We will say that u is an entropy solution to the obstacle problem if (2.1) holds.

To prove Theorem 2.1 we start by proving the result in some particular cases and then we proceed by approximation of g and f . Notice that since $1 < q \leq 2$, then $\frac{N}{2} \leq \frac{N}{q}$.

THEOREM 2.2. *Assume that $f, g \in L^r(\Omega)$ are positive functions with $r > \frac{N}{q}$, then for all $\lambda \geq 0$, there exists $u \in \mathcal{K}(\psi)$ a weak positive solution to problem*

$$(2.2) \quad \begin{cases} \int_{\Omega} \nabla u \nabla (v - u) \, dx + \int_{\Omega} |\nabla u|^q (v - u) \, dx \\ \geq \int_{\Omega} (\lambda g u + f)(v - u) \, dx \quad \text{for all } v \in \mathcal{K}(\psi), \end{cases}$$

PROOF. We divide the proof in several steps.

Step 1: Let $k > 0$ be fixed, then for all $n \in \mathbb{N}$, using classical results (see for instance [13] and [12]), there exists $w_n \in \mathcal{K}(\psi)$, a solution to the obstacle problem

$$(2.3) \quad \begin{aligned} & \int_{\Omega} \nabla w_n \nabla (v - w_n) \, dx + \int_{\Omega} \frac{|\nabla w_n|^q}{1 + \frac{1}{n} |\nabla w_n|^q} (v - w_n) \, dx \\ & \geq \int_{\Omega} (\lambda g T_k(w_n) + f)(v - w_n) \, dx \end{aligned}$$

for all $v \in \mathcal{K}(\psi)$.

For k fixed we pass to the limit in n . Let $v = T_m(w_n)$, since $\psi \in L^\infty(\Omega)$, then choosing m very large we conclude that v is an admissible test function in (2.3). Since $v - w_n = -G_m(w_n)$, it follows that

$$\int_{\Omega} |\nabla G_m(w_n)|^2 dx + \int_{\Omega} \frac{|\nabla G_m(w_n)|^q}{1 + \frac{1}{n}|G_m(w_n)|^q} G_m(w_n) \leq \int_{\Omega} (\lambda g T_k(w_n) + f) G_m(w_n) dx.$$

Thus

$$\int_{\Omega} |\nabla G_m(w_n)|^2 dx \leq \lambda k^2 \|g\|_1 + \int_{\Omega} f G_m(w_n) dx.$$

Using Poincaré inequality we get that $\int_{\Omega} |\nabla G_m(w_n)|^2 dx \leq C$ for all m .

Notice that choosing $m \gg k$ it follows that

$$(2.4) \quad \int_{\Omega} |\nabla G_m(w_n)|^2 dx \leq \int_{\Omega} f G_m(w_n) dx,$$

and then by using the classical Stampacchia estimates, see [15], we obtain that $\|w_n\|_{L^\infty} \leq C$ where C is independent of n .

We set now $v = \psi$, then

$$\begin{aligned} \int_{\Omega} \nabla w_n \nabla (w_n - \psi) dx + \int_{\Omega} \frac{|\nabla w_n|^q}{1 + \frac{1}{n}|\nabla w_n|^q} (w_n - \psi) dx \\ \leq \int_{\Omega} (\lambda g T_k(w_n) + f)(w_n - \psi) dx. \end{aligned}$$

Since $w_n \geq \psi$, we get

$$\int_{\Omega} |\nabla w_n|^2 dx \leq \int_{\Omega} \nabla w_n \nabla \psi dx + \int_{\Omega} (\lambda g T_k(w_n) + f)(w_n - \psi) dx.$$

Thus using Hölder inequality and the previous estimate we obtain that

$$\int_{\Omega} |\nabla w_n|^2 dx \leq C(f, g, \Omega, k) \quad \text{uniformly in } n,$$

therefore, up to a subsequence, $w_n \rightharpoonup u_k$ weakly in $W_0^{1,2}(\Omega)$. By weak-*convergence in $L^\infty(\Omega)$ we also have that $u_k \in W_0^{1,2} \cap L^\infty(\Omega)$ and $u_k \geq \psi$. Next we investigate the inequality satisfied by u_k . To do that we prove the following claim.

CONVERGENCE CLAIM. $w_n \rightarrow u_k$ strongly in $W_0^{1,2}(\Omega)$.

PROOF OF THE CONVERGENCE CLAIM. It is clear that for all $v \in \mathcal{H}(\psi)$,

$$(\lambda g T_k(w_n) + f)(w_n - v) \rightarrow (\lambda g T_k(u_k) + f)(u_k - v) \quad \text{strongly in } L^1(\Omega).$$

Let $v = w_n - (w_n - u_k)^+$, then $v \in \mathcal{H}(\psi)$ and $v - w_n = -(w_n - u_k)^+$, so we have

$$\begin{aligned} & \int_{\Omega} \nabla w_n \nabla (w_n - u_k)^+ dx + \int_{\Omega} \frac{|\nabla w_n|^q}{1 + \frac{1}{n} |\nabla w_n|^q} (w_n - u_k)^+ dx \\ & \leq \int_{\Omega} (\lambda g T_k(w_n) + f)(w_n - u_k) dx. \end{aligned}$$

It is clear that $\int_{\Omega} (\lambda g T_k(w_n) + f)(w_n - u_k) dx \rightarrow 0$ as $n \rightarrow \infty$. Hence we conclude that

$$\begin{aligned} & \int_{\Omega} |\nabla (w_n - u_k)^+|^2 dx + \int_{\Omega} \frac{|\nabla w_n|^q}{1 + \frac{1}{n} |\nabla w_n|^q} (w_n - u_k)^+ dx \\ & \leq - \int_{\Omega} \nabla u_k \nabla (w_n - u_k)^+ dx + o(1) = o(1). \end{aligned}$$

Thus $\int_{\Omega} |\nabla (w_n - u_k)^+|^2 dx = o(1)$ and then $(w_n - u_k)^+ \rightarrow 0$ strongly in $W_0^{1,2}(\Omega)$.

To complete the proof we follow closely the argument used in [6]. Consider $\phi(s) = se^{(1/4)s^2}$, which satisfies $\phi'(s) - |\phi(s)| \geq \frac{1}{2}$.

Let $v = w_n + \phi((w_n - u_k)^-)$, then $v \in \mathcal{H}(\psi)$ and $v - w_n = \phi((w_n - u_k)^-)$. It is clear that

$$\nabla(v - w_n) = \begin{cases} 0 & \text{if } w_n \geq u_k, \\ \phi'((w_n - u_k)^-)(\nabla u_k - \nabla w_n) & \text{if } w_n \leq u_k. \end{cases}$$

Using v as a test function in (2.3) we obtain that

$$\begin{aligned} (2.5) \quad & \int_{\Omega} \nabla w_n \phi'((w_n - u_k)^-) \nabla (w_n - u_k)^- dx + \int_{\Omega} H_n(\nabla w_n) \phi(w_n - u_k)^- dx \\ & \leq \lambda \int_{\Omega} g(x) T_k w_n \phi((w_n - u_k)^-) dx + \int_{\Omega} f(x) \phi((w_n - u_k)^-) dx. \end{aligned}$$

where $H_n(s) = \frac{|s|^q}{1 + \frac{1}{n}|s|^q}$. Therefore

$$\begin{aligned} (2.6) \quad & \int_{w_n \leq u_k} \phi'((w_n - u_k)^-) \nabla w_n \nabla (w_n - u_k)^- dx - \int_{\Omega} H_n(\nabla w_n) |\phi(w_n - u_k)^-| dx \\ & \leq \lambda \int_{w_n \leq u_k} g(x) T_k w_n \phi(u_k - w_n) dx + \int_{w_n \leq u_k} f(x) \phi(u_k - w_n) dx. \end{aligned}$$

Since $w_n \rightharpoonup u_k$ weakly in $W_0^{1,2}(\Omega)$, a direct computation shows that

$$\begin{aligned} & \int_{\Omega} \nabla w_n \phi'((w_n - u_k)^-) \nabla (w_n - u_k)^- dx \\ &= \int_{\Omega} |\nabla((w_n - u_k)^-)|^2 \phi'((w_n - u_k)^-) dx + o(1). \end{aligned}$$

As $q \leq 2$, $\forall \varepsilon > 0$ there exists a non negative constant C_ε such that

$$(2.7) \quad s^q \leq \varepsilon s^2 + C_\varepsilon, \quad s \geq 0.$$

Hence the second term in the left-hand side can be estimated in the following way,

$$\begin{aligned} & \int_{\Omega} H_n(\nabla w_n) \phi((w_n - u_k)^-) dx \\ & \leq \varepsilon \int_{\Omega} |\nabla w_n|^2 |\phi((w_n - u_k)^-)| dx + C(\varepsilon) \int_{\Omega} |\phi((w_n - u_k)^-)| dx \\ & = \varepsilon \int_{\Omega} |\nabla((w_n - u_k)^-)|^2 |\phi((w_n - u_k)^-)| dx - \varepsilon \int_{\Omega} |\nabla u_k|^2 |\phi((w_n - u_k)^-)| dx \\ & \quad + 2\varepsilon \int_{\Omega} \nabla w_n \nabla u_k |\phi((w_n - u_k)^-)| dx + C(\varepsilon) \int_{\Omega} |\phi((w_n - u_k)^-)| dx. \end{aligned}$$

Since $w_n \rightharpoonup u_k$ weakly in $W_0^{1,2}(\Omega)$ and $|\phi((w_n - u_k)^-)| \rightarrow 0$ almost everywhere and in $L^2(\Omega)$, it follows that,

- (i) $\int_{\Omega} |\nabla u_k|^2 |\phi((w_n - u_k)^-)| dx \rightarrow 0$ as $n \rightarrow \infty$,
- (ii) $\int_{\Omega} \nabla w_n \nabla u_k |\phi((w_n - u_k)^-)| dx \rightarrow 0$ as $n \rightarrow \infty$.

Therefore, passing to the limit as n tends to ∞ , we have

$$\int_{\Omega} H_n(\nabla w_n) \phi((w_n - u_k)^-) dx \leq \varepsilon \int_{\Omega} |\nabla w_n - \nabla u_k|^2 |\phi((w_n - u_k)^-)| dx + o(1).$$

Moreover, it is clear that the right-hand side in (2) goes to zero as $n \rightarrow \infty$. Since $\phi'(s) - |\phi(s)| > \frac{1}{2}$, choosing $\varepsilon \leq 1$ we conclude that

$$\begin{aligned} & \frac{1}{2} \int_{\Omega} |\nabla((w_n - u_k)^-)|^2 dx \\ & \leq \int_{\Omega} (\phi'((w_n - u_k)^-) - \varepsilon |\phi((w_n - u_k)^-)|) |\nabla((w_n - u_k)^-)|^2 dx \\ & \leq o(1), \end{aligned}$$

whence $w_n \rightarrow u_k$ in $W_0^{1,2}(\Omega)$ and the claim is proved. Moreover, from (2.7) it follows that

$$H_n(\nabla w_n) \leq c_1 |\nabla w_n|^2 + c_2.$$

By the claim, we have in particular the almost everywhere convergence of the gradients and therefore we conclude that

$$H_n(\nabla w_n) \rightarrow |\nabla u_k|^q \quad \text{in } L^1(\Omega).$$

Hence we find that $u_k \in \mathcal{H}(\psi)$ solves

$$(2.8) \quad \begin{aligned} \int_{\Omega} \nabla u_k \nabla (v - u_k) dx + \int_{\Omega} |\nabla u_k|^q (v - u_k) dx \\ \geq \int_{\Omega} (\lambda g T_k(u_k) + f)(v - u_k) dx \end{aligned}$$

for all $v \in \mathcal{H}(\psi)$.

Step 2: We claim the existence of a universal $M > 0$ that does not depend on k such that $\|u_k\|_{L^\infty(\Omega)} \leq M$. To prove the claim we use the fact that $f, g_0 \in L^r(\Omega)$ where $r > \frac{N}{2}$. Let $v = T_m(u_k)$, using v as a test function in (2.8) it follows that

$$\int_{\Omega} |\nabla G_m(u_k)|^2 dx + \int_{\Omega} |\nabla G_m(u_k)|^q G_m(u_k) \leq \lambda \int_{\Omega} g_0 G_m^2(u_k) dx + \int_{\Omega} f G_m(u_k) dx.$$

Notice that, using Poincaré inequality we get

$$\int_{\Omega} |\nabla G_m(u_k)|^q G_m(u_k) = \frac{1}{(1 + 1/q)^q} \int_{\Omega} |\nabla G_m^{1+1/q}(u_k)|^q dx \geq C \int_{\Omega} g_0 G_m^{1+q}(u_k) dx,$$

and

$$\begin{aligned} \lambda \int_{\Omega} g_0 G_m^2(u_k) dx + \int_{\Omega} f G_m(u_k) dx \\ \leq \varepsilon \int_{\Omega} g_0 G_m^{1+q}(u_k) dx + C(\varepsilon) \int_{u \geq m} g_0 dx + C \left(\int_{\Omega} G_m^{2^*}(u_k) dx \right)^{1/2^*}. \end{aligned}$$

Therefore we conclude that

$$\int_{\Omega} |\nabla G_m(u_k)|^2 dx + c \int_{\Omega} |\nabla G_m(u_k)|^q G_m(u_k) \leq \int_{u \geq m} g_0 dx + C \left(\int_{\Omega} G_m^{2^*}(u_k) dx \right)^{1/2^*}$$

where $C > 0$ is a positive constant that depends only on the data and is independent of m and k .

Recall that $f, g_0 \in L^r$ with $r > N/2$, then using Sobolev inequality,

$$\begin{aligned} & C \left(\int_{\Omega} G_m^{2^*}(u_k) dx \right)^{2/2^*} + \left(\int_{\Omega} G_m^{(1+1/q)q^*}(u_k) dx \right)^{q/q^*} \\ & \leq C \left(|u_k \geq m|^{1/r'} + C \int_{\Omega} G_m^{2^*}(u_k) dx \right)^{1/2^*} |u_k \geq m|^{1-1/r-1/2^*}. \end{aligned}$$

From Young's inequality there results that

$$\begin{aligned} & C \left(\int_{\Omega} G_m^{2^*}(u_k) dx \right)^{2/2^*} + \left(\int_{\Omega} G_m^{(1+1/q)q^*}(u_k) dx \right)^{q/q^*} \\ & \leq C(|u_k \geq m|^{1/r'} + C|u_k \geq m|^{2(1-1/r-1/2^*)}) \leq C|u_k \geq m|^\gamma, \end{aligned}$$

where $\gamma = \min\{2(1 - 1/r - 1/2^*), 1/r'\}$. By a direct computation we get easily that $2^*\gamma/2 > 1$.

We set $\beta(m) = |u_k \geq m|$, then for $m_1 < m_2$ we have

$$\begin{aligned} \beta^{1/2^*}(m_2)(m_2 - m_1) & \leq \left(\int_{u_k \geq m_2} |u_k - m_1|^{2^*} dx \right)^{1/2^*} \\ & \leq \left(\int_{u_k \geq m_1} |u_k - m_1|^{2^*} dx \right)^{1/2^*} \\ & \leq \beta^{\gamma/2}(m_1). \end{aligned}$$

Thus

$$\beta(m_2) \leq \frac{\beta^{2^*\gamma/2}(m_1)}{(m_2 - m_1)^{2^*}}.$$

Since $2^*\gamma/2 > 1$, using the Stampacchia classical result, (see [15]), there exists a universal constant $\bar{m} > 0$ such that $\beta(m) = 0$ if $m \geq \bar{m}$. Thus $u_k \leq \bar{m}$, and then choosing $k \gg \bar{m}$, we obtain that $u = u_k$ solves

$$(2.9) \quad \begin{cases} u \in \mathcal{K}(\psi) & \text{for all } v \in \mathcal{K}(\psi), \text{ we have} \\ \int_{\Omega} \nabla u \nabla(v - u) dx + \int_{\Omega} |\nabla u|^q(v - u) dx \geq \int_{\Omega} (\lambda g_0 u + f)(v - u_n) dx. \end{cases}$$

Hence we conclude the proof. \square

In the following result, we still consider a weight g_0 with the same summability condition as in Theorem 2.2, but now we assume $f \in L^1(\Omega)$.

THEOREM 2.3. *Assume that f, g_0 are positive functions, $f \in L^1(\Omega)$ and $g_0 \in L^r(\Omega)$ with $r > \frac{N}{q}$, then for all $\lambda \geq 0$ there exists $u \in W_0^{1,q}(\Omega)$ such that $u \geq \psi$ and for all $v \in \mathcal{K}(\psi)$ we have*

$$\int_{\Omega} \nabla u \nabla (T_k(v - u)) \, dx + \int_{\Omega} |\nabla u|^q (T_k(v - u)) \, dx \geq \int_{\Omega} (\lambda g_0 u + f) (T_k(v - u)) \, dx.$$

PROOF. Consider a sequence $f_n \in L^\infty(\Omega)$ such that $f_n \uparrow f$ in $L^1(\Omega)$. By Theorem 2.2, there exists a sequence of positive bounded functions $\{u_n\}$, solutions to problems,

$$(2.10) \quad \begin{cases} u_n \in \mathcal{K}(\psi), & \text{for all } v \in \mathcal{K}(\psi), \text{ we have} \\ \int_{\Omega} \nabla u_n \nabla (v - u_n) \, dx + \int_{\Omega} |\nabla u_n|^q (v - u_n) \, dx \geq \int_{\Omega} (\lambda g_0 u_n + f) (v - u_n) \, dx. \end{cases}$$

Consider the function

$$(2.11) \quad \Psi_k(s) = \begin{cases} 0 & \text{if } s \leq k \\ s - k & \text{if } k \leq s \leq k + 1 \\ 1 & \text{if } s \geq k + 1. \end{cases}$$

Define $v = u_n - \Psi_k(u_n)$, then $v \geq \psi$. Using v as a test function in (2.10) it follows that,

$$\begin{aligned} & \int_{k \leq u_n \leq k+1} |\nabla u_n|^2 \, dx + \int_{\Omega} |\nabla u_n|^q \Psi_k(u_n) \, dx \\ & \leq \lambda \int_{\Omega} g(x) u_n \Psi_k(u_n) \, dx + \int_{\Omega} f_n(x) \Psi_k(u_n) \, dx. \end{aligned}$$

Notice that,

$$\int_{\Omega} |\nabla u_n|^q \Psi_k(u_n) \, dx = \int_{\Omega} |\nabla \Theta(u_n)|^q \, dx,$$

where $\Theta(s) = \int_0^s \Psi_k^{1/q}(s) \, ds$. Using the hypothesis (1.3) on g we obtain,

$$\int_{\Omega} |\nabla \Theta(u_n)|^q \, dx \geq C(g_0, q) \left(\int_{\Omega} g \Theta(u_n) \, dx \right)^q.$$

Therefore, using the fact that $s \Psi_k(s) \leq \Theta(s) + C$, it follows that,

$$\begin{aligned} & \int_{k \leq u_n \leq k+1} |\nabla u_n|^2 \, dx + C(g_0, q) \left(\int_{\Omega} g_0 \Theta(u_n) \, dx \right)^q \\ & \leq \lambda \int_{\Omega} g_0 \Theta(u_n) \, dx + C\lambda \int_{\Omega} g_0(x) \, dx + \int_{\Omega} f_n(x) \, dx. \end{aligned}$$

Thus using Young's inequality there results

$$\int_{k \leq u_n \leq k+1} |\nabla u_n|^2 dx + C(g_0, q) \left(\int_{\Omega} g_0 \Theta(u_n) dx \right)^q \leq C,$$

and then

$$\int_{\Omega} |\nabla u_n|^q \Psi_k(u_n) dx \leq C, \quad \lambda \int_{\Omega} g_0(x) u_n \Psi_k(u_n) dx \leq C,$$

where C is a positive constant depending only on the data.

We set now $v = u_n - T_k(u_n - \psi)$. It is clear that $v \geq \psi$, using v as a test function in (2.10) we get

$$\begin{aligned} & \int_{\Omega} \nabla u_n \nabla T_k(u_n - \psi) dx + \int_{\Omega} |\nabla u_n|^q T_k(u_n - \psi) dx \\ & \leq \lambda \int_{\Omega} g_0 u_n T_k(u_n - \psi) dx + \int_{\Omega} f T_k(u_n - \psi) dx. \end{aligned}$$

Using the fact that $u_n \geq \psi$ and that $\lambda \int_{\Omega} g_0 u_n T_k(u_n - \psi) dx + \int_{\Omega} f T_k(u_n - \psi) dx \leq C$ for all n , it follows that

$$\int_{|u_n - \psi| \leq k} \nabla u_n \nabla (u_n - \psi) dx \leq C.$$

Then using Hölder's and Young's inequalities we get

$$\int_{|u_n - \psi| \leq k} |\nabla u_n|^2 dx \leq C.$$

Let $k > 0$, then

$$\int_{|u_n| \leq k} |\nabla u_n|^2 dx \leq \int_{|u_n - \psi| \leq k + \|\psi\|_{L^\infty}} |\nabla u_n|^2 dx \leq C.$$

Hence $\{T_k(u_n)\}$ is bounded in $W_0^{1,2}(\Omega)$ and $\{u_n\}$ is bounded in $W_0^{1,q}(\Omega)$. Thus we get the existence of u such that $u_n \rightharpoonup u$ weakly in $W_0^{1,q}(\Omega)$ and $T_k u_n \rightharpoonup T_k u$ weakly in $W_0^{1,2}(\Omega)$. It is clear by the assumption on g_0 that $g_0 u_n \rightarrow g_0 u$ strongly in $L^1(\Omega)$.

Define $\Phi_{k-1}(s) = T_1(G_{k-1}(s))$, then $\Phi_{k-1}(u_n) |\nabla u_n|^q \geq |\nabla u_n|^q \chi_{\{u_n \geq k\}}$.

Let $v = u_n - \Phi_{k-1}(u_n)$, then $v \geq \psi$. Using v as a test function in (2.10) there results

$$\int_{\Omega} |\nabla \Phi_{k-1}(u_n)|^2 dx + \int_{\Omega} \Phi_{k-1}(u_n) |\nabla u_n|^q dx \leq \int_{\Omega} (\lambda g_0(x) u_n + f_n(x)) \Phi_{k-1}(u_n) dx.$$

Since $\{u_n\}$ is uniformly bounded in $L^p(\Omega)$, $\forall p \leq q^*$, it follows that

$$\begin{aligned} |\{x \in \Omega, \text{ such that } k-1 < u_n(x) < k\}| &\rightarrow 0, \\ |\{x \in \Omega, \text{ such that } u_n(x) > k\}| &\rightarrow 0 \quad \text{as } k \rightarrow \infty, \end{aligned}$$

uniformly in n . Thus we conclude

$$(2.12) \quad \lim_{k \rightarrow \infty} \int_{\{u_n \geq k\}} |\nabla u_n|^q dx = 0, \quad \text{uniformly in } n.$$

We claim that $\nabla u_n \rightarrow \nabla u$, a.e. in Ω .

To prove the claim we follow the same arguments as in the proof of Theorem 2.2.

Let $v = u_n - (T_k(u_n) - T_k(u))^+$, then $v \in \mathcal{H}(\psi)$ and $v - u_n = -(T_k(u_n) - T_k(u))^+$, hence there result

$$\begin{aligned} &\int_{\Omega} \nabla u_n \nabla (T_k(u_n) - T_k(u))^+ dx + \int_{\Omega} |\nabla u_n|^q (T_k(u_n) - T_k(u))^+ dx \\ &\leq \int_{\Omega} (\lambda g_0 u_n + f_n) (T_k(u_n) - T_k(u))^+ dx. \end{aligned}$$

A direct computation shows that

$$\begin{aligned} &\int_{\Omega} \nabla u_n \nabla (T_k(u_n) - T_k(u))^+ dx \\ &= \int_{\Omega} |\nabla (T_k(u_n) - T_k(u))^+|^2 dx + \int_{\Omega} \nabla G_k(u_n) \nabla T_k(u) dx \\ &\quad + \int_{\Omega} \nabla T_k(u) \nabla (T_k(u_n) - T_k(u))^+ dx \\ &= \int_{\Omega} |\nabla (T_k(u_n) - T_k(u))^+|^2 dx + o(1). \end{aligned}$$

It is clear that $\int_{\Omega} (\lambda g_0 u_n + f_n) (T_k(u_n) - T_k(u))^+ dx \rightarrow 0$ as $n \rightarrow \infty$, therefore we conclude that

$$\int_{\Omega} |\nabla (T_k(u_n) - T_k(u))^+|^2 dx \rightarrow 0 \quad \text{as } n \rightarrow \infty.$$

Thus $(T_k(u_n) - T_k(u))^+ \rightarrow 0$ strongly in $W_0^{1,2}$. Take now $v = u_n + \phi((T_k(u_n) - T_k(u))^-)$, then $v \in \mathcal{H}(\psi)$. Using v as a test function in (2.10) we obtain that

$$\begin{aligned}
& \int_{\Omega} \nabla \phi'((w_n - u_k)^-) \nabla u_n \nabla (T_k(u_n) - T_k(u))^- dx \\
& \quad + \int_{\Omega} |\nabla u_n|^q \phi((T_k(u_n) - T_k(u))^-) dx \\
& \leq \lambda \int_{\Omega} g_0(x) u_n \phi((T_k(u_n) - T_k(u))^-) dx + \int_{\Omega} f_n(x) \phi((T_k(u_n) - T_k(u))^-) dx.
\end{aligned}$$

Thus

$$\begin{aligned}
& \int_{T_k(u_n) \leq T_k(u)} \phi'((T_k(u_n) - T_k(u))^-) \nabla u_n \nabla (u_n - u) dx \\
& \quad - \int_{\Omega} |\nabla u_n| \phi((T_k(u_n) - T_k(u))^-) dx \\
& \leq \lambda \int_{T_k(u_n) \leq T_k(u)} g_0(x) u_n \phi((T_k(u_n) - T_k(u))^-) dx \\
& \quad + \int_{T_k(u_n) \leq T_k(u)} f_n(x) \phi((T_k(u_n) - T_k(u))^-) dx.
\end{aligned}$$

Since $T_k u_n \rightharpoonup T_k u$ weakly in $W_0^{1,2}(\Omega)$, then

$$\begin{aligned}
& \int_{T_k(u_n) \leq T_k(u)} \phi'((T_k(u_n) - T_k(u))^-) \nabla u_n \nabla (u_n - u) dx \\
& = \int_{\Omega} |\nabla((T_k(u_n) - T_k(u))^-)|^2 \phi'((T_k(u_n) - T_k(u))^-) dx + o(1).
\end{aligned}$$

Since $q \leq 2$, as in the computation in the proof of Theorem 2.2 it follows that

$$\begin{aligned}
& \frac{1}{2} \int_{\Omega} |\nabla((T_k(u_n) - T_k(u))^-)|^2 dx \\
& \leq \int_{\Omega} (\phi'((T_k(u_n) - T_k(u))^-)) \\
& \quad - \varepsilon |\phi'((T_k(u_n) - T_k(u))^-)| |\nabla((T_k(u_n) - T_k(u))^-)|^2 dx \\
& \leq o(1),
\end{aligned}$$

whence $(T_k(u_n) - T_k(u))^- \rightarrow 0$ in $W_0^{1,2}(\Omega)$ and then $T_k(u_n) \rightarrow T_k(u)$ strongly in $W_0^{1,2}$. Hence the claim follows.

To finish the proof, we have to show that $|\nabla u_n|^q \rightarrow |\nabla u|^q$ strongly in $L^1(\Omega)$. Since the sequence of gradients converges a.e. in Ω , we have just to prove the

equi-integrability of the sequence $\{|\nabla u_n|^q\}$ and then apply Vitali's Theorem. Let $E \subset \Omega$ be a measurable set. Then,

$$\int_E |\nabla u_n|^q dx \leq \int_E |\nabla T_k u_n|^q dx + \int_{\{u_n \geq k\} \cap E} |\nabla u_n|^q dx.$$

Since $q \leq 2$, then for all $k > 0$, $T_k(u_n) \rightarrow T_k(u)$ strongly in $W_0^{1,q}(\Omega)$. Hence the integral $\int_E |\nabla T_k(u_n)|^q dx$ is uniformly small if $|E|$ is small enough. On the other hand, by (2.12) we obtain that

$$\int_{\{u_n \geq k\} \cap E} |\nabla u_n|^q dx \leq \int_{\{u_n \geq k\}} |\nabla u_n|^q dx \rightarrow 0 \quad \text{as } k \rightarrow \infty \text{ uniformly in } n.$$

The equi-integrability of $|\nabla u_n|^q$ follows immediately, and the proof is complete. \square

PROOF OF THEOREM 2.1. Consider $g_n(x) = \min\{g_0(x), n\} \in L^\infty(\Omega)$. It is clear that $g_n \rightarrow g_0$ strongly in $W^{-1,q'}$.

Using Theorem 2.3, we get the existence of a sequence of nonnegative functions $\{u_n\}$ such that u_n solves

$$(2.13) \quad \begin{cases} u_n \geq \psi & \text{for all } v \in \mathcal{K}(\psi), \text{ we have} \\ \int_{\Omega} \nabla u_n \nabla (T_k(v - u_n)) dx + \int_{\Omega} |\nabla u_n|^q (T_k(v - u_n)) dx \\ \geq \int_{\Omega} (\lambda g_n u_n + f)(T_k(v - u_n)) dx. \end{cases}$$

By setting $v = u_n - \Psi_k(u_n)$, where Ψ_k is defined in (2.11), and using the same computations as in the proof of Theorem 2.3, it follows that

$$\begin{aligned} & \int_{k \leq u_n \leq k+1} |\nabla u_n|^2 dx + C(g_n, q) \left(\int_{\Omega} g_n \Theta(u_n) dx \right)^q \\ & \leq \lambda \int_{\Omega} g_n \Theta(u_n) dx + C\lambda \int_{\Omega} g_0(x) dx + \int_{\Omega} f_n(x) dx, \end{aligned}$$

where

$$C(g_n, q) = \inf_{\phi \in W_0^{1,q}(\Omega) \setminus \{0\}} \frac{(\int_{\Omega} |\nabla \phi|^q dx)^{1/q}}{\int_{\Omega} g_n |\phi| dx}.$$

It is clear that $C(g_n, q) \uparrow C(g_0, q) > 0$. Then

$$\int_{k \leq u_n \leq k+1} |\nabla u_n|^2 dx + \left(\int_{\Omega} g_n \Theta(u_n) dx \right)^q \leq C,$$

and

$$\int_{\Omega} |\nabla u_n|^q \Psi_k(u_n) dx \leq C, \quad \lambda \int_{\Omega} g_n(x) u_n \Psi(u_n) dx \leq C,$$

where C is a positive constant depending only on the data. As in the proof of Theorem 2.3 we can prove that $\int_{\Omega} |\nabla u_n|^q dx \leq C$ and then $u_n \rightarrow u$ weakly in $W_0^{1,q}(\Omega)$. Since $g_n \rightarrow g$ strongly in $W^{-1,q'}$, then $\int_{\Omega} g_n u_n dx \rightarrow \int_{\Omega} g_0 u dx$ strongly in $L^1(\Omega)$. Now to complete the proof we follow closely the argument used in the proof of Theorem 2.3. \square

2.1. Partial uniqueness result. In this subsection we consider the case $q = 2$. We will prove a uniqueness result for positive solutions. We will use the next Comparison Principle that is a variation of the uniqueness result obtained in [1]. For the reader's convenience we include a short proof.

LEMMA 2.4. *Let $f \in L^1(\Omega)$ is a non negative function and suppose that g_0 is an admissible function in the sense of condition (1.3). Let $u_1, u_2 \in W_0^{1,2}(\Omega)$ be functions such that $u_1 > 0$ (resp. $u_2 > 0$) is a subsolution (resp. supesolution) to problem*

$$(2.14) \quad \begin{cases} -\Delta u + |\nabla u|^2 = \lambda g_0(x)u + f(x) & \text{in } \Omega, \\ u > 0 & \text{in } \Omega, \\ u = 0 & \text{on } \partial\Omega, \end{cases}$$

then $u_1 \leq u_2$ in Ω .

PROOF. For $i = 1, 2$ we set $v_i \equiv 1 - e^{-u_i}$, then $0 < v_i \leq 1$ in Ω and v_1 (resp. v_2) is a subsolution (resp. supersolution) to problem

$$(2.15) \quad \begin{cases} -\Delta v = \lambda g_0(x)(1-v) \log\left(\frac{1}{1-v}\right) + (1-v)f(x) & \text{in } \Omega, \\ v = 0 & \text{on } \partial\Omega. \end{cases}$$

Define

$$H(x, v) = \begin{cases} \lambda g(x)(1-v) \log\left(\frac{1}{1-v}\right) + (1-v)f(x), & \text{if } 0 < v < 1, \\ 0, & \text{if } v \geq 1. \end{cases}$$

By a direct computation we find that $\frac{H(v,x)}{v}$ is a non-increasing function in v for $v \geq 0$, then by similar arguments as in [1], we conclude that $v_1 \leq v_2$, therefore the result follows.

Then we can prove the following result about uniqueness.

THEOREM 2.5. *Assume $q = 2$, $\psi \equiv 0$ and that the hypotheses of Theorem 2.1 hold. Then problem (2.1) has a unique positive solution.*

PROOF. The existence of a nonnegative solution is a consequence of the Theorem 2.1.

Assume that u_1 and u_2 are two nonnegative solution to the obstacle problem (2.1). We claim that u_1 is strictly positive in Ω (the same conclusion holds for u_2). To prove the claim we consider $\phi \in \mathcal{C}_0^\infty(\Omega)$ a nonnegative function. Let $v = T_h(u) + \phi$, it is clear that for h large we have $v \geq \psi$ and $v \in L^\infty(\Omega) \cap W_0^{1,2}(\Omega)$. Using v as a test function in (2.1), it follows that

$$\begin{aligned} & \int_{\Omega} \nabla u_1 \nabla (T_k(T_h(u_1) - u_1 + \phi)) dx + \int_{\Omega} |\nabla u_1|^2 (T_k(T_h(u_1) - u_1 + \phi)) dx \\ & \geq \int_{\Omega} (\lambda g_0 u_1 + f) (T_k(T_h(u_1) - u_1 + \phi)) dx. \end{aligned}$$

Since $u_1 \in W_0^{1,2}(\Omega)$, then for fixed k , using the Dominated Convergence Theorem and a duality argument we can pass to the limit in h , hence

$$(2.16) \quad \int_{\Omega} \nabla u_1 \nabla (T_k(\phi)) dx + \int_{\Omega} |\nabla u_1|^2 (T_k(\phi)) dx \geq \int_{\Omega} (\lambda g_0 u_1 + f) (T_k(\phi)) dx.$$

Since $\phi \in \mathcal{C}_0^\infty(\Omega)$, then choosing k large enough we conclude that

$$\int_{\Omega} \nabla u_1 \nabla \phi dx + \int_{\Omega} |\nabla u_1|^2 \phi dx \geq \int_{\Omega} (\lambda g_0 u_1 + f) \phi dx.$$

Thus u_1 is a nonnegative supersolution to problem

$$(2.17) \quad \begin{cases} -\Delta w + |\nabla w|^2 = \lambda g_0(x)w + f & \text{in } \Omega, \\ w \geq 0 \text{ in } \Omega & \text{and } w \in W_0^{1,2}(\Omega). \end{cases}$$

From [1] we know that the above problem has a unique positive solution and that if w_1 is a supersolution to problem (2.17), then $w_1 \geq w$. Since u_1 is a supersolution, then the claim follows.

We follow now closely the argument used in [3]. Define

$$v_1 = u_1 - \delta\phi, \quad \text{and} \quad v_2 = u_2 + \delta\phi.$$

It is clear that $v_2 \geq 0$. We show that for δ small enough, depending on ϕ , then $v_1 \geq 0$. It is clear that $v_1 \geq 0$ in $\Omega \setminus \text{Supp}(\phi)$. Since $\text{Supp}(\phi) \llcorner \Omega$, then by using the strict positivity of u_1 there exists a positive constant c such that $u_1 \geq c$ in $\text{Supp}(\phi)$. Hence for $x \in \text{Supp}(\phi)$, we have

$$v_1(x) \geq C - \delta \|\phi\|_{L^\infty}.$$

Choosing $\delta \geq \frac{C}{2\|\phi\|_{L^\infty}}$, we conclude that $v_1(x) \geq C/2$ for all $x \in \text{Supp}(\phi)$. Hence $v_1 \geq 0$ in Ω . Notice that the same conclusion holds if we substitute v_1 by $v_1^h \equiv T_h(u_1) - \delta\phi$ where h is large enough.

It is clear that we cannot use v_1 and v_2 directly as a test function in the corresponding obstacle problem of u_1 and u_2 . Thus we use an approximation argument. Set

$$v_1^h = T_h(u_1) - \delta\phi, \quad v_2^h = T_h(u_2) + \delta\phi.$$

As above, for k fixed and passing to the limit in h , it follows that

$$\int_{\Omega} \nabla u_1 \nabla (T_k(\delta\phi)) \, dx + \int_{\Omega} |\nabla u_1|^2 (T_k(\delta\phi)) \, dx \leq \int_{\Omega} (\lambda g_0 u_1 + f) (T_k(\delta\phi)) \, dx$$

and

$$\int_{\Omega} \nabla u_2 \nabla (T_k(\delta\phi)) \, dx + \int_{\Omega} |\nabla u_2|^2 (T_k(\delta\phi)) \, dx \geq \int_{\Omega} (\lambda g_0 u_2 + f) (T_k(\delta\phi)) \, dx.$$

Letting $k \rightarrow \infty$ and using the fact that $\phi \in \mathcal{C}_0^\infty(\Omega)$, there results that

$$\int_{\Omega} \nabla u_1 \nabla \phi \, dx + \int_{\Omega} |\nabla u_1|^2 \phi \, dx \leq \int_{\Omega} (\lambda g_0 u_1 + f) \phi \, dx$$

and

$$\int_{\Omega} \nabla u_2 \nabla \phi \, dx + \int_{\Omega} |\nabla u_2|^2 \phi \, dx \geq \int_{\Omega} (\lambda g_0 u_2 + f) \phi \, dx.$$

Thus u_1 (resp. u_2) is a nonnegative subsolution (resp. supersolution) to (2.17), then by Lemma 2.4 we conclude that $u_1 \leq u_2$. Following the same argument as above we get easily that $u_2 \leq u_1$. Hence $u_1 = u_2$ and then the uniqueness result follows. \square

3. THE GENERAL RESULT

In this section we will consider the general case, that is, f and ψ can change sign. More precisely we will consider the following obstacle problem

$$(3.1) \quad \begin{cases} u \geq \psi \text{ a.e in } \Omega, & \text{for all } v \in \mathcal{K}(\psi), \text{ we have} \\ \int_{\Omega} \nabla u \nabla (v - u) \, dx + \int_{\Omega} h(u) |\nabla u|^q (v - u) \, dx \geq \int_{\Omega} (g(x, u) + f) (v - u) \, dx, \end{cases}$$

where $\psi \in W_0^{1,2}(\Omega) \cap L^\infty(\Omega)$ and $f \in L^1(\Omega)$. We will assume that h satisfies the next sign condition,

$$(3.2) \quad h(s) \text{ is an continuous increasing function such that, } h(s)s \geq 0, \quad \forall s \in \mathbb{R}$$

and

$$(3.3) \quad |g(x, s)| \leq \lambda g_0(x)|s| + g_1(x)$$

where g_0 satisfies the condition (1.3) and $g_1 \in L^1(\Omega)$.

We can formulate the general existence result as follows.

THEOREM 3.1. *Assume that conditions (3.2) and (1.2) hold. Then for all $f \in L^1(\Omega)$, there exists a function u such that $h(u)|\nabla u|^q \in L^1(\Omega)$, $T_k(u) \in W_0^{1,2}(\Omega)$ for all $k > 0$, and for all $v \in \mathcal{K}(\psi)$ we have*

$$(3.4) \quad \int_{\Omega} \nabla u \nabla (T_k(v - u)) \, dx + \int_{\Omega} h(u) |\nabla u|^q (T_k(v - u)) \, dx \\ \geq \int_{\Omega} (g(x, u) + f)(T_k(v - u)) \, dx.$$

We will say that u is an entropy solution to the obstacle problem (3.1).

As in section 2 we will prove the existence result for regular data and then we will pass to the limit.

For $n, m > 0$, we define,

$$h_n(s) = \frac{h(s)}{1 + \frac{1}{n}|h(s)|} \quad \text{and} \quad g_m(x, s) = \frac{g(x, s)}{1 + \frac{1}{m}|g(x, s)|}.$$

Then we have

THEOREM 3.2. *Assume that the above conditions (3.2) and (1.2) hold, then for all $f \in L^\infty(\Omega)$ and for all m fixed, there exists a $u_m \in W_0^{1,2}(\Omega)$ such that $h(x, u_m)|\nabla u_m|^q \in L^1(\Omega)$ and for all $v \in \mathcal{K}(\psi)$ we have*

$$(3.5) \quad \int_{\Omega} \nabla u_m \nabla (v - u_m) \, dx + \int_{\Omega} h(u_m) |\nabla u_m|^q (v - u_m) \, dx \\ \geq \int_{\Omega} (g_m(x, u_m) + f)(v - u_m) \, dx.$$

PROOF. Fixed $m > 0$, since $f \in L^\infty(\Omega)$, using the classical results in [13] or [12], there exists $u_n \in \mathcal{K}(\psi)$, a solution to the obstacle problem

$$(3.6) \quad \int_{\Omega} \nabla u_n \nabla (v - u_n) \, dx + \int_{\Omega} h_n(u_n) \frac{|\nabla u_n|^q}{1 + \frac{1}{n}|\nabla u_n|^q} (v - u_n) \, dx \\ \geq \int_{\Omega} (\lambda g_m(x, u_n) + f)(v - u_n) \, dx$$

for all $v \in \mathcal{K}(\psi)$. For fixed m , we want to pass to the limit in n . Let $v = T_k(w_n)$, since $\psi \in L^\infty(\Omega)$, then choosing k such that $k \gg \|\psi\|_\infty$, we conclude that v is an admissible test function. Notice that $v - u_n = -G_k(u_n)$, hence

$$\begin{aligned} & \int_{\Omega} |\nabla G_k(u_n)|^2 dx + \int_{\Omega} h_n(u_n) \frac{|\nabla G_k(u_n)|^q}{1 + \frac{1}{n} |G_k(u_n)|^q} G_k(u_n) \\ & \leq \int_{\Omega} (g_m(x, u_n) + f) G_k(u_n) dx. \end{aligned}$$

Since $h_n(u_n)G_k(u_n) \geq 0$, then

$$\int_{\Omega} |\nabla G_k(u_n)|^2 dx \leq \int_{\Omega} (m + |f|) |G_k(u_n)| dx.$$

Since $f \in L^\infty(\Omega)$, then by the classical Stampacchia result, see [15], the following L^∞ -estimate holds,

$$\|u_n\|_{L^\infty} \leq C(m, f, \Omega).$$

Take $v = \psi$ as test function in (3.6), then

$$\begin{aligned} & \int_{\Omega} \nabla u_n \nabla (u_n - \psi) dx + \int_{\Omega} h_n(u_n) \frac{|\nabla u_n|^q}{1 + \frac{1}{n} |\nabla u_n|^q} (u_n - \psi) dx \\ & \leq \int_{\Omega} (g_m(x, u_n) + f) (u_n - \psi) dx. \end{aligned}$$

Since $u_n \geq \psi$, we get

$$\int_{\Omega} |\nabla u_n|^2 dx \leq \int_{\Omega} \nabla u_n \nabla \psi dx + \int_{\Omega} g_m(x, u_n) + f) (u_n - \psi) dx.$$

Thus using Hölder inequality and the previous estimate we obtain that

$$\int_{\Omega} |\nabla u_n|^2 dx \leq C(f, m\Omega, k) \quad \text{uniformly in } n.$$

Therefore, up to a subsequence, $u_n \rightharpoonup u_m$ weakly in $W_0^{1,2}(\Omega)$ as $n \rightarrow \infty$. By weak-*convergence in $L^\infty(\Omega)$ we also have that $w_m \in W_0^{1,2} \cap L^\infty(\Omega)$ and $u_m \geq \psi$. Since $\{u_n\}$ is bounded in $L^\infty(\Omega)$, then as in the first section, following closely the argument used in [6], we get the strong convergence of u_n in $W_0^{1,2}(\Omega)$. Thus $u_m \in W_0^{1,2}(\Omega) \cap L^\infty(\Omega)$ and it solves the obstacle problem (3.5). Hence the result follows. \square

Now we can prove Theorem 3.1.

PROOF OF THEOREM 3.1. Let $\{f_n\}$ be a sequence of bounded functions such that $f_m \rightarrow f$ strongly in $L^1(\Omega)$ and consider u_m the solution to the obstacle problem (3.5) obtained above and let $k > 0$, then using $v = u_m - \Psi_k(u_m)$, where Ψ_k is defined in (2.11) with $\Psi(s) = -\Psi(-s)$ for $s < 0$, as a test function in (3.5) it follows that

$$\begin{aligned} & \int_{k \leq |u_m| \leq k+1} |\nabla u_m|^2 dx + \int_{\Omega} h(u_m) |\nabla u_m|^q \Psi_k(u_m) \\ & \leq \lambda \int_{\Omega} g_0(x) |u_m| |\Psi_k(u_m)| + \int_{\Omega} g_1(x) dx + \int_{\Omega} |f_n(x)| dx. \end{aligned}$$

Using the properties of h it follows that, for k large,

$$h(u_m) |\nabla u_m|^q \Psi_k(u_m) \geq C(k) |\nabla u_m|^q |\Psi_k(u_m)| \geq C |\nabla \Theta(u_m)|^q$$

where $\Theta(s) = \int_0^s |\Psi(s)|^{1/q} ds$. Thus

$$\int_{k \leq |u_m| \leq k+1} |\nabla u_m|^2 dx + \int_{\Omega} |\nabla \Theta(u_m)|^q \leq \lambda \int_{\Omega} g_0(x) |u_m| |\Psi_k(u_m)| + C.$$

Notice that $|s| |\Psi_k(s)| \leq \Theta(s) + C$, then

$$\int_{k \leq |u_m| \leq k+1} |\nabla u_m|^2 dx + \int_{\Omega} |\nabla \Theta(u_m)|^q \leq \lambda \int_{\Omega} g_0(x) \Theta(u_m) dx + C.$$

Since g_0 satisfies the condition (1.3), thus

$$\int_{k \leq |u_m| \leq k+1} |\nabla u_m|^2 dx + \int_{\Omega} |\nabla \Theta(u_m)|^q \leq C(k).$$

Therefore

$$\int_{\Omega} |\nabla u_n|^q |\Psi_k(u_m)| dx \leq C, \quad \lambda \int_{\Omega} g_0(x) |u_n| |\Psi(u_n)| dx \leq C$$

where C is a positive constant depending only on the data. As in the proof of Theorem 2.3 we can prove that $\int_{\Omega} |\nabla u_n|^q dx \leq C$ and then $u_n \rightarrow u$ weakly in

$W_0^{1,q}(\Omega)$. It is clear that $\int_{\Omega} g_0 u_n dx \rightarrow \int_{\Omega} g_0 u dx$ strongly in $L^1(\Omega)$, then using the Dominated Convergence Theorem we obtain that

$$g_m(x, u_m) \rightarrow g(x, u) \quad \text{strongly in } L^1(\Omega).$$

The rest of the proof follows exactly as in the proof of Theorem 2.3. \square

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Boumediene Abdellaoui, Sidi Mohamed Bouguima
Département de Mathématiques
Université Aboubekr Belkaïd
Tlemcen
Tlemcen 13000, Algeria
boumediene.abdellaoui@uam.es, bouguima@yahoo.fr

Ireneo Peral
Departamento de Matemáticas
U. Autónoma de Madrid and ICMAT
28049 Madrid, Spain
ireneo.peral@uam.es