

Asymptotic Behavior of Blowup Solutions of a Parabolic Equation with the p -Laplacian

By

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Abstract

We consider the blowup problem for $u_t = \Delta_p u + |u|^{p-2}u$ ($x \in \Omega$, $t > 0$) under the Dirichlet boundary condition and $p > 2$. We derive sufficient conditions on blowing up of solutions. In particular, it is shown that every non-negative and non-zero solution blows up in a finite time if the domain Ω is large enough. Moreover, we show that every blowup solution behaves asymptotically like a self-similar solution near the blowup time. The Rayleigh type quotient introduced in Lemma A plays an important role throughout this paper.

§1. Introduction and Results

In this paper we mainly consider the blowup problem for the following initial boundary value problem:

$$\begin{cases} u_t = \Delta_p u + |u|^{q-2}u, & x \in \Omega, \quad t > 0, \\ u(x, t) = 0, & x \in \partial\Omega, \quad t \geq 0, \\ u(x, 0) = u_0(x), & x \in \Omega, \end{cases} \quad (1.1)$$

where $p, q > 2$, $\Delta_p u = \operatorname{div}(|\nabla u|^{p-2} \nabla u)$ and Ω is a bounded domain in \mathbb{R}^N with smooth boundary $\partial\Omega$. Especially, we here study the case when $p = q$.

As for the existence and non-existence of global solutions of (1.1), the following results are well known (see [14], [9], [5], [11]):

- (i) When $p > q$, (1.1) has a global solution for any $u_0 \in W_0^{1,p}$.
- (ii) When $p < q$, for sufficiently small initial function $u_0 \in W_0^{1,p}$, (1.1) has a global solution, and if u_0 is large enough, the solution blows up in a finite time.

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(iii) When $p=q$, put $\lambda_1 = \inf \{ \|\nabla u\|_p / \|u\|_p : u \in W_0^{1,p} \setminus \{0\} \}$. If $\lambda_1 \geq 1$, (1.1) has a global solution for any $u_0 \in W_0^{1,p}$.

Here, $W_0^{1,p} \equiv W_0^{1,p}(\Omega)$ denotes the usual Sobolev space with the norm $\|u\|_{W_0^{1,p}} = \|\nabla u\|_p$, and $\|\cdot\|_p$ denotes the $L^p(\Omega)$ norm.

From the above results, we see that the case $p=q$ is critical for the existence of blowup solutions of (1.1). For the critical exponents of other equations and their role, we refer to the survey paper by Levine [8]. Here, we should note that little is known about the case when $p=q$ and $\lambda_1 < 1$. So, in what follows, we study (1.1) with the case when $p=q > 2$, that is, we consider the following problem:

$$\begin{cases} u_t = \Delta_p u + |u|^{p-2}u, & x \in \Omega, \quad t > 0, \\ u(x, t) = 0, & x \in \partial\Omega, \quad t \geq 0, \\ u(x, 0) = u_0(x), & x \in \Omega, \end{cases} \tag{P}$$

Our first purpose in this paper is to derive sufficient conditions on blowing up of solutions of (P) (Theorems B and C). The second purpose is to study the asymptotic behavior of solutions of (P). Here, we note that we consider not only the asymptotic behavior of blowup solutions but also that of global solutions. In both cases, we show that each solution of (P) behaves asymptotically like a self-similar solution of (P). First, we derive blowup rate and decay rate of solutions of (P) for each case (Theorem D). Next, we investigate the asymptotic profile of both blowup and global solutions of (P) near the maximal existence time (Theorem E). These results for the case $p > 2$ in (P) may be regarded as a natural extension of the linear case $p=2$ in (P).

To be more precise, we here recall the local existence results for (P). The local existence of strong solutions of (P) is already studied by many authors (see [5], [7], [10], [12]). Here, a function $u(x, t)$ is said to be a strong solution of (P) in $[0, T]$ if (i) $u \in C([0, T]; W_0^{1,p}(\Omega))$, (ii) $u_t, \Delta_p u$ and $|u|^{p-2}u \in L^2(0, T; L^2(\Omega))$, and (iii) u satisfies (P). Assume that $p > 2$, and $2(p-1) \leq Np/(N-p)$ if $p < N$. Then, for any $u_0 \in W_0^{1,p}$, there exists a positive number T such that (P) has a strong solution in $[0, T]$. Moreover, let T^* be the maximal existence time of the strong solution $u(t)$ of (P). Then, if $T^* < \infty$, it follows together with (1.6) below that

$$\lim_{t \rightarrow T^*} \|u(t)\|_2 = \lim_{t \rightarrow T^*} \|\nabla u(t)\|_p = \infty.$$

Furthermore, if we put $E(u) = \|\nabla u\|_p^p - \|u\|_p^p$, we have

$$\partial_t \|u(t)\|_2^2 = -2E(u(t)) \quad \text{a.e. in} \quad [0, T^*), \tag{1.2}$$

$$\partial_t E(u(t)) = -p \|u_t(t)\|_2^2 \quad \text{a.e. in} \quad [0, T^*), \quad (1.3)$$

We note that $E(\lambda u) = \lambda^p E(u)$ holds for any $\lambda > 0$ and $u \in W_0^{1,p}$, which is a special feature in the critical case. Our main idea in this paper is to introduce the Rayleigh type quotient $E(u) / \|u\|_2^2$. The following lemma is important in this paper.

Lemma A. *Assume that $u_0 \in W_0^{1,p}$, and let $u(t)$ be a strong solution of (P) in $[0, T^*)$. Then, we have*

$$\partial_t \frac{E(u(t))}{\|u(t)\|_2^2} \leq 0 \quad \text{a.e. in} \quad [0, T^*).$$

Lemma A follows immediately from (1.2) and (1.3), but it plays an essential role in the proofs of the following theorems. We should mention that a similar result to Lemma A is obtained by Berryman and Holland [1] for the fast diffusion $(u^{q-1})_t = \Delta u$ with $q > 2$. In [1] they study the asymptotic behavior of finite time extinction solutions of it.

First, we derive two sufficient conditions that the solution of (P) blows up in a finite time.

Theorem B. *Let $p > 2$ and $\lambda_1 < 1$. Assume that $u_0 \in W_0^{1,p}$ satisfies $E(u_0) < 0$. Then, the strong solution of (P) blows up in a finite time.*

Theorem C. *Let $p > 2$ and $\lambda_1 < 1$. Assume that $u_0 \in W_0^{1,p} \setminus \{0\}$ is non-negative in Ω . Then, the strong solution of (P) blows up in a finite time.*

Here, we recall that $\lambda_1 = \inf \{ \|\nabla u\|_2^2 / \|u\|_2^2 : u \in W_0^{1,p} \setminus \{0\} \}$, and if $\lambda_1 \geq 1$, every strong solution of (P) exists globally in time. Theorems B and C supplement the known results by many authors concerning the existence and non-existence of global solutions of (1.1) by giving information about the case of $p = q > 2$. In [2] Galaktionov showed a similar result to Theorem C for $u_t = \Delta u^m + u^m$ with $m > 1$ by using the so-called Kaplan method [6]. We should mention that this method is not applicable to our problem (P), and our proof of Theorem C is quite different from that of [2].

Next, we consider the asymptotic behavior of strong solutions of (P). We begin with deriving blowup rate and decay rate of strong solutions of (P).

Theorem D. *Assume $p > 2$ and $u_0 \in W_0^{1,p} \setminus \{0\}$. Let T^* be the maximal existence time of the strong solution $u(t)$ of (P). Put $\gamma_* = \lim_{t \rightarrow T^*} [E(u(t)) / \|u(t)\|_2^2]$.*

(i) If $T^* < \infty$, we have $\gamma_* < 0$ and

$$\lim_{t \rightarrow T^*} [-\gamma_* (p-2) (T^* - t)]^{\frac{1}{p-2}} \|u(t)\|_2 = 1. \tag{1.4}$$

(ii) If $T^* = \infty$ and $\gamma_* > 0$, we have

$$\lim_{t \rightarrow \infty} [\gamma_* (p-2) t]^{\frac{1}{p-2}} \|u(t)\|_2 = 1. \tag{1.5}$$

Remark 1.1. Put $\gamma_1 = \inf \{E(u) / \|u\|_2^p : u \in W_0^{1,p} \setminus \{0\}\}$. Then, we see that $\gamma_1 > -\infty$. In fact, by the Gagliardo-Nirenberg and the Young inequalities, there exist positive constants $\alpha \in (0, p)$, C_1 and C_2 such that

$$\|u\|_2^p \leq C_1 \|u\|_2^{p-\alpha} \|\nabla u\|_2^\alpha \leq \frac{1}{2} \|\nabla u\|_2^p + C_2 \|u\|_2^p, \quad u \in W_0^{1,p},$$

from which we have

$$\|\nabla u\|_2^p \leq 2E(u) + 2C_2 \|u\|_2^p, \quad u \in W_0^{1,p}, \tag{1.6}$$

and we have $\gamma_1 \geq -C_2$. So, it follows from Lemma A and this fact that the limit $\gamma_* = \lim_{t \rightarrow T^*} [E(u(t)) / \|u(t)\|_2^p]$ exists and $\gamma_* \geq \gamma_1$ holds for any strong solution $u(t)$ of (P). We also note that from Theorem B, if $T^* = \infty$, we have $\gamma_* \geq 0$. Moreover, we see that $\gamma_1 < 0$ [resp. $\gamma_1 = 0$, $\gamma_1 > 0$] if and only if $\lambda_1 < 1$ [resp. $\lambda_1 = 1$, $\lambda_1 > 1$].

Remark 1.2. A function $u(x, t) = v(t)w(x)$ of variable separation type is called a self-similar solution of (P) with $u_0(x) = v(0)w(x)$ if v and $w \in W_0^{1,p}$ satisfy

$$v_t = -\gamma|v|^{p-2}v \quad \text{in } \mathbb{R}, \tag{1.7}$$

$$-\Delta_p w - |w|^{p-2}w = \gamma w \quad \text{in } \mathcal{D}'(\Omega) \tag{1.8}$$

for some $\gamma \in \mathbb{R}$. From Theorem D, we see that the blowup rate and the decay rate of general strong solutions of (P) in Theorem D are the same as those of the self-similar solutions of (P).

Remark 1.3. In the case when $p < q$ in (1.1), the decay rate of small global solutions of (1.1) is given by H. Ishii [5]. However, it seems that in [5] there are no results for blowup rate of solutions of (1.1) when $2 < p < q$. For the semi-linear case $p=2 < q$, see Giga and Kohn [3] and references therein.

The following theorem states that the asymptotic profiles of solutions of (P) are given by the solutions of (1.8).

Theorem E. *Assume that $p > 2$ and $u_0 \in W_0^{1,p} \setminus \{0\}$. Let $T^* \in (0, \infty]$ be the maximal existence time of the strong solution $u(t)$ of (P). Then, for any sequence $\{t_j\}$ satisfying $t_j \rightarrow T^*$, there exist a subsequence $\{t_{j'}\}$ of $\{t_j\}$ and $w \in W_0^{1,p}$ such that*

$$\frac{u(t_{j'})}{\|u(t_{j'})\|_2} \rightarrow w \quad \text{in } W_0^{1,p}, \quad (1.9)$$

$$-\Delta_p w - |w|^{p-2} w = \gamma_* w \quad \text{in } \mathcal{D}'(\Omega), \quad \|w\|_2 = 1, \quad (1.10)$$

where $\gamma_* = \lim_{t \rightarrow T^*} [E(u(t)) / \|u(t)\|_2^p]$.

Remark 1.4. It is natural to ask in Theorem E whether the limit $u(t) / \|u(t)\|_2$ exists or not in $W_0^{1,p}$ as $t \rightarrow T^*$. At the present, we do not know the answer, even if the solution $u(t)$ of (P) is non-negative. Of course, if non-negative solution $w \in W_0^{1,p}$ of (1.10) is unique, then it follows immediately from Theorem E that $u(t) / \|u(t)\|_2 \rightarrow w$ in $W_0^{1,p}$ as $t \rightarrow T^*$ for any non-negative and non-zero solution $u(t)$ of (P). However, as we show in Section 3 for the case $N = 1$, non-negative solution of (1.10) is not unique in general.

The plan of this paper is as follows. In Section 2, we give the proofs of Lemma A and Theorems B, C, D and E. Lemma A will play an important role throughout this paper. Theorems B and D (ii) follow immediately from (1.2) and Lemma A. In order to prove Theorems D (i) and E, we use the rescaling arguments together with Lemma A. Theorem C is proved by contradiction, using Theorem E. In Section 3, we discuss the uniqueness and non-uniqueness of non-negative solutions of (1.10) for the case $N = 1$.

§2. Proofs of Theorems

In this section, we give the proofs of Lemma A and Theorems B, C, D and E. First, we give the proof of Lemma A.

Proof of Lemma A. From (1.2) and (1.3), we have

$$\partial_t \frac{E(u(t))}{\|u(t)\|_2^p} = \frac{\|u(t)\|_2^p \partial_t E(u(t)) - E(u(t)) \partial_t \|u(t)\|_2^p}{\|u(t)\|_2^{2p}}$$

$$\begin{aligned}
 &= \frac{-p\|u(t)\|_2^2\|u_t(t)\|_2^2 + (p/4)\|\partial_t\|u(t)\|_2^2\|u(t)\|_2^{-2}\|\partial_t\|u(t)\|_2^2}{\|u(t)\|_2^{2p}} \\
 &= \frac{p\{(\|\partial_t\|u(t)\|_2)^2 - 4\|u(t)\|_2^2\|u_t(t)\|_2^2\}}{4\|u(t)\|_2^{2+2}}
 \end{aligned}$$

a.e. in $[0, T^*)$. By the Cauchy-Schwarz inequality, we obtain Lemma A. \square

Next, we prove Theorems B and D (ii), using (1.2) and Lemma A only.

Proof of Theorem B. By Lemma A, we have

$$\frac{E(u(t))}{\|u(t)\|_2^2} \leq \frac{E(u_0)}{\|u_0\|_2^2}, \quad t \in [0, T^*).$$

Put $c_0 = -E(u_0)/\|u_0\|_2^2$. Then, from (1.2) and our assumption $E(u_0) < 0$, we have $c_0 > 0$ and

$$\|\partial_t\|u(t)\|_2^2 = -2E(u(t)) \geq 2c_0\|u(t)\|_2^2, \quad t \in [0, T^*). \tag{2.1}$$

Since we consider the case $p > 2$, it follows from (2.1) that $T^* < \infty$. \square

Proof of Theorem D (ii). From Lemma A, for any $\varepsilon > 0$ there exists a $T_\varepsilon > 0$ such that

$$\gamma_* \leq \frac{E(u(t))}{\|u(t)\|_2^2} \leq \gamma_* + \varepsilon, \quad t \in [T_\varepsilon, \infty). \tag{2.2}$$

By (1.2) and (2.2), we have

$$-2(\gamma_* + \varepsilon)\|u(t)\|_2^2 \leq \|\partial_t\|u(t)\|_2^2 \leq -2\gamma_*\|u(t)\|_2^2, \quad t \in [T_\varepsilon, \infty). \tag{2.3}$$

From (2.3), we get

$$\begin{aligned}
 &[\|u(T_\varepsilon)\|_2^{-(p-2)} + (\gamma_* + \varepsilon)(p-2)(t - T_\varepsilon)]^{-\frac{2}{p-2}} \\
 &\leq \|u(t)\|_2^2 \leq [\|u(T_\varepsilon)\|_2^{-(p-2)} + \gamma_*(p-2)(t - T_\varepsilon)]^{-\frac{2}{p-2}}, \quad t \in [T_\varepsilon, \infty),
 \end{aligned}$$

from which we have

$$[\gamma_*/(\gamma_* + \varepsilon)]^{\frac{1}{p-2}} \leq \liminf_{t \rightarrow \infty} [\gamma_*(p-2)t]^{\frac{1}{p-2}}\|u(t)\|_2$$

$$\leq \limsup_{t \rightarrow \infty} [\gamma_* (p-2)t]^{\frac{1}{p-2}} \|u(t)\|_2 \leq 1.$$

Since $\varepsilon > 0$ is arbitrary, we obtain (1.5). \square

Remark 2.1. When $T^* = \infty$, it follows from Theorem B that $\gamma_* \geq 0$. Conversely, if $\gamma_* \geq 0$, we have $T^* = \infty$. In fact, suppose that $\gamma_* \geq 0$. Then, it follows from the definition of γ_* that $E(u(t)) \geq 0$ for any $t \in [0, T^*)$. From (1.2), we see that $\|u(t)\|_2 \leq \|u_0\|_2$ for any $t \in [0, T^*)$, from which we have $T^* = \infty$. In the case when $\gamma_* = 0$, from the proof of Theorem D (ii), we see that there exists a positive constant C_1 such that $\|u(t)\|_2 \geq C_1(1+t)^{-1/(p-2)}$ for any $t \in [0, \infty)$.

Next, we prove Theorems D (i) and E, using the rescaling arguments.

Proof of Theorem D (i). First, from Remark 2.1, we see that $\gamma_* < 0$. In order to show (1.4), we introduce the rescaled function $\bar{u}(x, \tau)$ defined by

$$\begin{aligned} \bar{u}(x, \tau) &= (T^* - t)^{\frac{1}{p-2}} u(x, t), \quad t = T^* - e^{-\tau}, \\ &= \exp\left(-\frac{\tau}{p-2}\right) u(x, T^* - e^{-\tau}). \end{aligned}$$

Then, $\bar{u}(x, \tau)$ satisfies

$$\bar{u}_\tau = \Delta_p \bar{u} + |\bar{u}|^{p-2} \bar{u} - \frac{1}{p-2} \bar{u}, \quad \tau \in (-\log T^*, \infty). \quad (2.4)$$

Multiplying (2.4) by $\bar{u}(x, \tau)$ and integrating over Ω , we have

$$\partial_\tau \|\bar{u}(\tau)\|_2^2 = -2E(\bar{u}(\tau)) - \frac{2}{p-2} \|\bar{u}(\tau)\|_2^2. \quad (2.5)$$

Since we have $\lim_{\tau \rightarrow \infty} [E(\bar{u}(\tau)) / \|\bar{u}(\tau)\|_2^2] = \lim_{t \rightarrow T^*} [E(u(t)) / \|u(t)\|_2^2] = \gamma_*$, for any $\varepsilon > 0$ there exists $T_\varepsilon > 0$ such that

$$\gamma_* \leq \frac{E(\bar{u}(\tau))}{\|\bar{u}(\tau)\|_2^2} \leq \gamma_* + \varepsilon, \quad \tau \in [T_\varepsilon, \infty).$$

From (2.5), we have

$$f_\varepsilon(\|\bar{u}(\tau)\|_2^2) \leq \partial_\tau \|\bar{u}(\tau)\|_2^2 \leq f_0(\|\bar{u}(\tau)\|_2^2), \quad \tau \in [T_\varepsilon, \infty). \quad (2.6)$$

Here we put $f_\delta(s) = -2(\gamma_* + \delta)s^{p/2} - (2/(p-2))s$ for $\delta = 0$ and ε . To conclude the proof, we have only to show that

$$A_0 \leq \| \bar{u}(\tau) \|_2^2 \leq A_\varepsilon, \quad \tau \in [T_\varepsilon, \infty), \tag{2.7}$$

where $A_\delta = [- (\gamma_* + \delta) (p - 2)]^{-2/(p-2)}$ and $f_\delta(A_\delta) = 0$. In fact, since $\varepsilon > 0$ is arbitrary, (1.4) follows from (2.7) and the definition of $\bar{u}(x, \tau)$. We prove (2.7) by contradiction. First, suppose that there exists $\tau_0 \in [T_\varepsilon, \infty)$ such that $\| \bar{u}(\tau_0) \|_2^2 < A_0$. Then, from the second inequality of (2.6), we see that there exists a positive constant C_0 such that $\| \bar{u}(\tau) \|_2^2 \leq C_0 \exp \{ -2\tau / (p - 2) \}$ for any $\tau \geq \tau_0$. Since $\| \bar{u}(\tau) \|_2^2 = \exp \{ -2\tau / (p - 2) \} \| u(T^* - e^{-\tau}) \|_2^2$, we have $\| u(T^* - e^{-\tau}) \|_2^2 \leq C_0$ for any $\tau \geq \tau_0$. However, this contradicts the fact that $\lim_{t \rightarrow T^*} \| u(t) \|_2 = \infty$. Thus, we obtain the first inequality of (2.7). Next, suppose that there exists $\tau_1 \in [T_\varepsilon, \infty)$ such that $\| \bar{u}(\tau_1) \|_2^2 > A_\varepsilon$. From the first inequality of (2.6), we see that there exists $T_1 \in (\tau_1, \infty)$ such that $\lim_{\tau \rightarrow T_1} \| \bar{u}(\tau) \|_2^2 = \infty$. However, this contradicts the fact that $\bar{u}(\tau)$ exists for all $\tau \in (-\log T^*, \infty)$. Thus, we obtain the second inequality of (2.7), and the proof of Theorem D (i) is completed. \square

Proof of Theorem E. For the solution $u(x, t)$ of (P) in $[0, T^*)$, we define the rescaled function $\tilde{u}(x, \tau)$ as follows:

$$\tilde{u}(x, \tau) = \frac{u(x, t)}{\| u(t) \|_2}, \quad \tau(t) = \int_0^t \| u(s) \|_2^{p-2} ds.$$

Then, from Theorem D and Remark 2.1, we see that $\tau(T^*) = \infty$ and $\tilde{u}(\tau)$ satisfies

$$\tilde{u}_\tau = \Delta_p \tilde{u} + | \tilde{u} |^{p-2} \tilde{u} + E(\tilde{u}) \tilde{u}, \quad \tau \in [0, \infty). \tag{2.8}$$

First, we show that for any sequence $\{ \tau_j \}$ satisfying $\tau_j \rightarrow \infty$ there exist a subsequence $\{ \tau_{j'} \}$ of $\{ \tau_j \}$ and $w \in W_0^{1,p}$ such that

$$\tilde{u}(\tau_{j'}) \rightarrow w \quad \text{in } L^2(\Omega), \tag{2.9}$$

and w satisfies (1.10). Since $\| \tilde{u}(\tau) \|_2 = 1$ for $\tau \in [0, \infty)$, multiplying (2.8) by $\tilde{u}_\tau(x, \tau)$ and integrating over Ω , we have

$$\partial_\tau E(\tilde{u}(\tau)) = -p \| \tilde{u}_\tau(\tau) \|_2^2, \quad \tau \in [0, \infty). \tag{2.10}$$

From (2.10) and

$$\lim_{\tau \rightarrow \infty} E(\tilde{u}(\tau)) = \lim_{t \rightarrow T^*} \frac{E(u(t))}{\| u(t) \|_2^2} = \gamma_*, \tag{2.11}$$

we have $\int_0^\infty \|\tilde{u}_\tau(\tau)\|_2^2 d\tau < \infty$. Here, following the proof of Lemma 4 of Otani [11], we set $\tilde{u}_j(\sigma) = \tilde{u}(\tau_j + \sigma)$ for $0 \leq \sigma \leq 1$. Then, we see that $\{\tilde{u}_j\} \subset C([0,1]; W_0^{1,p}(\Omega))$, and \tilde{u}_j satisfies

$$\partial_\sigma \tilde{u}_j = \Delta_p \tilde{u}_j + |\tilde{u}_j|^{p-2} \tilde{u}_j + E(\tilde{u}_j) \tilde{u}_j, \quad \sigma \in [0,1]. \tag{2.12}$$

It follows from $\int_0^\infty \|\tilde{u}_\tau(\tau)\|_2^2 d\tau < \infty$ that

$$\|\partial_\sigma \tilde{u}_j\|_{L^2(0,1;L^2(\Omega))} \rightarrow 0. \tag{2.13}$$

Moreover, since $\|\tilde{u}_j(\sigma)\|_2 = 1$ for $\sigma \in [0, 1]$, it follows from (1.6) and (2.10) that

$$\sup_j \|\tilde{u}_j\|_{L^\infty(0,1;W_0^{1,p}(\Omega))} < \infty. \tag{2.14}$$

By (2.11) - (2.14), the monotonicity of $-\Delta_p$ and the standard compactness argument, we see that there exist a subsequence $\{\tilde{u}_{j'}\}$ of $\{\tilde{u}_j\}$ and $\tilde{w} \in L^\infty(0, 1; W_0^{1,p}(\Omega))$ such that

$$\tilde{u}_{j'} \rightarrow \tilde{w} \quad \text{in } C([0, 1]; L^2(\Omega)),$$

and $\tilde{w}(\sigma)$ satisfies (1.10) for each $\sigma \in [0, 1]$ (see the proofs of Theorem 1 of [14] and Lemma 4 of [11]). Putting $w = \tilde{w}(0)$, we see that there exists a subsequence $\{\tau_{j'}\}$ of $\{\tau_j\}$ satisfying (2.9) and w satisfies (1.10). Finally, we show that there exists a subsequence $\{\tau_{j''}\}$ of $\{\tau_{j'}\}$ such that

$$\tilde{u}(\tau_{j''}) \rightarrow w \quad \text{in } W_0^{1,p}(\Omega). \tag{2.15}$$

In fact, since $\{\tilde{u}(\tau_{j'})\}$ is bounded in $W_0^{1,p}$, it follows from (2.9) that there exists a subsequence $\{\tau_{j''}\}$ of $\{\tau_{j'}\}$ such that

$$\tilde{u}(\tau_{j''}) \rightarrow w \quad \text{weakly in } W_0^{1,p}(\Omega) \quad \text{and strongly in } L^p(\Omega). \tag{2.16}$$

Since w satisfies (1.10), it follows from (2.11) that $E(\tilde{u}(\tau_{j''})) \rightarrow \gamma_* = E(w)$. Moreover, it follows from (2.16) that $\|\tilde{u}(\tau_{j''})\|_p \rightarrow \|w\|_p$. Thus, we have

$$\|\nabla \tilde{u}(\tau_{j''})\|_p \rightarrow \|\nabla w\|_p. \tag{2.17}$$

Since $W_0^{1,p}$ is a uniformly convex Banach space, (2.15) follows from (2.16) and (2.17). This completes the proof of Theorem E. \square

Finally, we prove Theorem C. To prove it, we need to prepare one lemma.

Lemma 2.2. *Let $p > 2$, $\lambda_1 > 1$ and $\gamma \geq 0$. Suppose that $w \in W_0^{1,p}$ is non-negative in Ω , and satisfies $-\Delta_p w - |w|^{p-2}w = \gamma w$ in $\mathcal{D}'(\Omega)$. Then, we have $w \equiv 0$ in Ω .*

Proof of Lemma 2.2. Suppose that $w \not\equiv 0$ in Ω . Then, by the standard argument (see, e.g., [13, p.418]), we see that $w \in C^{1+\alpha}(\bar{\Omega})$ for some $\alpha \in (0,1)$ and w is positive in Ω . Let φ be a positive solution of $-\Delta_p \varphi = \lambda_1 |\varphi|^{p-2} \varphi$ in $\mathcal{D}'(\Omega)$. Since w satisfies $-\Delta_p w \geq |w|^{p-2} w$ in $\mathcal{D}'(\Omega)$, in the same way as in the proof of Theorem II of [4], we get $\varphi \equiv 0$ in Ω . This is a contradiction. Hence, we have $w \equiv 0$ in Ω . \square

Proof of Theorem C. We prove by contradiction. Let $u(t)$ be a global solution of (P) such that $u_0 \in W_0^{1,p} \setminus \{0\}$ is non-negative in Ω . Then, by the maximum principle as in [14], $u(t)$ is non-negative in Ω for $t \in [0, \infty)$. From Theorem B, we have $\gamma_* = \lim_{t \rightarrow \infty} [E(u(t)) / \|u(t)\|_2^2] \geq 0$. Moreover, from Theorem E, there exist a sequence $\{t_j\}$ satisfying $t_j \rightarrow \infty$ and $w \in W_0^{1,p}$ such that

$$\frac{u(t_j)}{\|u(t_j)\|_2} \rightarrow w \quad \text{in } W_0^{1,p}, \tag{2.18}$$

$$-\Delta_p w - |w|^{p-2}w = \gamma_* w \quad \text{in } \mathcal{D}'(\Omega).$$

Since $u(t)$ is non-negative in Ω for $t \in [0, \infty)$, from (2.18), we see that w is also non-negative in Ω . Thus, it follows from Lemma 2.2 that $w \equiv 0$ in Ω . However, this contradicts $\|w\|_2 = 1$. Hence, we obtain Theorem C. \square

§3. Eigenvalue Problem (1.10) for $N=1$

In this section, we consider the eigenvalue problem (1.10) for the case $N=1$. Especially, we are interested in the set of all non-negative solutions of (1.10) with $\gamma_* < 0$, which is related to the asymptotic profiles of non-negative blowup solutions of (P).

First, we consider the following boundary value problem:

$$\begin{cases} -(|u'|^{p-2}u')' - |u|^{p-2}u = -u(x), & x \in \Omega, \\ u \in W_0^{1,p}(\Omega), \quad u(x) \geq 0, \neq 0, & x \in \Omega. \end{cases} \tag{3.1}$$

Here, the symbol ' denotes the differentiation with respect to x . Let S_l be the set of all solutions of (3.1) for $\Omega = (-l, l)$. Then, the structure of S_l is determined as follows.

Proposition 3.1. *Let l_p be the positive number such that*

$$\lambda_1(-l_p, l_p) = \inf \left\{ \frac{\|u'\|_p^p}{\|u\|_p^p} : u \in W_0^{1,p}(-l_p, l_p), u \neq 0 \right\} = 1,$$

and $m_p = pl_p / (p-2)$.

- (1) *If $l \leq l_p$, then S_l is empty.*
- (2) *If $l_p < l \leq m_p$, then there exists a unique positive solution Φ_l of (3.1) and $S_l = \{\Phi_l\}$.*
- (3) *If $l > m_p$, then $S_l = \cup_{k=1}^{[l/m_p]} S_l^k$, where $[l/m_p]$ denotes the largest integer not exceeding l/m_p , and $S_l^k = \{\sum_{j=1}^k \Phi_{m_p}(\cdot - y_j) : -l \leq y_1 - m_p, y_j + 2m_p \leq y_{j+1}, j = 1, \dots, k-1, y_k + m_p \leq l\}$.*

As a corollary to Proposition 3.1, we have the main result in this section.

Theorem 3.2. *Let $\gamma < 0$ and $\Sigma(\gamma)$ be the set of all solutions of*

$$\begin{cases} -(|u|^{p-2}u'(x))' - |u|^{p-2}u(x) = \gamma u(x), & x \in (-l, l), \\ u \in W_0^{1,p}(-l, l), \quad \|u\|_2 = 1, \quad u(x) \geq 0, & x \in (-l, l). \end{cases}$$

- (1) *When $l \leq l_p$, $\Sigma(\gamma)$ is empty for any $\gamma < 0$.*
- (2) *When $l_p < l \leq m_p$, let $\gamma_1 = E(\Phi_l) / \|\Phi_l\|_2^2$. Then $\gamma_1 < 0$ and $\Sigma(\gamma_1) = \{\tilde{\Phi}_l\}$, where $\tilde{\Phi}_l = \Phi_l / \|\Phi_l\|_2$, and $\Sigma(\gamma)$ is empty if $\gamma \neq \gamma_1$.*
- (3) *When $l > m_p$, for $k=1, 2, \dots, [l/m_p]$, let $\gamma_k = k^{1-p/2} E(\Phi_{m_p}) / \|\Phi_{m_p}\|_2^2$. Then $\gamma_1 < \gamma_2 < \dots < \gamma_{[l/m_p]} < 0$ and $\Sigma(\gamma_k) = \{\sum_{j=1}^k \tilde{\Phi}_{m_p}(\cdot - y_j) : -l \leq y_1 - m_p, y_j + 2m_p \leq y_{j+1}, j=1, \dots, k-1, y_k + m_p \leq l\}$, and $\Sigma(\gamma)$ is empty if $\gamma \notin \{\gamma_1, \gamma_2, \dots, \gamma_{[l/m_p]}\}$.*

Theorem 3.2 follows immediately from Proposition 3.1. We note that γ_1 defined in Remark 1.1 coincides with that in Theorem 3.2 in this case. In order to prove Proposition 3.1, we consider the following initial value problem:

$$\begin{cases} (|u|^{p-2}u'(x))' = u(x) - |u|^{p-2}u(x), & x > 0, \\ u(0) = \alpha > 0, \quad u'(0) = 0. \end{cases} \tag{3.2}$$

Lemma 3.3. *Let $\alpha_p = (p/2)^{1/(p-2)}$ and $F(s) = (p/(p-1)) (|s|^{2/2} - |s|^{p/p})$, and let $x_\alpha = \infty$ if $\alpha < \alpha_p$, and $x_\alpha = \int_0^\alpha [F(s) - F(\alpha)]^{-1/p} ds$ if $\alpha \geq \alpha_p$. For $\alpha > 0$, there exists a unique solution φ_α of (3.2) in $(0, x_\alpha)$, and φ_α is positive in $(0, x_\alpha)$. Moreover, when $\alpha \geq \alpha_p$, $x_\alpha < \infty$ and φ_α satisfies $\varphi_\alpha(x_\alpha) = 0$, $\varphi'_\alpha(x_\alpha) < 0$ if $\alpha > \alpha_p$, and $\varphi'_\alpha(x_\alpha) = 0$ if $\alpha = \alpha_p$.*

Proof of Lemma 3.3. Let $u(x)$ be a solution of (3.2). Then, we have

$$|u'|^{p-2}u'(x) = \int_0^x [u(y) - |u|^{p-2}u(y)] dy, \quad x \geq 0. \tag{3.3}$$

When $\alpha = 1$, it follows from (3.3) that $u(x) = 1$ for $x \geq 0$. When $\alpha \neq 1$, from (3.3) we see that there exists $x_0 > 0$ such that $(\alpha - 1)u'(x) < 0$ for $0 < x < x_0$. Thus, $u(x)$ is twice differentiable in $(0, x_0)$. Multiplying the equation of (3.2) by u' and integrating over $(0, x)$ yields

$$|u'(x)|^p = F(u(x)) - F(\alpha), \quad x \geq 0. \tag{3.4}$$

From (3.3) and (3.4), we see that there exists a unique solution φ_α of (3.2) in $(0, x_\alpha)$, and φ_α is positive in $(0, x_\alpha)$. In particular, when $\alpha \geq \alpha_p$, $u = \varphi_\alpha(x)$ is given as the inverse function of $x = \int_u^\alpha [F(s) - F(\alpha)]^{-1/p} ds$. So, we see that $x_\alpha < \infty$ and φ_α satisfies $\varphi_\alpha(x_\alpha) = 0$, $\varphi'_\alpha(x_\alpha) < 0$ if $\alpha > \alpha_p$, and $\varphi'_\alpha(x_\alpha) = 0$ if $\alpha = \alpha_p$. □

Remark 3.4. By an elementary computation, we see that x_α is strictly decreasing with respect to $\alpha \geq \alpha_p$. It is known that $l_p = (p-1)^{1/p} B(1/p, 1-1/p)/p = [\pi (p-1)^{1/p}] / [p \sin(\pi/p)]$, where $B(\cdot, \cdot)$ is the beta function. Another elementary calculation yields $\lim_{\alpha \rightarrow \infty} x_\alpha = l_p$ and $x_{\alpha_p} = m_p$.

Proposition 3.1 follows from Lemma 3.3 and Remark 3.4. In particular, Φ_l is given by

$$\Phi_l(x) = \begin{cases} \varphi_{\alpha(l)}(x), & \text{for } 0 \leq x \leq l, \\ \varphi_{\alpha(l)}(-x), & \text{for } -l \leq x < 0, \end{cases}$$

where $\alpha(l) \in [\alpha_p, \infty)$ is the unique number such that $l = x_{\alpha(l)}$.

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