

**THE INFLUENCE OF DOMAIN GEOMETRY IN
THE BOUNDARY BEHAVIOR OF LARGE SOLUTIONS
OF DEGENERATE ELLIPTIC PROBLEMS**

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Abstract: In this paper we study the asymptotic boundary behavior of large solutions of the equation $\Delta u = d^\alpha u^p$ in a regular bounded domain Ω in \mathbb{R}^N , $N \geq 2$, where $d(x)$ denotes the distance from x to $\partial\Omega$, $p > 1$ and $\alpha > 0$. We precise the expansion which depends on the mean curvature of the boundary.

1 – Introduction: notations and main results

Let Ω be a regular bounded domain in \mathbb{R}^N , $N \geq 2$, $p > 1$ and $\alpha > 0$. We denote by $d(x)$ the distance from x to $\partial\Omega$, the boundary of Ω . In this paper we consider the semilinear degenerate equation

$$(1) \quad \Delta u = d^\alpha u^p \quad \text{in } \Omega$$

and we are interesting in the large solutions of (1), that is solutions of (1) which blow up at the boundary:

$$(2) \quad u(x) \rightarrow +\infty \quad \text{as } d(x) \rightarrow 0 .$$

Note already that the maximum principle implies that the solutions $u \in C^2(\Omega)$ of (1)–(2) are positive in Ω .

Equation (1) registers in problems of the form

$$(3) \quad \Delta u = p(x) f(u) \quad \text{in } \Omega .$$

Received: August 30, 2005; *Revised:* July 28, 2006.

AMS Subject Classification: 35J25, 35J70.

Keywords: elliptic equations; boundary behavior; blow-up.

Those problems were first studied by Bieberbach [4] for the case $p(x) = 1$, $f(u) = e^u$ and $N = 2$, in the context of Riemannian surfaces of constant negative curvature, and the theory of automorphic functions. The case $p(x) > 0$ for all $x \in \bar{\Omega}$ has been, largely dealt with in the literature (see [7], [12], [8], [5] for example).

Existence of solutions of (1)–(2) was established by Lair and Wood [9]. The question of the uniqueness of solutions of (1)–(2) is more delicate. When $\alpha = 0$ and $p > 1$, it is well known that problem (1)–(2) has a unique solution which satisfies

$$(4) \quad \lim_{d(x) \rightarrow 0} u(x) d(x)^{\frac{2}{p-1}} = \left(\frac{2(p+1)}{(p-1)^2} \right)^{\frac{1}{p-1}}.$$

This was first established by Loewner and Nirenberg [10] for the case $p = (N+2)/(N-2)$. Later we can find many extensions, see for example [1], [2] and [14] and the references cited there. The case $\alpha < 0$ and $p > 0$ is studied in [6]. In the general case $\alpha \geq 0$, Marcus and Véron proved the uniqueness of the solutions of (1)–(2) under the condition $1 < p < (N+1+\alpha)/(N-1)$. Our first theorem completes this result and gives the rate of the blow-up.

Theorem 1.1. *Let $u \in C^2(\Omega)$ be a solution of (1)–(2). Then it satisfies*

$$(5) \quad \lim_{d(x) \rightarrow 0} d(x)^{\frac{\alpha+2}{p-1}} u(x) = l$$

where l is given by

$$(6) \quad l = \left[\frac{(\alpha+2)(\alpha+p+1)}{(p-1)^2} \right]^{\frac{1}{p-1}}.$$

This theorem allows us to establish the uniqueness of solutions of (1)–(2) with no conditions on p and α .

Theorem 1.2. *Problem (1) possesses a unique large solution.*

In the second time we are interested in the influence of the geometry of Ω in the boundary behavior. When $\alpha = 0$, this problem was first studied by Bandle and Marcus [3] for the radially symmetric solutions of (1)–(2) in a ball. Later their result was extended by del Pino and Letelier [13] for general solutions. They proved that a lower-order term, still explosive, appears in the expansion

of u which depends linearly of the mean curvature of the boundary of Ω . More precisely, if $1 < p < 3$ and $\alpha = 0$, then on a sufficiently small neighborhood of $\partial\Omega$ we have the expansion

$$(7) \quad u(x) = \left(\frac{2(p+1)}{(p-1)^2} \right)^{\frac{1}{p-1}} d(x)^{-\frac{2}{p-1}} \left\{ 1 + \frac{N+1}{p+3} H(\bar{x}) d(x) + o(d(x)) \right\} .$$

Here, for all x in a neighborhood of $\partial\Omega$, \bar{x} denotes the unique point of the boundary such that $d(x) = |x - \bar{x}|$ and $H(\bar{x})$ the mean curvature of the boundary at that point. Estimate (7) generalizes to our case $\alpha \geq 0$ in the following way.

Theorem 1.3. *Let $u \in C^2(\Omega)$ a large solution of (1). Then, on a sufficiently small neighborhood of $\partial\Omega$:*

$$(8) \quad u(x) = l d(x)^{-\frac{2+\alpha}{p-1}} \left\{ 1 + \frac{N-1}{\alpha+p+3} H(\bar{x}) d(x) + o(d(x)) \right\} .$$

This theorem implies that on a sufficiently small neighborhood of $\partial\Omega$:

$$(9) \quad u(x) - l d(x)^{-\frac{2+\alpha}{p-1}} = \frac{N-1}{\alpha+p+3} H(\bar{x}) d(x)^{-\frac{\alpha+3-p}{p-1}} + o\left(d(x)^{-\frac{\alpha+3-p}{p-1}}\right) .$$

Therefore, we obtain that

- if $p > \alpha + 3$, then the first member of (9) tends to 0 at the boundary,
- if $p = \alpha + 3$, then $u(x) - l d(x)^{-\frac{2+\alpha}{p-1}} = \frac{N-1}{\alpha+p+3} H(\bar{x}) + o(1)$,
- if $p < \alpha + 3$, then the first member of (9) is not bounded and the blow-up depends on the mean curvature. Roughly, the “more curved” or “sharper” towards the exterior of Ω is around a given point of $\partial\Omega$, the higher the explosion rate at that point is.

That is a generalization of the results of Bandle and Marcus [3] for the radially symmetric solutions of (1)–(2) in a ball $\Omega = B(0, R)$:

- if $p > 3$, then $u(r) - l(R-r)^{-\frac{2}{p-1}} \rightarrow 0$ when $r \rightarrow R$,
- if $p = 3$, then $u(r) - l(R-r)^{-\frac{2}{p-1}} \rightarrow \frac{C(N)}{R}$ when $r \rightarrow R$, which represents the mean curvature of the ball,
- if $p < 3$, then $u(r) - l(R-r)^{-\frac{2}{p-1}} \rightarrow \infty$ when $r \rightarrow R$.

Our paper is organized as follows:

1. Introduction
2. Asymptotic behavior and uniqueness
3. Boundary influence in the explosion rate.

2 – Asymptotic behavior and uniqueness

We begin this section by proving a classical estimate for all solution u of (1) (see [12]).

Proposition 2.1 (Osserman estimate). *There exist two positive constants $a = a(\partial\Omega)$ and $C = C(\Omega, \alpha, p)$ such that for all solution $u \in C^2(\Omega)$ of equation (1), we have:*

$$(10) \quad u(x) \leq C d(x)^{-\frac{2+\alpha}{p-1}}$$

for all $x \in \Omega$ such that $d(x) < a$.

Proof: Since Ω is regular there exist $\tilde{a} = \tilde{a}(\Omega) > 0$ and $M = M(\Omega) > 0$ such that

$$(11) \quad |\Delta d(x)| \leq M, \quad |\nabla d(x)| = 1$$

for all $x \in \Omega$ such that $d(x) < \tilde{a}$. Set $a = \min(1, \frac{\tilde{a}}{2})$. Let $x_0 \in \Omega$ such that $d(x_0) < a$ and $r_0 = d(x_0)/2$. We denote by B_0 the ball centered at x_0 of radius r_0 and we define the function w in B_0 as follows:

$$(12) \quad w(x) = \lambda d(x)^{-\frac{\alpha}{p-1}} \left(r_0^2 - |x - x_0|^2 \right)^{-\frac{2}{p-1}}$$

with $\lambda > 0$ to determine such that

$$(13) \quad -\Delta w + d^\alpha w^p \geq 0 \quad \text{in } B_0.$$

A straightforward computation gives:

$$\begin{aligned} -\Delta w + d^\alpha w^p &= \lambda d^{-\frac{\alpha}{p-1}} \left(r_0^2 - |x - x_0|^2 \right)^{-\frac{2p}{p-1}} \times \\ &\quad \left[-\frac{\alpha(\alpha+p-1)}{(p-1)^2} \left(r_0^2 - |x - x_0|^2 \right)^2 d^{-2} + \frac{\alpha}{p-1} \left(r_0^2 - |x - x_0|^2 \right)^2 d^{-1} \Delta d \right. \\ &\quad + \frac{8\alpha}{(p-1)^2} \left(r_0^2 - |x - x_0|^2 \right) d^{-1} \nabla d \cdot (x - x_0) - \frac{8(p+1)}{(p-1)^2} |x - x_0|^2 \\ &\quad \left. - \frac{4N}{p-1} \left(r_0^2 - |x - x_0|^2 \right) + \lambda^{p-1} \right]. \end{aligned}$$

Since $|x - x_0| < d(x_0) \leq 1$, $d(x) \geq d(x_0)/2$ and $r_0^3 < r_0^2$, there exists a constant $L = L(\alpha, p, M) > 0$ such that

$$-\Delta w + d^\alpha w^p \geq \lambda d^{-\frac{\alpha}{p-1}} \left(r_0^2 - |x - x_0|^2 \right)^{-\frac{2p}{p-1}} (-Lr_0^2 + \lambda^{p-1})$$

in B_0 . Therefore, we choose $\lambda = L^{\frac{1}{p-1}} r_0^{\frac{2}{p-1}}$ and we obtain (13). Note that $w(x) = +\infty$ if $x \in \partial B_0$ because $-2/(p-1) < 0$. The comparison principle implies $u \leq w$ in B_0 and in particular

$$u(x_0) \leq w(x_0) = L^{\frac{1}{p-1}} \left(\frac{d(x_0)}{2} \right)^{\frac{2}{p-1}} (d(x_0))^{-\frac{\alpha}{p-1}} \left(\frac{d(x_0)}{2} \right)^{-\frac{4}{p-1}}$$

which gives inequality (10). ■

We now establish an estimate from below for the solutions of (1)–(2). The results of [1] and [2] can't be used because the distance function d is not positive in $\bar{\Omega}$. Nevertheless we can adapt them as follows.

Proposition 2.2. *Let $u \in C^2(\Omega)$ be a solution of (1)–(2). Then*

$$(14) \quad \liminf_{d(x) \rightarrow 0} d(x)^{\frac{\alpha+2}{p-1}} u(x) \geq l$$

where l is defined in (6).

Proof: Let $\varepsilon > 0$, \tilde{a} be as the proof of Proposition 2.1 and $\beta \in (0, 1)$. We define

$$\underline{u}(x) = \beta l \left((d(x) + \varepsilon)^{-\frac{\alpha+2}{p-1}} - (\bar{a} + \varepsilon)^{-\frac{\alpha+2}{p-1}} \right)$$

where \bar{a} will be determined such that $\bar{a} < \tilde{a}$. We have $\underline{u} > 0$ on $\partial\Omega$ and $\underline{u}(x) = 0$ for all x such that $d(x) = \bar{a}$. Moreover a straightforward computation yields

$$-\Delta \underline{u} + d^\alpha \underline{u}^p = \beta \left[\Delta d \left(\frac{\alpha+2}{p-1} \right) l (d + \varepsilon)^{-\frac{\alpha+p+1}{p-1}} - l^p (d + \varepsilon)^{-\frac{\alpha+2p}{p-1}} + d^\alpha \beta^{p-1} l^p \left((d + \varepsilon)^{-\frac{\alpha+2}{p-1}} - (\bar{a} + \varepsilon)^{-\frac{\alpha+2}{p-1}} \right)^p \right]$$

in $0 < d(x) < \bar{a}$. Using inequality (11), we obtain

$$-\Delta \underline{u} + d^\alpha \underline{u}^p \leq \beta l^p (d + \varepsilon)^{-\frac{\alpha+2p}{p-1}} \left[M \left(\frac{\alpha+2}{p-1} \right) l^{1-p} (d + \varepsilon) - 1 + \beta^{p-1} \left(\frac{d}{d + \varepsilon} \right)^\alpha \right]$$

which implies

$$-\Delta \underline{u} + d^\alpha \underline{u}^p \leq \beta l^p (d + \varepsilon)^{-\frac{\alpha+2p}{p-1}} \left[\overline{M} (d + \varepsilon) - (1 - \beta^{p-1}) \right]$$

with $\overline{M} = M \left(\frac{\alpha+2}{p-1} \right) l^{1-p}$. We now choose $\bar{a} = \frac{1}{2} \min \left(\tilde{a}, \frac{1-\beta^{p-1}}{\overline{M}} \right)$ and impose $\varepsilon < \frac{1}{2} \left(\frac{1-\beta^{p-1}}{\overline{M}} \right)$.

Then \bar{u} is a subsolution of (1) in $0 < d(x) < \bar{a}$. By the maximum principle we derive $\underline{u} \leq u$ in $0 < d(x) < \bar{a}$. Letting ε tend to 0, this implies for all $\beta \in (0, 1)$ and x such that $d(x) < \bar{a}$:

$$\beta l \left[1 - \left(\frac{d(x)}{\bar{a}} \right)^{\frac{\alpha+2}{p-1}} \right] \leq d(x)^{\frac{\alpha+2}{p-1}} u(x).$$

Therefore for all $\beta \in (0, 1)$:

$$\beta l \leq \liminf_{d(x) \rightarrow 0} d(x)^{\frac{\alpha+2}{p-1}} u(x)$$

which ends the proof. ■

Because of Proposition 2.2, we can describe the asymptotic behavior of radially symmetric solutions of (1)–(2).

Proposition 2.3. *Let $R > 0$ and $v \in C^2(0, R)$ a solution of*

$$(15) \quad -v'' - \frac{N-1}{r} v' + (R-r)^\alpha v^p = 0$$

in $(0, R)$ such that

$$\lim_{r \rightarrow R} v(r) = +\infty.$$

Then

$$(16) \quad \lim_{r \rightarrow R} (R-r)^{\frac{\alpha+2}{p-1}} v(r) = l$$

where l is defined in (6). ■

We omit the proof of this proposition because it follows the idea of [14]: the function $w(t) = (R-r)^{\frac{\alpha+2}{p-1}} v(r)$ with $R-r = e^{-t}$ is bounded and satisfies a second order differential equation in a neighborhood of infinity and the ω -limit set of a trajectory of that equation is $\{0\}$ or $\{l\}$. Therefore Proposition 2.2 implies Proposition 2.3. Those results allows us to prove Theorem 1.1.

Proof of Theorem 1.1: In view of (14) we must only prove that

$$(17) \quad \limsup_{d(x) \rightarrow 0} d(x)^{\frac{\alpha+2}{p-1}} u(x) \leq l .$$

Still the results of [1], [2] or [14] don't apply directly but we can adapt them. Let $y \in \partial\Omega$. Since $\partial\Omega$ is smooth, there exists a ball B_y centered at a point Y of radius R_y such that $B_y \subset \Omega$ and $\overline{B_y} \cap \partial\Omega = \{y\}$. We introduce the function V defined by $V(x) = v(|x|)$ for all $x \in B_{R_y}$ where v is a function as in Proposition 2.3 with $R = R_y$. The function v exists because it is the radial solution of (1)–(2) for $\Omega = B$ (see [9]). Let $k > 1$. Finally we introduce the function V_k defined by $V_k(x) = k^{\frac{2}{p-1}} V(k(x - Y))$ for all $x \in B(Y, \frac{R_y}{k})$. Note that $B(Y, \frac{R_y}{k}) \subset B_y$ and V_k is solution of

$$-\Delta V_k + \left(R_y - k|x - Y|\right)^\alpha V_k^p = 0$$

in $B(y, \frac{R_y}{k})$ and satisfies

$$\lim_{|x-Y| \rightarrow \frac{R_y}{k}} V_k(x) = +\infty .$$

Since $x \in B(Y, \frac{R_y}{k})$ implies $d(x) \geq R_y - |x - Y| \geq R_y - k|x - Y|$, the comparison principle involves $u \leq V_k$ in $B(Y, \frac{R_y}{k})$. Letting k tend to 1, we obtain

$$(18) \quad u(x) \leq v(|x - Y|) \quad \text{in } B_y .$$

Because of Proposition 2.3, for all $\varepsilon > 0$ there exists $\eta > 0$ such that

$$(19) \quad \left| s^{\frac{\alpha+2}{p-1}} v(R_y - s) - l \right| < \varepsilon \quad \forall s \in (0, \eta) .$$

Let $\tilde{\eta} > 0$ be sufficiently small so that for all $x \in \Omega$ with $d(x) < \tilde{\eta}$ there exists a unique $y \in \partial\Omega$ such that $|x - y| = d(x)$. Then for all $x \in \Omega$ such that $d(x) < \min(\eta, \tilde{\eta})$, both inequalities (18) and (19) imply

$$d(x)^{\frac{\alpha+2}{p-1}} u(x) \leq d(x)^{\frac{\alpha+2}{p-1}} v(R_y - d(x)) < l + \varepsilon$$

and inequality (17) holds. ■

Proof of Theorem 1.2: Large solutions of (1) satisfy (5). Then two large solutions u_1 and u_2 of (1) are such that

$$\lim_{d(x) \rightarrow 0} \frac{u_1(x)}{u_2(x)} = 1 .$$

and the result follows as in [1] or [11]. ■

3 – Boundary influence in the explosion rate

In this section we prove Theorem 1.3. As in [13] we construct suitable sub- and supersolutions of (1) in a neighborhood of $\partial\Omega$ which are inspired of the radial study that we omit here.

Since Ω is regular there exists $\bar{b} > 0$ such that d is a function of class C^2 in $\{x \in \Omega / d(x) < \bar{b}\}$, $|\nabla d(x)| = 1$ and

$$(20) \quad \Delta d(x) = -(N-1)H(\bar{x}) + o(1) \quad \text{as } d(x) \rightarrow 0.$$

Let $b_0 \in (0, \bar{b})$, $b \in (0, b_0)$ and $\varepsilon > 0$. We introduce the function Ψ defined in $E_{b,b_0} = \{x \in \Omega / b < d(x) < b_0\}$ by

$$\Psi(x) = l(d(x) - b)^{-\frac{\alpha+2}{p-1}} + \frac{l}{\alpha+p+3} \left((N-1)H(\bar{x}) + \varepsilon \right) (d(x) - b)^{-\frac{\alpha+3-p}{p-1}}.$$

We claim that if b_0 is chosen sufficiently small, independently of ε and b , then Ψ is a supersolution in E_{b,b_0} . Indeed, a straightforward computation using (20) gives:

$$\begin{aligned} \Delta \Psi &= l^p (d(x) - b)^{-\frac{\alpha+2p}{p-1}} + l(d(x) - b)^{-\frac{\alpha+p+1}{p-1}} \left[\frac{\alpha+2}{p-1} (N-1)H(\bar{x}) \right. \\ &\quad \left. + \frac{(\alpha+3-p)(\alpha+2)}{(\alpha+p+3)(p-1)^2} \left((N-1)H(\bar{x}) + \varepsilon \right) + o(1) \right. \\ &\quad \left. + \frac{\alpha+3-p}{(p-1)(\alpha+p+3)} \left((N-1)H(\bar{x}) + \varepsilon \right) \left((N-1)H(\bar{x}) + o(1) \right) (d(x) - b) \right]. \end{aligned}$$

On the other hand, we have

$$\begin{aligned} d(x)^\alpha \Psi^p &\geq (d(x) - b)^\alpha \Psi^p \\ &\geq l^p (d(x) - b)^{-\frac{\alpha+2p}{p-1}} \times \\ &\quad \left[1 + \frac{p}{\alpha+p+3} \left((N-1)H(\bar{x}) + \varepsilon \right) (d(x) - b) + o(d(x) - b) \right]. \end{aligned}$$

Then

$$\begin{aligned}
 -\Delta\Psi + d^\alpha\Psi^p &\geq \\
 &\geq l\left(d(x) - b\right)^{-\frac{\alpha+p+1}{p-1}} \times \\
 &\quad \left[-\frac{\alpha+2}{p-1}(N-1)H(\bar{x}) - \frac{(\alpha+3-p)(\alpha+2)}{(\alpha+p+3)(p-1)^2} \left((N-1)H(\bar{x}) + \varepsilon\right) \right. \\
 &\quad - \frac{\alpha+3-p}{(p-1)(\alpha+p+3)} \left((N-1)H(\bar{x}) + \varepsilon\right) \left((N-1)H(\bar{x}) + o(1)\right) \left(d(x) - b\right) \\
 &\quad \left. + \frac{l^{p-1}p}{\alpha+p+3} \left((N-1)H(\bar{x}) + \varepsilon\right) + o(1) \right].
 \end{aligned}$$

Since

$$-\frac{\alpha+2}{p-1} - \frac{(\alpha+3-p)(\alpha+2)}{(\alpha+p+3)(p-1)^2} + \frac{l^{p-1}p}{\alpha+p+3} = 0,$$

and since the coefficient of ε is $(\alpha+2)/(p-1)$, it implies that there exists $b_0 = b_0(\varepsilon) \in (0, \bar{b})$ such that for all $0 < b < b_0$:

$$-\Delta\Psi + d^\alpha\Psi^p \geq 0 \quad \text{in } E_{b,b_0}.$$

Consider the solution u of (1)–(2). We claim that there exists a positive number K independent of $b \in (0, b_0)$ such that:

$$(21) \quad \Psi(x) + K \geq u(x)$$

for all $x \in \Omega$ with $d(x) = b_0$. In fact, if we define

$$M_0 = \max_{d(x)=b_0} u(x),$$

we can compute for all x such that $d(x) = b_0$:

$$\Psi(x) = l(b_0 - b)^{-\frac{\alpha+2}{p-1}} + \frac{l}{\alpha+p+3} \left((N-1)H(\bar{x}) + \varepsilon\right) (b_0 - b)^{-\frac{\alpha+3-p}{p-1}}.$$

Since $\partial\Omega$ is regular, there exists a real $b_1 \in (0, b_0)$ such that

$$\left| \frac{1}{\alpha+p+3} \left((N-1)H(\bar{x}) + \varepsilon\right) (b_0 - b) \right| \leq \frac{1}{2}$$

for all $b \in (b_1, b_0)$, where \bar{x} is such that $d(x) = |x - \bar{x}|$. Therefore

$$1 + \frac{1}{\alpha+p+3} \left((N-1)H(\bar{x}) + \varepsilon\right) (b_0 - b) \geq \frac{1}{2}$$

and then

$$\Psi(x) \geq \frac{l}{2}(b_0 - b)^{-\frac{\alpha+2}{p-1}} \geq \frac{l}{2}(b_0 - b_1)^{-\frac{\alpha+2}{p-1}}$$

for all $b \in (b_1, b_0)$, where \bar{x} is such that $d(x) = |x - \bar{x}|$. On the other hand, for all $b \in (0, b_1]$ and $d(x) = b_0$:

$$\begin{aligned} \Psi(x) &= l(b_0 - b)^{-\frac{\alpha+2}{p-1}} + \frac{l}{\alpha + p + 3} \left((N-1)H(\bar{x}) + \varepsilon \right) (b_0 - b)^{-\frac{\alpha+3-p}{p-1}} \\ &\geq lb_0^{-\frac{\alpha+2}{p-1}} - C(b_0 - b_1)^{-\frac{\alpha+3-p}{p-1}} \end{aligned}$$

with $C > 0$ and because the assumption if we assume $\alpha + 3 - p > 0$ (we omit the proof in the case $\alpha + 3 - p \leq 0$ which is simpler). Finally we obtain for all $b \in (0, b_0)$:

$$\Psi(x) \geq L = \min\left(\frac{l}{2}(b_0 - b_1)^{-\frac{\alpha+2}{p-1}}, lb_0^{-\frac{\alpha+2}{p-1}} - C(b_0 - b_1)^{-\frac{\alpha+3-p}{p-1}}\right),$$

then, for all x such that $d(x) = b_0$,

$$u \leq M_0 \leq \max(1, M_0 - L) + L \leq \max(1, M_0 - L) + \psi$$

which implies (21).

On the other hand the function $\Psi + K$ is itself a supersolution of equation (1) in E_{b,b_0} . Therefore the comparison principle implies (21) in E_{b,b_0} . Letting b tend to 0, we obtain

$$u(x) \leq ld(x)^{-\frac{\alpha+2}{p-1}} + \frac{l}{\alpha + p + 3} \left((N-1)H(\bar{x}) + \varepsilon \right) d(x)^{-\frac{\alpha+3-p}{p-1}} + K$$

for all $x \in \Omega$ such that $0 < d(x) < b_0$. In the same way, by considering subsolutions in the form

$$\phi(x) = l(d(x) + b)^{-\frac{2+\alpha}{p-1}} + \frac{l}{\alpha + p + 3} \left((N-1)H(\bar{x}) - \varepsilon \right) (d(x) + b)^{-\frac{\alpha+3-p}{p-1}} - \bar{K}$$

we obtain expansion (8). ■

ACKNOWLEDGEMENTS – The authors are pleased to acknowledge Catherine Bandle for their profitable discussions.

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