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PFH spectral invariants on the two-sphere and the large scale geometry of Hofer's metric

Abstract. We resolve three longstanding questions related to the large scale geometry of the group of Hamiltonian diffeomorphisms of the two-sphere, equipped with Hofer's metric. Namely: (1) we resolve the Kapovich–Polterovich question by showing that this group is not quasi-isometric to the real line; (2) more generally, we show that the kernel of Calabi over any proper open subset is unbounded; and (3) we show that the group of area and orientation preserving homeomorphisms of the two-sphere is not a simple group. We also find, as a corollary, that the group of area-preserving diffeomorphisms of the open disc, equipped with an area form of finite area, is not perfect. Central to all of our proofs are new sequences of spectral invariants over the two-sphere, defined via periodic Floer homology.

Keywords. Area-preserving homeomorphisms, symplectic topology, periodic Floer homology, Hofer's metric

Contents

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1. Introduction

It is a remarkable fact that the group of Hamiltonian diffeomorphisms of a symplectic manifold admits a bi-invariant Finsler metric, known as *Hofer's metric*. The existence of such a metric on an infinite-dimensional Lie group is highly unusual, due to the lack of compactness, and stands in contrast to the fact that a simple[1] finite-dimensional Lie group admits a bi-invariant Finsler metric only if it is compact; see [43, Prop. 1.3.15].

The theme of this article is the large-scale geometry of Hofer's metric, on $\mathrm{Ham}(\mathbb{S}^2, \omega)$, the Hamiltonian diffeomorphisms of the two-sphere.[2] Our first result, Theorem 1.4, settles two longstanding questions, presented below, about the quasi-isometry type of $\mathrm{Ham}(\mathbb{S}^2, \omega)$. The first of the two questions was posed by Kapovich and Polterovich in 2006.

Question 1.1. *Is* $\mathrm{Ham}(\mathbb{S}^2, \omega)$ *quasi-isometric to the real line* $\mathbb{R}$*?*

The second question is due to Polterovich and dates back to the 2000s. To state it, consider a connected, proper open set $U \subset \mathbb{S}^2$ and denote by $\mathrm{Ham}_U(\mathbb{S}^2, \omega)$ the subgroup of $\mathrm{Ham}(\mathbb{S}^2, \omega)$ consisting of Hamiltonian diffeomorphisms supported in U. This subgroup carries a well-known group homomorphism called the *Calabi homomorphism*:

$$\mathrm{Cal} : \mathrm{Ham}_U(\mathbb{S}^2, \omega) \to \mathbb{R},$$

whose definition we recall in Section 2.1 (see (7)).

Question 1.2. *Suppose*[3] *that* $\mathrm{Area}(U) \leq \frac{1}{2} \mathrm{Area}(\mathbb{S}^2)$. *Is the kernel of* $\mathrm{Cal} :$ $\mathrm{Ham}_U(\mathbb{S}^2, \omega) \to \mathbb{R}$ *an unbounded subset of* $\mathrm{Ham}(\mathbb{S}^2, \omega)$*?*

The Hofer geometry of the two-sphere has long remained mysterious, and these two basic questions have received much attention over the past years. This is especially the

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[1] $\mathrm{Ham}(M, \omega)$ is simple for closed M, by a theorem of Banyaga [2].

[2] It is known that the group $\mathrm{Ham}(\mathbb{S}^2, \omega)$ is in fact the set of all diffeomorphisms of $\mathbb{S}^2$ which preserve the area form ω.

[3] If $\mathrm{Area}(U) > \frac{1}{2} \mathrm{Area}(\mathbb{S}^2)$, then the question is known to have an affirmative answer by Polterovich [41].

case for Question 1.1, which appears as Problem 21 on the list of open problems of McDuff–Salamon [36, Sec. 14.2]; it is mentioned as one of the motivations behind the influential article of Polterovich and Shelukhin [44, Sec. 1.3]; and it is highlighted in several articles such as [4, 13, 29, 46].

We also continue the direction of research initiated in our recent article [9]. In particular, we answer the following question from the 1980 article of Fathi [15] on the algebraic structure of $\mathrm{Homeo}_0(\mathbb{S}^2, \omega)$, the group of all area and orientation preserving homeomorphisms of the two-sphere.[4]

Question 1.3. *Is the group* $\mathrm{Homeo}_0(\mathbb{S}^2, \omega)$ *simple?*

Although at first glance this question might appear unrelated to Hofer's geometry, we will see that the large scale geometry of Hofer's metric plays a crucial role in the solution. The two-sphere is the only closed manifold for which the question of simplicity of the component of the identity in the group of volume-preserving homeomorphisms remained open; for other closed manifolds this was settled by Fathi in the late 1970s.

1.1. The large-scale geometry of the kernel of Calabi

Let d_H denote the Hofer metric on $\mathrm{Ham}(M, \omega)$, the group of Hamiltonian diffeomorphisms of a closed and connected symplectic manifold (M, ω); we will review the definition of d_H, and other basic notions from symplectic geometry, in Section 2.1.

A fundamental notion in large-scale geometry is that of *quasi-isometry*, which we now recall. A *quasi-isometric embedding* is a mapping $\Phi : (X_1, d_1) \to (X_2, d_2)$ of metric spaces for which there exist constants $A \geq 1, B \geq 0$ such that

$$\frac{1}{A} d_1(x, y) - B \leq d_2(\Phi(x), \Phi(y)) \leq A d_1(x, y) + B. \tag{1}$$

The map Φ satisfying the above is said to be a *quasi-isometry* if it is *quasi-surjective*, i.e. there exists a constant $C > 0$ such that every point in X_2 is within distance C of the image $\Phi(X_1)$.

The large-scale geometry of Hofer's metric, on general symplectic manifolds, has been studied extensively ever since Hofer's discovery of the metric in 1990 [18]; see for example [1, 10, 20, 27, 28, 38, 44, 46, 49, 52]. Usually, (Ham, d_H) is a "large" metric space. For example, it is conjectured to be always unbounded, and this has been proven for many manifolds [10, 31, 35, 38, 41, 47, 52]. Moreover, Usher [52] has proven that, for a large class of manifolds, including closed surfaces of positive genus,[5] it admits a

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[4]The group $\mathrm{Homeo}_0(\mathbb{S}^2, \omega)$ can alternatively be described as the connected component of the group of area-preserving homeomorphisms of $\mathbb{S}^2$. For any transformation group, the simplicity question is only interesting for the component of the identity because it forms a normal subgroup of the ambient group.

[5]As observed in [46], such results for surfaces of positive genus can be deduced from the arguments in [31, 41].

quasi-isometric embedding of infinite-dimensional normed vector spaces; see also Py's article [46].

Despite all the above progress, a famous case that has been difficult to understand is that of the two-sphere. All that is known is that $\mathrm{Ham}(\mathbb{S}^2, \omega)$, and the subgroup $\mathrm{Ham}_U(\mathbb{S}^2, \omega)$, are unbounded and admit a quasi-isometric embedding of the real line $\mathbb{R}$; this was proven by Polterovich [41]. As for the kernel of $\mathrm{Cal} : \mathrm{Ham}_U(\mathbb{S}^2, \omega) \to \mathbb{R}$ with $\mathrm{Area}(U) \leq \frac{1}{2}\mathrm{Area}(\mathbb{S}^2)$, it is not even known if it is unbounded, i.e. whether it is quasi-isometric to the point. It is our understanding that when Questions 1.1 and 1.2 were posed, there were not even clear conjectures about what their answers should be.

Our first point in the present work is that the kernel of Calabi is indeed rather big, which we illustrate in two different ways.

Theorem 1.4. *Let $U \subsetneq \mathbb{S}^2$. Then:*

(a) *For any $n \in \mathbb{N}$, there exists a quasi-isometric embedding of $\mathbb{R}^n$ into $(\mathrm{Ham}(\mathbb{S}^2, \omega), d_H)$ whose image is contained in the kernel of the Calabi homomorphism $\mathrm{Cal} : \mathrm{Ham}_U(\mathbb{S}^2, \omega) \to \mathbb{R}$.*

(b) *The kernel of $\mathrm{Cal} : \mathrm{Ham}_U(\mathbb{S}^2, \omega) \to \mathbb{R}$ is not coarsely proper.*

To review the terminology here, recall that a metric space (X, d) is said to be *coarsely proper* if there exists $R_0 > 0$ such that every bounded subset of (X, d) can be covered by finitely many balls of radius R_0; see [8, Def. 3.D.10]. Examples of coarsely proper spaces include the Euclidean space $\mathbb{R}^n$ or any bounded spaces – in particular, part (b) of Theorem 1.4 resolves Questions 1.1 and 1.2 – but on the other hand, an infinite-dimensional Banach space is not coarsely proper. Recall also that a *quasi-flat* in a metric space (X, d) is the image of a quasi-isometric embedding of $\mathbb{R}^n$; moreover, the *quasi-flat rank* of a metric space (X, d) is the supremum of all n such that there exists a quasi-isometric embedding of $\mathbb{R}^n$ into X. Thus, part (a) of Theorem 1.4 is equivalent to the statement that the metric space $(\mathrm{Ham}(\mathbb{S}^2, \omega), d_H)$ and the subset given by the kernel of the Calabi homomorphism $\mathrm{Cal} : \mathrm{Ham}_U(\mathbb{S}^2, \omega) \to \mathbb{R}$ have infinite quasi-flat rank. Now, it is known that the quasi-flat rank of $\mathbb{R}^n$ is n and so we see that part (a) of Theorem 1.4 also answers Questions 1.1 and 1.2. In fact, we will see in Example 1.5 below that Theorem 1.4 tells us quite a bit more about the quasi-isometry type of the metric spaces in question.

Example 1.5. Let (G, d) be a finite-dimensional connected Lie group, with a left invariant Finsler metric induced from a norm on its Lie algebra; we call such a d a *compatible metric*. As was explained above, the existence of Hofer's metric dramatically contrasts with the situation for finite-dimensional Lie groups; one might hope that the large-scale geometry also sees this. Indeed, it is known that any such (G, d) both has finite quasi-flat rank and is coarsely proper. So, our main theorem precludes this as a quasi-isometry type for $(\mathrm{Ham}(\mathbb{S}^2), d)$ or for the kernel of Calabi. Similarly, any finitely generated group, or more generally, any locally compact and compactly generated group (here we refer the reader to [8] for the precise definition) is coarsely proper [8, Prop. 3.D.29]. It would be interesting to understand to what degree the quasi-isometry type of

$\text{Ham}(\mathbb{S}^2)$ is unique, for example whether it differs[6] from that of $\text{Ham}(S)$ for other surfaces S.

Remark 1.6. Contemporaneously with our work, Polterovich–Shelukhin [45] have shown, using very different methods, that there is an isometric embedding of the space of even compactly supported functions on $(-1/8, 1/8)$ into $\text{Ham}(\mathbb{S}^2, \omega)$. This clearly answers the Kapovich–Polterovich question, and moreover implies that $\text{Ham}(\mathbb{S}^2, \omega)$ is neither coarsely proper nor of finite quasi-flat rank. It would be very interesting to relate our methods here to the methods in [45].

1.2. Nonsimplicity of $\text{Homeo}_0(\mathbb{S}^2, \omega)$

We turn now to continuous symplectic geometry.

In our recent article [9], we proved that the group of compactly supported area-preserving homeomorphisms of the disc is not simple. Our next theorem settles the simplicity question for the sphere. Recall that $\text{Homeo}_0(\mathbb{S}^2, \omega)$ denotes the identity component in the group of area-preserving homeomorphisms of the two-sphere.

Theorem 1.7. $\text{Homeo}_0(\mathbb{S}^2, \omega)$ *is not simple.*

In fact, as in our previous article [9], this theorem implies a stronger statement by appealing to a beautiful argument of Epstein and Higman [14, 17]. Recall that a group is *perfect* if it is equal to its commutator subgroup.

Corollary 1.8. $\text{Homeo}_0(\mathbb{S}^2, \omega)$ *is not perfect.*

Theorem 1.7 answers a question of Fathi[7] from the 70s [15, Appendix A.6], whose history we now briefly review. The question of simplicity of groups of homeomorphisms and diffeomorphisms was studied extensively in the 50s, 60s, and 70s and is fairly well-understood in most scenarios. However, area-preserving homeomorphisms of surfaces have remained mysterious. For example, in the case of closed manifolds, the simplicity question had been answered by the late 70s for all of the following groups: homeomorphisms, diffeomorphisms[8], volume-preserving diffeomorphisms and symplectomorphisms. And in the case of volume-preserving homeomorphisms it was answered by Fathi [15] for every closed manifold other than the two-sphere. Fathi asked the aforementioned question answered by Theorem 1.7 in the work [15].

We remark that nonsimplicity of $\text{Homeo}_0(\mathbb{S}^2, \omega)$ is surprising as it stands in dramatic contrast to the fact that on closed simply connected manifolds, such as spheres of dimension greater than 1, this is the only example of the "usual" transformation

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[6]We have learned in recent conversation with Polterovich that this question is wide open.

[7]In fact, Theorem 5.2 below answers Fathi's question for all compact genus zero surfaces; see Remark 5.6.

[8]We are considering C^∞ diffeomorphisms here. For C^k diffeomorphisms, simplicity is known for all k except when $k = \dim(M) + 1$, which remains open to date.

groups known to be nonsimple. For example, it is known that for simply connected manifolds the identity component in any of the groups mentioned in the previous paragraph is simple except, of course, in our case of area-preserving homeomorphisms of the sphere.

The simplicity of the aforementioned groups was established through the works of a long list of mathematicians who studied the question from the 30s to the late 70s. For a summary of the long history of the simplicity question, we refer the interested reader to [9, Sec. 1].

1.2.1. The perfectness question for volume-preserving diffeomorphisms of $\mathbb{R}^n$. We now explain an application to the study of the algebraic structure of diffeomorphism groups.

Let Ω be a volume form on $\mathbb{R}^n$ and denote by $\mathrm{Diff}(\mathbb{R}^n, \Omega)$ the group of all diffeomorphisms of $\mathbb{R}^n$ which preserve Ω. McDuff proved in 1980 that although $\mathrm{Diff}(\mathbb{R}^n, \Omega)$ is nonsimple,[9] it is always perfect for $n \geq 3$; see [34]. There are two distinct cases of McDuff's theorem, namely the finite volume case, which is the same as the case of an open ball with its standard volume form, and the infinite volume case. In both cases, however, the 2-dimensional case has remained open. Theorem 1.7 allows us to settle this question in the finite area case.

Corollary 1.9. *Assume that $\int_{\mathbb{R}^2} \Omega < +\infty$. Then $\mathrm{Diff}(\mathbb{R}^2, \Omega)$ is not perfect.*

We prove this corollary in Section 5.3.

1.3. New spectral invariants

We now discuss the main tools that we use and develop here for proving the above theorems. We henceforth view $\mathbb{S}^2$ as the unit sphere in standard $\mathbb{R}^3$ and equip it with the symplectic form $\omega := \frac{1}{4\pi} d\theta \wedge dz$, where (θ, z) are cylindrical coordinates. Note that this gives the sphere a total area of 1.

Periodic Floer homology and spectral invariants

To prove our results we use a version of Floer homology for area-preserving diffeomorphisms, called *periodic Floer homology* (PFH), which was introduced by Hutchings [23]; we will review PFH in Section 2.3. As will be reviewed in Section 3, one can use PFH to define a collection of invariants of Hamiltonians on the sphere,

$$c_{d,k} : C^\infty(\mathbb{S}^1 \times \mathbb{S}^2) \to \mathbb{R},$$

which are indexed by $d \in \mathbb{N}$ and $k \in \mathbb{Z}$ with k having the same parity as d.

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[9]In this case, nonsimplicity follows from the fact that the compactly supported volume-preserving diffeomorphisms form a proper normal subgroup.

We show in Section 3 that these invariants have various useful properties; see Proposition 3.2. In particular, we show that they can be used to define invariants

$$c_{d,k} : \widetilde{\mathrm{Ham}}(\mathbb{S}^2, \omega) \to \mathbb{R}, \quad c_d : \widetilde{\mathrm{Ham}}(\mathbb{S}^2, \omega) \to \mathbb{R},$$

where $c_d := c_{d,-d}$, which are well-defined on the universal cover of $\mathrm{Ham}(\mathbb{S}^2, \omega)$. Moreover, we show in Proposition 3.5 that if d is even then $c_{d,k} : \widetilde{\mathrm{Ham}}(\mathbb{S}^2, \omega) \to \mathbb{R}$ descends to $\mathrm{Ham}(\mathbb{S}^2, \omega)$ and so in particular we obtain

$$c_d : \mathrm{Ham}(\mathbb{S}^2, \omega) \to \mathbb{R},$$

defined for even d.

Homogenization

As is evident from the works of Entov–Polterovich [10, 12], for the purposes of applications to Hofer's geometry, it is often beneficial to homogenize spectral invariants. This is true in our work as well, and in fact we prove Theorem 1.4 using the homogenizations of the invariants c_d which we now introduce. More precisely, we can define, for $\varphi \in \mathrm{Ham}(\mathbb{S}^2, \omega)$ and for all $d \in \mathbb{N}$,

$$\mu_d(\varphi) := \limsup_{n \to \infty} \frac{c_d(\tilde{\varphi}^n)}{n}, \tag{2}$$

where $\tilde{\varphi} \in \widetilde{\mathrm{Ham}}(\mathbb{S}^2, \omega)$ is any lift of φ; we show in Proposition 3.6 that the above $\limsup$ is well-defined and that $\mu_d(\varphi)$ does not depend on the choice of $\tilde{\varphi} \in \widetilde{\mathrm{Ham}}(\mathbb{S}^2, \omega)$. We also define the related invariant $\zeta_d : C^\infty(\mathbb{S}^2) \to \mathbb{R}$ by

$$\zeta_d(H) := \limsup_{n \to \infty} \frac{c_d(nH)}{n}. \tag{3}$$

We will see that these two homogenized invariants are related by the formula

$$\mu_d(\varphi_H^1) = \zeta_d(H) - d \int_{\mathbb{S}^2} H\omega.$$

A useful property of any μ_d is that it coincides with (a multiple of) the Calabi invariant for Hamiltonian diffeomorphisms with small supports. More precisely, suppose that $\mathrm{supp}(\varphi)$, the support of $\varphi \in \mathrm{Ham}(\mathbb{S}^2, \omega)$, is contained in a topological disc D with $\mathrm{Area}(D) < \frac{1}{d+1}$. Then

$$\frac{1}{d} \mu_d(\varphi) = -\mathrm{Cal}(\varphi). \tag{4}$$

The above properties of μ_d, ζ_d will be proven in Section 3.

Remark 1.10. The properties of the μ_d are reminiscent of the *Calabi quasi-morphism* of Entov–Polterovich [10]. It is an open question whether $\mathrm{Ham}(\mathbb{S}^2, \omega)$ admits any Hofer continuous (homogeneous) quasi-morphisms other than the one constructed by Entov–Polterovich [10]. We plan to investigate in future work whether the invariants μ_d are quasi-morphisms.

The Hofer Lipschitz property and monotone twists A critical fact which we will show, and which is at the heart of all applications to Hofer's geometry, is the Hofer Lipschitz *property. For the invariants μ_d this means that*

$$|\mu_d(\varphi) - \mu_d(\psi)| \leq C_d d_H(\varphi, \psi)$$

for all $\varphi, \psi \in \mathrm{Ham}(\mathbb{S}^2, \omega)$. The Lipschitz constant is $C_d = d$. In particular, these invariants can be used to bound the Hofer distance from below.

In view of the Hofer Lipschitz property, to prove our results, we will have to produce examples of Hamiltonian diffeomorphisms whose invariants we can compute. This will be done by studying *monotone twist Hamiltonians*, that is, autonomous Hamiltonians $H : \mathbb{S}^2 \to \mathbb{R}$ of the form

$$H(z, \theta) = \tfrac{1}{2}h(z),$$

where $h' \geq 0$, $h'' \geq 0$, $h(-1) = h'(-1) = 0$; we developed a combinatorial model in our previous work [9] which can be used to compute the c_d for Hamiltonians like this under the additional technical assumption that $h'(1) \in \mathbb{N}$. For monotone twist Hamiltonians, the invariant ζ_d has a beautiful expression.

Proposition 1.11. *For any Hamiltonian H as above we have*

$$\zeta_d(H) = \tfrac{1}{2} \sum_{i=1}^{d} h\left(-1 + \frac{2i}{d+1}\right).$$

In other words, ζ_d is the sum of the values of H on d equally distributed horizontal circles. We learn from the above proposition that $\zeta_d(H)$ is at least as large as the value H takes on each of the d circles $\mathcal{C}_i = \{(z, \theta) : z = -1 + \frac{2i}{d+1}\}$, where $i \in \{1, \ldots, d\}$. This bears some resemblance to the notion of *heaviness* of equators introduced in the work of Entov–Polterovich [12]. What is surprising is that the circles $\mathcal{C}_i$ are all displaceable for $d \geq 2$, while heaviness of a set, as defined in [12], implies that the set is not displaceable by Hamiltonian diffeomorphisms. Sensitivity to the displaceable circles $\mathcal{C}_i$ is the distinguishing feature of our invariants μ_d, ζ_d which powers our applications to the Hofer geometry of the kernel of Calabi.

C^0-continuity and nonsimplicity of $\mathrm{Homeo}_0(\mathbb{S}^2, \omega)$

To prove Theorem 1.7, we need invariants which are continuous with respect to the C^0-topology. The invariants c_d and μ_d, while useful, are not in general C^0-continuous. We remedy this by taking certain linear combinations of the c_d to define C^0-continuous invariants

$$\eta_d : \mathrm{Ham}(\mathbb{S}^2, \omega) \to \mathbb{R}.$$

Not only are these invariants C^0-continuous, but also they extend continuously to $\mathrm{Homeo}_0(\mathbb{S}^2, \omega)$. Moreover, they are Hofer Lipschitz. We summarize the properties of the η_d in Proposition 3.9.

1.4. Relationship to previous work

As mentioned above, in our previous work we used PFH to define spectral invariants for compactly supported area-preserving diffeomorphisms and homeomorphisms of the two-disc. For all of the applications discussed here, we need to rework this theory over the two-sphere. In the disc case, we could assume that the maps were generated by a Hamiltonian that vanishes near the boundary of the disc. This is no longer possible, so new ideas are needed.

One idea here, familiar to specialists (see for example [10, 37, 47]), is to work with mean-normalized Hamiltonians. A careful analysis shows that this gives invariants which are well-defined on $\widetilde{\mathrm{Ham}}$; then, after homogenization as in the previous section, we can obtain invariants of Ham. These invariants would be enough to prove the theorems of Section 1.1. However, as stated above, they are not C^0-continuous, and so cannot be used to study the algebraic structure of the homeomorphism group. This is where the η_d, defined by taking a difference of spectral invariants, come in. The crucial insight for this, which was initially surprising to us, is that the c_d for even d descend from $\widetilde{\mathrm{Ham}}$ to Ham.

2. Preliminaries

In this section we fix our notation and introduce the necessary background on symplectic geometry and periodic Floer homology.

2.1. Recollections

Here we recall some basic facts about symplectic geometry and the Hofer distance.

Let (M, ω) be a symplectic manifold. Let $H \in C^\infty(\mathbb{S}^1 \times M)$ be a Hamiltonian; if M happens to be noncompact, then we consider only compactly supported Hamiltonians. We can think of H as a family of functions H_t on M, depending on time; we think of $\mathbb{S}^1$ as parametrized by $0 \le t \le 1$. Such an H gives rise to a possibly time-varying vector field X_{H_t} on M, called the *Hamiltonian vector field*, defined by

$$\omega(X_{H_t}, \cdot) = dH_t.$$

The flow of X_{H_t} is called the *Hamiltonian flow* and is denoted φ_H^t. The set of time-1 maps of Hamiltonian flows is called the set of *Hamiltonian diffeomorphisms* of M and denoted $\mathrm{Ham}(M, \omega)$; it forms a subgroup of the symplectomorphisms of (M, ω). We can define the *Hofer norm* $\|\varphi\|$ of any $\varphi \in \mathrm{Ham}(M, \omega)$ as follows. First, to a Hamiltonian $H \in C^\infty(\mathbb{S}^1 \times M)$, we associate the norm

$$\|H\|_{1,\infty} := \int_0^1 \left(\max_M(H_t) - \min_M(H_t) \right) dt.$$

We then define

$$\|\varphi\| := \inf \{\|H\|_{1,\infty} : \varphi = \varphi_H^1\}.$$

The above quantity is invariant under conjugation, i.e. $\|\psi^{-1}\varphi\psi\| = \|\varphi\|$. This follows from the fact that $\varphi^t_{H\circ\psi} = \psi^{-1}\varphi^t_H\psi$; see [19, Sec. 5.1, Prop. 1], for example.

Finally, we can define a metric on $\mathrm{Ham}(M, \omega)$, the *Hofer metric*, by

$$d_H(\varphi, \psi) = \|\varphi^{-1} \circ \psi\|.$$

As mentioned above, this yields a nondegenerate, bi-invariant metric, which is quite remarkable given the noncompactness of Ham. Nondegeneracy is what is difficult to prove and it was established by Hofer for $\mathbb{R}^{2n}$ [18], by Polterovich for rational symplectic manifolds [40], and by Lalonde–McDuff in full generality [30].

The bi-invariance of Hofer's distance also implies the following inequalities:

$$d_H(\varphi_1\varphi_2, \psi_1\psi_2) \leq d_H(\varphi_1, \psi_1) + d_H(\varphi_2, \psi_2), \tag{5}$$

$$d_H(\varphi, \psi^{-1}\varphi\psi) \leq 2d_H(\psi, \mathrm{Id}). \tag{6}$$

Indeed, (5) follows from (9) below, and (6) is proved as follows:

$$d_H(\varphi, \psi^{-1}\varphi\psi) = \|\varphi^{-1}\psi^{-1}\varphi\psi\| \leq \|\varphi^{-1}\psi^{-1}\varphi\| + \|\psi\| = 2d_H(\psi, \mathrm{Id}).$$

Now let $M = \mathbb{S}^2 = \{(x, y, z) \in \mathbb{R}^3 : x^2 + y^2 + z^2 = 1\}$. This has a symplectic form $\omega := \frac{1}{4\pi}d\theta \wedge dz$, where (θ, z) are cylindrical coordinates. We let $\mathrm{Diff}(S^2, \omega)$ denote the set of smooth diffeomorphisms φ such that $\varphi^*\omega = \omega$. In fact, $\mathrm{Diff}(S^2, \omega) = \mathrm{Ham}(\mathbb{S}^2, \omega)$. The Hofer geometry of $\mathrm{Diff}(S^2, \omega)$, with this identification implied, will be the topic of study in the present work. We recall, for later use, that the fundamental group of $\mathrm{Ham}(\mathbb{S}^2, \omega)$ is $\mathbb{Z}/2\mathbb{Z}$ and is generated by Rot, the full rotation around the North-South axis of the sphere; for a proof of this see, for example, [42]; the Hamiltonian $H(\theta, z) = \frac{1}{2}z$ generates this full rotation.

We will denote the universal cover of $\mathrm{Ham}(\mathbb{S}^2, \omega)$ by $\widetilde{\mathrm{Ham}}(\mathbb{S}^2, \omega)$. This can be described as the set of Hamiltonian paths, considered up to homotopy relative to endpoints; here, by a *Hamiltonian path*, we mean a path of Hamiltonian diffeomorphisms $\{\varphi^t, 0 \leq t \leq 1\}$. This is a two-fold covering, by the discussion in the previous paragraph.

We next recall the *displacement energy* of a subset $A \subset \mathbb{S}^2$. This is by definition the quantity

$$e(A) := \inf\{\|\phi\| : \phi(A) \cap \bar{A} = \emptyset\}.$$

It is known that for a disjoint union of closed discs, each with area a and whose union covers less than half the area of the sphere, the displacement energy is a. We will need the following lemma in Section 4.3.

Lemma 2.1. *Let $D, D' \subset \mathbb{S}^2$ be disjoint closed discs of equal area. Then*

$$\inf\{\|\phi\| : \phi(D) = D'\} = \mathrm{Area}(D).$$

Proof. Let us denote $a := \mathrm{Area}(D)$ and $E := \inf\{\|\phi\| : \phi(D) = D'\}$. It follows from the above discussion of displacement energy that $E \geq a$.

For the reverse inequality, note that the same discussion also implies that for any $\varepsilon > 0$, there exists $\psi \in \mathrm{Ham}(\mathbb{S}^2, \omega)$ with $\|\psi\| < a + \varepsilon$ and $\psi(D) \cap D = \emptyset$. Since $\psi(D)$

and D' have the same area and are both contained in $\mathbb{S}^2 \setminus D$, there exists a Hamiltonian diffeomorphism χ, supported in $\mathbb{S}^2 \setminus D$, which maps $\psi(D)$ onto D'. The assumption on the support implies that $\chi^{-1}(D) = D$. We now pick $\phi = \chi \psi \chi^{-1}$. We see that $\phi(D) = D'$ and by conjugation invariance of the Hofer norm we have $\|\phi\| = \|\psi\| < a + \varepsilon$. Since such a diffeomorphism ϕ may be found for any $\varepsilon > 0$, this shows the reverse inequality $E \leq a$. ∎

Next, we review the definition of the *Calabi homomorphism*

$$\mathrm{Cal} : \mathrm{Ham}_U(\mathbb{S}^2, \omega) \to \mathbb{R},$$

alluded to in the introduction. Recall that, for open $U \subsetneq \mathbb{S}^2$, we denote by $\mathrm{Ham}_U(\mathbb{S}^2, \omega)$ the subgroup of $\mathrm{Ham}(\mathbb{S}^2, \omega)$ consisting of Hamiltonian diffeomorphisms which are supported in U. Given $\varphi \in \mathrm{Ham}_U(\mathbb{S}^2, \omega)$, define

$$\mathrm{Cal}(\varphi) = \int_{\mathbb{S}^1} \int_{\mathbb{S}^2} H(t, \cdot) \omega \, dt, \tag{7}$$

where $H \in C^\infty(\mathbb{S}^1 \times \mathbb{S}^2)$ is any Hamiltonian supported in U whose time-1 flow is φ. It is well-known that $\mathrm{Cal}(\varphi)$ does not depend on the choice of H, and moreover $\mathrm{Cal} : \mathrm{Ham}_U(\mathbb{S}^2, \omega) \to \mathbb{R}$ is a group homomorphism; see [5, 36] for further details.

In Section 5, we will also consider the group $\mathrm{Homeo}_0(\mathbb{S}^2, \omega)$ of *area and orientation preserving homeomorphisms* of $\mathbb{S}^2$. This is defined to be the group of homeomorphisms of $\mathbb{S}^2$ preserving the measure induced by ω in the component of the identity. It has a distance d_{C^0}, called the C^0-*distance*, defined by picking a Riemannian metric d on $\mathbb{S}^2$, and defining

$$d_{C^0}(\varphi, \psi) = \sup_{x \in M} d(\varphi(x), \psi(x)).$$

We remark for later use that $\mathrm{Diff}(\mathbb{S}^2, \omega)$ sits densely in the C^0-distance in $\mathrm{Homeo}_0(\mathbb{S}^2, \omega)$.

2.2. The spectrum

We now recall the action spectrum, defined in [9, Sec. 2.5]. Let $H \in C^\infty(\mathbb{S}^1 \times \mathbb{S}^2)$. Recall the *action functional* associated to H,

$$\mathcal{A}_H(z, u) = \int_0^1 H(t, z(t)) \, dt + \int_{D^2} u^* \omega, \tag{8}$$

defined for capped loops (z, u). The critical points of $\mathcal{A}_H$ are pairs (z, u), where z is a 1-periodic orbit of φ_H^t, and the set of associated critical values is called the *action spectrum* $\mathrm{Spec}(H)$ of H. The forthcoming PFH spectral invariants will take values in the *order d action spectrum* of H, defined by

$$\mathrm{Spec}_d(H) := \bigcup_{k_1 + \cdots + k_j = d} (\mathrm{Spec}(H^{k_1}) + \cdots + \mathrm{Spec}(H^{k_j})),$$

where H^k denotes the k-fold composition of H with itself. Here, the *composition* is defined by

$$(G \# H)(t, x) = \begin{cases} 2\rho'(2t)H_{\rho(2t)}(x) & \text{if } t \in [0, 1/2], \\ 2\rho'(2t-1)G_{\rho(2t-1)}(x) & \text{if } t \in [1/2, 1], \end{cases}$$

where $\rho : [0, 1] \to [0, 1]$ is a fixed nondecreasing smooth function which is equal to 0 near 0 and to 1 near 1. Note that we do not need H and G to be 1-periodic to define the composition, and even if they are not 1-periodic, $G \# H$ will still be, since it is zero for t close to 0 and 1. The time-1 map of $G \# H$ is $\varphi_G^1 \circ \varphi_H^1$. Note that for any Hamiltonians G_1, G_2, H_1, H_2, we have

$$\|G_1 \# H_1 - G_2 \# H_2\|_{1,\infty} = \|G_1 - G_2\|_{1,\infty} + \|H_1 - H_2\|_{1,\infty}. \tag{9}$$

We state here some of the properties of the order d action spectrum which will be used in the following sections. Recall that $H \in C^\infty(\mathbb{S}^1 \times \mathbb{S}^2)$ is said to be *mean-normalized* if $\int_{\mathbb{S}^2} H(t, \cdot)\omega = 0$ for all $t \in \mathbb{S}^2$. Two Hamiltonians H_0, H_1 are said to be *homotopic* if there exists a smooth path of Hamiltonians connecting H_0 to H_1 such that $\varphi_{H_0}^1 = \varphi_{H_s}^1 = \varphi_{H_1}^1$ for all $s \in [0, 1]$. In other words, the Hamiltonian paths $\{\varphi_{H_0}^t\}$ and $\{\varphi_{H_1}^t\}$, for $0 \le t \le 1$, coincide as elements of the universal cover $\widetilde{\text{Ham}}(\mathbb{S}^2, \omega)$. Here is a list of properties of Spec_d which will be needed.

(i) *Symplectic invariance:*

$$\text{Spec}_d(H \circ \psi) = \text{Spec}_d(H)$$

for all $H \in C^\infty(\mathbb{S}^1 \times \mathbb{S}^2)$ and $\psi \in \text{Ham}(\mathbb{S}^2, \omega)$.

(ii) *Homotopy invariance:* If H_0, H_1 are mean-normalized and homotopic, then $\text{Spec}_d(H_0) = \text{Spec}_d(H_1)$.

(iii) *Measure zero:* $\text{Spec}_d(H)$ is of measure zero.

The above properties are well-known in the case of $\text{Spec}(H)$, that is, when $d = 1$; see for example [37]. It is not difficult to see that the first two properties follow from the case $d = 1$: symplectic invariance follows from the identity $(H \circ \psi)^k = H^k \circ \psi$ for any $k \in \mathbb{N}$, and homotopy invariance is a consequence of the fact that H_0^k, H_1^k are mean-normalized and homotopic, for any $k \in \mathbb{N}$, if H_0 and H_1 are. As we will now explain, the third property also follows from the $d = 1$ case. As a consequence of the definition of $\text{Spec}_d(H)$, it is sufficient to prove that the set $\text{Spec}(H^{k_1}) + \cdots + \text{Spec}(H^{k_j})$ is of measure zero for any $k_1, \ldots, k_j$ with $k_1 + \cdots + k_j = d$. To that end, let $(M, \omega \oplus \cdots \oplus \omega)$ be the symplectic manifold obtained by taking the j-fold product of $(\mathbb{S}^2, \omega)$ and consider the Hamiltonian $F : \mathbb{S}^1 \times M \to \mathbb{R}$ defined by

$$F(t, x_1, \ldots, x_j) = H^{k_1}(t, x_1) + \cdots + H^{k_j}(t, x_j).$$

We conclude that $\text{Spec}(H^{k_1}) + \cdots + \text{Spec}(H^{k_j})$ has measure zero by observing that it coincides with $\text{Spec}(F)$ which we know has measure zero.

2.3. Definition of PFH

We now recall the definition of periodic Floer homology (PFH), for example from [23], which is a tool that will be central in our work. While PFH can be defined over any surface, for simplicity we consider the case where our surface is $\mathbb{S}^2$, which is the only case that is relevant for the present work.

We start with some preliminaries. Let $\varphi \in \mathrm{Diff}(\mathbb{S}^2, \omega)$. Given φ, we can define the *mapping torus*

$$Y_\varphi := \mathbb{S}^2 \times [0,1]_t / \sim, \quad (x, 1) \sim (\varphi(x), 0).$$

This has a natural vector field $R := \partial_t$, which we call the *Reeb vector field*, a natural one-form dt, and a natural two-form ω_φ induced from the area form ω. The pair (dt, ω_φ) is a *stable Hamiltonian structure* in the sense of for example [3, 7, 26, 53]. The manifold Y_φ has a plane field ξ defined to be the vertical tangent bundle for the fibration $\pi : Y_\varphi \to \mathbb{S}^1$.

We will be interested in closed integral curves

$$\alpha : \mathbb{R}/T\mathbb{Z} \to Y_\varphi$$

of R, modulo reparametrization of the domain, which we call *closed orbits*; we can identify an embedded closed orbit with its image. A closed orbit α has an integral *degree* $d(\alpha) := \pi_*[\alpha] \in H_1(\mathbb{S}^1) = \mathbb{Z}$. The *linearized return map* P_α for a closed orbit α is defined for any $p \in \alpha$ as the linearization of the time-T flow of R on $\xi|_p$. A closed orbit is called *nondegenerate* if 1 is not an eigenvalue of the linearized return map; a nondegenerate closed orbit is called *hyperbolic* if the eigenvalues of P_α are real, and *elliptic* if the eigenvalues lie on the unit circle; these definitions do not depend on the choice of p.

Define an *orbit set* $\alpha := \{(\alpha_i, m_i)\}$ to be a finite set, where the α_i are distinct embedded closed orbits of R, and the m_i are positive integers. The *degree* of the orbit set α is the sum of the degrees of the α_i. The map φ is d-*nondegenerate* if every closed orbit with degree at most d is nondegenerate; this is a generic condition. A degree d orbit set for a d-nondegenerate φ is called *admissible* if $m_i = 1$ whenever α_i is hyperbolic.

Let $X = \mathbb{R}_s \times Y_\varphi$. This has a natural symplectic form

$$\omega = ds \wedge dt + \omega_\varphi.$$

The pair (X, ω) is called the *symplectization* of Y_φ. Recall that an *almost complex structure* on X is a smooth bundle map $J : TX \to TX$ such that $J^2 = -1$. A J-*holomorphic curve* in X is a map $u : (\Sigma, j) \to (X, J)$ satisfying the equation

$$du \circ j = J \circ du.$$

Here, Σ is a closed (possibly disconnected) Riemann surface minus a finite number of punctures, and the map u is assumed to be asymptotic to Reeb orbits near the punctures, see for example [22] for the precise definition.

The *periodic Floer homology* $PFH(\mathbb{S}^2, \varphi, d)$ is the homology of a chain complex $PFC(\mathbb{S}^2, \varphi, d)$. The chain complex $PFC(\mathbb{S}^2, \varphi, d)$ is freely generated over $\mathbb{Z}_2$ by admissible orbit sets α of degree $d > 0$. The chain complex differential ∂ counts J-holomorphic

curves in X, for generic admissible J; here, an almost complex structure is called *admissible* if it preserves ξ, is $\mathbb{R}$-invariant, sends ∂_s to R, and its restriction to ξ is tamed by ω_φ. More precisely,

$$\langle \partial\alpha, \beta \rangle = \#\mathcal{M}_J^{I=1}(\alpha, \beta),$$

where I denotes the ECH index, defined below, we are considering curves in X up to equivalence of currents and modulo translation in the $\mathbb{R}$ direction, and $\#$ denotes the mod 2 count.

It is shown in [24, 25] that[10] $\partial^2 = 0$, so the homology is well-defined; it is shown in [33] that it agrees with a version of Seiberg–Witten Floer cohomology and in particular is independent of φ.

To define spectral invariants, we will want to use a *twisted* version of PFH, denoted $\widetilde{PFH}(\mathbb{S}^2, \varphi, d)$; as we will see in §3.1.1, the twisted PFH carries a natural action filtration which we will use to define the spectral invariants. To define twisted PFH, let γ be any degree 1 cycle in Y_φ, transverse to ξ; choose a homotopy class of trivializations τ_0 on $\xi|_\gamma$. The twisted PFH chain complex $\widetilde{PFC}$ is generated by pairs (α, Z), called *twisted PFH generators*, where α is a degree d admissible orbit set, and $Z \in H_2(Y_\varphi, \alpha, d\gamma)$. The differential counts $I = 1$ curves C from (α, Z) to (β, Z'), namely curves $C \in \mathcal{M}_J^{I=1}(\alpha, \beta)$, such that

$$[C] + Z' = Z.$$

For each d, there is a grading, defined below, which we call the k-*grading*. The homology is an invariant, and so can be computed, with the result that for $d \geq 0$ we have

$$\widetilde{PFH}_*(\mathbb{S}^2, \varphi, d) = \begin{cases} \mathbb{Z}_2 & \text{if } * = d \bmod 2, \\ 0 & \text{otherwise.} \end{cases} \tag{10}$$

The above identity can be proven via a direct computation when φ is taken to be an irrational rotation of the sphere; for more details see, for example, [9, Sec. 3.3]. We now define the ECH index I, and the grading k.

The ECH index I depends only on the relative homology class $A \in H_2(Y_\varphi, \alpha, \beta)$ between two orbit sets. We have

$$I(A) = c_\tau(A) + Q_\tau(A) + CZ_\tau^I(A), \tag{11}$$

where τ denotes a homotopy class of trivializations of ξ over all Reeb orbits, $c_\tau(A)$ denotes the *relative Chern class* of ξ restricted to A, $Q_\tau(A)$ denotes the *relative self-intersection*, and CZ_τ^I denotes the total Conley–Zehnder index. We will not need the precise definitions of these terms in the present work, so we omit them for brevity, referring the reader to [21] for the details.

We can define the promised k-grading. The definitions of the relative Chern class and relative self-intersection extend verbatim to relative homology classes $A \in H_2(Y_\varphi, \alpha, d\gamma)$,

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[10]More precisely, [25] proves that the differential in embedded contact homology squares to zero. As pointed out in [24, 33] this proof carries over, nearly verbatim, to our setting.

once a trivialization τ over the simple orbits in α and a trivialization τ_0 over γ have been chosen. With the preceding understood, we now define

$$k(\alpha, Z) := c_{\tau,\tau_0}(Z) + Q_{\tau,\tau_0}(Z) + CZ_\tau^I(\alpha).$$

To simplify the notation, we will denote $k(\alpha, Z)$ by $I(Z)$ below.

3. The spectral invariants

We now use the twisted PFH to define various invariants. We begin by summarizing for the reader what will be done in this section.

To set the stage for what is coming, it is helpful to recall what was done in [9, Sec. 3.4]. There, we defined spectral invariants $c_{d,k}(H)$ for $H \in \mathcal{H}$ where

$$\mathcal{H} := \{H \in C^\infty(\mathbb{S}^1 \times \mathbb{S}^2) : \varphi_H^t(p_-) = p_-, \; H(t, p_-) = 0, \; \forall t \in [0, 1],$$
$$-1/4 < \text{rot}(\{\varphi_H^t\}, p_-) < 1/4\},$$

where $\text{rot}(\{\varphi_H^t\}, p_-)$ is the rotation number of the isotopy $\{\varphi_H^t\}_{t\in[0,1]}$ at p_-. It was shown in addition that these invariants depend only on the time-1 map. Spectral invariants for compactly supported disc maps were then defined by identifying the disc with the northern hemisphere.

Our goal now is to define spectral invariants for all $H \in C^\infty(\mathbb{S}^1 \times \mathbb{S}^2)$ and to find invariants that depend only on $\varphi \in \text{Ham}(\mathbb{S}^2, \omega)$, rather than on a choice of generating Hamiltonian. Here is how we do this. First we extend the procedure in [9] from $H \in \mathcal{H}$ to arbitrary H to get invariants $c_{d,k}$, defined when k and d have the same parity. These $c_{d,k}$ extend the $c_{d,k}$ from our previous work: that is, if $H \in \mathcal{H} \subset C^\infty(\mathbb{S}^1 \times \mathbb{S}^2)$, then the definition of $c_{d,k}(H)$ here agrees with that in [9]. Similarly to our previous work, we can then define $c_d := c_{d,-d}$. This choice of $k = -d$ is not quite canonical (see Remark 3.4), but is convenient and suffices for our purposes: what is crucial is that $c_d(0) = 0$.

As alluded to in the introduction, these c_d are in general *not* invariants of the time-1 map, and so are not well-suited on their own for proving our main theorems. However, we can use the c_d to form new invariants. First, we show that the c_d for even d are invariants when we restrict to mean-normalized Hamiltonians; similarly, the homogenizations μ_d, ζ_d are also invariants restricted to mean-normalized Hamiltonians. None of these invariants are C^0-continuous, so we use a linear combination of the c_d for d even to define another sequence η_d.

Thus, to summarize for the ease of the reader, the main product of this section are invariants c_d and η_d defined for d even, and μ_d, ζ_d defined for all d, together with proofs of their properties that we will need. The μ_d and ζ_d are related by formula (25). The μ_d are used to prove Theorem 1.4, while the η_d are used to prove Theorem 1.7; the c_d are used to construct the μ_d and the η_d.

3.1. Invariants for Hamiltonians

We begin by introducing PFH spectral invariants $c_{d,k}(H)$ for Hamiltonians $H \in C^\infty(\mathbb{S}^1 \times \mathbb{S}^2)$. This requires first recalling a construction of Hutchings, assigning a spectral invariant to every nonzero twisted PFH class.

3.1.1. The nondegenerate case.
A Hamiltonian $H \in C^\infty(\mathbb{S}^1 \times \mathbb{S}^2)$ is called d-*nondegenerate* if its time-1 flow $\varphi = \varphi_H^1$ is d-nondegenerate. We now explain how to define PFH spectral invariants for d-nondegenerate Hamiltonians by extending the definition in [9] in a natural way.

We begin by explaining the aforementioned construction of Hutchings for assigning a spectral invariant to a nonzero twisted PFH class. A twisted PFH generator has an *action* defined by

$$\mathcal{A}(\alpha, Z) = \int_Z \omega_\varphi.$$

The differential decreases the action (see for example [9, Sec. 3.3]), so the action induces a filtration on the twisted PFH chain complex: we can define $\widetilde{PFC}^L$ to be the subcomplex generated by twisted PFH generators with action no more than L. Denote the homology of this complex by $\widetilde{PFH}^L$. For any nonzero class $\sigma \in \widetilde{PFH}(\mathbb{S}^2, \varphi, d)$, we can now define $c_\sigma(\varphi, \gamma, \tau_0)$ to be the smallest L such that σ is in the image of the inclusion induced map

$$\widetilde{PFH}^L \to \widetilde{PFH}.$$

We can think of this as the minimum action required to represent σ.

The number $c_\sigma(\varphi, \gamma, \tau_0)$ depends on the choice of reference cycle γ and trivialization τ_0 over γ; we will now define the *PFH spectral invariants* associated to a d-nondegenerate Hamiltonian H by using the Hamiltonian flow to fix a natural reference cycle.

To make this precise, let H be a d-nondegenerate Hamiltonian and write $\varphi = \varphi_H^1$. Consider the trivialization

$$\Psi_H : \mathbb{S}^1 \times \mathbb{S}^2 \to Y_\varphi, \quad (t, x) \mapsto ((\varphi_H^t)^{-1}(x), t). \tag{12}$$

Define $\gamma_H = \Psi_H(\mathbb{S}^1 \times \{p_-\})$. This is trivialized by the pushforward τ_H of an $\mathbb{S}^1$-invariant trivialization over p_-. We will now use the twisted PFH chain complex for Y_φ, with respect to the reference cycle γ_H, to define the spectral invariants.

Assume first that H vanishes at p_- for all time. For each $d \in \mathbb{N}$, we define

$$c_{d,k}(H) := c_\sigma(\varphi_H^1, \gamma_H, \tau_H), \quad d \equiv k \bmod 2,$$

where σ is the unique nonzero class in $\widetilde{PFH}_k(\mathbb{S}^2, \varphi, d)$. We emphasize that, even fixing the Hamiltonian diffeomorphism, this can and will depend on H, since the trivialized reference cycle γ_H does. We note that for such an H,

$$c_{d,k}(H) = \mathcal{A}(\alpha, Z) \tag{13}$$

for some twisted PFH generator (α, Z). Indeed, as explained in [9, Sec. 3.3], this follows from the fact that the subset $\{\mathcal{A}(\alpha, Z) : (\alpha, Z) \in \widehat{PFC}(\varphi, d)\} \subset \mathbb{R}$ is discrete, since under our nondegeneracy assumption there are only finitely many orbit sets of degree d.

Finally, for arbitrary H we reduce to the case of H vanishing at p by demanding that the *shift property*, stated in Proposition 3.2 below, hold. This says that

$$c_{d,k}(H + h) = c_{d,k}(H) + d \int_0^1 h(t)\, dt, \tag{14}$$

when $h : \mathbb{S}^1 \to \mathbb{R}$ is any function.

In principle, $c_{d,k}(H)$ could depend on the choice of admissible J, but we will see by the monotonicity property below that it does not.

3.1.2. Key properties. We now prove that the PFH spectral invariants have the following key properties and extend to all, possibly degenerate, Hamiltonians.

Theorem 3.1. *The PFH spectral invariant $c_{d,k}(H)$ admits a unique extension to all $H \in C^\infty(\mathbb{S}^1 \times \mathbb{S}^2)$ such that the extended spectral invariant*

$$c_{d,k} : C^\infty(\mathbb{S}^1 \times \mathbb{S}^2) \to \mathbb{R}$$

has the following properties:

(1) Continuity: *For any $H, G \in C^\infty(\mathbb{S}^1 \times \mathbb{S}^2)$, we have*

$$d \int_{\mathbb{S}^1} \min(H_t - G_t)\, dt \leq c_{d,k}(H) - c_{d,k}(G) \leq d \int_{\mathbb{S}^1} \max(H_t - G_t)\, dt.$$

(2) Spectrality: $c_{d,k}(H) \in \operatorname{Spec}_d(H)$.

Before giving the proof, we note that the second item of the theorem implies that if H, G vanish at p_-, then

$$|c_{d,k}(H) - c_{d,k}(G)| \leq d \|H - G\|_{1,\infty}, \tag{15}$$

which is an alternative variant of the Hofer continuity property.

Proof of Theorem 3.1. The proof proceeds along similar lines to [9, Thm. 3.6].

Step 1: Reducing to the d-nondegenerate case. We now assume that the theorem has been proved for d-nondegenerate H, and explain how this implies the result for all H. Given any H, take any sequence of d-nondegenerate H_i which C^2-converges to H, and define

$$c_{d,k}(H) = \lim_{i \to \infty} c_{d,k}(H_i). \tag{16}$$

This limit exists, and does not depend on the choice of approximating H_i, due to the continuity property with $H = H_i$ and $G = H_j$. The same inequality implies that the extension from d-nondegenerate H is unique as claimed; the continuity and shift properties for d-nondegenerate H imply these properties for all H. Spectrality for d-nondegenerate H implies spectrality for all H by Arzelà–Ascoli.

Step 2: Reducing to Hamiltonians that vanish at p_-. It remains to prove continuity and spectrality in the nondegenerate case.

We now show that by using the shift property (14), it suffices to prove these properties for Hamiltonians vanishing at p_-. We begin with continuity. Consider arbitrary H, G. Then we can write

$$H = \tilde{H} + h, \quad G = \tilde{G} + g, \tag{17}$$

where h and g are defined as the restrictions of H, G to p_-, and $\tilde{H}, \tilde{G}$ vanish on p_-. Then, by the shift property,

$$c_{d,k}(H) - c_{d,k}(G) = c_{d,k}(\tilde{H}) - c_{d,k}(\tilde{G}) + d \int_{\mathbb{S}^1} (h(t) - g(t)) \, dt.$$

Thus, if continuity holds for $\tilde{H}$ and $\tilde{G}$, then we have

$$c_{d,k}(H) - c_{d,k}(G) \le d \int_{\mathbb{S}^1} \max(\tilde{H}_t - \tilde{G}_t) \, dt + d \int_{\mathbb{S}^1} (h(t) - g(t)) \, dt.$$

Now, since h, g only depend on t, we have

$$\max(\tilde{H}_t - \tilde{G}_t) = \max(H_t - G_t) + g(t) - h(t).$$

Combining this equality with the previous inequality proves the rightmost inequality required for continuity. Similarly, if continuity holds for $\tilde{H}$ and $\tilde{G}$, then we have

$$c_{d,k}(H) - c_{d,k}(G) \ge d \int_{\mathbb{S}^1} \min(\tilde{H}_t - \tilde{G}_t) \, dt + d \int_{\mathbb{S}^1} (h(t) - g(t)) \, dt,$$

and we know that

$$\min(\tilde{H}_t - \tilde{G}_t) = \min(H_t - G_t) + g(t) - h(t),$$

which yields the leftmost inequality required for continuity to hold.

Similarly, if spectrality holds for $\tilde{H}$ in (17), then it holds for H by the shift property, because the addition of h does not change the set of critical points of $\mathcal{A}_H$, hence by (8), $\mathrm{Spec}_d(H) = \mathrm{Spec}_d(\tilde{H}) + d \int_{\mathbb{S}^1} h(t) \, dt$.

Thus, we can assume H and G vanish at p_-.

Step 3: Continuity when H and G vanish at p_-. Under (12), the stable Hamiltonian structure (dt, ω_φ) is of the form $(dt, \omega + dH \wedge dt)$, and $R = \partial_t + X_H$. The natural symplectic form on the symplectization $X = \mathbb{R} \times Y_\varphi$ under (12) is

$$\omega_H = ds \wedge dt + \omega + dH \wedge dt,$$

where s is the coordinate on $\mathbb{R}$. We henceforth identify Y_φ with $\mathbb{S}^1 \times \mathbb{S}^2$ using (12), we implicitly identify orbit sets on Y_φ with the corresponding orbit sets on $\mathbb{S}^1 \times \mathbb{S}^2$, and we identify the trivialized reference cycle (γ_H, τ_H) with the $\mathbb{S}^1$-invariant trivialized cycle γ over p_-.

Given H and G, we pick a function β which is 0 for sufficiently small s, is 1 for s sufficiently large, and satisfies $1 + \beta'(H - G) > 0$, we define $K = G + \beta(s)(H - G)$, and we consider the form

$$\omega_X = ds \wedge dt + \omega + d(K dt),$$

which is symplectic and agrees with ω_H for sufficiently positive s and with ω_G for sufficiently negative s.

The general theory of (twisted) PFH cobordism maps, as developed by Chen [6], guarantees a chain map $\Psi_{H,G}$ between the twisted PFH chain complexes for H and G, counting ECH index zero J_X-holomorphic buildings from (α, Z) to (β, Z'), and inducing an isomorphism, where J_X is a fibration compatible almost complex structure on X, in the sense that it preserves the vertical tangent bundle and its ω_X-orthogonal complement.

So, given $d \geq 1$ and $k \in \mathbb{Z}$ of the same parity, let $(\alpha_1, Z_1) + \cdots + (\alpha_m, Z_m)$ be a cycle in $\widetilde{PFC}(\varphi_H^1, d)$ representing $\sigma_{d,k}$ with

$$c_{\sigma_{d,k}}(\varphi_H^1) = \mathcal{A}(\alpha_1, Z_1) \geq \cdots \geq \mathcal{A}(\alpha_m, Z_m)$$

and let (β, Z') be a generator in $\widetilde{PFC}(\varphi_G^1, d)$ with maximal action among the elements of the support of $\Psi_{H,G}((\alpha_1, Z_1) + \cdots + (\alpha_m, Z_m))$.

Thus, we have a J_X-holomorphic building C from some (α_i, Z_i), which we will denote by (α, Z), to (β, Z'). Since, just as in [9], our argument only involves action and index considerations, we can assume that C consists of a single level, and we know that

$$Z' + [C] = Z$$

as elements of $H_2(\mathbb{S}^1 \times \mathbb{S}^2, \alpha, d\gamma)$. Hence, as $I([C]) = 0$, we must have $I(Z) = I(Z') = k$, so that

$$c_{d,k}(\varphi_H^1) - c_{d,k}(\varphi_G^1) \geq \mathcal{A}(\alpha, Z) - \mathcal{A}(\beta, Z'). \tag{18}$$

We now claim that

$$\mathcal{A}(\alpha, Z) - \mathcal{A}(\beta, Z') = \int_C (\omega + dK \wedge dt + K' ds \wedge dt), \tag{19}$$

where K' denotes the derivative with respect to s, and for the rest of this section dK denotes the derivative in the $\mathbb{S}^2$ direction.

The proof of this is just as in [9, Lem. 3.8]. Indeed, as in that proof, we have

$$\mathcal{A}(\alpha, Z) = \int_Z (\omega + d(H dt)), \quad \mathcal{A}(\beta, Z') = \int_{Z'} (\omega + d(G dt)),$$

and

$$\int_C \omega = \int_Z \omega - \int_{Z'} \omega.$$

Moreover, $\int_C d(K dt) = \int_Z d(H dt) - \int_{Z'} d(G dt)$, since H and G vanish on γ. So, putting all this together, we have

$$\mathcal{A}(\alpha, Z) - \mathcal{A}(\beta, Z') = \int_C (\omega + d(K dt)),$$

proving (19).

Moreover, we have $\int_C (\omega + dK \wedge dt) \geq 0$, since as in the proof of [9, Lem. 3.8], the form $\omega + dK \wedge dt$ is pointwise nonnegative along C, and so in fact we obtain

$$\mathcal{A}(\alpha, Z) - \mathcal{A}(\beta, Z') \geq \int_C K' ds \wedge dt. \tag{20}$$

The argument in [9, Lem. 3.8] also shows that $ds \wedge dt$ is pointwise nonnegative on C.

Now we have

$$\int_C K' ds \wedge dt = \int_C \beta'(s)(H - G) ds \wedge dt \geq \int_C \beta'(s) \min(H_t - G_t) ds \wedge dt,$$

since $ds \wedge dt$ is pointwise nonnegative along C. We can evaluate the rightmost integral above by projecting to the (s, t) plane; this projection has degree d, and $\int \beta' = 1$, so the above inequality combined with (20) and (18) gives the leftmost inequality required for continuity.

To prove the other inequality, we switch the roles of H and G in the above argument, and again invoke the corresponding versions of (18) and (20) to get

$$c_{d,k}(\varphi_G^1) - c_{d,k}(\varphi_H^1) \geq \int_C \beta'(s)(G - H) ds \wedge dt,$$

hence

$$c_{d,k}(\varphi_H^1) - c_{d,k}(\varphi_G^1) \leq \int_C \beta'(s)(H - G) ds \wedge dt \leq \int_C \beta'(s) \max(H_t - G_t) ds \wedge dt,$$

where in the rightmost inequality we have used the fact that $ds \wedge dt$ is pointwise non-negative. We then project to the (s, t) plane as above to obtain the rightmost inequality required for continuity.

Step 4: Spectrality when H vanishes at p_-. Since H vanishes at p_-, we know by (13) that any $c_{d,k}(H)$ equals $\mathcal{A}(\alpha, Z)$ for some twisted PFH generator (α, Z).

Recall that $\mathcal{A}(\alpha, Z)$ is the action of some relative homology class. We first construct a particular homology class Z_α from a periodic orbit α, and show that the action of this class lies in the action spectrum. More precisely, let x be a q-periodic point of $\varphi = \varphi_H^1$, and pick a capping disc u for the orbit $\gamma(t) = (\varphi_H^t(x))_{t \in [0,q]}$ such that $u(0, 0) = p_-$. Equip the disc with polar coordinates (θ, ρ) with $\theta \in \mathbb{R}/q\mathbb{Z}$, $\rho \leq 1$, and then consider the homology class Z_α represented by

$$\mathbb{R}/q\mathbb{Z} \times [0, 1] \to \mathbb{S}^1 \times \mathbb{S}^2, \quad (\theta, \rho) \mapsto (\theta \bmod 1, u(\theta, \rho)).$$

We now compute

$$\mathcal{A}(\beta, Z_\alpha) = \int_{Z_\alpha} (\omega + dH \wedge dt) = \int_{Z_\alpha} (\omega + d(Hdt))$$

$$= \int u^*\omega + \int_{\partial Z_\alpha} Hdt = \int u^*\omega + \int_0^q H_t(\gamma(t))\, dt$$

$$= \mathcal{A}_H(\gamma, u) \in \mathrm{Spec}_q(H),$$

where, in the third equality above, we have used the fact that H vanishes at p_-.

Now, given an arbitrary (α, Z), write $\alpha = \{(\alpha_i, q_i)\}$. We can write $Z = \sum Z_{\alpha_i} + y[\mathbb{S}^2]$. Then

$$\mathcal{A}(\alpha, Z) = y + \sum \mathcal{A}(\alpha_i, Z_{\alpha_i}).$$

The right hand side is an element of $\mathrm{Spec}_d(H)$, since we can for example absorb the y into the capping of any particular orbit. Hence, $c_{d,k}(H) \in \mathrm{Spec}_d(H)$. ∎

We now collect some additional useful properties of the $c_{d,k}$.

Proposition 3.2. *The spectral invariant $c_{d,k} : C^\infty(\mathbb{S}^1 \times \mathbb{S}^2) \to \mathbb{R}$ satisfies:*

(1) Normalization: $c_{d,k}(0) = 0$ *for* $-d \le k \le d$.

(2) Monotonicity: *Suppose that* $H \le G$. *Then*

$$c_{d,k}(H) \le c_{d,k}(G).$$

(3) Shift: *Let* $h : \mathbb{S}^1 \to \mathbb{R}$ *be a function of time. Then*

$$c_{d,k}(H + h) = c_{d,k}(H) + d \int_{\mathbb{S}^1} h(t)\, dt.$$

(4) Symplectic invariance: $c_{d,k}(H \circ \psi) = c_{d,k}(H)$ *for any* $\psi \in \mathrm{Ham}(\mathbb{S}^2, \omega)$.

(5) Homotopy invariance: *If* H_0 *and* H_1 *are mean-normalized and homotopic, then* $c_{d,k}(H_0) = c_{d,k}(H_1)$.

(6) Support control: *If the support of H is contained in a topological disc D with* $\mathrm{Area}(D) < \frac{1}{d+1}$, *and* $-d \le k \le d$, *then* $|c_{d,k}(H)| \le 2d\,\mathrm{Area}(D)$.

Proof. Normalization follows from our previous work [9, Thm. 3.6], since as mentioned previously the $c_{d,k}$ extend the spectral invariants we defined there. The shift property is immediate from the definition. The monotonicity property follows formally from continuity: indeed, by continuity we have

$$c_{d,k}(H) - c_{d,k}(G) \le d \int_{\mathbb{S}^1} \max(H_t - G_t)\, dt,$$

and so if $H \le G$ then the integrand in the above inequality is nonpositive, so that we obtain monotonicity.

To prove symplectic invariance, let ψ_t be a Hamiltonian isotopy such that $\psi_0 = \mathrm{Id}$ and $\psi_1 = \psi$. It is sufficient to show that the function $t \mapsto c_{d,k}(H \circ \psi_t)$ is constant. To see this, recall from Section 2.2 that $\mathrm{Spec}_d(H \circ \psi_t) = \mathrm{Spec}_d(H)$ and so the function $t \mapsto c_{d,k}(H \circ \psi_t)$, which is continuous by the continuity property of Theorem 3.1, takes values in the measure-zero set $\mathrm{Spec}_d(H)$ and so it must be constant.

The proof of homotopy invariance is analogous. Let $H_s, 0 \leq s \leq 1$, be a smooth path of mean-normalized Hamiltonians connecting H_0 to H_1. Note that, by the homotopy invariance of the action spectrum from Section 2.2, we have $\mathrm{Spec}_d(H_s) = \mathrm{Spec}_d(H_0)$ for all $d \in \mathbb{N}$ and $s \in [0,1]$. Then the continuous function $s \mapsto c_{d,k}(H_s)$ is constant because it takes values in the measure-zero set $\mathrm{Spec}_d(H_0)$. We conclude that $c_{d,k}(H_0) = c_{d,k}(H_1)$.

It therefore remains to prove support control. The proof will rely on the following lemma. We will say that a set U is *d-displaced* by a map Ψ if the sets $U, \Psi(U), \ldots, \Psi^d(U)$ are pairwise disjoint.

Lemma 3.3. *Let F be a Hamiltonian and let B be an open topological disc which is d-displaced by φ_F^1. Then, for any Hamiltonian G which is supported in B, we have* $c_{d,k}(G \,\#\, F) = c_{d,k}(F)$.

A similar lemma was established in [9, Lem. 4.4] but only for maps supported in the northern hemisphere. The argument presented here is essentially the same and so we will be rather brief.

Proof of Lemma 3.3. Let $(K^s)_{s \in [0,1]}$ be a smooth one-parameter family of Hamiltonians such that for any $s \in [0,1]$, the time-1 map of K^s is $\varphi_G^s \varphi_F^1$ and the isotopy of K^s consists in following first the isotopy generated by F and then that generated by $sG(st, x)$. More precisely, we may take $K^s = G^s \,\#\, F$, where $G^s(t, x) := sG(st, x)$. It generates the isotopy

$$\varphi_{K^s}^t = \begin{cases} \varphi_F^{\rho(2t)} & \text{if } t \in [0, 1/2], \\ \varphi_G^{s\rho(2t-1)} \varphi_F^1 & \text{if } t \in [1/2, 1]. \end{cases}$$

Then, for all $s \in [0,1]$, $\mathrm{Spec}_d(K^s) = \mathrm{Spec}_d(F)$: the argument for this is exactly the same as the argument[11] in [9, Lem. 4.4] and so we will omit it. This implies that for any (d, k) the continuous map $s \mapsto c_{d,k}(K^s)$ takes values in $\mathrm{Spec}_d(F)$; the fact that this map is continuous is a consequence of Hofer continuity of $c_{d,k}$ (see (15)). Since $\mathrm{Spec}_d(F)$ is totally disconnected, we deduce that these functions are all constant, and so $c_{d,k}(K^0) = c_{d,k}(K^1)$.

Since $K^1 = G \,\#\, F$, it is sufficient to show that $c_{d,k}(K^0) = c_{d,k}(F)$ to finish the proof. To see this, note that the Hamiltonian flows $\varphi_{K^0}^t$ and φ_F^t are homotopic rel. endpoints, and moreover $\int_{\mathbb{S}^1} \int_{\mathbb{S}^2} K^0 \omega \, dt = \int_{\mathbb{S}^1} \int_{\mathbb{S}^2} F\omega \, dt$. It then follows from the homotopy invariance and shift properties that $c_{d,k}(K^0) = c_{d,k}(F)$. ∎

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[11]To orient a reader who reads [9, Lem. 4.4], note that the argument there refers to the spectrum of the time-1 maps of K^s and F rather than to the Hamiltonians themselves; this is because in that proof, the Hamiltonians are all assumed to be zero on the southern hemisphere so we can refer to the spectrum in terms of the time-1 map; however, the argument for that lemma extends to the case here with no changes.

We will now use Lemma 3.3 to establish the support-control inequality.

Proof of the support-control inequality. Fix $d > 0$. Let H be a Hamiltonian supported in a disc D of area smaller than $\frac{1}{d+1}$. This area condition implies that we can find a Hamiltonian F such that the disc D is d-displaced by φ_F^1.

Furthermore, for any $\varepsilon > 0$, we may assume that $\|F\|_{1,\infty} \leq \mathrm{Area}(D) + \varepsilon$. To see this, note that we can find an area-preserving diffeomorphism ψ such that $\psi(D)$ is sandwiched between two meridians (that is, curves with $\theta = $ constant) of $\mathbb{S}^2$ which enclose a region of area $\mathrm{Area}(D) + \frac{\varepsilon}{2}$. Suppose that ε is so small that $\mathrm{Area}(D) + \varepsilon < \frac{1}{d+1}$. Then, consider the Hamiltonian $K = \frac{\mathrm{Area}(D)+\varepsilon}{2} z$ whose time-1 map φ_K^1 is the horizontal rotation of angle $\mathrm{Area}(D) + \varepsilon$, which d-displaces $\psi(D)$. Then we may set $F = K \circ \psi$, whose time-1 map, $\psi^{-1} \varphi_K^1 \psi$, d-displaces the disc D. Clearly, $\|F\|_{1,\infty} = \mathrm{Area}(D) + \varepsilon$.

By Lemma 3.3, we have $c_{d,k}(H \# F) = c_{d,k}(F)$. Using this and (15), we obtain

$$|c_{d,k}(H) - c_{d,k}(F)| = |c_{d,k}(H \# 0) - c_{d,k}(H \# F)| \leq d \|H \# 0 - H \# F\|_{1,\infty}$$
$$= d \|F\|_{1,\infty}.$$

Hence,

$$|c_{d,k}(H)| \leq |c_{d,k}(F)| + d\|F\|_{1,\infty} \leq 2d\|F\|_{1,\infty} = 2d\,\mathrm{Area}(D) + 2\varepsilon.$$

This completes the proof of the support-control inequality.   ∎

We have now completed the proof of Proposition 3.2.   ∎

3.2. Invariants for Hamiltonian diffeomorphisms

The goal of this section is to introduce PFH spectral invariants, and other related invariants, for Hamiltonian diffeomorphisms.

3.2.1. The c_d. We begin by noting that we can now define the PFH spectral invariants on the universal cover $\widetilde{\mathrm{Ham}}(\mathbb{S}^2, \omega)$ as follows. Given $\tilde{\varphi} \in \widetilde{\mathrm{Ham}}(\mathbb{S}^2, \omega)$, let H be a mean-normalized Hamiltonian such that the Hamiltonian path $\{\varphi_H^t\}, 0 \leq t \leq 1$, is a representative for $\tilde{\varphi}$. Define

$$c_{d,k}(\tilde{\varphi}) := c_{d,k}(H). \tag{21}$$

The mapping

$$c_{d,k} : \widetilde{\mathrm{Ham}}(\mathbb{S}^2, \omega) \to \mathbb{R}$$

is well-defined as a consequence of the homotopy invariance property in Proposition 3.2. Note that for any (not necessarily normalized) Hamiltonian H, the shift property yields

$$c_{d,k}(\tilde{\varphi}) = c_{d,k}(H) - d \int_{\mathbb{S}^1} \int_{\mathbb{S}^2} H_t \omega \, dt, \tag{22}$$

where $\tilde{\varphi}$ is the lift of φ given by the isotopy $(\varphi_H^t)_{t \in [0,1]}$.

However, to prove our main theorems, we will want invariants of $\mathrm{Ham}(\mathbb{S}^2, \omega)$ rather than $\widetilde{\mathrm{Ham}}(\mathbb{S}^2, \omega)$.

To produce such invariants, we start by showing that, as mentioned above, we can use the $c_{d,k}$ to define invariants that are independent of the choice of mean-normalized Hamiltonian. To begin, let $\tilde{\varphi} \in \widetilde{\mathrm{Ham}}(\mathbb{S}^2, \omega)$. Define

$$c_d(\tilde{\varphi}) := c_{d,-d}(\tilde{\varphi}).$$

Next, we will prove that, for even d, the map $c_d : \widetilde{\mathrm{Ham}}(\mathbb{S}^2, \omega) \to \mathbb{R}$ descends to $\mathrm{Ham}(\mathbb{S}^2, \omega)$. In other words, we will show that there is a well-defined map

$$c_d : \mathrm{Ham}(\mathbb{S}^2, \omega) \to \mathbb{R}.$$

This is the content of Proposition 3.5.

Remark 3.4. The c_d as defined here are not canonical. We could equally well define

$$c_d(\varphi) := c_{d,k}(\varphi_H^1)$$

for any $-d \leq k \leq d$ with the same parity as d. What is important for the applications in our paper is to choose a k such that the $c_{d,k}$ have the normalization property. It is also instructive to note that because addition of the homology class of a fiber of the map $Y_\varphi \to S^1$ induces a canonical bijection of the twisted PFH chain complex, we have

$$c_{d,k+2d+2}(H) = c_{d,k}(H) + 1. \tag{23}$$

In particular, as a function on $C^\infty(\mathbb{S}^1 \times \mathbb{S}^2)$, any $c_{d,k'}$ differs by a constant function from some $c_{d,k}$ with $-d \leq k \leq d$.

To summarize, then, there are essentially $d + 1$ possible spectral invariants corresponding to degree d, and we have made a noncanonical choice of one of them moving forward, with the main goal of simplifying the notation.

For future use, we also define

$$c_d(H) := c_{d,-d}(H)$$

for any $H \in C^\infty(\mathbb{S}^1 \times \mathbb{S}^2)$.

Proposition 3.5. *For any positive even integer d and any even integer k, the invariant $c_{d,k} : \widetilde{\mathrm{Ham}}(\mathbb{S}^2, \omega) \to \mathbb{R}$ descends to $\mathrm{Ham}(\mathbb{S}^2, \omega)$. In other words, it does not depend on the choice of mean-normalized H.*

In particular, we obtain a well-defined invariant $c_d : \mathrm{Ham}(\mathbb{S}^2, \omega) \to \mathbb{R}$ for any positive even integer d.

Proof. Let H be any Hamiltonian and $K = \frac{1}{2}(z + 1)$. Note that the Hamiltonian K vanishes at p_- and its time-1 flow is rotation by 2π about the z-axis. We will show below that for any positive integer d and integer k of the same parity as d,

$$c_{d,k}(H \# K) = c_{d,k}(H) + d/2. \tag{24}$$

This implies the conclusin of Proposition 3.5, by the following argument. Let H_1 and H_2 be mean-normalized Hamiltonians generating the same time-1 map. We can assume that H_1 and H_2 are not homotopic, or else the conclusion holds by Proposition 3.2 (5). Then $H_1 \# (K - 1/2)$ and H_2 are homotopic, and $H_1 \# (K - 1/2)$ is mean-normalized, hence

$$c_{d,k}(H_1 \# (K - 1/2)) = c_{d,k}(H_2).$$

On the other hand, by the shift property,

$$c_{d,k}(H_1 \# (K - 1/2)) = c_{d,k}(H_1 \# K) - d/2,$$

so that the conclusion follows from (24).

It remains to prove (24). To do this, we first note that

$$c_{d,k}(H \# 0) = c_{d,k}(H).$$

Indeed, $H \# 0$ and H are homotopic, with the same mean. Thus, it suffices to show that

$$c_{d,k}(H \# K) = c_{d,k}(H \# 0) + d/2.$$

To prove this, by the shift property we can assume that H vanishes at p_-. Then $H \# K$ and $H \# 0$ have the same time-1 map φ, and the same reference cycle $\gamma \subset Y_\varphi$. The only difference between them is the trivialization of V over γ; more precisely, if τ' denotes the trivialization over γ induced by $H \# K$ and τ denotes the trivialization induced by $H \# 0$, then $\tau' = \tau - 1$. Thus, since in this case the identity map is an isomorphism of the twisted PFH chain complexes, which shifts the grading by $d^2 + d$ by [21, (6), Lem. 2.5 (b)], we have

$$c_{d,k}(H \# K) = c_{d,k+d^2+d}(H \# 0) = c_{d,k}(H \# 0) + d/2$$

as desired; here, we have used (23) for the second equality above. ∎

3.2.2. Homogenized invariants. We introduced the homogenizations μ_d and ζ_d in the introduction (see (2) and (3)). The next proposition states that they are well-defined.

Proposition 3.6. *There are well-defined maps* $\mu_d : \mathrm{Ham}(\mathbb{S}^2, \omega) \to \mathbb{R}$ *and* $\zeta_d : C^\infty(\mathbb{S}^2) \to \mathbb{R}$ *given by*

$$\mu_d(\varphi) = \limsup_{n \to \infty} \frac{c_d(\tilde{\varphi}^n)}{n} \quad \text{and} \quad \zeta_d(H) = \limsup_{n \to \infty} \frac{c_d(nH)}{n}$$

for any $\varphi \in \mathrm{Ham}(\mathbb{S}^2, \omega)$ *and* $H \in C^\infty(\mathbb{S}^2)$.

Remark 3.7. One can more generally define $\zeta_d(H) := \limsup_{n \to \infty} c_d(H^n)/n$ for any (not necessarily autonomous) Hamiltonian $H \in C^\infty(\mathbb{S}^1 \times \mathbb{S}^2)$. However, we choose to restrict ζ_d to $C^\infty(\mathbb{S}^2)$ in analogy with [11], where a similar map ζ was defined and was proved to satisfy the properties of a symplectic quasi-state. It would be interesting to see if our ζ_d also has these properties.

Proof of Proposition 3.6. The fact that both the lim sup exist follows directly from the continuity property of $c_d := c_{d,-d}$ in Theorem 3.1. This shows that ζ_d is well-defined on $C^\infty(\mathbb{S}^2)$ and that μ_d is well-defined on the universal cover $\widetilde{\mathrm{Ham}}(\mathbb{S}^2, \omega)$. It remains to show that μ_d descends to $\mathrm{Ham}(\mathbb{S}^2, \omega)$.

To see this, let $\varphi \in \mathrm{Ham}(\mathbb{S}^2, \omega)$. Let $\tilde{\varphi}, \tilde{\varphi}'$ be two lifts of φ to $\widetilde{\mathrm{Ham}}(\mathbb{S}^2, \omega)$. Recall that $\pi_1(\mathrm{Ham}(\mathbb{S}^2, \omega))$ has only two elements and the nontrivial element is represented by the isotopy $\{\varphi_R^t\}$, where $R(\theta, z) = z/2$ is the Hamiltonian which generates a full 2π rotation around the z-axis. As a consequence, for any $n \in \mathbb{N}$, since $\tilde{\varphi}^n$ and $\tilde{\varphi}'^n$ are both lifts of the same diffeomorphism φ^n, we have either $\tilde{\varphi}^n = \tilde{\varphi}'^n$ or $\tilde{\varphi}^n = \tilde{\varphi}'^n \tilde{\varphi}_R$. In both cases, the continuity property of Theorem 3.1 gives an upper bound which does not depend on n:

$$|c_d(\tilde{\varphi}^n) - c_d(\tilde{\varphi}'^n)| \le d \|R\|_{1,\infty}.$$

It then follows that

$$\limsup_{n\to\infty} \frac{c_d(\tilde{\varphi}^n)}{n} = \limsup_{n\to\infty} \frac{c_d(\tilde{\varphi}'^n)}{n},$$

which proves that μ_d descends to $\mathrm{Ham}(\mathbb{S}^2, \omega)$. ∎

We next state some of the properties of μ_d which will be used in our arguments.

Proposition 3.8. *The invariant* $\mu_d : \mathrm{Ham}(\mathbb{S}^2, \omega) \to \mathbb{R}$ *has the following properties:*

(1) Normalization: $\mu_d(\mathrm{Id}) = 0$.

(2) Hofer continuity: *For all* φ, ψ *we have*

$$|\mu_d(\varphi) - \mu_d(\psi)| \le d\, d_H(\varphi, \psi).$$

(3) Calabi property: *Suppose that* $\mathrm{supp}(\varphi)$ *is contained in a topological disc D. If* $\mathrm{Area}(D) < \frac{1}{d+1}$, *then*

$$\frac{1}{d}\mu_d(\varphi) = -\mathrm{Cal}(\varphi),$$

where $\mathrm{Cal} : \mathrm{Ham}_c(D, \omega) \to \mathbb{R}$ *denotes the Calabi invariant.*

(4) Relationship to ζ_d: *For any* $H \in C^\infty(\mathbb{S}^2)$,

$$\mu_d(\varphi_H^t) = \zeta_d(tH) - td \int_{\mathbb{S}^2} H\omega \tag{25}$$

for all $t \in \mathbb{R}$.

Proof. The first item follows immediately from the definition of μ_d combined with the fact that $c_d(0) = 0$.

To prove the second item, let $\varphi, \psi \in \mathrm{Ham}(\mathbb{S}^2, \omega)$ and H, K be mean-normalized Hamiltonians with $\varphi_H^1 = \varphi$ and $\varphi_K^1 = \psi$. We also denote by $\tilde{\varphi}, \tilde{\psi} \in \widetilde{\mathrm{Ham}}(\mathbb{S}^2, \omega)$ the lifts of φ, ψ respectively given by H, K. Then, by definition,

$$\tfrac{1}{d}|c_d(\tilde{\varphi}^n) - c_d(\tilde{\psi}^n)| = \tfrac{1}{d}|c_d(H^{\#n}) - c_d(K^{\#n})|$$

for any $n > 0$. Here $H^{\#n}$ denotes the n-fold composition $H \# \cdots \# H$. By the continuity property of c_d, we have

$$\tfrac{1}{d}|c_d(H^n) - c_d(K^n)| \le \|H^{\#n} - K^{\#n}\|_{1,\infty} = n\|H - K\|_{1,\infty}.$$

Note that this last equality follows from (9). From this inequality, we deduce

$$|\mu_d(\varphi) - \mu_d(\psi)| \le d\|H - K\|_{1,\infty}.$$

Since this holds for any choices of Hamiltonians H, K, and since we can restrict to mean-normalized Hamiltonians in computing the Hofer norm of $\varphi^{-1}\psi$, the Hofer continuity property follows.

The Calabi property is a consequence of the support-control property from Proposition 3.2. Indeed, given any Hamiltonian H with support in D, (22) yields

$$c_d(\tilde{\varphi}) = c_d(H) - d \int_{\mathbb{S}^1} \int_{\mathbb{S}^2} H_t \omega \, dt,$$

where $\tilde{\varphi}$ is the lift of φ given by the isotopy $(\varphi_H^t)_{t \in [0,1]}$. The integral in the second term on the right hand side above is precisely $-\mathrm{Cal}(\varphi)$, while the first term is bounded from above by $2d\,\mathrm{Area}(D)$, by support control. Applying this to φ^n, for any $n > 0$, we get

$$\left|\tfrac{1}{n}c_d(\tilde{\varphi}^n) + d\,\mathrm{Cal}(\varphi^n)\right| = \tfrac{1}{n}|c_d(\tilde{\varphi}^n) + d\,\mathrm{Cal}(\varphi^n)| = \tfrac{1}{n}|c_d(H^n)| \le \tfrac{1}{n}2d\,\mathrm{Area}(D).$$

The Calabi property follows from this inequality.

As for the last item, it follows from the definitions of μ_d and ζ_d, and (22), that

$$\mu_d(\varphi_H^t) = \mu_d(\varphi_{tH}^1) = \zeta_d(tH) - d \int_{\mathbb{S}^1} \int_{\mathbb{S}^2} tH\omega \, dt. \qquad \blacksquare$$

3.2.3. The invariants η_d. Although we can use the invariants c_d or μ_d to get invariants of the time-1 map, these invariants will not in general be C^0-continuous, as they require mean-normalizing the Hamiltonian. We obtain C^0-continuous invariants by defining, for even $d \in \mathbb{N}$, the numbers

$$\eta_d : \ \mathrm{Ham}(\mathbb{S}^2, \omega) \to \mathbb{R}, \quad \varphi \mapsto c_d(\varphi) - \frac{d}{2}c_2(\varphi). \tag{26}$$

The fact that η_d is well-defined is an immediate consequence of Proposition 3.5. Observe that, by Proposition 3.5 and the shift property of Proposition 3.2, we have

$$\eta_d(\varphi) = c_d(H) - \frac{d}{2}c_2(H), \tag{27}$$

where H is any Hamiltonian with time-1 flow φ.

We now prove that η_d has various properties, most notably C^0-continuity.

Proposition 3.9. *The invariant η_d is well-defined and has the following properties for all $\varphi, \psi \in \mathrm{Ham}(\mathbb{S}^2, \omega)$:*

(1) Normalization: $\eta_d(\mathrm{Id}) = 0$.

(2) Hofer continuity: $|\eta_d(\varphi) - \eta_d(\psi)| \leq d\, d_H(\varphi, \psi)$.

(3) Support control: *If the support of $\varphi \in \mathrm{Ham}(\mathbb{S}^2, \omega)$ is contained in a topological disc D with $\mathrm{Area}(D) < \frac{1}{d+1}$, then $|\eta_d(\varphi)| \leq 2d\,\mathrm{Area}(D)$.*

(4) C^0-continuity: *The mapping $\eta_d : \mathrm{Ham}(\mathbb{S}^2, \omega) \to \mathbb{R}$ is continuous with respect to the C^0-topology on $\mathrm{Ham}(\mathbb{S}^2, \omega)$, and furthermore it extends continuously to $\mathrm{Homeo}_0(\mathbb{S}^2, \omega)$.*

We recall in relation to the support-control inequality that the total area of the sphere is assumed to be 1.

Proof of Proposition 3.9. The first and third properties are immediate consequences of the same properties of the invariants $c_{d,k}$; see Theorem 3.1 and Proposition 3.2. To prove the second, note first of all that if H and G are any Hamiltonians, then by the continuity property of Theorem 3.1, we have

$$(c_d(H) - c_d(G)) - \frac{d}{2}(c_2(H) - c_2(G)) \leq d\,\|H - G\|.$$

Now let K be any Hamiltonian generating $\varphi^{-1}\psi$ and let G generate φ. Then $H := G \# K$ generates ψ and hence by the above inequality,

$$\eta_d(\psi) - \eta_d(\varphi) \leq d\,\|K\|, \quad \eta_d(\varphi) - \eta_d(\psi) \leq d\,\|K\|,$$

which yields the Hofer continuity property, since K was arbitrary.

We only have to establish C^0-continuity. This takes up the remainder of this subsection. Our proof will follow the lines (and use some of the intermediate steps) of [9, Sec. 4], which established a similar result for the invariant c_d restricted to maps supported in the northern hemisphere. We fix some degree $d > 0$. The result will follow from the next proposition.

Proposition 3.10. *For any $h \in \mathrm{Homeo}_0(\mathbb{S}^2, \omega)$ and $\varepsilon > 0$, there exists $\delta > 0$ such that for all $f, g \in \mathrm{Ham}(\mathbb{S}^2, \omega)$ satisfying $d_{C^0}(f, h) < \delta$ and $d_{C^0}(g, \mathrm{Id}) < \delta$, we have $|\eta_d(gf) - \eta_d(f)| < \varepsilon$.*

Let us temporarily assume this proposition and explain how it implies the C^0-continuity of η_d. Let $\{f_i\}_{i \in \mathbb{N}} \subset \mathrm{Ham}(\mathbb{S}^2, \omega)$ be a sequence which C^0-converges to $h \in \mathrm{Ham}(\mathbb{S}^2, \omega)$. We may write f_i in the form $g_i h$ with $d_{C^0}(g_i, \mathrm{Id}) \to 0$ as $i \to \infty$. By the proposition we have $|\eta_d(g_i h) - \eta_d(h)| \to 0$, which proves the C^0-continuity. To prove extension to $\mathrm{Homeo}_0(\mathbb{S}^2, \omega)$ let $h \in \mathrm{Homeo}_0(\mathbb{S}^2, \omega)$ and let $\{f_i\}_{i \in \mathbb{N}} \subset \mathrm{Ham}(\mathbb{S}^2, \omega)$ be a sequence which C^0-converges to h. Then $d_{C^0}(f_i, f_j) = d_{C^0}(f_i f_j^{-1}, \mathrm{Id})$ becomes arbitrarily small when i, j are large enough and so Proposition 3.10 implies that $|\eta_d(f_i) - \eta_d(f_j)| = |\eta_d((f_i f_j^{-1})f_j) - \eta_d(f_j)|$ becomes arbitrarily small for i, j large enough so that $\eta_d(f_i)$ converges. Proposition 3.10 also similarly implies that if $\{f_i'\}_{i \in \mathbb{N}}$ is another sequence converging to h, then $|\eta_d(f_i) - \eta_d(f_i')| \to 0$, hence the limit does

not depend on the choice of limiting sequence. This allows us to consistently define $\eta_d(h)$ for any $h \in \mathrm{Homeo}_0(\mathbb{S}^2, \omega)$ by setting

$$\eta_d(h) := \lim_{i \to \infty} \eta_d(f_i)$$

for any sequence f_i which C^0-converges to h. ∎

We now prove Proposition 3.10.

Proof of Proposition 3.10. We give the proof of this proposition in two steps.

Step A: Continuity in the non-finite-order case. We first assume that h is not of finite order in the group $\mathrm{Homeo}_0(\mathbb{S}^2, \omega)$. Then there exists a point $x \in \mathbb{S}^2$ such that $h^{d!}(x) \neq x$. For such a point and for any integers $0 \leq p < q \leq d$, we have $h^{q-p}(x) \neq x$. By composing with h^p, we also have $h^q(x) \neq h^p(x)$. Therefore, the points $x, h(x), \ldots, h^d(x)$ are pairwise distinct. Let B be a small ball centered at x such that the closure $\bar{B}$ of B is d-displaced by h.

Let $\varepsilon > 0$. We choose $\delta' > 0$ so small that any map f such that $d_{C^0}(f, h) < \delta'$ must also d-displace $\bar{B}$.

The next lemma says roughly that a C^0-small element of $\mathrm{Ham}(\mathbb{S}^2, \omega)$ is Hofer-close to being supported in B.

Lemma 3.11. *Let B be any open topological disc. For all $\varepsilon' > 0$, there exists $\delta > 0$ such that for all $g \in \mathrm{Ham}(\mathbb{S}^2, \omega)$ with $d_{C^0}(g, \mathrm{Id}) < \delta$, there is $\phi \in \mathrm{Ham}(\mathbb{S}^2, \omega)$ supported in B such that $d_H(\phi, g) \leq \varepsilon'$.*

A similar result was proved in [9, Lem. 4.6] but only for maps g supported in the northern hemisphere. By conjugating by an appropriate area-preserving map, this particular case implies that Lemma 3.11 holds for maps g supported in any topological disc of area $1/2$. In fact, the factor $1/2$ here is not essential to the proof of [9, Lem. 4.6]: the same argument, which we omit[12] for brevity, shows that it also holds for maps supported in an embedded disc of any area. The general case then immediately follows from the next fragmentation lemma: indeed, given the lemma below, and given g, we can first fragment g into maps supported on embedded discs and then approximate each of these maps by maps supported in B.

Lemma 3.12 ([48, Prop. 3.1]). *There exist open topological embedded discs D_1, D_2 which cover $\mathbb{S}^2$ such that for any $\alpha > 0$, there exists $\delta > 0$ such that for any $g \in \mathrm{Ham}(S^2, \omega)$ satisfying $d_{C^0}(g, \mathrm{Id}) < \delta$, there exist $g_1, g_2 \in \mathrm{Ham}(\mathbb{S}^2, \omega)$ with $\mathrm{supp}(g_i) \subset D_i$ and $d_{C^0}(g_i, \mathrm{Id}) < \alpha$ for $i = 1, 2$ such that $g = g_1 \circ g_2$.*

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[12]To help the reader who reads the argument in [9, Lem. 4.6], we note that the only change is that the factors of $1/2$, which come from the fact that the northern hemisphere has area $1/2$ (see the end of the second paragraph of the proof there) must be changed to some number $\ell < 1$. This change can be accommodated by choosing what are called N and m in the proof to be such that $\ell/N < \mathrm{Area}(B), \ell/m < 1/2$ and $4\ell\frac{N+1}{m} < \epsilon$. With these changes understood, the argument can be repeated verbatim.

Having established Lemma 3.11, we can now continue the proof of Proposition 3.10. Let $\delta > 0$ be as provided by Lemma 3.11 for $\varepsilon' = \varepsilon/d$. We may assume without loss of generality that $\delta < \delta'$. Then, by Lemma 3.11, for any Hamiltonian diffeomorphism g with $d_{C^0}(g, \mathrm{Id}) < \delta$, there exists $\phi \in \mathrm{Ham}(\mathbb{S}^2, \omega)$ supported in B satisfying $d_H(\phi, g) \leq \varepsilon/d$. Now let $f, g \in \mathrm{Ham}(\mathbb{S}^2, \omega)$ be such that $d_{C^0}(f, h) < \delta$ and $d_{C^0}(g, \mathrm{Id}) < \delta$. By Hofer continuity of η_d and Lemma 3.3, it follows that

$$\begin{aligned} |\eta_d(gf) - \eta_d(f)| &\leq |\eta_d(gf) - \eta_d(\phi f)| + |\eta_d(\phi f) - \eta_d(f)| \\ &\leq d\, d_H(gf, \phi f) + 0 = d\, d_H(g, \phi) \leq \varepsilon. \end{aligned}$$

This concludes the proof in the case where h is not of finite order.

Step B: The finite-order case. We will now conclude the argument by reducing the case where h is of finite order to the case where it is not. Let h be of finite order and let $\varepsilon > 0$. We may pick a Hamiltonian diffeomorphism ψ such that $\|\psi\| < \varepsilon/3$ and $h\psi$ is not of finite order.[13] Then, by Step A, there exists δ such that for any f', g' satisfying $d_{C^0}(f', h\psi) < \delta$ and $d_{C^0}(g', h) < \delta$, we have $|\eta_d(g'f') - \eta_d(f')| < \varepsilon/3$.

Now take f, g as in the statement of the proposition. We apply the triangle inequality to obtain

$$\begin{aligned} |\eta_d(gf) - \eta_d(f)| \\ \leq |\eta_d(gf) - \eta_d(gf\psi)| + |\eta_d(gf\psi) - \eta_d(f\psi)| + |\eta_d(f\psi) - \eta_d(f)|. \end{aligned}$$

By Hofer continuity and the above estimate on $\|\psi\|$, we have

$$|\eta_d(gf) - \eta_d(gf\psi)| \leq \varepsilon/3, \quad |\eta_d(f\psi) - \eta_d(f)| \leq \varepsilon/3.$$

Thus, to finish the proof of the proposition, it remains to show that

$$|\eta_d(gf\psi) - \eta_d(f\psi)| \leq \varepsilon/3.$$

This follows from the previous paragraph, since if $d_{C^0}(f, h) < \delta$, then $d_{C^0}(f\psi, h\psi) < \delta$. ∎

4. The quasi-isometry type of the kernel of Calabi

Equipped with our new spectral invariants, we now prove Theorem 1.4.

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[13] For the benefit of the reader, we briefly sketch why such a ψ exists. Since h has finite order, all $x \in \mathbb{S}^2$ are periodic and we let ℓ be the maximal period of a point. Then the set of points of period ℓ is open, because it is $\{x \in \mathbb{S}^2 : h^k(x) \neq x, \forall k = 1, \ldots, \ell - 1\}$. It follows that if we fix a point x of period ℓ, there exists an open set U containing x such that $h^\ell|_U = \mathrm{Id}_U$ and $U, \ldots, h^{\ell-1}(U)$ are pairwise disjoint. Now, let ψ be a C^1-small (hence Hofer-small) map supported in U which coincides with an irrational rotation around x in a smaller open subset $V \subset U$. Then $h\psi$ does not have finite order. Indeed, for any $n \in \mathbb{N}$ and $y \in V \setminus \{x\}$, we have $(h\psi)^n(y) = \psi^{n/\ell}(y)$ if ℓ divides n and $(h\psi)^n(y) \notin U$ otherwise. In both cases, $(h\psi)^n(y) \neq y$. Thus, such a ψ suits our needs.
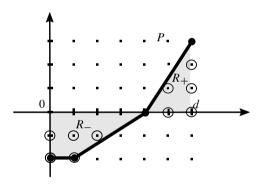
4.1. Homogenization and monotone twists

4.1.1. A combinatorial model. We begin by recalling the combinatorial model of [9, Thm. 6.1], which gives an explicit formula for $c_d(H)$ where $H : \mathbb{S}^2 \to \mathbb{R}$ is a monotone twist, and which we will need below and for the proof of Theorem 1.7 as well.

Here and below we use the notations of [9, Sec. 5.2]. To summarize, recall that a *lattice path* P is the graph of a piecewise linear function $Y : [0, d] \to \mathbb{R}_{\geq 0}$ such that the vertices of P are at integer lattice points; the number d is called the *degree* of the path and is assumed to be an integer below. A lattice path is called *concave* if it never crosses any of its tangent lines. We can think of a lattice path as a collection of maximal line segments, called *edges*, joined end to end. We regard any edge as an integer multiple $m_{p,q}$ of a primitive vector (q, p); these vectors are directed with the convention that $q \geq 1$.

To any concave lattice path, we associate a number $j(P)$ as follows. We form the region R_+ bounded by the x-axis, the line $x = d$, and the part of P above the x-axis, we form the region R_- bounded by the axes and the part of P below the x-axis, we define j_+ to be the number of lattice points in R_+, not including lattice points on P, and we define j_- to be the number of lattice points in R_-, not including the lattice points on the x-axis. Finally, we define $j(P) := j_+ - j_-$. See Figure 1. We define the *combinatorial index* $I(P)$ by

$$I(P) := 2j(P) - d.$$

Fig. 1. The lattice points which contribute to the count for $j(P)$ are circled. Here, $j_+(P) = 5$, $j_-(P) = 5$, $d = 6$, thus $j(P) = 0$ and $I(P) = -6$.

Now let $H = \frac{1}{2}h$ be a monotone twist. Assume that $h(-1) = h'(-1) = 0, h'' > 0$, and $h'(1)$ is an integer. We call such a monotone twist *nice*. We call a lattice path P *compatible* with h if for every edge $m_{p,q}(q, p)$, there exists some $z_{p,q}$ such that $h'(z_{p,q}) = p/q$. If P is a concave lattice path, compatible with a nice monotone twist, then we define the *action* $\mathcal{A}(P)$ by setting

$$\mathcal{A}(q, p) := \tfrac{1}{2}\big(p(1 - z_{p,q}) + qh(z_{p,q})\big), \tag{28}$$

and extending by linearity.

We can now state the formula from [9, Thm. 6.1] for the invariants $c_{d,k}$: for all degrees $d \geq 1$ and all gradings k, when H is a nice monotone twist we have

$$c_{d,k}(H) = \max \{ \mathcal{A}(P) : 2j(P) - d = k \}, \tag{29}$$

where the maximum is over concave lattice paths that are compatible with h. Invoking this formula is justified because our nice monotone twists are in $\mathcal{H}$, and the $c_{d,k}$ defined here extend the definition from [9]; see the discussion at the beginning of Section 3.

4.1.2. ζ_d of monotone twists. We now apply the combinatorial model from the previous section to compute the invariants ζ_d for monotone twists. In particular, we can now give the promised proof of Proposition 1.11.

Proof of Proposition 1.11. First note that ζ_d is Lipschitz continuous with respect to the uniform norm on $C^\infty(\mathbb{S}^2)$; this follows from Hofer continuity of c_d. Thus, both sides of the equation in Proposition 1.11 are continuous with respect to uniform norm. Since any monotone twist H can be approximated uniformly by nice monotone twists, we deduce that it is sufficient to prove Proposition 1.11 for such twists. For the rest of the proof, we therefore assume that H is a nice monotone twist.

We note that the index $I(P) = 2j(P) - d$ from §4.1.1 is equivalently given by

$$I(P) = 2A + y + w - e, \tag{30}$$

where y and w are respectively the minimal and maximal vertical coordinates of P, e is the number of edges in P, and A is the (signed) area of the region enclosed by P, the x-axis and the vertical line $\{d\} \times \mathbb{R}$. Indeed, by shifting the path if necessary, it suffices to prove this when $y = 0$, in which case it follows from Pick's formula that this corresponds to the definition given in §4.1.1.

Let us introduce some notation. For any $i \in \{1, \ldots, d\}$, we set $a_i = Y(i) - Y(i-1)$, where Y is the function $[0, d] \to \mathbb{R}$ such that $P = \mathrm{graph}(Y) = \{(x, Y(x)) : x \in [0, d]\}$. Then

$$2A = 2dy + a_1 + (2a_1 + a_2) + \cdots + (2a_1 + \cdots + 2a_{d-1} + a_d)$$
$$= 2dy + (2d-1)a_1 + (2d-3)a_2 + \cdots + a_d.$$

Using (30) and the relation $w = y + a_1 + \cdots + a_d$, the condition $I(P) = -d$ becomes

$$2dy + 2y + 2da_1 + (2d-2)a_2 + \cdots + 2a_d - e = -d.$$

Therefore, under this condition, we may express y in terms of the a_i as

$$y = -\frac{1}{d+1}\left(da_1 + (d-1)a_2 + \cdots + a_d + \frac{d-e}{2} \right)$$
$$= \frac{e-d}{2(d+1)} + \sum_{i=1}^{d}\left(-a_i + \frac{i}{d+1}a_i \right). \tag{31}$$

Let us now turn our attention to the action. It is given by (28): we state here a reformulated version with the a_i, namely

$$\mathcal{A}(P) = y + \sum_{i=1}^{d} \frac{1}{2}(a_i(1 - z_i) + h(z_i)),$$

where z_i is the unique point such that $h'(z_i) = a_i$. Using (31), we obtain

$$\mathcal{A}(P) = \frac{e - d}{2(d + 1)} + \sum_{i=1}^{d}\left(\frac{i}{d + 1}a_i - a_i\right) + \sum_{i=1}^{d}\frac{1}{2}(a_i(1 - z_i) + h(z_i))$$

$$= \frac{e - d}{2(d + 1)} + \frac{1}{2}\sum_{i=1}^{d}\left(h'(z_i)\left(-1 + \frac{2i}{d + 1} - z_i\right) + h(z_i)\right). \tag{32}$$

Note that the term $\frac{e-d}{2(d+1)}$ belongs to $(-1/2, 0]$. It vanishes when all edges in P have horizontal displacement 1.

By (29), $c_d = c_{d,-d}$ is obtained by maximizing the value of $\mathcal{A}(P)$ over all possible paths with $I(P) = -d$.

To compute this maximum, consider the function

$$F(t_1, \ldots, t_d) = \sum_{i=1}^{d}\left(h'(t_i)\left(-1 + \frac{2i}{d + 1} - t_i\right) + h(t_i)\right),$$

defined on the set E of tuples $(t_1, \ldots, t_d)$ such that $-1 \le t_1 \le \cdots \le t_d \le 1$. We have

$$\mathcal{A}(P) = \frac{1}{2}F(z_1, \ldots, z_d) + \frac{e - d}{2(d + 1)}. \tag{33}$$

We may compute the partial derivatives of F,

$$\frac{\partial F}{\partial t_i} = \left(-1 + \frac{2i}{d + 1} - t_i\right)h''(t_i),$$

and we see that it is positive for $t_i < -1 + \frac{2i}{d+1}$ and negative for $t_i > -1 + \frac{2i}{d+1}$. This implies that F attains its maximum at $(t_1, \ldots, t_n)$ such that $t_i = -1 + \frac{2i}{d+1}$ for all i. Thus,

$$\max_{E} F = \sum_{i=1}^{d} h\left(-1 + \frac{2i}{d + 1}\right). \tag{34}$$

Now it follows from the versions of (29), (33), (34) for nH, and the fact that $e - d \le 0$, that for all n,

$$\frac{1}{n}c_d(nH) \le \frac{1}{2}\sum_{i=1}^{d} h\left(-1 + \frac{2i}{d + 1}\right). \tag{35}$$

To proceed, let $a_1^n \leq \cdots \leq a_d^n$ be sequences in $\frac{1}{n}\mathbb{N}$ which converge respectively to $-1 + \frac{2}{d+1}, -1 + \frac{4}{d+1}, \ldots, 1 - \frac{2}{d+1}$. Let z_i^n be such that $h'(z_i^n) = a_i^n$. Then

$$F(z_1^n, \ldots, z_d^n) \xrightarrow{n \to \infty} \max_E F = \frac{1}{2} \sum_{i=1}^{d} h\left(-1 + \frac{2i}{d+1}\right).$$

Moreover, we can construct a lattice path, for nH, such that (33) holds with the $z_i = z_i^n$ and $e = d$. We therefore deduce $\frac{1}{n} c_d(nH) \geq \frac{1}{2} F(z_1^n, \ldots, z_d^n)$ and we conclude in combination with (35) that

$$\lim_{n \to \infty} \frac{1}{n} c_{d,-d}(nH) = \frac{1}{2} \sum_{i=1}^{d} h\left(-1 + \frac{2i}{d+1}\right),$$

as desired. ∎

Remark 4.1. (i) It follows from the previous proposition that we have the following composition property for monotone twists: for any monotone twists ϕ, ψ we have

$$\mu_d(\phi\psi) = \mu_d(\phi) + \mu_d(\psi). \tag{36}$$

Indeed, it follows from Proposition 1.11 that $\zeta_d(H_1 + H_2) = \zeta_d(H_1) + \zeta_d(H_2)$, which gives (36).

(ii) For any monotone twist $\varphi \in \mathrm{Ham}(\mathbb{S}^2, \omega)$, it can easily be shown that $\mu_d(\varphi) = \lim_{n \to \infty} c_d(\tilde{\varphi}^n)/n$, i.e. the lim sup in formula (2) is in fact a limit for such φ.

4.2. A family of Hamiltonians

We will now construct a certain family of Hamiltonian diffeomorphisms which will be used to establish Theorem 1.4. Let $U \subset \mathbb{S}^2$ be a proper open set containing the north pole p_+ and let $\iota > 0$ be an integer. For all $i \in \mathbb{N}$, we denote

$$D_i := \{(z, \theta) \in \mathbb{S}^2 : 1 - 2/d_i < z \leq 1\}, \quad \text{where } d_i = 2^{\iota + i + 1}.$$

We choose ι large enough so that all the D_i's are contained in U. Note that each D_i is an embedded disc and

$$\mathrm{Area}(D_i) = 1/d_i.$$

The next lemma states the properties of our family of maps.

Lemma 4.2. *There exist autonomous Hamiltonians* $(H_i)_{i \in \mathbb{N}}$ *such that* H_i *is supported in* D_i *and the following properties are satisfied for all* $t \geq 0$, $i \in \mathbb{N}$:

(1) $\varphi_{H_i}^t$ *is a monotone twist for all* $t \in \mathbb{R}$,

(2) $\mathrm{Cal}(\varphi_{H_i}^t) = t/2$,

(3) $d_H(\varphi_{H_i}^t, \mathrm{Id}) \leq 2t + 2$,

(4) *if* $j > i$, *then* $\mu_{d_i}(\varphi_{H_j}^t) = -t d_i/2$,

(5) $\mu_{d_i}(\varphi^t_{H_i}) > -t d_i/2$ and if $j < i$, then $\mu_{d_i}(\varphi^t_{H_j}) \geq -t\frac{5d_i}{16}$,

(6) $\varphi^t_{H_i}\varphi^s_{H_j} = \varphi^s_{H_j}\varphi^t_{H_i}$ for all $t, s \in \mathbb{R}$.

Proof. Consider the functions $f_i : [-1, 1] \to \mathbb{R}$ defined by

$$f_i(z) := \begin{cases} 0, & z \in [-1, 1 - 2/d_i], \\ d_i^2(z - (1 - 2/d_i)), & z \in [1 - 2/d_i, 1]. \end{cases}$$

These functions are nonsmooth but to ensure that our future Hamiltonians are smooth, we approximate them by smooth functions h_i satisfying the following conditions:

  (i) $h'_i(z), h''_i(z) \geq 0$ and $|f_i(z) - h_i(z)| \leq 1/d_i$ for all $z \in [-1, 1]$,

 (ii) the support of h_i is contained in the interior of $[1 - 2/d_i, 1]$.

(iii) $\int_{-1}^1 h_i(z)\, dz = 2$.

Let $H_i(z, \theta) = \frac{1}{2}h_i(z)$ and observe that H_i is supported in D_i and that $\varphi^t_{H_i}$ is a monotone twist for any $t > 0$ (proving item (1) of the lemma).

The fact that $\mathrm{Cal}(\varphi^t_{H_i}) = \frac{t}{2}$ readily follows from property (iii).

To prove item (3), we will need the following lemma, whose proof we postpone to the end of this section. The idea behind this lemma goes back to Sikorav, who implemented it in the case of $\mathbb{R}^{2n}$ in [50]; see also [19, §5.6].

Lemma 4.3. *Let $H : \mathbb{S}^2 \to \mathbb{R}$ denote an autonomous Hamiltonian such that the support of H is contained in a disc D with $\mathrm{Area}(D) < 1/N$. Then $d_H(\varphi^1_H, \mathrm{Id}) \leq \frac{1}{N}\max(H) + 2$.*

Since the support of H_i is a disc of area less than $1/d_i$, this lemma leads to

$$d_H(\varphi^t_{H_i}, \mathrm{Id}) \leq \frac{1}{d_i}\max(t H_i) + 2 \leq \left(1 + \frac{1}{2d_i^2}\right)t + 2 \leq 2t + 2,$$

which implies item (3).

Item (4) is a consequence of the Calabi property from Proposition 3.8 (3), because $\varphi^t_{H_j}$ is supported in D_j and $\mathrm{Area}(D_j) < \frac{1}{d_i+1}$ for $j > i$. The last item of Lemma 4.2 is also easy to check. Indeed, since the Hamiltonians H_i are functions of z, they all Poisson commute, hence their flows commute.

It remains to prove item (5). Since

$$\zeta_{d_i}(t H_j) = \mu_{d_i}(\varphi^t_{H_j}) + d_i \int_{\mathbb{S}^2} t H_j \omega = \mu_{d_i}(\varphi^t_{H_j}) + t\frac{d_i}{2},$$

we just need to prove that $\zeta_{d_i}(H_i) > 0$ and $\zeta_{d_i}(H_j) \geq 3d_i/16$ when $i > j$.

As already mentioned, the above conditions (i) and (ii) ensure that $\varphi^t_{H_j}$ is a monotone twist for all $t > 0$. By Proposition 1.11,

$$\zeta_{d_i}(H_j) = \frac{1}{2}\sum_{m=1}^{d_i} h_j\left(-1 + \frac{2m}{d_i + 1}\right).$$

We can rewrite the above sum as

$$\zeta_{d_i}(H_j) = \frac{1}{2} \sum_{k=0}^{N} h_j \left(-1 + \frac{2(d_i - k)}{d_i + 1} \right),$$

where N is the largest integer such that $-1 + \frac{2(d_i - N)}{d_i + 1}$ is in the support of h_j, that is, $-1 + \frac{2(d_i - N)}{d_i + 1} > 1 - \frac{2}{d_j}$. A simple computation reveals that

$$N + 1 = d_i / d_j. \tag{37}$$

Consider the (nonsmooth) function $f_j(z)$. By the definition of h_j, we have $h_j \geq f_j - 1/d_j$ and so

$$\zeta_{d_i}(H_j) \geq \left[\frac{1}{2} \sum_{k=0}^{N} f_j \left(-1 + \frac{2(d_i - k)}{d_i + 1} \right) \right] - \frac{1}{2} \frac{N+1}{d_j}. \tag{38}$$

The first term on the right hand side in the above equation may be computed explicitly. Indeed, $f_j(z)$ is linear in z and so the above is just an arithmetic sum. First, note that $-1 + \frac{2(d_i - k)}{d_i + 1} = 1 - \frac{2}{d_i + 1}(k + 1)$, and so

$$\sum_{k=0}^{N} f_j \left(-1 + \frac{2(d_i - k)}{d_i + 1} \right) = \sum_{k=0}^{N} f_j \left(1 - \frac{2}{d_i + 1}(k + 1) \right).$$

Next, one can easily check that $f_j(1 - \frac{2}{d_i+1}(k+1)) = 2d_j - \frac{2d_j^2}{d_i+1}(k+1)$. Thus,

$$\sum_{k=0}^{N} f_j \left(1 - \frac{2}{d_i+1}(k+1) \right) = \sum_{k=0}^{N} \left[2d_j - \frac{2d_j^2}{d_i+1}(k+1) \right]$$

$$= 2d_j(N+1) - \frac{2d_j^2}{d_i+1} \frac{(N+1)(N+2)}{2}$$

$$= d_i \left(2 - \frac{d_i + d_j}{d_i + 1} \right). \tag{39}$$

For $i = j$, (37)–(39) yield

$$\zeta_{d_i}(H_i) \geq d_i \left(1 - \frac{d_i}{d_i + 1} - \frac{1}{2d_i^2} \right) > 0,$$

as desired, which implies the first part of item (5).

For $i > j$, (37)–(39) give

$$\zeta_{d_i}(H_i) \geq \frac{1}{2} d_i \left(2 - \frac{d_i + d_j}{d_i + 1} - \frac{1}{d_j^2} \right)$$

$$\geq \frac{1}{2} d_i \left(2 - \frac{d_i + d_j}{d_i} - \frac{1}{8} \right) = \frac{1}{2} d_i \left(\frac{7}{8} - \frac{d_j}{d_i} \right) \geq \frac{3d_i}{16},$$

where we have used $1/d_j^2 \leq 1/8$ for the second inequality and $d_j/d_i \leq 1/2$ for the last inequality. This concludes the proof of item (5). ∎

We end this section (and the proof of Lemma 4.2) with the proof of Lemma 4.3.

Proof of Lemma 4.3. Since $\text{Area}(D) < 1/N$, by Lemma 2.1, we can find $\psi_1, \ldots, \psi_N \in \text{Ham}(\mathbb{S}^2, \omega)$ such that

$$\psi_i(D) \cap \psi_j(D) = \emptyset \quad \text{for } i \neq j, \quad d_H(\psi_i, \text{Id}) \leq 1/N.$$

Define the Hamiltonians $F_i = \frac{1}{N} H \circ \psi_i^{-1}$ and note that $\varphi_{F_i}^1 = \psi_i \varphi_H^{1/N} \psi_i^{-1}$. Let $F = \sum_{i=1}^{N} F_i$. Observe that since F_i is supported in $\psi_i(D)$, the F_i's have disjoint supports, and so $\max(F) = \max(F_i) = \frac{1}{N} \max(H)$. Hence,

$$d_H(\varphi_F^1, \text{Id}) \leq \frac{1}{N} \max(H),$$

from the definition of the Hofer norm. Therefore, to prove the claim it is sufficient to show that $d_H(\varphi_H^1, \varphi_F^1) \leq 2$. This can be proved using inequalities (5) and (6) as follows:

$$d_H(\varphi_H^1, \varphi_F^1) = d_H\left(\prod_{i=1}^{N} \varphi_H^{1/N}, \prod_{i=1}^{N} \psi_i \varphi_H^{1/N} \psi_i^{-1}\right)$$

$$\leq \sum_{i=1}^{N} d_H(\varphi_H^{1/N}, \psi_i \varphi_H^{1/N} \psi_i^{-1}) \leq \sum_{i=1}^{N} 2 d_H(\psi_i, \text{Id}) \leq 2. \qquad \blacksquare$$

4.3. Quasi-flats in the kernel of Calabi

We are now ready to present a proof of Theorem 1.4. First note that without loss of generality we may assume that the open set U contains the North Pole p_+. This allows us to use the constructions of the preceding section.

We begin by proving the first part of the theorem, regarding the quasi-flat rank.

Proof of Theorem 1.4 (a). Let $\mathbb{R}_+^n := \{(t_1, \ldots, t_n) : t_i \geq 0 \text{ for all } i\}$. We equip $\mathbb{R}_+^n$ with the distance induced by the sup norm, that is, we define the distance between $(t_1, \ldots, t_n), (s_1, \ldots, s_n) \in \mathbb{R}_+^n$ to be

$$\|(t_1, \ldots, t_n) - (s_1, \ldots, s_n)\|_\infty = \max\{|t_i - s_i| : i = 1, \ldots, n\}.$$

The mapping

$$\Phi : \mathbb{R}_+^n \to \text{Ham}(\mathbb{S}^2, \omega),$$

$$(t_1, \ldots, t_n) \mapsto \varphi_{H_1}^{t_1} \circ \cdots \circ \varphi_{H_n}^{t_n} \circ \varphi_{H_{n+1}}^{-(t_1 + \cdots + t_n)},$$

takes values in $\text{Ham}_U(\mathbb{S}^2, \omega)$. Moreover, since the H_i all have the same integral over $\mathbb{S}^2$, the mapping Φ takes values in the kernel of the Calabi homomorphism.

We will show that there exists an invertible $n \times n$ matrix A such that

$$\|A(\mathbf{t} - \mathbf{s})\|_\infty \leq d_H(\Phi(\mathbf{t}), \Phi(\mathbf{s})) \leq 2n(2\|\mathbf{t} - \mathbf{s}\|_\infty + 1 + 1/n), \tag{40}$$

where $\mathbf{t}, \mathbf{s}$ stand for $(t_1, \ldots, t_n)$ and $(s_1, \ldots, s_n)$, respectively. As a consequence,

$$\frac{1}{\|A^{-1}\|_{\mathrm{op}}} \|\mathbf{t} - \mathbf{s}\|_\infty \le d_H(\Phi(\mathbf{t}), \Phi(\mathbf{s})) \le 4n\|\mathbf{t} - \mathbf{s}\|_\infty + 2n + 2,$$

where $\|A^{-1}\|_{\mathrm{op}}$ is the operator norm of A^{-1} as a linear map of the normed space $(\mathbb{R}^n, \|\cdot\|_\infty)$.

The above clearly implies that Φ is a quasi-isometric embedding of $(\mathbb{R}^n_+, \|\cdot\|_\infty)$ into $(\mathrm{Ham}(\mathbb{S}^2, \omega), d_H)$. This establishes Theorem 1.4 (a) because $(\mathbb{R}^n, \|\cdot\|_\infty)$ quasi-isometrically embeds into $(\mathbb{R}^{2n}_+, \|\cdot\|_\infty)$; an explicit formula for such an embedding is

$$\mathbf{L} : \mathbb{R}^n \to \mathbb{R}^{2n}_+,$$
$$(x_1, \ldots, x_n) \mapsto (L(x_1), \ldots, L(x_n)),$$

where $L : \mathbb{R} \to \mathbb{R}^2_+$ is defined as

$$L(x) := \begin{cases} (0, -x), & x \le 0, \\ (x, 0), & x \ge 0. \end{cases}$$

For a proof of the fact that $\mathbf{L}$ is a quasi-isometric embedding see [51, Lem. 8.12].

We now turn our attention to (40), beginning with the proof of the right inequality:

$$d_H(\Phi(\mathbf{t}), \Phi(\mathbf{s})) = \|\Phi(\mathbf{t}) \circ \Phi(\mathbf{s})^{-1}\| = \|\varphi_{H_1}^{t_1 - s_1} \cdots \varphi_{H_n}^{t_n - s_n} \varphi_{H_{n+1}}^{-\sum_i (t_i - s_i)}\|$$

$$\le \|\varphi_{H_{n+1}}^{-\sum_i (t_i - s_i)}\| + \sum_{i=1}^{n} \|\varphi_{H_i}^{t_i - s_i}\| \le \sum_{i=1}^{n} (4|t_i - s_i| + 2) + 2$$

$$\le 2n \left(2\|\mathbf{t} - \mathbf{s}\|_\infty + 1 + \frac{1}{n} \right).$$

Above, the second equality on the first line is a consequence of Lemma 4.2 (6), the first inequality on the second line follows from the triangle inequality, and the second inequality in the second line is a consequence of Lemma 4.2 (3).

It remains to prove the left inequality of (40). Let us consider the following two families of monotone twists:

$$\alpha(\mathbf{t}) = \varphi_{H_1}^{t_1} \cdots \varphi_{H_n}^{t_n}, \quad \beta(\mathbf{t}) = \varphi_{H_{n+1}}^{t_1 + \cdots + t_n}.$$

By definition $\Phi(\mathbf{t}) = \alpha(\mathbf{t})\beta(\mathbf{t})^{-1}$ and since monotone twists commute,

$$d_H(\Phi(\mathbf{t}), \Phi(\mathbf{s})) = \|\beta(\mathbf{t})\alpha(\mathbf{t})^{-1}\alpha(\mathbf{s})\beta(\mathbf{s})^{-1}\| = \|(\alpha(\mathbf{t})\beta(\mathbf{s}))^{-1}\alpha(\mathbf{s})\beta(\mathbf{t})\|$$
$$= d_H(\alpha(\mathbf{t})\beta(\mathbf{s}), \alpha(\mathbf{s})\beta(\mathbf{t})).$$

Combining this with Proposition 3.8 (2) gives

$$\max_{i=1,\ldots,n} \left| \frac{\mu_{d_i}(\alpha(\mathbf{t})\beta(\mathbf{s}))}{d_i} - \frac{\mu_{d_i}(\alpha(\mathbf{s})\beta(\mathbf{t}))}{d_i} \right| \le d_H(\Phi(\mathbf{t}), \Phi(\mathbf{s})). \tag{41}$$

By Proposition 3.8 (4) and (36), we can write

$$\mu_d(\alpha(\mathbf{t})\beta(\mathbf{s})) - \mu_d(\alpha(\mathbf{s})\beta(\mathbf{t})) = \sum_{j=1}^{n} \left(\mu_d(\varphi_{H_j}^1) - \mu_d(\varphi_{H_{n+1}}^1) \right)(t_j - s_j)$$

for any d. It follows that the left hand side in (41) equals $\|A(\mathbf{t} - \mathbf{s})\|_\infty$, where A is the matrix whose ij entry (for $1 \le i, j \le n$) is

$$A_{ij} = \frac{\mu_{d_i}(\varphi_{H_j}^1) - \mu_{d_i}(\varphi_{H_{n+1}}^1)}{d_i}.$$

Lemma 4.2 (4) tells us that $\mu_{d_i}(\varphi_{H_{n+1}}^1) = \mu_{d_i}(\varphi_{H_j}^1) = -d_i/2$ for $j > i$. It follows that $A_{ij} = 0$ for $j > i$, i.e. the matrix A is lower triangular. From Lemma 4.2 (5) we deduce that the diagonal entries of A are nonzero. Hence, A is invertible, which proves (40). We have completed the proof of Theorem 1.4 (a). ∎

4.4. The kernel of Calabi is not coarsely proper

In this section, we prove the remainder of Theorem 1.4, that the kernel of the Calabi homomorphism defined on $\mathrm{Ham}_U(\mathbb{S}^2, \omega)$ is not coarsely proper. Recall from the introduction that a metric space (X, d) is said to be coarsely proper if there exists $R_0 > 0$ such that every bounded subset of (X, d) can be covered by finitely many balls of radius R_0.

Proof of Theorem 1.4 (b). For any fixed $r > 0$ consider the set

$$X_r := \{\varphi_{H_1}^{-r}\varphi_{H_{2i}}^r : i \ge 1\},$$

where the H_i are the Hamiltonians provided by Lemma 4.2.

Claim 4.4. *The set X_r is $\frac{3r}{16}$-separated, i.e. for $i \ne j$ we have*

$$3r/16 \le d_H(\varphi_{H_1}^{-r}\varphi_{H_{2i}}^r, \varphi_{H_1}^{-r}\varphi_{H_{2j}}^r).$$

The above claim implies that the set X_r, which is bounded by Lemma 4.2, cannot be covered by finitely many balls of radius $r/16$. Since this holds for every value of r, and since X_r is included in $\ker(\mathrm{Cal})$, we conclude that $\ker(\mathrm{Cal})$ is not coarsely proper, which proves the theorem.

Proof of Claim 4.4. Suppose that $i < j$, pick $k \in \mathbb{N}$ such that $2i < k < 2j$ and consider d_k as in Lemma 4.2. By the Hofer continuity property of μ_{d_k} from Proposition 3.8, we have

$$\frac{1}{d_k}|\mu_{d_k}(\varphi_{H_{2i}}^r) - \mu_{d_k}(\varphi_{H_{2j}}^r)| \le d_H(\varphi_{H_{2i}}^r, \varphi_{H_{2j}}^r) = d_H(\varphi_{H_1}^{-r}\varphi_{H_{2i}}^r, \varphi_{H_1}^{-r}\varphi_{H_{2j}}^r).$$

By Lemma 4.2, we have $\mu_{d_k}(\varphi_{H_{2j}}^r) = -rd_k/2$ and $\mu_{d_k}(\varphi_{H_{2i}}^r) \ge -r\frac{5d_k}{16}$. Thus, $\frac{1}{d_k}|\mu_{d_k}(\varphi_{H_{2i}}^r) - \mu_{d_k}(\varphi_{H_{2j}}^r)| \ge 3r/16$, which completes the proof. ∎

We have now proved Theorem 1.4 (b). ∎

5. Nonsimplicity of $\mathrm{Homeo}_0(\mathbb{S}^2, \omega)$

We conclude by proving Theorem 1.7.

5.1. Outline of the argument

To prove nonsimplicity of $\mathrm{Homeo}_0(\mathbb{S}^2, \omega)$, we explicitly construct a proper normal subgroup which we call the group of *finite-energy homeomorphisms* and denote by $\mathrm{FHomeo}(\mathbb{S}^2, \omega)$. We introduced these homeomorphisms in [9] where we proved that they form a proper normal subgroup of the compactly supported area-preserving homeomorphisms of the disc. Here, we will give a slight variant of the definition in [9] which is more natural from the point of view of Hofer's geometry.

Definition 5.1. We say $\varphi \in \mathrm{Homeo}_0(\mathbb{S}^2, \omega)$ is a *finite-energy homeomorphism* if there exists a sequence $\{\varphi_i\}_{i \in \mathbb{N}}$ of Hamiltonian diffeomorphisms which is bounded with respect to Hofer's distance and which converges uniformly to φ. We denote by $\mathrm{FHomeo}(\mathbb{S}^2, \omega)$ the set of all finite-energy homeomorphisms.

Theorem 1.7 follows immediately from the following result, whose proof will occupy the remainder of this section.

Theorem 5.2. $\mathrm{FHomeo}(\mathbb{S}^2, \omega)$ *is a proper normal subgroup of* $\mathrm{Homeo}_0(\mathbb{S}^2, \omega)$.

We prove the above using arguments similar to those given in [9]. Here is a brief outline. As we shall see, it is not hard to show that $\mathrm{FHomeo}(\mathbb{S}^2, \omega)$ forms a normal subgroup of $\mathrm{Homeo}_0(\mathbb{S}^2, \omega)$; the main challenge is proving that it is proper.

To do this, we use the invariant $\eta_d : \mathrm{Ham}(\mathbb{S}^2, \omega) \to \mathbb{R}$. We showed above that it is continuous with respect to the C^0-topology on $\mathrm{Ham}(\mathbb{S}^2, \omega)$, and moreover it extends continuously to $\mathrm{Homeo}_0(\mathbb{S}^2, \omega)$; see Proposition 3.9. A straightforward argument shows that for any $\varphi \in \mathrm{FHomeo}(\mathbb{S}^2, \omega)$ there exists a constant C, depending on φ, such that for all (even) d we have

$$\frac{\eta_d(\varphi)}{d} \leq C. \tag{42}$$

We will prove Theorem 5.2 by showing that certain so-called *infinite twist* homeomorphisms $\psi \in \mathrm{Homeo}_0(\mathbb{S}^2, \omega)$ satisfy

$$\lim_{d \to \infty} \frac{\eta_d(\psi)}{d} = \infty, \tag{43}$$

so fail to satisfy (42). This last step requires estimating $\eta_d(\psi)$, for which we rely on the combinatorial model from Section 4.1.1.

We end this section by highlighting the differences between our proof, in this article, of nonsimplicity of $\mathrm{Homeo}_0(\mathbb{S}^2, \omega)$ and the proof of nonsimplicity of $\mathrm{Homeo}_c(\mathbb{D}^2, \omega)$ given in [9]. In both articles we use PFH spectral invariants $c_d : C^\infty(\mathbb{S}^1 \times \mathbb{S}^2) \to \mathbb{R}$. Given an arbitrary Hamiltonian H, the value of $c_d(H)$ depends on H and so c_d does

not yield a well-defined invariant of Hamiltonian diffeomorphisms. However, in [9] we overcome this problem by restricting the domain of c_d to a certain class of Hamiltonians which is suitable for the purposes of that article; see [9, Sec. 3.4]. In the current article, we do not have the possibility of restricting the domain of c_d. Instead, we work with η_d, which is well-defined for all Hamiltonian diffeomorphisms of the sphere as proved in Section 3.2.

Another difference between the two proofs is the manner in which we show properness of FHomeo. In both articles this is achieved by exhibiting area-preserving homeomorphisms ψ satisfying (43). The proof of this given in [9] involves verifying for certain smooth twist maps a conjecture of Hutchings, concerning recovering the Calabi invariant from the asymptotics of PFH spectral invariants, whereas our proof here, which is shorter, uses the forthcoming Claim 5.4. The proof of this claim, however, relies on the combinatorial model for PFH developed in [9, Sec. 5]. We should remark that part of the motivation for the somewhat longer argument in [9] was that Hutchings' conjecture is of independent interest, hence useful to verify for twist maps.

5.2. An infinite twist is not a finite-energy homeomorphism

We now carry out the above outline. We begin by describing the infinite twist homeomorphisms ψ.

Denote by $p_+ \in \mathbb{S}^2$ the North Pole of the sphere, i.e. the point whose z-coordinate is 1, in the cylindrical coordinate system introduced in Section 2.1. We say a function $F : \mathbb{S}^2 \setminus \{p_+\} \to \mathbb{R}$ is an *infinite twist Hamiltonian* if it is of the form

$$F(z, \theta) = \tfrac{1}{2} f(z), \tag{44}$$

where $f : [-1, 1) \to \mathbb{R}$ is a smooth function such that $f' \geq 0$, $f'' \geq 0$ and

$$\lim_{d \to \infty} \frac{1}{d} f\left(1 - \frac{2}{d+1}\right) = \infty. \tag{45}$$

Observe that F defines a smooth Hamiltonian on $\mathbb{S}^2 \setminus \{p_+\}$ whose flow is given by

$$\varphi_F^t(\theta, z) = (\theta + 2\pi t f'(z), z).$$

We extend the flow φ_F^t to $\mathbb{S}^2$ by defining $\varphi_F^t(p_+) = p_+$; this yields an area-preserving flow on $\mathbb{S}^2$ which is nonsmooth at the point p_+. We say $\psi \in \mathrm{Homeo}_0(\mathbb{S}^2, \omega)$ is an *infinite twist homeomorphism* if it is of the form

$$\psi := \varphi_F^1 \tag{46}$$

for some F. We will call ψ an *adapted* infinite twist if the corresponding f satisfies the following technical hypothesis:

$$f'\left(1 - \frac{2}{d+1}\right) \in (d+1)\mathbb{N}$$

for all $d \geq 2$.

We can now give the promised proof of the remaining theorem.

Proof of Theorem 5.2. We begin by noting that the argument in [9, Prop. 2.1], repeated verbatim, shows that $\mathrm{FHomeo}(\mathbb{S}^2, \omega)$ forms a normal subgroup of $\mathrm{Homeo}_0(\mathbb{S}^2, \omega)$. It remains to show that it is proper.

Step 1: Linear growth in FHomeo. We first show that for any $\varphi \in \mathrm{FHomeo}(\mathbb{S}^2, \omega)$ the linear growth condition (42) holds. This is an immediate consequence of the properties in Proposition 3.9. Indeed, let $\varphi \in \mathrm{FHomeo}(\mathbb{S}^2, \omega)$ and choose a sequence $\varphi_i \xrightarrow{C^0} \varphi$ that is uniformly bounded with respect to Hofer's distance. Since the φ_i are bounded and $\eta_d(\mathrm{Id}) = 0$, the Hofer continuity property ensures a bound of the form $\eta_d(\varphi_i) \leq d \cdot C$ for some uniform constant C; then, by C^0-continuity, the same bound holds for φ.

Step 2: Superlinear growth of some infinite twists. It remains to prove that FHomeo is proper. The structure of the remainder of our argument will now be to prove properness, assuming the technical Claim 5.3 below which makes use of the adapted condition, and then prove the claim. From now until the end of the paper, we therefore assume that F is an adapted infinite twist Hamiltonian whose support is contained in the interior of the disc $\{(z, \theta) : z \geq 7/8\}$, which is of area $1/16$. Imposing this assumption enables us to apply the following promised technical claim. Recall below that the η_d are defined only for even d.

Claim 5.3. *Fix $d \geq 4$, define $z_0 := 1 - \frac{2}{d+1}$. Let H be a smooth monotone twist Hamiltonian supported in a disc of area at most $1/12$. Assume that $p := h'(z_0) \in (d + 1)\mathbb{N}$. Then*

$$\eta_d(\varphi_H^1) \geq H(z_0) - d/6.$$

We defer the proof for the moment. Assuming it, we can produce superlinear growth of the η_d as follows.

Claim 5.4. $\eta_d(\varphi_F^1) \geq \frac{1}{2}f(1 - \frac{2}{d+1}) - \frac{d}{6}$ *for $d \geq 4$.*

Proof of Claim 5.4. For every $i \in \mathbb{N}$, let $F_i : \mathbb{S}^2 \to \mathbb{R}$ be a sequence of smooth Hamiltonians of the form

$$F_i(z, \theta) = \tfrac{1}{2} f_i(z),$$

where $f_i : [-1, 1] \to \mathbb{R}$ is a smooth function such that $f_i' \geq 0$, $f_i'' \geq 0$ and $f_i(z) = f(z)$ for $z \in [-1, 1 - 1/i]$.

Observe that $\varphi_{F_i}^1 \xrightarrow{C^0} \varphi_F^1$, because $(\varphi_{F_i}^1)^{-1} \circ \varphi_F^1$ is supported in the disc $\{(z, \theta) : z \geq 1 - 1/i\}$. Hence, by the C^0-continuity of η_d established in Proposition 3.9, we have

$$\eta_d(\varphi_F^1) = \lim_{i \to \infty} \eta_d(\varphi_{F_i}^1)$$

for every d. Applying Claim 5.3 to F_i, for i sufficiently large with respect to d, yields

$$\eta_d(\varphi_{F_i}^1) \geq F_i\left(1 - \frac{2}{d+1}\right) - \frac{d}{6} = \frac{1}{2}f_i\left(1 - \frac{2}{d+1}\right) - \frac{d}{6} = \frac{1}{2}f\left(1 - \frac{2}{d+1}\right) - \frac{d}{6},$$

for $d > 3$, as claimed. ∎

It follows immediately from the previous claim that $\psi := \varphi_F^1$ satisfies (43), which, as explained in Step 1, implies that an adapted infinite twist φ_F^1 is not a finite-energy homeomorphism.

To complete the proof of Theorem 5.2, it therefore remains to prove Claim 5.3.

Proof of Claim 5.3. Recall from (27) that $\eta_d(\varphi) = c_d(H) - \frac{d}{2}c_2(H)$. Hence, to prove the claim, it is sufficient to show that the following two inequalities hold:

$$c_2(H) \le \frac{1}{3}, \tag{47}$$

$$c_d(H) \ge H\left(1 - \frac{2}{d+1}\right). \tag{48}$$

To prove (47), we invoke the support-control inequality of Proposition 3.2, which gives $c_2(H) \le 2 \cdot 2 \cdot \frac{1}{12}$, since the area of the support of H is bounded by $1/12$.

Next, we prove (48). By the continuity property of c_d from Theorem 3.1, we may perform a small perturbation of h, near $z = 1$, and assume that $h'(1) \in \mathbb{N}$, in other words that our twist is nice. This allows us to apply Theorem 6.1 of [9] whose statement we recalled in Section 4.1.1.

Recall the notation $z_0 = 1 - \frac{2}{d+1}$. By [9, Thm. 6.1] we have

$$c_d(H) \ge \mathcal{A}(P)$$

for any degree d lattice path P of combinatorial index $I(P) = -d$; see Section 4.1.1.

Recall the notation $p := h'(z_0)$, which is by assumption an integer. By assumption, there exists an integer $a > 0$ such that $p = a(d+1)$. Take P to be the lattice path obtained by joining the lattice points $(0, -a)$, $(d-1, -a)$ and $(d, p-a)$. This is a concave lattice path made up of two edges. It satisfies

$$\mathcal{A}(P) = \frac{p}{2}(1-z_0) + \frac{1}{2}h(z_0) - a,$$
$$I(P) = 2j(P) - d = 2((p-a) - da) - d = -d.$$

Hence,

$$c_d(H) \ge \frac{p}{2}(1-z_0) + \frac{1}{2}h(z_0) - a = \frac{p}{2}\left(1 - z_0 - \frac{2}{d+1}\right) + \frac{1}{2}h(z_0)$$
$$= \frac{1}{2}h(z_0) = H\left(1 - \frac{2}{d+1}\right). \qquad \blacksquare$$

We have thus completed the proof of Theorem 5.2. $\qquad \blacksquare$

Remark 5.5. The infinite twist Hamiltonian F, introduced above, generates a 1-parameter subgroup φ_F^t of $\mathrm{Homeo}_0(\mathbb{S}^2, \omega)$. It follows immediately from the above proof that $\varphi_F^t \notin \mathrm{FHomeo}(\mathbb{S}^2, \omega)$ for $t \ne 0$. This yields an injective group homomorphism from the real line $\mathbb{R}$ into the quotient $\mathrm{Homeo}_0(\mathbb{S}^2, \omega)/\mathrm{FHomeo}(\mathbb{S}^2, \omega)$. One can show that this injection is not a surjection; see [45]. However, we have not been able to determine whether $\mathrm{Homeo}_0(\mathbb{S}^2, \omega)/\mathrm{FHomeo}(\mathbb{S}^2, \omega)$ is isomorphic to $\mathbb{R}$ as an abelian group.

Remark 5.6. The group $\text{FHomeo}(\Sigma, \omega)$ of finite-energy homeomorphisms can be defined on any surface Σ; it forms a normal subgroup of $\overline{\text{Ham}}_c(\Sigma, \omega)$, the group of Hamiltonian homeomorphisms of (Σ, ω) which coincide with the identity near the boundary, if $\partial \Sigma \neq \emptyset$. Recall that a Hamiltonian homeomorphism is a homeomorphism which can be written as the C^0-limit of Hamiltonian diffeomorphisms. It is well-known that $\overline{\text{Ham}}_c(\Sigma, \omega)$ coincides with the kernel of the mass-flow homomorphism of Fathi [15].

Suppose now that (Σ, ω) is any compact surface of genus zero, with boundary, and view it as embedded into $(\mathbb{S}^2, \omega)$. There is an inclusion $\text{FHomeo}(\Sigma, \omega) \subset \text{FHomeo}(\mathbb{S}^2, \omega)$. The infinite twist ψ can be placed on (Σ, ω) and the fact that it is not a finite-energy homeomorphism of the sphere implies that $\psi \notin \text{FHomeo}(\Sigma, \omega)$. We conclude that $\overline{\text{Ham}}_c(\Sigma, \omega)$ is not simple. Moreover, as in the case of Corollary 1.8, one can conclude that $\overline{\text{Ham}}_c(\Sigma, \omega)$ is not perfect either. This answers a question of Fathi [15, Appendix A.6], concerning the simplicity of the kernel of the mass-flow homomorphism, for compact genus-zero surfaces.

We remark that the infinite twist and FHomeo can be defined on any symplectic manifold. However, our methods for proving properness of finite-energy homeomorphisms are currently applicable to dimension 2 only.

5.3. Proof of Corollary 1.9

We may assume without loss of generality that $\int_{\mathbb{R}^2} \Omega = 1$. As alluded to in the introduction, by a version of Moser's argument for noncompact manifolds, due to Greene and Shiohama [16], there exists a smooth diffeomorphism $\psi : \mathbb{R}^2 \to \mathbb{S}^2 \setminus \{p\}$ such that $\psi^* \omega = \Omega$. Here p denotes the North Pole in $\mathbb{S}^2$. This gives rise to an injective group homomorphism $\Psi : \text{Diff}(\mathbb{R}^2, \Omega) \to \text{Homeo}_0(\mathbb{S}^2, \omega)$, defined for any $h \in \text{Diff}(\mathbb{R}^2, \Omega)$ by $\Psi(h)(x) := \psi h \psi^{-1}(x)$ for $x \neq p$ and $\Psi(h)(p) = p$. The image of Ψ is the set of elements of $\text{Homeo}_0(\mathbb{S}^2, \omega)$ that fix p and are smooth in the complement of p.

In particular, the image of Ψ contains an adapted infinite twist homeomorphism τ, which we showed above is not in $\text{FHomeo}(\mathbb{S}^2, \omega)$. By the Epstein–Higman argument cited in the introduction, the commutator subgroup of Homeo_0 is contained in any nontrivial normal subgroup. In particular, τ is not a product of commutators. This implies that $\Psi^{-1}(\tau)$ is not a product of commutators in $\text{Diff}(\mathbb{R}^2, \Omega)$, which is therefore not perfect. ∎

Remark 5.7. The above argument similarly shows that $\text{Homeo}(\mathbb{R}^2, \Omega)$, the group of area-preserving homeomorphisms of the plane, is not perfect if Ω has finite total area. This holds more generally if Ω is only assumed to be a good measure and not necessarily a smooth area form; being good means that Ω is nonatomic and positive on nonempty open sets. In this case, one can repeat the above argument, using the classical Oxtoby–Ulam theorem [39] instead of [16].

5.4. Remarks on Hofer's geometry

We close by briefly discussing the large scale geometry of FHomeo.

It is possible to define Hofer's distance for area-preserving homeomorphisms as follows. Given $\varphi \in \text{FHomeo}(\mathbb{S}^2, \omega)$, we define its Hofer distance from the identity by

$$\tilde{d}_H(\varphi, \text{Id}) := \inf d_H(\varphi_i, \text{Id}), \tag{49}$$

where the infimum is taken over all sequences $\{\varphi_i\} \subset \text{Ham}(\mathbb{S}^2, \omega)$ which converge uniformly to φ. Define $\tilde{d}_H(\varphi, \psi) := \tilde{d}_H(\varphi^{-1}\psi, \text{Id})$.

We leave it to the reader to check that this defines a bi-invariant distance on $\text{FHomeo}(\mathbb{S}^2, \omega)$.

It is a natural question to try to better understand this space. For example, one could ask if FHomeo has infinite quasi-flat rank. We strongly suspect that the answer is, in fact, positive as our tools are robust with respect to the C^0-topology and so one can adapt the proof of Theorem 1.4 to prove that FHomeo does have infinite quasi-flat rank. Similarly, it can be shown that FHomeo is not coarsely proper.

One could define $\tilde{d}_H(\varphi, \text{Id})$, via (49), for arbitrary $\varphi \in \text{Homeo}_0(\mathbb{S}^2, \omega)$.

If φ is not a finite energy homeomorphism, i.e. if $\varphi \notin \text{FHomeo}(\mathbb{S}^2, \omega)$, then we get

$$\tilde{d}_H(\varphi, \text{Id}) = \infty.$$

Hence, we may view homeomorphisms which are not finite-energy as those which are infinitely far from diffeomorphisms, in Hofer's distance. This is the point of view expressed in Le Roux's article [32, Question 1]. Theorem 5.2 tells us that such homeomorphisms do exist.

A question which arises immediately as a consequence of our definition of $\tilde{d}_H$ is whether $\tilde{d}_H(\varphi, \psi)$ coincides with the usual Hofer distance $d_H(\varphi, \psi)$ when $\varphi, \psi \in \text{Ham}(\mathbb{S}^2, \omega)$. We do not know the answer to this question. Note that this is equivalent to asking if the (usual) Hofer distance is lower semicontinuous with respect to the C^0-topology; this was posed as an open question by Le Roux [32].

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