



Partial Differential Equations — *On a (p, q) -Laplacian problem with parametric concave term and asymmetric perturbation*, by SALVATORE A. MARANO, SUNRA J. N. MOSCONI and NIKOLAOS S. PAPAGEORGIOU, communicated on April 21, 2017.

ABSTRACT. — A Dirichlet problem driven by the (p, q) -Laplace operator and an asymmetric concave reaction with positive parameter is investigated. Four nontrivial smooth solutions (two positive, one negative, and the remaining nodal) are obtained once the parameter turns out to be sufficiently small. Under an oddness condition near the origin for the perturbation, a whole sequence of sign-changing solutions, which converges to zero, is produced.

KEY WORDS: (p, q) -Laplacian, asymmetric perturbation, concave term, extremal constant-sign and nodal solution

MATHEMATICS SUBJECT CLASSIFICATION: 35J20, 35J92, 58E05

1. INTRODUCTION

Let Ω be a bounded domain in \mathbb{R}^N with a C^2 -boundary $\partial\Omega$, let $1 < s < q \leq p < +\infty$, and let $\mu \in \mathbb{R}_0^+$. Consider the Dirichlet problem

$$(1.1) \quad \begin{cases} -\Delta_p u - \mu \Delta_q u = \lambda |u|^{s-2} u + f(x, u) & \text{in } \Omega, \\ u = 0 & \text{on } \partial\Omega, \end{cases}$$

where Δ_r , $r > 1$, denotes the r -Laplace operator, namely

$$\Delta_r u := \operatorname{div}(|\nabla u|^{r-2} \nabla u) \quad \forall u \in W_0^{1,r}(\Omega),$$

$p = q$ iff $\mu = 0$, λ is a real parameter, while $f : \Omega \times \mathbb{R} \rightarrow \mathbb{R}$ satisfies Carathéodory's conditions.

The non-homogeneous differential operator $Au := \Delta_p u + \Delta_q u$ that appears in (1.1) is usually called (p, q) -Laplacian. It stems from a wide range of important applications, including biophysics [9], plasma physics [27], reaction-diffusion equations [2, 6], as well as models of elementary particles [3, 5, 8].

This paper treats the existence of multiple solutions, with a precise sign information, to (1.1) when, roughly speaking,

- 1) $\lambda > 0$ is suitably small, and
- 2) $t \mapsto f(x, t)$ exhibits an asymmetric behavior as t goes from $-\infty$ to $+\infty$.

We will assume that, for an appropriate constant $C > 0$,

$$\begin{aligned} -C &\leq \liminf_{t \rightarrow -\infty} \frac{f(x, t)}{|t|^{p-2}t} \leq \limsup_{t \rightarrow -\infty} \frac{f(x, t)}{|t|^{p-2}t} \leq \lambda_{1,p} \\ &\leq \liminf_{t \rightarrow +\infty} \frac{f(x, t)}{t^{p-1}} \leq \limsup_{t \rightarrow +\infty} \frac{f(x, t)}{t^{p-1}} \leq C \end{aligned}$$

uniformly in $x \in \Omega$, where $\lambda_{1,p}$ indicates the first eigenvalue of $(-\Delta_p, W_0^{1,p}(\Omega))$. Hence, $f(x, \cdot)$ grows $(p-1)$ -linearly at $\pm\infty$ and only a partial interaction with $\lambda_{1,p}$ is allowed (nonuniform non-resonance).

Since $s < q \leq p$, the term $t \mapsto \lambda|t|^{s-2}t$ represents a parametric ‘concave’ contribution inside the reaction of (1.1).

Under 1), 2), and a further hypothesis involving the rate of $f(x, \cdot)$ near zero, Problem (1.1) admits four nontrivial $C_0^1(\bar{\Omega})$ -solutions, two positive, one negative, and the remaining nodal; see Theorem 4.2. If, moreover, $t \mapsto f(x, t)$ turns out to be odd in a neighborhood of zero then there exists a whole sequence $\{u_n\}$ of nodal solutions such that $u_n \rightarrow 0$ in $C^1(\bar{\Omega})$; cf. Theorem 4.3.

The adopted approach exploits variational methods, truncation techniques, as well as results from Morse theory. Regularity is a standard matter.

Many recent papers have been devoted to elliptic problems with either

- p -Laplacian and asymmetric nonlinearity (see, e.g., [7, 19, 20, 21, 22] and the references therein), or
- (p, q) -Laplacian and symmetric reaction (see for instance [4, 17, 18] and the references given there).

On the contrary, to the best of our knowledge, few articles treat equations driven by the (p, q) -Laplace operator and an asymmetric nonlinearity. Actually, we can only mention [24], where $\mu := 1$, $q := 2$, the parametric concave term does not appear, f satisfies somewhat different assumptions, and a complete sign information on the solutions is not performed. A wider bibliography on these topics can be found in the survey paper [16].

2. PRELIMINARIES

Let $(X, \|\cdot\|)$ be a real Banach space and let X^* be its topological dual, with duality brackets $\langle \cdot, \cdot \rangle$. An operator $A : X \rightarrow X^*$ is called of type $(S)_+$ provided

$$x_n \rightharpoonup x \quad \text{in } X, \quad \limsup_{n \rightarrow +\infty} \langle A(x_n), x_n - x \rangle \leq 0 \quad \Rightarrow \quad x_n \rightarrow x \quad \text{in } X.$$

For $\varphi \in C^1(X)$ and $c \in \mathbb{R}$, put

$$\begin{aligned} \varphi^c &:= \{x \in X : \varphi(x) \leq c\}, & K_\varphi &:= \{x \in X : \varphi'(x) = 0\}, \\ K_\varphi^c &:= \{x \in K_\varphi : \varphi(x) = c\}. \end{aligned}$$

Given an isolated critical point $\bar{x} \in K_\varphi^c$, we define the k -th critical group of φ at \bar{x} as

$$C_k(\varphi, \bar{x}) := H_k(\varphi^c \cap U, \varphi^c \cap U \setminus \{\bar{x}\}), \quad k \in \mathbb{N}_0,$$

where U is any neighborhood of \bar{x} such that $K_\varphi \cap \varphi^c \cap U = \{\bar{x}\}$ and $H_k(A, B)$ denotes the k -th relative singular homology group for the pair (A, B) with integer coefficients. The excision property of singular homology ensures that this definition does not depend on the choice of U ; see [23] for details.

We say that φ satisfies the Cerami condition when

- (C) *Every sequence $\{x_n\} \subseteq X$ such that $\{\varphi(x_n)\}$ is bounded and $(1 + \|x_n\|)\varphi'(x_n) \rightarrow 0$ in X^* admits a strongly convergent subsequence.*

The following version [23] of the mountain pass theorem will be employed.

THEOREM 2.1. *If $\varphi \in C^1(X)$ satisfies (C), $x_0, x_1 \in X$, $0 < \rho < \|x_0 - x_1\|$,*

$$\max\{\varphi(x_0), \varphi(x_1)\} < m_\rho := \inf_{\|x-x_0\|=\rho} \varphi(x),$$

and

$$c := \inf_{\gamma \in \Gamma} \max_{t \in [0,1]} \varphi(\gamma(t)), \quad \text{where } \Gamma := \{\gamma \in C^0([0,1], X) : \gamma(0) = x_0, \gamma(1) = x_1\},$$

then: $m_\rho \leq c$; K_φ^c is nonempty; $C_1(\varphi, \bar{x}) \neq 0$ provided $\bar{x} \in K_\varphi^c$ turns out to be isolated.

Hereafter, Ω will denote a fixed bounded domain in \mathbb{R}^N with a C^2 -boundary $\partial\Omega$. Let $u, v : \Omega \rightarrow \mathbb{R}$ be measurable and let $t \in \mathbb{R}$. The symbol $u \leq v$ means $u(x) \leq v(x)$ for almost every $x \in \Omega$, $t^\pm := \max\{\pm t, 0\}$, $u^\pm(\cdot) := u(\cdot)^\pm$. If $p \in [1, +\infty)$ then $p' := p/(p-1)$ is the conjugate exponent of p and p^* indicates the Sobolev critical exponent in dimension N , namely

$$p^* = \begin{cases} \frac{Np}{N-p} & \text{when } p < N, \\ \text{any } q > 1 & \text{for } p = N, \\ +\infty & \text{otherwise.} \end{cases}$$

Set, provided $r \in [1, +\infty]$,

$$L_+^r(\Omega) = \{u \in L^r(\Omega) : u \geq 0 \text{ a.e. in } \Omega\}.$$

If $r < +\infty$ then, as usual,

$$\|u\|_r := \left(\int_\Omega |\nabla u|^r dx \right)^{1/r}, \quad u \in W_0^{1,r}(\Omega), \quad \text{and}$$

$$|u|_r := \left(\int_\Omega |u|^r dx \right)^{1/r}, \quad u \in L^r(\Omega).$$

$W^{-1,r'}(\Omega)$ denotes the dual space of $W_0^{1,r}(\Omega)$ while $A_r : W_0^{1,r}(\Omega) \rightarrow W^{-1,r'}(\Omega)$ is the nonlinear operator stemming from the negative r -Laplacian, i.e.,

$$\langle A_r(u), v \rangle := \int_{\Omega} |\nabla u|^{r-2} \nabla u \cdot \nabla v \, dx \quad \forall u, v \in W_0^{1,r}(\Omega).$$

It is known [10, Section 6.2] that A_r turns out to be bounded, continuous, strictly monotone, as well as of type $(S)_+$.

Given $\xi \in L_+^{\infty}(\Omega) \setminus \{0\}$, we define

$$(2.1) \quad \lambda_{1,r}(\xi) := \inf \left\{ \frac{\int_{\Omega} |\nabla u|^r \, dx}{\int_{\Omega} \xi |u|^r \, dx} : u \in W_0^{1,r}(\Omega), u \neq 0 \right\}.$$

When no confusion can arise, simply write $\lambda_{1,r} := \lambda_{1,r}(1)$. Some basic properties of $\lambda_{1,r}(\xi)$ and its eigenfunctions are listed below.

PROPOSITION 2.1. *Let $1 < r < +\infty$ and let $\xi \in L_+^{\infty}(\Omega) \setminus \{0\}$. Then:*

1. $\lambda_{1,r}(\xi)$ is positive and attained on a positive function $\hat{u}_{1,r} \in W_0^{1,r}(\Omega)$, which fulfills $|\hat{u}_{1,r}|_r = 1$ as well as

$$(2.2) \quad A_r(u) = \lambda_{1,r}(\xi) \xi |u|^{r-2} u.$$

2. Solutions to (2.2) coincide with minima of (2.1) and form a one-dimensional linear space.
3. The function $\xi \mapsto \lambda_{1,r}(\xi)$ is monotone (strictly) decreasing with respect to the a.e. ordering of $L_+^{\infty}(\Omega)$.

Through the compactness of the embedding $W_0^{1,r}(\Omega) \hookrightarrow L^r(\Omega)$ one can verify [25, p. 356] the next result.

PROPOSITION 2.2. *If $\xi \in L_+^{\infty}(\Omega) \setminus \{\lambda_{1,r}\}$ and $\xi \leq \lambda_{1,r}$ then there exists a constant $c(\xi) > 0$ such that*

$$\|u\|_r^r - \int_{\Omega} \xi |u|^r \, dx \geq c(\xi) \|u\|_r^r \quad \forall u \in W_0^{1,r}(\Omega).$$

We will also employ the linear space

$$C_0^1(\bar{\Omega}) := \{u \in C^1(\bar{\Omega}) : u|_{\partial\Omega} = 0\},$$

which is complete with respect to the standard $C^1(\bar{\Omega})$ -norm. Its positive cone

$$C_+ := \{u \in C_0^1(\bar{\Omega}) : u(x) \geq 0 \text{ in } \bar{\Omega}\}$$

has a nonempty interior given by

$$\text{int}(C_+) = \left\{ u \in C_+ : u(x) > 0 \ \forall x \in \Omega, \frac{\partial u}{\partial n}(x) < 0 \ \forall x \in \partial\Omega \right\}.$$

Here, $n(x)$ denotes the outward unit normal to $\partial\Omega$ at x .

Suppose $g : \Omega \times \mathbb{R} \rightarrow \mathbb{R}$ is a Carathéodory function growing sub-critically, i.e.,

$$|g(x, t)| \leq c(1 + |t|^{r-1}) \quad \text{in } \Omega \times \mathbb{R},$$

where $c > 0$, $1 \leq r < p^*$. Write, as usual, $G(x, t) := \int_0^t g(x, \tau) \, d\tau$ and consider the C^1 -functional $\varphi : W_0^{1,p}(\Omega) \rightarrow \mathbb{R}$ defined by

$$\varphi(u) := \frac{1}{p} \|u\|_p^p + \frac{\mu}{q} \|u\|_q^q - \int_{\Omega} G(x, u(x)) \, dx \quad \forall u \in W_0^{1,p}(\Omega),$$

with $1 < q \leq p$ and $\mu \geq 0$. The next result [11] establishes a relation between local minimizers of φ in $C_0^1(\overline{\Omega})$ and in $W_0^{1,p}(\Omega)$.

PROPOSITION 2.3. *If $u_0 \in W_0^{1,p}(\Omega)$ is a local $C_0^1(\overline{\Omega})$ -minimizer of φ , then $u_0 \in C_0^{1,\alpha}(\overline{\Omega})$ for some $\alpha \in (0, 1)$ and u_0 turns out to be a local $W_0^{1,p}(\Omega)$ -minimizer of φ .*

3. SOLUTIONS OF CONSTANT SIGN

In this section we will construct three nontrivial constant-sign solutions to Problem (1.1) provided the parameter is small enough. From now on, everywhere in Ω stands for almost everywhere and $q = p$ iff $\mu = 0$.

The hypotheses on the reaction f will be as follows.

(h₀) $f : \Omega \times \mathbb{R} \rightarrow \mathbb{R}$ is a Carathéodory function such that

$$|f(x, t)| \leq C(1 + |t|^{p-1}) \quad \forall (x, t) \in \Omega \times \mathbb{R},$$

where $C \in \mathbb{R}^+$.

(h₁) There exists $\xi_1 \in L_+^\infty(\Omega) \setminus \{\lambda_{1,p}\}$ satisfying

$$\lambda_{1,p} \leq \xi_1(x) \leq \liminf_{t \rightarrow +\infty} \frac{f(x, t)}{t^{p-1}} \quad \text{uniformly in } x \in \Omega.$$

(h₂) There is $\xi_2 \in L_+^\infty(\Omega) \setminus \{\lambda_{1,p}\}$ such that

$$\limsup_{t \rightarrow -\infty} \frac{f(x, t)}{|t|^{p-2}t} \leq \xi_2(x) \leq \lambda_{1,p} \quad \text{uniformly with respect to } x \in \Omega.$$

(h₃) There exist $\delta_0, \theta_0 \in (0, 1)$ fulfilling

$$0 \leq f(x, t)t \leq \mu\lambda_{1,q}|t|^q + \theta_0\lambda_{1,p}|t|^p \quad \forall (x, t) \in \Omega \times [-\delta_0, \delta_0].$$

REMARK 3.1. It should be noted that (h_0) – (h_2) entail

$$-C \leq \liminf_{t \rightarrow -\infty} \frac{f(x, t)}{|t|^{p-2}t} \leq \xi_2(x) \leq \lambda_{1,p} \leq \xi_1(x) \leq \limsup_{t \rightarrow +\infty} \frac{f(x, t)}{t^{p-1}} \leq C.$$

If $1 < s < q$ and $\lambda \in \mathbb{R}^+$, we put

$$g_\lambda(x, t) := \lambda|t|^{s-1}t + f(x, t),$$

which still satisfies a growth condition like (h_0) , but with a different positive constant depending on λ , say C_λ , and

$$G_\lambda(x, t) := \int_0^t g_\lambda(x, \tau) d\tau.$$

The energy functional $\varphi_\lambda \in C^1(W_0^{1,p}(\Omega))$ that stems from Problem (1.1) is defined by

$$\varphi_\lambda(u) := \frac{1}{p} \|u\|_p^p + \frac{\mu}{q} \|u\|_q^q - \int_\Omega G_\lambda(x, u(x)) dx \quad \forall u \in W_0^{1,p}(\Omega).$$

Suitable truncations of it will be employed. With this aim, set

$$g_\lambda^+(x, t) := g_\lambda(x, t^+), \quad g_\lambda^-(x, t) := g_\lambda(x, -t^-), \quad G_\lambda^\pm(x, t) := \int_0^t g_\lambda^\pm(x, \tau) d\tau.$$

Evidently, $G_\lambda^+(x, t) = G_\lambda(x, t^+)$, $G_\lambda^-(x, t) = G_\lambda(x, -t^-)$, and the associated functionals

$$\varphi_\lambda^\pm(u) := \frac{1}{p} \|u\|_p^p + \frac{\mu}{q} \|u\|_q^q - \int_\Omega G_\lambda^\pm(x, u(x)) dx, \quad u \in W_0^{1,p}(\Omega),$$

turn out to be C^1 as well. Likewise the proof of [18, Theorem 4.1], using the nonlinear regularity theory developed in [13, 14], the strong maximum principle, and the Hopf boundary point lemma [26, pp. 111 and 120], yields

PROPOSITION 3.1. *Under (h_0) and (h_3) , nontrivial critical points for φ_λ^+ (resp., φ_λ^-) actually are critical points of φ_λ and belong to $\text{int}(C_+)$ (resp., $-\text{int}(C_+)$).*

LEMMA 3.1. *If (h_0) – (h_2) hold true, then*

1. φ_λ^+ satisfies Condition (C).
2. φ_λ^- is coercive (hence it fulfills the Cerami condition too).

PROOF. 1. Let $\{u_n\} \subseteq W_0^{1,p}(\Omega)$ be such that $\{\varphi_\lambda^+(u_n)\}$ is bounded and

$$(1 + \|u_n\|_p)(\varphi_\lambda^+)'(u_n) \rightarrow 0 \quad \text{in } W^{-1,p'}(\Omega).$$

Since the embedding $W_0^{1,p}(\Omega) \hookrightarrow L^p(\Omega)$ is compact, while $A_p + \mu A_q$ enjoys property $(S)_+$, it suffices to show that $\{u_n\}$ is bounded. One has

$$(3.1) \quad \left| \langle A_p(u_n), v \rangle + \mu \langle A_q(u_n), v \rangle - \int_{\Omega} g_{\lambda}^+(x, u_n) v dx \right| \leq \varepsilon_n \frac{\|v\|_p}{1 + \|u_n\|_p} \quad \forall v \in W_0^{1,p}(\Omega),$$

where $\varepsilon_n \rightarrow 0^+$. Letting $v := -u_n^-$ yields

$$\|u_n^-\|_p + \mu \|u_n^-\|_q \leq \varepsilon_n,$$

so that $u_n^- \rightarrow 0$ and

$$(3.2) \quad \left| \langle A_p(u_n^+), v \rangle + \mu \langle A_q(u_n^+), v \rangle - \int_{\Omega} g_{\lambda}^+(x, u_n^+) v dx \right| \leq \varepsilon'_n \|v\|_p$$

for some $\varepsilon'_n \rightarrow 0^+$. Suppose $\|u_n^+\|_p \rightarrow +\infty$ and put $w_n := u_n^+ / \|u_n^+\|_p$. From $\|w_n\|_p \equiv 1$ it follows, up to subsequences,

$$(3.3) \quad w_n \rightharpoonup w \quad \text{in } W_0^{1,p}(\Omega), \quad w_n \rightarrow w \quad \text{in } L^p(\Omega), \quad w \geq 0.$$

Moreover,

$$|g_{\lambda}^+(\cdot, u_n^+)|_{p'}^{p'} \leq C_{\lambda} \int_{\Omega} (1 + |u_n^+|^p) dx \leq C_{\lambda} (|\Omega| + \lambda_{1,p}^{-1} \|u_n^+\|_p^p)$$

because f satisfies (h_0) . This implies

$$\left| \frac{g_{\lambda}^+(\cdot, u_n^+)}{\|u_n^+\|_p^{p-1}} \right|_{p'} \leq C_{\lambda}^{1/p'} \left(\frac{|\Omega|}{\|u_n^+\|_p^p} + \frac{1}{\lambda_{1,p}} \right)^{1/p'}.$$

Since the right-hand side is bounded, we may suppose

$$\frac{g_{\lambda}^+(\cdot, u_n^+)}{\|u_n^+\|_p^{p-1}} \rightharpoonup h \quad \text{in } L^{p'}(\Omega).$$

Recalling that $s < p$, namely (h_1) holds true for g_{λ}^+ , and proceeding as in [23, pp. 317–318] produces

$$(3.4) \quad h = \eta w^{p-1} \quad \text{for some } \eta \in L^{\infty}(\Omega) \text{ with } \xi_1(x) \leq \eta(x) \leq C.$$

Through (3.2) we then have

$$(3.5) \quad \left| \langle A_p(w_n), v \rangle + \mu \|u_n^+\|_p^{q-p} \langle A_q(w_n), v \rangle - \int_{\Omega} \frac{g_{\lambda}^+(x, u_n^+)}{\|u_n^+\|_p^{p-1}} v dx \right| \leq \varepsilon'_n \frac{\|v\|_p}{\|u_n^+\|_p^{p-1}}.$$

Now, choose $v := w_n - w$ and use (3.3) to arrive at

$$\lim_{n \rightarrow +\infty} \langle A_p(w_n), w_n - w \rangle = 0.$$

Therefore, by [23, Proposition 2.72], $w_n \rightarrow w$ in $W_0^{1,p}(\Omega)$, whence $\|w\|_p = 1$. Via (3.5) we thus obtain, letting $n \rightarrow +\infty$,

$$\langle A_p(w), v \rangle = \int_{\Omega} \eta w^{p-1} v \, dx \quad \forall v \in W_0^{1,p}(\Omega),$$

i.e., $\lambda = 1$ is an eigenvalue for the problem

$$-\Delta_p u = \lambda \eta |u|^{p-2} u \quad \text{in } \Omega, \quad u = 0 \quad \text{on } \partial\Omega$$

associated with the eigenfunction w . However, due to Item 4) of Proposition 2.1, (h₁), and (3.4),

$$1 = \lambda_{1,p}(\lambda_{1,p}) > \lambda_{1,p}(\xi_1) \geq \lambda_{1,p}(\eta).$$

Point 3) in the same result ensures that w changes sign, contradicting $w \geq 0$.

2. By (h₀) and (h₂), for every $\varepsilon > 0$ there exists a constant $C_\varepsilon > 0$ such that

$$F(x, t) \leq \frac{1}{p} (\xi_2(x) + \varepsilon) |t|^p + C_\varepsilon \quad \forall (x, t) \in \Omega \times (-\infty, 0].$$

Thus, on account of Proposition 2.2,

$$\begin{aligned} \varphi_\lambda^-(u) &\geq \frac{1}{p} \|u\|_p^p + \frac{\mu}{q} \|u\|_q^q - \frac{\lambda}{s} |u|_s^s - \int_{\Omega} \frac{1}{p} (\xi_2 + \varepsilon) |u^-|^p \, dx - C_\varepsilon |\Omega| \\ &\geq \frac{1}{p} \|u^+\|_p^p + \frac{1}{p} \left(\|u^-\|_p^p - \int_{\Omega} \xi_2 |u^-|^p \, dx \right) - \frac{\varepsilon}{p} \|u^-\|_p^p - \frac{\lambda}{s} |u|_s^s - C_\varepsilon |\Omega| \\ &\geq \frac{1}{p} \|u^+\|_p^p + \frac{1}{p} \left(c(\xi_2) - \frac{\varepsilon}{\lambda_{1,p}} \right) \|u^-\|_p^p - \frac{\lambda}{s} |u|_s^s - C_\varepsilon |\Omega|. \end{aligned}$$

Choosing $\varepsilon := \lambda_{1,p} c(\xi_2) / 2$ and recalling that $s < p$ finally provides the desired coercivity property. \square

With slight modifications one can verify the next lemma.

LEMMA 3.2. *Under (h₀)–(h₂), the functional φ_λ satisfies Condition (C) for all $\lambda > 0$.*

PROOF. Fix $\lambda > 0$. Let $\{u_n\} \subseteq W_0^{1,p}(\Omega)$ be such that $\{\varphi_\lambda(u_n)\}$ is bounded and

$$(1 + \|u_n\|_p)(\varphi_\lambda)'(u_n) \rightarrow 0 \quad \text{in } W^{-1,p'}(\Omega).$$

Then (3.1) holds with g_λ instead of g_λ^+ . Choosing $v := -u_n^-$, it furnishes

$$(3.6) \quad \|u_n^-\|_p^p + \mu \|u_n^-\|_q^q - \lambda |u_n^-|_s^s - \int_{\Omega} f(x, -u_n^-)(-u_n^-) dx \leq \varepsilon_n, \quad n \in \mathbb{N},$$

where $\varepsilon_n \rightarrow 0^+$. Thanks to (h_0) and (h_2) , for every $\varepsilon > 0$ there exists a constant $C_\varepsilon > 0$ such that

$$f(x, t)t \leq (\xi_2(x) + \varepsilon)|t|^p + C_\varepsilon \quad \forall (x, t) \in \Omega \times]-\infty, 0].$$

So, the proof of Conclusion 2 in the previous lemma carries over, giving the coerciveness of the functional

$$u \mapsto \|u^-\|_p^p + \mu \|u^-\|_q^q - \lambda |u^-|_s^s - \int_{\Omega} f(x, -u^-)(-u^-) dx, \quad u \in W_0^{1,p}(\Omega).$$

Hence, due to (3.6), the sequence $\{u_n^-\}$ has to be bounded. To check that the same holds for $\{u_n^+\}$, suppose on the contrary $\|u_n^+\|_p \rightarrow +\infty$ and put $w_n := u_n / \|u_n^+\|_p$. Obviously, $\{w_n\}$ turns out to be bounded, because so is $\{u_n^-\}$. Moreover, $w_n^- \rightarrow 0$ while, along a subsequence when necessary,

$$w_n \rightharpoonup w \quad \text{in } W_0^{1,p}(\Omega), \quad w_n \rightarrow w \quad \text{in } L^p(\Omega).$$

As before, via (h_0) we see that $\{g_\lambda(\cdot, u_n) \|u_n^+\|_p^{1-p}\}$ is bounded in $L^{p'}(\Omega)$. Now, divide the present version of (3.1) by $\|u_n^+\|_p^{p-1}$, test with $v := w_n - w$, use the inequality $q < p$, and let $n \rightarrow +\infty$ to achieve

$$\lim_{n \rightarrow +\infty} \langle A_p(w_n), w_n - w \rangle = 0,$$

which implies $w_n \rightarrow w$ in $W_0^{1,p}(\Omega)$. Consequently, $w \geq 0$ and $\|w\|_p \geq 1$. Since

$$\left| \frac{g_\lambda(\cdot, -u^-)}{\|u_n^+\|_p^{p-1}} \right|_{p'} \leq C_\lambda^{1/p'} \left(\frac{|\Omega|}{\|u_n^+\|_p^p} + \frac{1}{\lambda_{1,p}} \frac{\|u_n^-\|_p^p}{\|u_n^+\|_p^p} \right)^{1/p'} \rightarrow 0,$$

we have

$$(3.7) \quad \frac{g_\lambda(\cdot, u_n)}{\|u_n^+\|_p^{p-1}} - \frac{g_\lambda(\cdot, u_n^+)}{\|u_n^+\|_p^{p-1}} \rightarrow 0 \quad \text{in } L^{p'}(\Omega).$$

The same arguments of [23, pp. 317–318] yield here

$$\frac{g_\lambda(\cdot, u_n^+)}{\|u_n^+\|_p^{p-1}} \rightarrow \eta w^{p-1} \quad \text{in } L^{p'}(\Omega),$$

with appropriate $\eta \in L^\infty(\Omega)$ fulfilling $\xi_1(x) \leq \eta(x) \leq C$. Thanks to (3.7), this holds true also for $\{g_\lambda(\cdot, u_n) \|u_n^+\|_p^{1-p}\}$. Hence, from

$$\left| \langle A_p(w_n), v \rangle + \mu \|u_n^+\|_p^{q-p} \langle A_q(w_n), v \rangle - \int_\Omega \frac{g_\lambda(x, u_n)}{\|u_n^+\|_p^{p-1}} v \, dx \right| \leq \varepsilon'_n \frac{\|v\|_p}{\|u_n^+\|_p^{p-1}}$$

(cf. (3.5)) it follows, when $n \rightarrow +\infty$,

$$\langle A_p(w), v \rangle = \int_\Omega \eta w^{p-1} v \, dx \quad \forall v \in W_0^{1,p}(\Omega).$$

Now the proof goes on exactly as the one of Item 1 in Lemma 3.1. \square

LEMMA 3.3. *If (h_0) and (h_3) are satisfied, then there exists a constant $\lambda^* > 0$ such that to every $\lambda \in (0, \lambda^*)$ corresponds a $\rho_\lambda \in \mathbb{R}^+$ complying with*

$$m_\lambda := \inf_{\|u\|_p = \rho_\lambda} \varphi_\lambda^+(u) > 0.$$

PROOF. Fix any $r \in (p, p^*)$. Through (h_0) and (h_3) we obtain

$$f(x, t) \leq \mu \lambda_{1,q} t^{q-1} + \theta_0 \lambda_{1,p} t^{p-1} + C_r t^{r-1}, \quad (x, t) \in \Omega \times [0, +\infty),$$

which, when integrated, entails

$$F(x, t) \leq \mu \frac{\lambda_{1,q}}{q} t^q + \theta_0 \frac{\lambda_{1,p}}{p} t^p + \frac{C_r}{r} t^r \quad \text{in } \Omega \times [0, +\infty).$$

Here, $C_r \in \mathbb{R}^+$. By the Sobolev, Hölder, and Poincaré inequalities one has

$$\begin{aligned} \varphi_\lambda^+(u) &\geq \frac{1}{p} \|u\|_p^p + \frac{\mu}{q} \|u\|_q^q - \frac{\lambda}{s} |u^+|_s^s - \mu \frac{\lambda_{1,q}}{q} |u|_q^q - \theta_0 \frac{\lambda_{1,p}}{p} |u|_p^p - \frac{C_r}{r} |u|_r^r \\ &\geq \frac{1 - \theta_0}{p} \|u\|_p^p - \frac{\lambda}{s} |u^+|_s^s - \frac{C_r}{r} |u|_r^r \\ &\geq \frac{1 - \theta_0}{p} \|u\|_p^p - C_1(\Omega) \frac{\lambda}{s} |u^+|_{p^*}^s - \frac{C_2(\Omega) C_r}{r} |u|_{p^*}^r \\ &\geq \left[\frac{1 - \theta_0}{p} - \bar{C}_1 \lambda \|u\|_p^{s-p} - \bar{C}_2 \|u\|_p^{r-p} \right] \|u\|_p^p \end{aligned}$$

for appropriate positive constants \bar{C}_1, \bar{C}_2 . Letting $\|u\|_p = \lambda^{\frac{1}{r-s}}$ yields

$$\varphi_\lambda^+(u) \geq \left[\frac{1 - \theta_0}{p} - \bar{C}_1 \lambda^{1 - \frac{p-s}{r-s}} - \bar{C}_2 \lambda^{\frac{r-p}{r-s}} \right] \lambda^{\frac{1}{r-s}} = \left[\frac{1 - \theta_0}{p} - (\bar{C}_1 + \bar{C}_2) \lambda^{\frac{r-p}{r-s}} \right] \lambda^{\frac{1}{r-s}}.$$

This immediately brings the conclusion, because $s < p < r$ and $0 < \theta_0 < 1$. \square

From now on, λ^* will denote the real number just found.

LEMMA 3.4. *Suppose (h_0) – (h_1) hold true. Then*

$$\lim_{\tau \rightarrow +\infty} \varphi_\lambda^+(\tau \hat{u}_{1,p}) = -\infty,$$

with $\hat{u}_{1,p}$ as in Proposition 2.1.

PROOF. Thanks to (h_0) – (h_1) , for every $\varepsilon > 0$ there exists a constant $C_\varepsilon > 0$ such that

$$(3.8) \quad F(x, t) \geq \frac{\xi_1(x) - \varepsilon}{p} t^p - C_\varepsilon \quad \forall (x, t) \in \Omega \times [0, +\infty).$$

The properties of $\hat{u}_{1,p}$ and ξ_1 produce

$$\int_\Omega (\xi_1 - \lambda_{1,p}) \hat{u}_{1,p}^p dx > 0.$$

Choose $\varepsilon > 0$ fulfilling

$$\theta := \int_\Omega (\xi_1 - \lambda_{1,p}) \hat{u}_{1,p}^p dx - \varepsilon \int_\Omega \hat{u}_{1,p}^p dx > 0.$$

Since $\|\hat{u}_{1,p}\|_p^p = \lambda_{1,p} |\hat{u}_{1,p}|_p^p = \lambda_{1,p}$, via (3.8) we get

$$\begin{aligned} \varphi_\lambda^+(\tau \hat{u}_{1,p}) &\leq \frac{\tau^p}{p} \|\hat{u}_{1,p}\|_p^p + \mu \frac{\tau^q}{q} \|\hat{u}_{1,p}\|_q^q - \frac{\tau^p}{p} \int_\Omega (\xi_1 - \varepsilon) \hat{u}_{1,p}^p dx - \lambda \frac{\tau^s}{s} |\hat{u}_{1,p}|_s^s + C_\varepsilon |\Omega| \\ &\leq -\theta \frac{\tau^p}{p} + \mu \frac{\tau^q}{q} \|\hat{u}_{1,p}\|_q^q - \lambda \frac{\tau^s}{s} |\hat{u}_{1,p}|_s^s + C_\varepsilon |\Omega| \end{aligned}$$

for all $\tau > 0$. The conclusion follows from $q < p$. □

Now, critical point arguments will provide three constant-sign solutions.

THEOREM 3.1. *Let (h_0) – (h_3) be satisfied. Then:*

1. *For every $\lambda \in (0, \lambda^*)$, Problem (1.1) admits two positive solutions $u_0, u_1 \in \text{int}(C_+)$.*
2. *For every $\lambda > 0$ there exists a negative solution $u_2 \in -\text{int}(C_+)$ to (1.1).*

PROOF. 1. Pick $\lambda \in (0, \lambda^*)$. Lemma 3.4 gives a $\tau \in \mathbb{R}^+$ so large that $\varphi_\lambda^+(\tau \hat{u}_{1,p}) < 0$. On account of Lemmas 3.1 and 3.3, Theorem 2.1 applies to φ_λ^+ . Thus, there is $u_0 \in W_0^{1,p}(\Omega)$ fulfilling

$$(\varphi_\lambda^+)'(u_0) = 0, \quad \varphi_\lambda^+(u_0) \geq m_\lambda > 0,$$

whence $u_0 \neq 0$. By Proposition 3.1, the function u_0 turns out to be a solution of (1.1) lying in $\text{int}(C_+)$. Next, define

$$B_{\rho_\lambda} := \{u \in W_0^{1,p}(\Omega) : \|u\|_p < \rho_\lambda\},$$

where ρ_λ comes from Lemma 3.3. A standard procedure based on the weak sequential lower semicontinuity of φ_λ^+ ensures that this functional attains its minimum at some $u_1 \in \bar{B}_{\rho_\lambda}$. Fix $w \in \text{int}(C_+)$ and choose $\tau_1 > 0$ complying with

$$\|\tau_1 w\|_p < \rho_\lambda, \quad \tau_1 \sup_{x \in \Omega} w(x) \leq \delta_0.$$

Thanks to (h₃) we have

$$f(x, \tau w(x)) \geq 0 \quad \forall \tau \in (0, \tau_1),$$

which easily entails

$$\varphi_\lambda^+(\tau w) \leq \frac{\tau^p}{p} \|w\|_p^p + \mu \frac{\tau^q}{q} \|w\|_q^q - \lambda \frac{\tau^s}{s} |w|_s^s < 0$$

provided τ is sufficiently small (recall that $s < q < p$). Hence, a fortiori,

$$\varphi_\lambda^+(u_1) < 0.$$

The above inequality brings both $u_1 \neq u_0$ and $u_1 \in B_{\rho_\lambda} \setminus \{0\}$. On account of [15, Lemma 4.3] we thus arrive at $(\varphi_\lambda^+)'(u_1) = 0$. Finally, due to Proposition 3.1, the function u_1 lies in $\text{int}(C_+)$ and solves (1.1).

2. φ_λ^- is coercive (cf. Lemma 3.1) and weakly sequentially lower semicontinuous. So, it attains its minimum at some $u_2 \in W_0^{1,p}(\Omega)$. As before, we see that $\varphi_\lambda^-(u_2) < 0$, whence $u_2 \neq 0$. Since $(\varphi_\lambda^-)'(u_2) = 0$, Proposition 3.1 applies to get the conclusion. \square

4. NODAL SOLUTIONS

Let us first show that (1.1) admits extremal constant-sign, namely a smallest positive and a biggest negative, solutions. Indeed, (h₀) and (h₃) yield a real number $c_0 > 0$ fulfilling

$$f(x, t) \geq -c_0 |t|^p \quad \forall (x, t) \in \Omega \times \mathbb{R}.$$

The same arguments exploited to prove [18, Lemma 2.2] ensure here that, given $\lambda > 0$, the auxiliary problem

$$-\Delta_p u - \mu \Delta_q u = \lambda |u|^{s-2} u - c_0 |u|^{p-2} u \quad \text{in } \Omega, \quad u = 0 \quad \text{on } \partial\Omega$$

has only one positive solution $\bar{u}_\lambda \in \text{int}(C_+)$, while, by oddness, $\bar{v}_\lambda := -\bar{u}_\lambda$ turns out to be its unique negative solution. Reasoning as made for [17, Lemma 3.3] we next achieve

LEMMA 4.1. *Under (h_0) – (h_3) , any positive (resp., negative) solution u of (1.1) satisfies the inequality $\bar{u}_\lambda \leq u$ (resp., $u \leq \bar{v}_\lambda$).*

These facts give rise to the following result; cf. the proof of [18, Lemma 4.2].

LEMMA 4.2. *Assume (h_0) – (h_3) . Then, for every $\lambda \in (0, \lambda^*)$, Problem (1.1) possesses a smallest positive solution $u_\lambda^+ \in \text{int}(C_+)$ and a greatest negative solution $u_\lambda^- \in -\text{int}(C_+)$.*

We are in a position now to produce a nodal solution through a mountain pass procedure. Set

$$[u_\lambda^-, u_\lambda^+] := \{u \in W_0^{1,p}(\Omega) : u_\lambda^- \leq u \leq u_\lambda^+ \text{ a.e. in } \Omega\}.$$

THEOREM 4.1. *If (h_0) – (h_3) hold true and $\lambda \in (0, \lambda^*)$ then there exists a sign-changing solution $u_3 \in C_0^1(\bar{\Omega}) \cap [u_\lambda^-, u_\lambda^+]$ to (1.1).*

PROOF. The proof is similar to that of [17, Theorem 3.8]; so, we only sketch it. Define, for every $(x, t) \in \Omega \times \mathbb{R}$,

$$\hat{g}_\lambda(x, t) := \begin{cases} \lambda |u_\lambda^-(x)|^{s-2} u_\lambda^-(x) + f(x, u_\lambda^-(x)) & \text{if } t < u_\lambda^-(x), \\ \lambda |t|^{s-2} t + f(x, t) & \text{if } u_\lambda^-(x) \leq t \leq u_\lambda^+(x), \\ \lambda |u_\lambda^+(x)|^{s-2} u_\lambda^+(x) + f(x, u_\lambda^+(x)) & \text{if } t > u_\lambda^+(x), \end{cases}$$

as well as $\hat{G}_\lambda(x, t) := \int_0^t \hat{g}_\lambda(x, \tau) d\tau$. The associated energy functional

$$\hat{\phi}_\lambda(u) := \frac{1}{p} \|u\|_p^p + \frac{\mu}{q} \|u\|_q^q - \int_\Omega \hat{G}_\lambda(x, u(x)) dx, \quad u \in W_0^{1,p}(\Omega),$$

clearly fulfills

$$(4.1) \quad K_{\hat{\phi}_\lambda} \subseteq [u_\lambda^-, u_\lambda^+],$$

while using Proposition 2.3 one verifies that both u_λ^+ and u_λ^- turn out to be local $W_0^{1,p}(\Omega)$ -minimizers of $\hat{\phi}_\lambda$; see [18, 17] for details. We may suppose $K_{\hat{\phi}_\lambda}$ finite, otherwise $\hat{\phi}_\lambda$ (and so ϕ_λ , by standard nonlinear regularity theory) would have infinitely many critical points in $[u_\lambda^-, u_\lambda^+] \setminus \{u_\lambda^-, u_\lambda^+\}$, which brings the conclusion. Consequently, u_λ^+ and u_λ^- are strict local minimizers. Without loss of generality, assume $\hat{\phi}_\lambda(u_\lambda^+) \leq \hat{\phi}_\lambda(u_\lambda^-)$. The reasoning employed to establish [1, Proposition 29] produces here $\rho \in (0, 1)$ such that

$$(4.2) \quad \|u_\lambda^- - u_\lambda^+\|_p > \rho, \quad \max\{\hat{\phi}_\lambda(u_\lambda^+), \hat{\phi}_\lambda(u_\lambda^-)\} < m_\rho := \inf_{\|u - u_\lambda^-\|_p = \rho} \hat{\phi}_\lambda(u).$$

Moreover, $\hat{\phi}_\lambda$ satisfies Condition (C), because it evidently is coercive. Hence, Theorem 2.1 applies, and there exists $u_3 \in K_{\hat{\phi}_\lambda}$ such that $m_\rho \leq \hat{\phi}_\lambda(u_3)$. Combining

(4.2) with (4.1) entails $u_3 \in [u_\lambda^-, u_\lambda^+] \setminus \{u_\lambda^-, u_\lambda^+\}$ while, by Theorem 2.1 again,

$$(4.3) \quad C_1(\hat{\varphi}_\lambda, u_3) \neq 0.$$

Now, through $u_\lambda^+, -u_\lambda^- \in \text{int}(C_+)$, $\hat{\varphi}_\lambda|_{[u_\lambda^-, u_\lambda^+]} = \varphi_\lambda|_{[u_\lambda^-, u_\lambda^+]}$, and $0 \in \text{int}_{C_0^1(\bar{\Omega})}([u_\lambda^-, u_\lambda^+])$, we infer

$$C_k(\hat{\varphi}_\lambda|_{C_0^1(\bar{\Omega})}, 0) = C_k(\varphi_\lambda|_{C_0^1(\bar{\Omega})}, 0), \quad k \in \mathbb{N}_0.$$

Since $C_0^1(\bar{\Omega})$ is dense in $W_0^{1,p}(\Omega)$, one has

$$C_k(\hat{\varphi}_\lambda|_{C_0^1(\bar{\Omega})}, 0) = C_k(\hat{\varphi}_\lambda, 0), \quad C_k(\varphi_\lambda|_{C_0^1(\bar{\Omega})}, 0) = C_k(\varphi_\lambda, 0).$$

Let us verify that

$$(4.4) \quad C_k(\hat{\varphi}_\lambda, 0) = 0 \quad \forall k \in \mathbb{N}_0,$$

which will force $u_3 \neq 0$ thanks to (4.3). Pick any $\theta \in (s, q)$. Assumption (h₃) directly yields

$$g_\lambda(x, t)t = \lambda|t|^s + f(x, t)t \geq \lambda|t|^\theta \quad \text{in } \Omega \times [-\delta_0, \delta_0].$$

If $\delta_1 < \delta_0$ fulfills

$$|t| \leq \delta_1 \quad \Rightarrow \quad \left(\frac{\theta}{s} - 1\right)\lambda|t|^s \geq \mu\lambda_{1,q}|t|^q + \theta_0\lambda_{1,p}|t|^p$$

then from (h₃) and the obvious inequality $F(x, t) \geq 0$ for $|t| \leq \delta_1$ we deduce

$$g_\lambda(x, t)t \leq \lambda|t|^s + \mu\lambda_{1,q}|t|^q + \theta_0\lambda_{1,p}|t|^p \leq \lambda\frac{\theta}{s}|t|^s \leq \theta\left(\frac{\lambda}{s}|t|^s + F(x, t)\right) \leq \theta G_\lambda(x, t).$$

Consequently,

$$\lambda|t|^\theta \leq g_\lambda(x, t)t \leq \theta G_\lambda(x, t), \quad (x, t) \in \Omega \times [-\delta_1, \delta_1].$$

Recalling that $\hat{\varphi}_\lambda$ is coercive, [17, Theorem 3.6] can be used to achieve (4.4), as desired. Finally, u_3 is nodal by extremality of u_λ^- and u_λ^+ (cf. Lemma 4.2), while standard nonlinear regularity results give $u_3 \in C_0^1(\bar{\Omega})$. \square

Theorems 3.1 and 4.1 together produce the following

THEOREM 4.2. *Let (h₀)–(h₃) be satisfied. Then there exists a real number $\lambda^* > 0$ such that for every $\lambda \in (0, \lambda^*)$ Problem (1.1) admits four nontrivial solutions: one smallest positive, $u_\lambda^+ \in \text{int}(C_+)$; a further positive solution $\bar{u}_\lambda \in \text{int}(C_+)$; one greatest negative, $u_\lambda^- \in -\text{int}(C_+)$; a nodal solution $u_\lambda \in C_0^1(\bar{\Omega}) \cap [u_\lambda^-, u_\lambda^+]$.*

Our next goal is to find a whole sequence of sign-changing solutions that converges to zero. With this aim, the local symmetry condition below involving $f(x, \cdot)$ will be posited whatever $x \in \Omega$.

(h₄) The function $t \mapsto f(x, t)$ is odd on $[-\delta_1, \delta_1]$ for some $\delta_1 \in \mathbb{R}^+$.

THEOREM 4.3. *If (h₀)–(h₄) hold true then to every $\lambda > 0$ corresponds a sequence $\{u_n\}$ of nodal solutions to (1.1) such that $u_n \rightarrow 0$ in $C^1(\overline{\Omega})$.*

PROOF. Pick $\lambda > 0$ and define, provided $(x, t) \in \Omega \times \mathbb{R}$,

$$(4.5) \quad h_\lambda(x, t) := \begin{cases} -\lambda\delta_1^{s-1} - f(x, \delta_1) & \text{when } t < -\delta_1, \\ \lambda|t|^{s-2}t + f(x, t) & \text{when } |t| \leq \delta_1, \\ \lambda\delta_1^{s-1} + f(x, \delta_1) & \text{when } t > \delta_1, \end{cases}$$

as well as $H_\lambda(x, t) := \int_0^t h_\lambda(x, \tau) d\tau$. By (h₄), the associated energy functional

$$\psi_\lambda(u) := \frac{1}{p} \|u\|_p^p + \frac{\mu}{q} \|u\|_q^q - \int_\Omega H_\lambda(x, u(x)) dx, \quad u \in W_0^{1,p}(\Omega),$$

turns out to be even, besides C^1 and coercive.

Let $V \subseteq C_0^1(\overline{\Omega}) \subseteq W_0^{1,p}(\Omega)$ be any finite dimensional space. Bearing in mind that all norms on V are equivalent, we have

$$u \in V, \|u\|_p \leq \rho \quad \Rightarrow \quad |u(x)| \leq \min\{\delta_0, \delta_1\}, \quad x \in \Omega,$$

for some constant $\rho > 0$ depending on V . Hence, $F(x, u(x)) \geq 0$ thanks to (h₃), which entails

$$\psi_\lambda(u) \leq \frac{1}{p} \|u\|_p^p + \frac{\mu}{q} \|u\|_q^q - \frac{\lambda}{s} |u|_s^s \quad \text{whenever } u \in V, \|u\|_p \leq \rho.$$

Since $s < q \leq p$, there exists $\rho' \in (0, \rho)$ such that $\psi_\lambda(u) < 0$ provided $u \in V$, $\|u\|_p = \rho'$. So, Theorem 1 of [12] furnishes a sequence

$$(4.6) \quad \{u_n\} \subseteq K_{\psi_\lambda} \cap \psi_\lambda^{-1}((-\infty, 0))$$

converging to zero in $W_0^{1,p}(\Omega)$. The nonlinear regularity theory ensures a uniform $C^{1,\alpha}(\overline{\Omega})$ -bound on $\{u_n\}$. Through Ascoli-Arzelà's theorem we thus obtain $u_n \rightarrow 0$ in $C^1(\overline{\Omega})$. Consequently, if \bar{u}_λ is the barrier given by Lemma 4.1 then

$$(4.7) \quad \sup_{x \in \Omega} |u_n(x)| < \min \left\{ \delta_1, \sup_{x \in \Omega} \bar{u}_\lambda(x) \right\}$$

for any n large enough. These u_n evidently solve Problem (1.1), because of (4.5)–(4.7). Due to Lemma 4.1 and (4.7) again, they must be nodal. □

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