# New R Matrices Associated with Cyclic Representations of $U_q(A_2^{(2)})$

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#### **Abstract**

New R matrices are constructed as intertwiners of N-dimensional representations of  $U_q(A_2^{(2)})$  at  $q^N = 1$ . Analogous construction for  $U_q(A_1^{(1)})$  reproduces the chiral Potts model.

#### §1. Introduction

Let V be a finite dimensional vector space and  $R(\xi, \eta)$  a linear operator acting on  $V \otimes V$  with parameters  $\xi, \eta \in S$  (to be called spectral parameters). We call  $R(\xi, \eta)$  an R matrix if the Yang-Baxter equation

$$(R(\eta, \lambda) \otimes 1)(1 \otimes R(\xi, \lambda))(R(\xi, \eta) \otimes 1)$$

$$= (1 \otimes R(\xi, \eta))(R(\xi, \lambda)) \otimes 1)(1 \otimes R(\eta, \lambda)) \tag{1.1}$$

holds.

A scheme of constructing R matrices is as follows [11]. Consider a Hopf algebra U. Suppose that a family of representations  $(V, \pi_{\xi})_{\xi \in S}$  of U is given in such a way that

- (i)  $(V \otimes V \otimes V, \pi_{\xi} \otimes \pi_{\eta} \otimes \pi_{\lambda})$  is indecomposable for generic  $\xi$ ,  $\eta$ ,  $\lambda$ , i.e., if  $F \in \text{End}(V \otimes V \otimes V)$  satisfies  $[F, (\pi_{\xi} \otimes \pi_{\eta} \otimes \pi_{\lambda})(g)] = 0$  for any  $g \in U$  then F is a scalar.
- (ii) there exists an intertwiner  $R: V \otimes V \cong V \otimes V$  such that

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$$R(\xi, \eta)(\pi_{\xi} \otimes \pi_{\eta})(g) = (\pi_{\eta} \otimes \pi_{\xi})(g)R(\xi, \eta) \qquad (g \in U)$$

for any  $\xi$ ,  $\eta \in S$ ,

(iii)  $R(\xi, \xi) = 1$ .

Under these conditions  $R(\xi, \eta)$  satisfies (1.1).

For example, take  $U=U_q(A_1^{(1)})$ ,  $V=\mathbb{C}^2$ ,  $S=\mathbb{C}^\times=\mathbb{C}\setminus\{0\}$ . Using  $\xi=x$  as a coordinate of S, we set

$$\pi_{\xi}(e_0) = \begin{pmatrix} 0 & 0 \\ x & 0 \end{pmatrix}, \quad \pi_{\xi}(f_0) = \begin{pmatrix} 0 & x^{-1} \\ 0 & 0 \end{pmatrix}, \quad \pi_{\xi}(t_0) = \begin{pmatrix} q^{-1} & 0 \\ 0 & q \end{pmatrix},$$

$$\pi_{\xi}(e_1) = \begin{pmatrix} 0 & x \\ 0 & 0 \end{pmatrix}, \quad \pi_{\xi}(f_1) = \begin{pmatrix} 0 & 0 \\ x^{-1} & 0 \end{pmatrix}, \quad \pi_{\xi}(t_1) = \begin{pmatrix} q & 0 \\ 0 & q^{-1} \end{pmatrix}.$$

Here  $e_i$ ,  $f_i$ ,  $t_i$  (i=0, 1) are the Chevalley generators. By this choice the conditions (i), (ii), (iii) are satisfied and  $R(\xi, \eta)$  gives the Boltzmann weights of the 6 vertex model.

Bazhanov and Stroganov [5] found that the Boltzmann weights of the chiral Potts model [2], [4] are obtained by a certain algebraic procedure starting from the R matrix of the 6 vertex model. In this paper we reformulate their discovery in the above scheme by taking U to be a certain central extension of  $U_q(A_1^{(1)})$  with q a root of 1. This Hopf algebra is known as the quantum double of a 'Borel' subalgebra of  $U_q(A_1^{(1)})$  [6]. Let us denote it by  $\tilde{U}_q(A_1^{(1)})$ . In [8] we have shown that the Fateev-Zamolodchikov model, the trigonometric limit of the chiral Potts model, is obtained by the scheme (i), (ii), (iii) for  $U = U_q(A_1^{(1)})$ . The central extension enables us to reproduce the whole of the chiral Potts model.

The story goes as follows. If q is a primitive N-th root of 1,  $U_q(A_1) = U_q(\mathfrak{S}I(2,\mathbb{C}))$  admits a 3 parameter family of N-dimensional irreducible representations [12], [13]. (See [7], [9] for general results. See also [1], [3], [8].) It is extended to a 6 parameter family of N-dimensional irreducible representations of  $\widetilde{U}_q(A_1^{(1)})$ . The requirement (ii) restricts the parameters  $\xi$ ,  $\eta$  to be on an algebraic surface  $\mathcal{S}$ . In fact,  $\mathcal{S}$  factorizes essentially into two identical curves:  $\mathcal{S} = \mathcal{C} \times \mathcal{C}$ . Accordingly,  $R(\xi, \eta)$  factorizes into 4 pieces. They are the Boltzmann weights of the chiral Potts model.

The main achievement of this paper is the construction of new R matrices corresponding to the case  $A_2^{(2)}$ , as opposed to  $A_1^{(1)}$  for the chiral Potts model. (For generic q the intertwiner for  $U_q(A_2^{(2)})$  gives the Izergin-Korepin model [10], [11] in the simplest case.) We again start from the N-dimensional irreducible representations of  $U_q(A_1)$ . In this case only those representations which send

the Casimir element of  $U_q(A_1)$  to zero can be extended to the representations of  $\widetilde{U}_q(A_2^{(2)})$ . This restriction effects that the set  $\mathcal{S}$  is only a curve, given explicitly by

$$c^{N}+c^{-N}=\Gamma_{1}(a^{N}+a^{-N})$$

where  $\Gamma_1$  is a modulus. This is essentially the same curve as  $\mathcal{C}$  for the chiral Potts model. The R matrix is given in Theorem 3.4. Unlike the case  $A_1^{(1)}$ , however, we have not found a basis of  $V=\mathbb{C}^N$  for which the matrix elements of  $R(\xi, \eta)$  factorize.

The plan of this paper is as follows. In Section 2 we construct a family of representations of  $\widetilde{U}_q(A_2^{(2)})$ . In Section 3 we solve the equation for the intertwiner and give new R matrices. We show also the indecomposability of  $V \otimes V \otimes V$  to prove the validity of the Yang-Baxter equation. The intertwiner for  $\widetilde{U}_q(A_1^{(1)})$  and the chiral Potts model are discussed in Section 4.

## §2. Algebra $\tilde{U}_o(A_2^{(2)})$ and its Cyclic Representations

Let us first recall the definition of the quantized enveloping algebra of type  $A_2^{(2)}$ . We consider only the case when the deformation parameter q is a root of 1.

Throughout this paper we fix a positive odd integer  $N \ge 3$ , and a primitive N-th root of unity q. We set  $\omega = q^2$ .

Let  $(a_{ij}) = \begin{pmatrix} 2 & -1 \\ -4 & 2 \end{pmatrix}$  be the generalized Cartan matrix of type  $A_2^{(2)}$ , and set  $q_0 = q^4$ ,  $q_1 = q$ . The quantized enveloping algebra  $U_q$  of type  $A_2^{(2)}$  is a Calgebra generated by  $e_i$ ,  $f_i$ ,  $t_i$ ,  $t_i^{-1}$  (i = 0, 1), subject to the following defining relations.

$$t_i t_i^{-1} = t_i^{-1} t_i = 1, \quad t_i t_j = t_j t_i,$$
 (2.1a)

$$t_i e_j t_i^{-1} = q_i^{a_{ij}} e_j$$
,  $t_i f_j t_i^{-1} = q_i^{-a_{ij}} f_j$ ,  $[e_i, f_j] = \delta_{ij} \frac{t_i - t_i^{-1}}{q_i - q_i^{-1}}$ , (2.1b)

$$\sum_{l=0}^{1-a_{ij}} (-1)^l {1-a_{ij} \brack l}_i e_i^l e_j e_i^{1-a_{ij}-l} = 0, \qquad i \neq j,$$
 (2.1c)

$$\sum_{l=0}^{1-a_{ij}} (-1)^l \begin{bmatrix} 1-a_{ij} \\ l \end{bmatrix}_i f_i^l f_j f_i^{1-a_{ij}-l} = 0, \quad i \neq j.$$
 (2.1d)

Here we use the following notations.

$$\begin{bmatrix} m \\ n \end{bmatrix}_{i} = \frac{[m]_{i} \cdots [m-n+1]_{i}}{[n]_{i} \cdots [1]_{i}} , \quad [m]_{i} = \frac{q_{i}^{m} - q_{i}^{-m}}{q_{i} - q_{i}^{-1}} .$$

For our purposes it is necessary to enlarge this algebra by adding 2 central elements  $z_i$  (i=0, 1). Let us denote the enlarged algebra by  $\tilde{U}_q = \tilde{U}_q(A_2^{(2)})$ .

This algebra can be endowed with a structure of a Hopf algebra. We shall need only the comultiplication  $\Delta\colon \widetilde{U}_q \to \widetilde{U}_q \otimes \widetilde{U}_q$  defined by

$$\Delta(e_i) = e_i \otimes 1 + z_i t_i \otimes e_i \tag{2.2a}$$

$$\Delta(f_i) = f_i \otimes t_i^{-1} + z_i^{-1} \otimes f_i , \qquad (2.2b)$$

$$\Delta(t_i) = t_i \otimes t_i \,, \tag{2.2c}$$

$$\Delta(z_i) = z_i \otimes z_i . \tag{2.2d}$$

Note that this comultiplication coincides with that of  $U_q$  if we set  $z_i=1$ . As in the case of  $U_q$  we form tensor products of representations via this comultiplication.

Let us denote by  $U_q(A_1)$  the subalgebra of  $\tilde{U}_q$  generated by  $e_1$ ,  $f_1$ ,  $t_1$ ,  $t_1^{-1}$ . As  $q^N = 1$ , the powers  $e_1^N$ ,  $f_1^N$ ,  $t_1^{\pm N}$  belong to the center of  $U_q(A_1)$ . We consider finite dimensional irreducible representations such that  $e_1^N$  and  $f_1^N$  are non zero constants. We call such representations cyclic representations.

The cyclic representation of  $U_q(A_1)$  is N-dimensional and depend on 3 continuous parameters [12], [13]. This is described as follows. Let V be an N-dimensional vector space over  $\mathbb{C}$ . Choose two linear operators X, Z on V satisfying

$$ZX = \omega XZ$$
,  $X^N = Z^N = 1$ .

Proposition 2.1. An N-dimensional cyclic representation

$$\pi_{a_0a_1x_1}: U_q(A_1) \to \operatorname{End}(V) \qquad (a_0, a_1, x_1 \in \mathbb{C}^{\times})$$

is given as follows.

$$\begin{split} \pi_{a_0a_1x_1}(e_1) &= x_1 X \;, \\ \pi_{a_0a_1x_1}(f_1) &= -\frac{x_1^{-1}}{(q-q^{-1})^2} \Big( \frac{\omega a_1}{a_0} Z + \frac{a_0}{\omega a_1} Z^{-1} - a_0 a_1 - \frac{1}{a_0 a_1} \Big) X^{-1} \;, \\ \pi_{a_0a_1x_1}(t_1) &= \frac{qa_1}{a_0} Z \;. \end{split}$$

The Casimir element

$$qt_1+q^{-1}t_1^{-1}+(q-q^{-1})^2f_1e_1$$

takes the value  $a_0a_1+(a_0a_1)^{-1}$  for this representation.

We will extend this representation to that of  $\tilde{U}_q$ .

**Proposition 2.2.** Let  $\pi_{a_0a_1x_1}$  be as in Proposition 2.1. There exists a representation  $\pi \colon \widetilde{U}_q(A_2^{(2)}) \to \operatorname{End}(V)$  such that  $\pi \mid_{U_q(A_1)} = \pi_{a_0a_1x_1}$  if and only if  $(a_0a_1)^2 = -1$ . In this case  $\pi$  is given by

$$\pi(e_0) = \mp \frac{x_0}{(q^4 - q^{-4})^2} ((qaZ + q^{-1}a^{-1}Z^{-1})X^{-1})^2, \quad \pi(e_1) = x_1X, \quad (2.3a)$$

$$\pi(f_0) = x_0^{-1} X^2$$
,  $\pi(f_1) = -\frac{x_1^{-1}}{(q-q^{-1})^2} (qaZ + q^{-1}a^{-1}Z^{-1})X^{-1}$ , (2.3b)

$$\pi(t_0) = \pm a^{-2}Z^{-2}, \quad \pi(t_1) = aZ,$$
 (2.3c)

$$\pi(z_0) = c_0, \quad \pi(z_1) = c_1,$$
 (2.3d)

with some  $c_0$ ,  $c_1$ ,  $x_0 \in \mathbb{C}^{\times}$ , and  $a = qa_1/a_0$ .

*Proof.* From the requirements  $t_0e_1t_0^{-1}=q^{-4}e_1$ ,  $[t_0, t_1]=0$ , we have

$$\pi(t_0)=bZ^{-2}\,,$$

where  $b \in \mathbb{C}^{\times}$  is to be determined. The conditions  $t_1 f_0 t_1^{-1} = q^4 f_0$ ,  $t_0 f_0 t_0^{-1} = q^{-8} f_0$ ,  $[e_1, f_0] = 0$  fix the form of  $\pi(f_0)$  to be

$$\pi(f_0) = x_0^{-1} X^2$$
,

where  $x_0 \in \mathbb{C}^{\times}$  is some constant. From  $t_1 e_0 t_1^{-1} = q^{-4} e_0$ ,  $t_0 e_0 t_0^{-1} = q^8 e_0$  we know that  $\pi(e_0)$  has the following form

$$\pi(e_0) = \varphi(Z)X^{-2}$$

with some function  $\varphi$ . The condition  $[e_0, f_0] = (t_0 - t_0^{-1})/(q^4 - q^{-4})$  implies that  $\varphi$  must have the form

$$\varphi(Z) = -\frac{x_0}{(q^4 - q^{-4})^2} (q^{-4}bZ^{-2} + q^4b^{-1}Z^2 + c),$$

where  $c \in \mathbb{C}$  is to be determined. Finally the requirement  $[e_0, f_1] = 0$  fixes the constants b and c. Namely we have

$$a_0a_1+a_0^{-1}a_1^{-1}=0$$
 ,  $b=\pm\left(rac{a_0}{qa_1}
ight)^{\!\!2}\!\!, \quad c=\pm(q^2+q^{-2})$  .

The first relation means the Casimir element vanishes. By setting  $a=qa_1/a_0$ , we have (2.3). The Serre relations (2.1c), (2.1d) can be checked by using the vanishing of the Casimir element and the fact that  $\pi(e_0)$  and  $\pi(f_0)$  are proportional to  $\pi(f_1)^2$  and  $\pi(e_1)^2$ , respectively.  $\square$ 

Hereafter we consider only representations (2.3) such that  $\pi(t_0 t_1^2) = 1$ . We denote this representation by  $\pi_{\xi}$  where  $\xi = (a, c_0, c_1, x_0, x_1) \in (\mathbb{C}^{\times})^5$ .

The following will be used in Section 3. Let us denote the representation with  $a=c_0=c_1=1$  and  $x_0=x_1=x$  by  $\pi_x$ .

Lemma 2.3. Assume  $N \neq 3$ . Let  $(V', \pi')$  be a representation of  $U_q(A_2^{(2)})$ , and consider the equations for  $F(x) \in \text{End}(V \otimes V')$ 

$$[(\pi_x \otimes \pi')(f_i), F(x)] = 0 (i = 0, 1),$$
  

$$(\pi_x \otimes \pi')(t_1)F(x) = \omega^m F(x)(\pi_x \otimes \pi')(t_1). (2.4)$$

Then for generic x any solution has the form  $F(x)=Z^m \otimes F'(x)$ , where  $F'(x) \in \text{End}(V')$  satisfies

$$[\pi'(f_i), F'(x)] = 0 (i = 0, 1),$$
  

$$\pi'(t_1)F'(x) = \omega^m F'(x)\pi'(t_1). (2.5)$$

*Proof.* Clearly  $Z^m \otimes F'(x)$  with F'(x) satisfying (2.5) is a solution of (2.4). The coefficients of the linear equations (2.4) are polynomials in x. Therefore it is sufficient to prove the assertion for F(x) which are polynomials in x. In terms of Z and X, the equations (2.4) are

$$[T_i, F(x)] = 0, (i = 1, 2),$$
 (2.6a)

$$(Z \otimes \pi'(t_1))F(x) = \omega^m F(x) (Z \otimes \pi'(t_1)), \qquad (2.6b)$$

where

$$T_1 = X^2 \otimes \pi'(t_1)^2 + x_1 \otimes \pi'(f_0),$$
  

$$T_2 = Y \otimes \pi'(t_1)^{-1} + x_1 \otimes \pi'(f_1),$$

and

$$Y = x\pi_x(f_1) = -(qZ+q^{-1}Z^{-1})X^{-1}/(q-q^{-1})^2$$
.

It follows from (2.6b) that F(x) commutes with  $1 \otimes \pi'(t_1)^N$ . Then using

$$\begin{split} T_1' &= (1 \otimes \pi'(t_1)^{-N}) T_1^{(N+1)/2} \\ &= X \otimes \pi'(t_1) + x \frac{\omega^2}{\omega^2 + 1} X^{-1} \otimes \pi'(t_1^{-1}f_0) + O(x^2) \;, \end{split}$$

we find that F(x) commutes with

$$egin{align} T_3 &= -q^{-1} rac{q-q^{-1}}{q+q^{-1}} (\omega T_1' T_2 - T_2 T_1') \ &= Z^{-1} \otimes 1 + x igg( \omega^3 Z + rac{\omega^2}{\omega^2 + 1} Z^{-1} igg) X^{-2} \otimes \pi'(t_1^{-2} f_0) + O(x^2) \ \end{aligned}$$

and

$$\begin{split} T_4 &= (\omega T_2 T_3 - T_3 T_2) / x (\omega - 1) \\ &= Z^{-1} \otimes \pi'(f_1) + \omega^6 \frac{q + q^{-1}}{(q - q^{-1})^2} Z^2 X^{-3} \otimes \pi'(t_1^{-3} f_0) + O(x) \; . \end{split}$$

Putting x=0, F(0) commutes with  $Z\otimes 1$ ,  $X\otimes \pi'(t_1)$  and

$$Z^{-1} \otimes \pi'(f_1) + \omega^6 \frac{q + q^{-1}}{(q - q^{-1})^2} Z^2 X^{-3} \otimes \pi'(t_1^{-3} f_0) \; .$$

The commutativity with the first two operators and (2.6b) show that F(0) is of the form

$$F(0) = Z^m \otimes F'$$

where  $F' \in \text{End}(V')$  satisfies

$$\pi'(t_1)F'=\omega^m F'\pi'(t_1).$$

In the case  $N \neq 3$ ,  $Z^{-1}$  and  $Z^2X^{-3}$  are linearly independent. From the commutativity with the last operator, it then follows that F' commutes with  $\pi'(f_i)$  (i=0, 1). Therefore (F(x)-F(0))/x satisfies (2.4). By repeating this we have the conclusion.  $\square$ 

**Proposition 2.4.** Assume  $N \neq 3$ . For generic  $x_i$ , if  $F \in \text{End}(V^{\otimes n})$  satisfies

$$\begin{split} &[(\pi_{x_1} \otimes \cdots \otimes \pi_{x_n})(f_i), \, F] = 0 \qquad (i = 0, \, 1) \,, \\ &(\pi_{x_1} \otimes \cdots \otimes \pi_{x_n})(t_1)F = \omega^m F(\pi_{x_1} \otimes \cdots \otimes \pi_{x_n})(t_1) \,, \end{split}$$

then F is a scalar. Moreover, if  $m \equiv 0 \mod N$  then F=0.

*Proof.* Thanks to the Lemma 2.3, the problem reduces to the case n=1. In this case, the equations are

$$[X^2, F] = [Y, F] = 0, \quad ZF = \omega^m FZ.$$

From this, the assertion follows.  $\square$ 

We have also proved the following directly by using computer.

**Lemma 2.5.** Suppose that N=3 and consider the representation  $\pi_{\xi}$  such that  $x_0=x_1=1$ ,  $c_0=c^{-2}$  and  $c_1=c$ . For generic a, c the tensor product  $(V \otimes V \otimes V, \pi_{\xi} \otimes \pi_{\xi} \otimes \pi_{\xi})$  is indecomposable.

## §3. The Intertwiner for $\tilde{U}_q(A_2^{(2)})$

Now we shall solve

$$R(\xi, \tilde{\xi})(\pi_{\xi} \otimes \pi_{\tilde{\xi}})(g) = (\pi_{\tilde{\xi}} \otimes \pi_{\xi})(g)R(\xi, \tilde{\xi}) \qquad (g \in \tilde{U}_{q}(A_{2}^{(2)})). \tag{3.1}$$

As a result we obtain a new R matrix whose spectral parameters live on a curve. Firstly we derive necessary conditions for the existence of a solution. The following will be used frequently.

Lemma 3.1. Let  $\varepsilon$  be a primitive N-th root of unity. If A, B are elements of a  $\mathbb{C}$ -algebra satisfying  $AB = \varepsilon BA$ , then we have  $(A+B)^N = A^N + B^N$ .

For convenience we shall call an expression *invariant* if it remains the same under the exchange of  $\xi$  and  $\tilde{\xi}$ .

**Proposition 3.2.** Let  $a^{2N} \neq -1$  and set  $c = c_1$ ,  $d = c_0 c_1^2$ . For the existence of an intertwiner (3.1), it is necessary that the following are invariants:

$$\Gamma_1 = \frac{c^N + c^{-N}}{a^N + a^{-N}}, \quad \Gamma_2 = x_0^N (a^{2N} - c^{2N} d^{-N}), \quad \Gamma_3 = \frac{1 - a^N c^N}{x_1^N}, \quad (3.2a)$$

$$\Gamma_4 = \frac{c^{2N}d^{-N} + c^{-2N}d^N + 2}{(a^N + a^{-N})^2}, \quad \Gamma_5 = d + d^{-1}.$$
 (3.2b)

For generic values of the parameters, the Jacobian of the map  $G:(a, c, x_0, x_1, d) \mapsto (\Gamma_i)_{1 \le i \le 5}$  is nonzero.

*Proof.* If  $R(\pi_{\xi} \otimes \pi_{\xi})(g) = (\pi_{\xi} \otimes \pi_{\xi})(g)R(g \in \widetilde{U}_{q}(A_{2}^{(2)}))$  and R is invertible, then tr  $(\pi_{\xi} \otimes \pi_{\xi})(g) = \text{tr } (\pi_{\xi} \otimes \pi_{\xi})(g)$ . Apply this to  $g = e_{i}^{N}$ ,  $f_{i}^{N}$ ,  $e_{0}e_{1}^{2}$  and  $f_{0}f_{1}^{2}$ . Using (2.2) and Lemma 3.1, we obtain the following invariants  $r_{\xi}$ .

$$egin{aligned} au_1 &= rac{1 - a^{-2N} \, c_0^N}{x_0^N (a^N + a^{-N})^2}, \quad r_2 &= x_0^N (a^{2N} - c_0^{-N}) \,, \ au_3 &= rac{1 - a^N \, c_1^N}{x_1^N}, \quad r_4 &= rac{x_1^N (a^{-N} - c_1^{-N})}{a^N + a^{-N}}, \ au_5 &= rac{1 - c_0 \, c_1^2}{x_0 \, x_1^2}, \quad r_6 &= rac{x_0 \, x_1^2 (1 - c_0 \, c_1^2)}{c_0 \, c_1^2}. \end{aligned}$$

The  $\Gamma_i$  are obtained by setting  $\Gamma_1=1-r_3r_4$ ,  $\Gamma_2=r_2$ ,  $\Gamma_3=r_3$ ,  $\Gamma_4=1-r_1r_2$  and  $\Gamma_5=r_5r_6+2$ . The Jacobian of the map G is found to be

$$4N^4\frac{1-d^2}{acx_0x_1d^2}\frac{a^N-a^{-N}}{(a^N+a^{-N})^4}\frac{(d^N-1)(c^{2N}+d^N)}{c^Nd^N}\varGamma_2\varGamma_3.$$

This completes the proof.

In view of the above proposition, we must impose some condition on  $\xi$  and  $\tilde{\xi}$  in order to obtain an R matrix depending continuously on them. Here-

after we shall assume

$$c_0 c_1^2 = 1$$

and denote the parameters  $(a, c, x_0, x_1) \in (\mathbb{C}^{\times})^4$  by the same letter  $\xi$ . This choice makes  $\Gamma_5$  trivial and  $\Gamma_4 = (\Gamma_1)^2$ . Now (3.2a) reduces to the following invariants:

$$\Gamma_1 = \frac{c^N + c^{-N}}{a^N + a^{-N}}, \quad \Gamma_2 = x_0^N (a^{2N} - c^{2N}), \quad \Gamma_3 = \frac{1 - a^N c^N}{x_1^N}.$$
 (3.3)

This defines a family of algebraic curves  $C_{\Gamma} \subset \{\xi = (a, c, x_0, x_1) \in (\mathbb{C}^{\times})^4\}$  parametrized by  $\Gamma = (\Gamma_1, \Gamma_2, \Gamma_3)$ .

We shall show that (3.1) has a solution under (3.3). First we prepare a lemma. Let  $\varepsilon$  be a primitive N-th root of unity, and let  $\mathscr{W}_{\varepsilon}$  be the C-algebra generated by z and x satisfying  $zx = \varepsilon xz$ .

**Lemma 3.3.** Let  $\sigma$  be a representation of  $\mathcal{W}_{\epsilon}$  on a vector space V such that  $\sigma(z)^N = a$ ,  $\sigma(x)^N = b$ , a,  $b \in \mathbb{C}$ . Set  $Z = \sigma(z)$  and  $X = \sigma(x)$ . Let  $\alpha$ ,  $\beta$ ,  $\gamma$ ,  $\delta$  be complex numbers satisfying  $\alpha^N a + \beta^N = \gamma^N a + \delta^N$ . Then if we define P(Z) by

$$P(Z) = \sum_{k=0}^{N-1} p_k Z^k, \quad p_k = \prod_{l=1}^k \left( \varepsilon \gamma - \varepsilon^l \alpha \right) \prod_{l=k+1}^{N-1} \left( \varepsilon^l \beta - \delta \right), \tag{3.4}$$

it satisfies

$$P(Z)(\alpha Z + \beta)X = (\gamma Z + \delta)XP(Z). \tag{3.5}$$

Let  $\tilde{P}(Z) = \sum_{k=0}^{N-1} \tilde{p}_k Z^k$  be P(Z) with  $\alpha$  and  $\gamma$ , and  $\beta$  and  $\delta$  being interchanged. Then

$$P(Z)\widetilde{P}(Z) = \rho$$
,

with

$$ho = aNrac{r-lpha}{eta-\delta}rac{lpha^N\delta^N-eta^N\gamma^N}{lpha\delta-eta\gamma}\,.$$

*Proof.* Without loss of generality we may assume that  $ab \neq 0$  and that  $\sigma$  is an N-dimensional irreducible representation. Then  $Z^k X^l$   $(0 \leq k, l \leq N-1)$  are linearly independent. Therefore (3.5) is equivalent to

$$\begin{aligned} p_{k+1}(\varepsilon^{k+1}\beta - \delta) &= p_k(\varepsilon \gamma - \varepsilon^{k+1}\alpha) & (0 \le k \le N - 2), \\ p_0(\beta - \delta) &= p_{N-1}a(\varepsilon \gamma - \alpha). \end{aligned} \tag{3.6}$$

This recursion relation is satisfied by (3.4). Since  $P(Z)\tilde{P}(Z)$  satisfies (3.5) with

 $\alpha = r$  and  $\beta = \delta$ , (3.6) implies that it is proportional to the identity. Therefore, using the formula

$$\sum_{k=0}^{N-1} \frac{y}{\varepsilon^k x - y} = \frac{N y^N}{x^N - y^N}$$

we find  $P(Z)\tilde{P}(Z) = \rho$  with

$$\begin{split} \rho &= p_0 \tilde{p}_0 + a \sum_{k=1}^{N-1} p_k \tilde{p}_{N-k} \\ &= -\frac{a}{(\beta - \delta)^2} \sum_{k=0}^{N-1} \prod_{l=1}^k \left( \varepsilon \gamma - \varepsilon^l \alpha \right) \prod_{l=k+1}^N \left( \varepsilon^l \beta - \delta \right) \prod_{l=1}^{N-k} \left( \varepsilon \alpha - \varepsilon^l \gamma \right) \prod_{l=N-k+1}^N \left( \varepsilon^l \delta - \beta \right) \\ &= a \frac{(\gamma - \alpha)(\alpha^N - \gamma^N)(\beta^N - \delta^N)}{\beta - \delta} \sum_{k=0}^{N-1} \frac{\varepsilon^k}{(\varepsilon^k \alpha - \gamma)(\varepsilon^k \beta - \delta)} \\ &= a N \frac{\gamma - \alpha}{\beta - \delta} \frac{\alpha^N \delta^N - \beta^N \gamma^N}{\alpha \delta - \beta \gamma} . \quad \Box \end{split}$$

**Theorem 3.4.** For generic values of the parameters  $\xi$  and  $\tilde{\xi}$  satisfying (3.3), the equation (3.1) has a unique invertible solution up to a scalar multiple. It is explicitly given by

$$R(\xi, \tilde{\xi}) = P^{1}(C_{1})P^{2}(C_{2})P^{3}(C_{3}) = P^{3}(\tilde{C}_{3})P^{1}(C_{1})P^{2}(C_{2})$$
(3.7)

where

$$\begin{split} C_1 &= X^2 \otimes Z^2 X^{-2} \,, \quad C_2 = Z X^{-1} \otimes X \,, \\ C_3 &= (\pi_{\xi} \otimes \pi_{\xi}) (e_1)^2 (\pi_{\xi} \otimes \pi_{\xi}) (f_0)^{-1} / x_0 \tilde{x}_0 \,, \\ \tilde{C}_3 &= (\pi_{\xi} \otimes \pi_{\xi}) (e_1)^2 (\pi_{\xi} \otimes \pi_{\xi}) (f_0)^{-1} / x_0 \tilde{x}_0 \,, \\ P^i(C) &= \sum_{k=0}^{N-1} p_k^i C^k \qquad (i=1,2,3) \,, \\ p_k^1 &= \prod_{l=1}^k \left( \omega^4 a^2 x_0 - \omega^{4l} \tilde{a}^2 \tilde{x}_0 \right) \prod_{l=k+1}^{N-1} \left( \omega^{4l} c^2 x_0 - \tilde{c}^2 \tilde{x}_0 \right) \,, \\ p_k^2 &= \prod_{l=1}^k \left( \omega \tilde{a} \tilde{c} x_1 - \omega^l a c \tilde{x}_1 \right) \prod_{l=k+1}^{N-1} \left( \omega^l x_1 - \tilde{x}_1 \right) \,, \\ p_k^3 &= \prod_{l=1}^k \left( \omega^2 a c x_0 - \omega^{2l} \tilde{a} \tilde{c} \tilde{x}_0 \right) \prod_{l=k+1}^{N-1} \omega \left( \omega^{2l} \tilde{a} \tilde{c} x_1^2 - a c \tilde{x}_1^2 \right) \,, \end{split}$$

Proof. Set

$$\mathcal{Q}_1=1\!\otimes\! X^2,\quad \mathcal{Q}_2=X\!\otimes\! 1,\quad \mathcal{Q}_3=C_1^{(N+1)/2}=q^{-1}X\!\otimes\! ZX^{-1}\,.$$

Note that

$$C_3 = ((\tilde{a}^2 \tilde{x}_0 C_1 + c^2 x_0) \mathcal{Q}_1)^{-1} ((ac \tilde{x}_1 C_2 + x_1) \mathcal{Q}_2)^2$$
,

and that

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$$A_2^{(2)}$$
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$$C_3^N = \frac{\tilde{a}^N \tilde{c}^N x_1^{2N} - a^N c^N \tilde{x}_1^{2N}}{a^N c^N x_0^N - \tilde{a}^N \tilde{c}^N \tilde{x}_0^N}.$$
 (3.8)

We have the following commutation relations.

$$C_1 \mathcal{Q}_1 = \omega^4 \mathcal{Q}_1 C_1$$
,  $C_2 \mathcal{Q}_2 = \omega \mathcal{Q}_2 C_2$ ,  $C_3 \mathcal{Q}_3 = \omega^2 \mathcal{Q}_3 C_3$ , (3.9a)

$$[C_i, C_i] = [\mathcal{Q}_i, \mathcal{Q}_i] = [C_i, \mathcal{Q}_i] = 0 \qquad (1 \le i \ne j \le 2), \qquad (3.9b)$$

$$[C_3, C_1^{(N+1)/2}C_2] = 0$$
,  $[\Omega_3, C_1] = [\Omega_3, C_2] = 0$ . (3.9c)

Thanks to Lemma 3.4  $P^{i}(C_{i})$  (i=1, 2, 3) satisfies

$$P^{1}(C_{1})(\tilde{a}^{2}\tilde{x}_{0}C_{1}+c^{2}x_{0})\Omega_{1}=(a^{2}x_{0}C_{1}+\tilde{c}^{2}\tilde{x}_{0})\Omega_{1}P^{1}(C_{1}), \qquad (3.10a)$$

$$P^{2}(C_{2})(ac\tilde{x}_{1}C_{2}+x_{1})\Omega_{2}=(\tilde{a}\tilde{c}x_{1}C_{2}+\tilde{x}_{1})\Omega_{2}P^{2}(C_{2}), \qquad (3.10b)$$

$$P^{3}(C_{3})\tilde{a}\tilde{c}(\tilde{x}_{0}C_{3}+\omega x_{1}^{2})\Omega_{3}=ac(x_{0}C_{3}+\omega \tilde{x}_{1}^{2})\Omega_{3}P^{3}(C_{3}). \qquad (3.10c)$$

Now we shall show that  $R=R(\xi, \tilde{\xi})$  given by (3.7) satisfies (3.1) for  $g=t_i$  (i=0,1),  $f_0$  and  $e_1$ . In terms of  $C_i$  and  $Q_i$  (i=1,2), the equations become

$$\begin{split} RC_1^{(N+1)/2}C_2 &= C_1^{(N+1)/2}C_2R \;, \\ R(\tilde{a}^2\tilde{x}_0C_1 + c^2x_0)\mathcal{Q}_1 &= (a^2x_0C_1 + \tilde{c}^2\tilde{x}_0)\mathcal{Q}_1R \;, \\ R(ac\tilde{x}_1C_2 + x_1)\mathcal{Q}_2 &= (\tilde{a}\tilde{c}x_1C_2 + \tilde{x}_1)\mathcal{Q}_2R \;. \end{split}$$

They follow immediately from (3.9), (3.10a), (3.10b).

Next we shall turn to (3.1) for  $g=f_1$ . It is sufficient to check (3.1) for  $g=f_1e_1$ . After some calculations we obtain

$$-(q-q^{-1})^{2}(\pi_{\xi}\otimes\pi_{\xi})(f_{1}e_{1})$$

$$=q(\pi_{\xi}\otimes\pi_{\xi})(t_{1})+q^{-1}(\pi_{\xi}\otimes\pi_{\xi})(t_{1}^{-1})$$

$$+\frac{1}{x_{1}\tilde{x}_{1}}\left(\frac{\tilde{a}}{c}(\tilde{x}_{0}C_{3}+\omega x_{1}^{2})\mathcal{Q}_{3}+\frac{c}{\tilde{a}}(x_{0}C_{3}+\omega^{-1}\tilde{x}_{1}^{2})\mathcal{Q}_{3}^{-1})\right). \tag{3.11}$$

Note that

$$P^{1}(C_{1})P^{2}(C_{2})C_{3} = \tilde{C}_{3}P^{1}(C_{1})P^{2}(C_{2}).$$
 (3.12)

Using (3.9c), (3.10c), (3.11) and (3.12) we obtain

$$R(\pi_{\mathfrak{k}} \otimes \pi_{\widetilde{\mathfrak{k}}})(f_1 e_1) = (\pi_{\widetilde{\mathfrak{k}}} \otimes \pi_{\mathfrak{k}})(f_1 e_1)R$$
.

Finally we shall consider (3.1) for  $g=e_0$ . This equation can be checked directly. In the case  $N \pm 3$ , it can be shown also by the following argument. Let

$$F = R^{-1}((\pi_{\tilde{\epsilon}} \otimes \pi_{\tilde{\epsilon}})(e_0)R - R(\pi_{\tilde{\epsilon}} \otimes \pi_{\tilde{\epsilon}})(e_0)).$$

We can easily show that F satisfies (2.4) with m=-2. From Proposition 2.4, it vanishes. Therefore R satisfies (3.1) for  $g=e_0$ . Clearly R satisfies (3.1) for  $g=z_i$  (i=0, 1). This completes the proof.  $\square$ 

Remark. If we set

$$\kappa = \tilde{a}^{2N}\tilde{x}_0^N + c^{2N}x_0^N = a^{2N}x_0^N + \tilde{c}^{2N}\tilde{x}_0^N$$

then  $\kappa R(\xi, \tilde{\xi})$  is holomorphic on the curve  $\mathcal{C}_{\Gamma}$ .

Remark. When  $\Gamma_1=1$  and  $\Gamma_2=\Gamma_3=0$ ,  $C_{\Gamma}$  degenerates to a rational curve. Letting  $a, c \to 1$  we find that  $x_0/x_1$  is an invariant. The R matrix becomes a polynomial in the single variable  $x=x_0\tilde{x}_1/x_1\tilde{x}_0$ . We call this the trigonometric case.

Proposition 3.5. The obtained R matrices satisfy the following inversion relation

$$R(\xi,\,\tilde{\xi})R(\tilde{\xi},\,\xi) = \rho(\xi,\,\tilde{\xi})$$

where

$$\begin{split} \rho(\tilde{\varepsilon},\,\tilde{\tilde{\varepsilon}}) &= \omega^{-2} N^{\,3} (ac\tilde{a}\tilde{c})^{N-1} \frac{a^2 x_0 - \tilde{a}^2 \tilde{x}_0}{c^2 x_0 - \tilde{c}^2 \tilde{x}_0} \frac{\tilde{a}\tilde{c} x_1 - ac\tilde{x}_1}{x_1 - \tilde{x}_1} \frac{a^N c^N x_0^N + \tilde{a}^N \tilde{c}^N \tilde{x}_0^N}{ac x_0 + \tilde{a}\tilde{c}\tilde{x}_0} \\ &\times \left( \frac{\tilde{a}^N \tilde{c}^N x_1^{2N} - a^N c^N \tilde{x}_1^{2N}}{\tilde{a}\tilde{c} x_1^2 - ac\tilde{x}_1^2} \right)^2 \frac{x_0^N x_1^{2N} - \tilde{x}_0^N \tilde{x}_1^{2N}}{x_0 x_1^2 - \tilde{x}_0 \tilde{x}_1^2} \;. \end{split}$$

*Proof.* Let  $\tilde{P}^i$  be  $P^i$  with  $\mathcal{E}$  and  $\tilde{\mathcal{E}}$  being interchanged. Thanks to Lemma 3.3,  $P^i(C_i)\tilde{P}^i(C_i)$  (i=1, 2, 3) are proportional to the identity. Therefore, noting  $[C_1, C_2] = 0$ , we find

$$\begin{split} R(\xi,\,\tilde{\xi})R(\tilde{\xi},\,\xi) &= P^{1}(C_{1})P^{2}(C_{2})P^{3}(C_{3})\tilde{P}^{3}(C_{3})\tilde{P}^{1}(C_{1})\tilde{P}^{2}(C_{2}) \\ &= P^{1}(C_{1})\tilde{P}^{1}(C_{1})P^{2}(C_{2})\tilde{P}^{2}(C_{2})P^{3}(C_{3})\tilde{P}^{3}(C_{3}) \,. \end{split}$$

Using (3.8) and Lemma 3.3 we obtain the expression for  $\rho(\xi, \tilde{\xi})$ .  $\square$ 

In order to show that the R satisfies the Yang-Baxter equation, it remains to prove the indecomposability of the tensor products of three cyclic representations of  $\tilde{U}_q(A_2^{(2)})$ . Let  $CV = \bigcup_{\Gamma} C_{\Gamma} \times C_{\Gamma} \times C_{\Gamma} \subset \mathbb{C}^{15}$ 

Proposition 3.6. For generic  $(\Gamma, \xi, \eta, \lambda) \in \mathcal{V}$ , if  $F \in \text{End}(V \otimes V \otimes V)$  commutes with  $(\pi_{\xi} \otimes \pi_{\eta} \otimes \pi_{\lambda})(g)$  for any  $g \in \widetilde{U}_{q}$ , then F is a scalar operator.

*Proof.* Since  $C_{\Gamma}$  is irreducible for generic  $\Gamma$ , the variety CV is irreducible. Therefore we can show the assertion by specialization argument. This is already done in Proposition 2.4 and Lemma 2.5.  $\square$ 

*Remark.* From the R matrix of Theorem 3.4, one can get a local Hamiltonian  $\sum_{i} H_{ij+1}$  by a standard procedure. More precisely, set

$$H = \frac{d}{d\varepsilon} \log R(\xi, \, \tilde{\xi})|_{\varepsilon=0} \,,$$

where  $\tilde{\xi} = \tilde{\xi}(\varepsilon) \to \xi$  as  $\varepsilon \to 0$ . As usual let  $X_j$ ,  $Z_j$  be the operators acting as X, Z on the j-th component in the tensor product of  $V^{\otimes L}$ , and similarly for  $H_{jj+1}$ . Up to a scalar multiple and a term proportional to the identity, we have

$$\begin{split} H_{jj+1} &= H_{jj+1}^1 + H_{jj+1}^2 + H_{jj+1}^3 \,, \\ H_{jj+1}^1 &= 2c^{2N} \sum_{k=1}^{N-1} \frac{1}{1 - \omega^{4k}} \left( -\frac{\omega^4 a^2}{c^2} \right)^k C_1^k \,, \\ H_{jj+1}^2 &= (1 + a^N c^N) \sum_{k=1}^{N-1} \frac{1}{1 - \omega^k} (-\omega a c)^k C_2^k \,, \\ H_{jj+1}^3 &= (a^{2N} + c^{2N}) \sum_{k=1}^{N-1} \frac{1}{1 - \omega^{2k}} (-\omega)^k C_3^k \,, \end{split}$$

where

$$C_1 = X_j^2 Z_{j+1}^2 X_{j+1}^{-2}$$
,  $C_2 = Z_j X_j^{-1} X_{j+1}$ ,  
 $C_3 = (X_j + ac Z_j X_{j+1})^2 (a^2 X_j^2 Z_{j+1}^2 + c^2 X_{j+1}^2)^{-1}$ .

## §4. Intertwiners for $\tilde{U}_q(A_1^{(1)})$ and the Chiral Potts Model

In this section  $\tilde{U}_q$  means the algebra  $\tilde{U}_q(A_1^{(1)})$ . It is defined by the same formulas (2.1), (2.2), wherein  $(a_{ij}) = \begin{pmatrix} 2 & -2 \\ -2 & 2 \end{pmatrix}$  and  $q_0 = q_1 = q$ . As in Section 2 we shall consider N-dimensional representations obtained as the extensions of Proposition 2.1 for  $U_q(A_1)$ , the subalgebra generated by  $e_1$ ,  $f_1$  and  $t_1^{\pm 1}$ . We find it more convenient to make the change of variables

$$X = r^{-1}(a_1^2 Z - 1)X'$$
,  $x_1 = \frac{r}{q - q^{-1}}x_1'$ .

Here  $r^N = a_1^{2N} - 1$ , so that  $X'^N = 1$  and  $ZX' = \omega X'Z$ . Dropping primes we thus have  $\pi : \widetilde{U}_q \to \operatorname{End}(V)$   $(V = \mathbb{C}^N)$ , where

$$\begin{split} \pi\left(e_{0}\right) &= x_{0} X^{-1} \frac{a_{0}^{2} Z^{-1} - 1}{q - q^{-1}} \;, \quad \pi(e_{1}) = x_{1} \frac{a_{1}^{2} Z - 1}{q - q^{-1}} X \;, \\ \pi\left(f_{0}\right) &= (a_{0} a_{1} x_{0})^{-1} \frac{a_{1}^{2} Z - 1}{q - q^{-1}} X \;, \quad \pi(f_{1}) = (a_{0} a_{1} x_{1})^{-1} X^{-1} \frac{a_{0}^{2} Z^{-1} - 1}{q - q^{-1}} \;, \end{split}$$

$$\pi(t_0) = \frac{a_0}{qa_1} Z^{-1}, \quad \pi(t_1) = \frac{qa_1}{a_0} Z, \quad \pi(z_2) = c_0, \quad \pi(z_1) = c_1.$$

The six parameters  $\xi = (a_0, a_1, c_0, c_1, x_0, x_1) \in (\mathbb{C}^{\times})^6$  entering  $\pi$  will be exhibited as  $\pi_{\xi}$ .

Given  $\xi$ ,  $\tilde{\xi} \in (\mathbb{C}^{\times})^6$  we now look for an intertwiner  $R: V \otimes V \xrightarrow{\sim} V \otimes V$  such that

$$R(\xi, \tilde{\xi})(\pi_{\xi} \otimes \pi_{\xi})(g) = (\pi_{\xi} \otimes \pi_{\xi})(g)R(\xi, \tilde{\xi}) \qquad (g \in \tilde{U}_{g}). \tag{4.1}$$

By a similar reasoning as given in Section 3, we have

**Proposition 4.1.** Assume  $a_i^{2N} \neq 1$ ,  $(a_0a_1)^2 \neq -1$ ,  $\tilde{a}_i^{2N} \neq 1$ ,  $(\tilde{a}_0\tilde{a}_1)^2 \neq -1$ . For an intertwiner to exist it is necessary that the following quantities are invariants:

$$r_1 = \frac{1 - c_0^N a_0^N a_1^{-N}}{x_0^N (a_0^{2N} - 1)}, \quad r_2 = \frac{1 - c_1^N a_0^{-N} a_1^N}{x_1^N (a_1^{2N} - 1)},$$
 (4.2a)

$$r_3 = \frac{(a_0 a_1 x_0)^N (a_0^{-N} a_1^N - c_0^{-N})}{a_1^{2N} - 1} , \quad r_4 = \frac{(a_0 a_1 x_1)^N (a_0^N a_1^{-N} - c_1^{-N})}{a_0^{2N} - 1} , \quad (4.2b)$$

$$r_5 = \frac{1 - c_0 c_1}{x_0 x_1 ((a_0 a_1)^2 + 1)}, \quad r_6 = \frac{x_0 x_1 (1 - c_0^{-1} c_1^{-1})}{1 + (a_0 a_1)^{-2}}. \tag{4.2c}$$

For generic  $\xi$  the Jacobian of the map  $\xi \mapsto (\gamma_i)$  is nonzero.

Hereafter we shall assume

$$c_0c_1=1$$
.

This makes (4.2c) trivial, and  $r_1r_3=r_2r_4$ . Eqs. (4.2a), (4.2b) then define a surface  $S=S_{\gamma_1\gamma_2\gamma_3}$  written in the coordinates  $(a_0, a_1, c_0, x_0, x_1)$ . As it turns out, S is essentially a product of two curves. To see this, consider the curve  $C_k$  in the coordinates  $(x, y, \mu) \in (\mathbb{C}^{\times})^3$  [4]

$$C_k: x^N + y^N = k(1 + x^N y^N), \quad \mu^N = \frac{k'}{1 - kx^N} = \frac{1 - ky^N}{k'}.$$
 (4.3)

Here the parameter k is a modulus and  $k^2+k'^2=1$ . Set

$$k^2 = -r_1 r_3$$
 ,  $\kappa_0^N = -k/r_1$  ,  $\kappa_1^N = -k/r_2$  .

Then the following gives an algebraic correspondence:

$$\mathcal{S}_{\gamma_1\gamma_2\gamma_3} \longleftrightarrow \mathcal{C}_k \times \mathcal{C}_k$$

$$(a_0, a_1, c_0, x_0, x_1) \mapsto (x, y, \mu, x', y', \mu')$$

$$x_0 = \kappa_0 x'$$
,  $x_1 = \kappa_1 y'$ ,  $a_0^2 = \frac{y}{x'\mu\mu'}$ ,  $a_1^2 = \frac{x\mu\mu'}{y'}$ ,  $c_0 a_0 a_1 = \frac{qx}{y'}$ .

Hereafter we shall use the letter  $r=(x_r, y_r, \mu_r)$  to denote a point on  $C_k$ . To describe the intertwiners, define matrices

$$S_{rs} = \sum_{a=0}^{N-1} \widehat{W}_{rs}(a) (X^{-1} \otimes X)^{a},$$

$$T_{rs} = \sum_{a=0}^{N-1} \overline{W}_{rs}(a) Z^{a}.$$

Here  $r, s \in \mathcal{C}_k$ , and the coefficients  $\widehat{W}_{rs}(a)$  and  $\overline{W}_{rs}(a)$  are defined via the recursion relations

$$\begin{split} \frac{\widehat{W}_{rs}(a)}{\widehat{W}_{rs}(a-1)} &= \frac{\mu_s y_r - \mu_r y_s \omega^{a-1}}{\mu_s x_s - \mu_r x_r \omega^a}, \\ \frac{\overline{W}_{rs}(a)}{\overline{W}_{rs}(a-1)} &= \mu_r \mu_s \frac{x_r \omega - x_s \omega^a}{y_s - y_r \omega^a}. \end{split}$$

Note also that  $W_{rs}(l) = \sum_{a=0}^{N-1} \hat{W}_{rs}(a) \omega^{-al}$  are given by

$$\frac{W_{rs}(a)}{W_{rs}(a-1)} = \frac{\mu_r}{\mu_s} \frac{y_s - x_r \omega^a}{y_r - x_s \omega^a}.$$

These are the Boltzmann weights of the chiral Potts model [4].

**Theorem 4.2.** Consider  $\xi = (r, r')$ ,  $\tilde{\xi}' = (\tilde{r}, \tilde{r}') \in C_k \times C_k$  where  $C_k$  is given by (4.3). Then up to a scalar multiple the intertwiner (4.1) is given by

$$R(\xi,\,\tilde{\xi})=S_{r'\tilde{r}}(T_{r'\tilde{r}'}\otimes T_{r\tilde{r}})S_{r\tilde{r}'}.$$

*Proof.* Set  $K=\mathbb{Q}(k, k')$ . Let  $\mathcal{O}_N(q)$  be the N-th cyclotomic polynomial in q where q is an indeterminate variable. Set  $\widetilde{K}=K[q]/K\mathcal{O}_N(q)$ . Let  $A=\widetilde{K}[x, x^{-1}, y, y^{-1}, \mu, \mu^{-1}]$  be the coordinate ring of  $\mathcal{C}_k$  over  $\widetilde{K}$ , and  $B=A\otimes_{\widetilde{K}}A$  the coordinate ring of  $\mathcal{C}_k\times\mathcal{C}_k$ . We consider a B-algebra  $\mathcal{W}$  generated by Z and X with the defining relations  $ZX=\omega XZ$  ( $\omega=q^2$ ),  $Z^N=X^N=1$ . We may regard

$$\frac{q-q^{-1}}{\kappa_1}(\pi_{\xi}\otimes\pi_{\widetilde{\xi}})(e_1)$$

$$= x_r\mu_r\mu_{r'}ZX\otimes 1 - y_{r'}X\otimes 1 + x_r^{r}\mu_r\mu_{r'}\mu_r^{r}\mu_{r'}Z\otimes ZX - y_r^{r}\mu_r\mu_{r'}Z\otimes X$$

as an element of  $W \otimes_{\tilde{k}} W$ , and (4.1) for  $g = e_1$  as an equation in  $W \otimes_{\tilde{k}} W$ .

This equation is shown by using the following and similar identities in  $\mathcal{W}$  (see Lemma 3.3).

$$S_{rs}(y_s(X^{-1} \otimes X) - \omega x_r)(Z \otimes 1) = \frac{\mu_s}{\mu_r}(y_r(X^{-1} \otimes X) - \omega x_s)(Z \otimes 1)S_{rs},$$

$$T_{rs}(\mu_r \mu_s x_s Z - y_s)X = (\mu_r \mu_s x_r Z - y_s)XT_{rs}.$$

In order to prove the case  $g=e_0$  we can use the following K-linear antiinvolution \* of  $\mathcal{W}$ :

$$x_r^* = y_r$$
,  $\mu_r^* = \mu_r^{-1}$ ,  $x_s^* = y_s$ ,  $\mu_s^* = \mu_s^{-1}$ ,  $Z^* = Z^{-1}$ ,  $X^* = X^{-1}$ ,  $q^* = q^{-1}$ .

We have

$$\left(\frac{q-q^{-1}}{\kappa_1}(\pi_{\xi}\otimes\pi_{\widetilde{\xi}})(e_1)\right)^* = \frac{q-q^{-1}}{\kappa_0}(\pi_{\xi}\otimes\pi_{\widetilde{\xi}})(e_0).$$

We also have

$$S_{rs}((y_s(X^{-1} \otimes X) - \omega x_r)(Z \otimes 1))^* = \left(\frac{\mu_s}{\mu_r}(y_r(X^{-1} \otimes X) - \omega x_s)(Z \otimes 1)\right)^* S_{rs}, \quad (4.4a)$$

$$T_{rs}((\mu_r \mu_s x_s Z - y_r)X)^* = ((\mu_r \mu_s x_r Z - y_s)X)^* T_{rs}. \quad (4.4b)$$

Therefore we obtain (4.1) for  $g=e_0$ .

For  $g=f_0$ ,  $f_1$ , we use another anti-involution  $^{\wedge}$  of  $\mathcal{W}$ :

$$\hat{x}_r = \frac{1}{qx_r}, \quad \hat{y}_r = \frac{q}{y_r}, \quad \hat{\mu}_r = -\frac{qx_r\mu_r}{y_r},$$
 $\hat{x}_s = \frac{1}{qx_s}, \quad \hat{y}_s = \frac{q}{y_s}, \quad \hat{\mu}_s = -\frac{qx_s\mu_s}{y_s},$ 
 $\hat{Z} = Z, \quad \hat{X} = X^{-1}, \quad \hat{q} = q.$ 

Then we have

$$egin{aligned} &(\kappa_0 c_0 ilde{c}_0 (\pi_{m{\xi}} \otimes \pi_{\widetilde{m{\xi}}})(f_0))^\wedge = -rac{q^2}{\kappa_0 c_0 ilde{c}_0} (\pi_{m{\xi}} \otimes \pi_{\widetilde{m{\xi}}})(t_0^{-1} e_0) \ , \ &\left(rac{\kappa_1}{c_0 ilde{c}_0} (\pi_{m{\xi}} \otimes \pi_{\widetilde{m{\xi}}})(f_1)
ight)^\wedge = -rac{c_0 ilde{c}_0}{\kappa_1} (\pi_{m{\xi}} \otimes \pi_{\widetilde{m{\xi}}})(t_1^{-1} e_1) \ . \end{aligned}$$

The identities (4.4) with \* replaced by  $^{\wedge}$  are also valid. Therefore we obtain (4.1) for  $g=f_0, f_1$ .  $\square$ 

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#### References

- [1] Arnaudon, D., Periodic and flat irreducible representations of SU(3)<sub>q</sub>, Comm. Math. Phys., 134 (1990), 523-537.
- [2] Au-Yang, H., McCoy, B.M., Perk, J.H.H., Tang, S. and Yan, M.-L., Commuting transfer matrices in the chiral Potts models: Solutions of star-triangle equations with genus >1, Phys. Lett. A., 123 (1987), 219–223.
- [3] Bazhanov, V.V. and Kashaev, R.M., Cyclic L operators related with 3-state R-matrix, Comm. Math. Phys., 136 (1990), 607-624.
- [4] Baxter, R.J., Perk, J.H.H. and Au-Yang, H., New solutions of the star-triangle relations for the chiral Potts model, *Phys. Lett. A.*, 128 (1988), 138-142.
- [5] Bazhanov, V.V. and Stroganov, Yu. G., Chiral Potts models as a descendant of the six-vertex models, J. Stat. Phys., 51 (1990), 799-817.
- [6] Drinfeld, V.G., Quantum groups, Proc. ICM Berkeley (1987), 798-820.
- [7] Date, E., Jimbo, M., Miki, K. and Miwa, T., Cyclic representations of  $U_q(\widehat{\mathfrak{sl}}(n+1,\mathbb{C}))$  at  $q^M=1$ , Publ. RIMS Kyoto Univ., 27 (1991), 347-366.
- [8] ——, R-matrix for cyclic representations of  $U_q(\widehat{\mathfrak{sl}}(3,\mathbb{C}))$  at  $q^3=1$ , Phys. Lett. A. 148 (1990), 45–49.
- [9] De Concini, C. and Kac, V.G., Representations of quantum groups at roots of 1, in 'Operator Algebras, Unitary Representation, Enveloping algebras and Invariant Theory', Actes du Colloque en l'honneur de Jacques Dixmier, eds. A. Connes, M. Duflo, A. Joseph and R. Rentschler, *Prog in Math.*, 92 Birkhäuser, (1990), 471-506.
- [10] Izergin, A.G. and Korepin, V.E., The inverse scattering method approach to the quantum Shabat-Mikhailov model, Comm. Math. Phys., 79 (1981), 303-316.
- [11] Jimbo, M., Quantum R matrix for the generalized Toda system, Comm. Math. Phys., 102 (1986), 537-547.
- [12] Roche, P. and Arnaudon, D., Irreducible representations of the quantum analogue of SU(2), Lett. Math. Phys., 17 (1989), 295-300.
- [13] Sklyanin, E.K., Some algebraic structures connected with the Yang-Baxter equation. Representations of quantum algebras, Funct. Anal. Appl., 17 (1984), 273–284.