

Global weak solutions to a time-periodic body-liquid interaction problem

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Abstract. We prove existence of time-periodic weak solutions to the coupled liquid-structure problem constituted by an incompressible Navier–Stokes fluid interacting with a rigid body of finite size, subject to an *undamped* linear restoring force. The fluid flow is generated by a uniform, time-periodic velocity field V far from the body. We emphasize that our result is global, in the sense that no restriction is imposed on the magnitude of V and, rather remarkably, the frequency of V is entirely arbitrary. Thus, in particular, it can coincide with any multiple of a natural frequency of vibration of the body so that, with this model, resonance cannot occur. Although based on the classical “invading domains” technique, our approach requires several new ideas. Indeed, due to the lack of sufficient dissipation, it appears quite unfeasible to show the existence of a fixed point of the Poincaré map at the finite-dimensional level along the Galerkin approximant. Therefore, unlike the usual strategy, such a result must be proven directly in a class of weak solutions, and therefore in the infinite-dimensional framework.

Introduction

One among the many, captivating problems concerning the interaction of a liquid with an elastic structure arises when the liquid is in a time-periodic regime so as to generate a similar motion of the structure. Of particular interest is the case when the frequency of the flow, ω , approaches or even coincides with a multiple of a natural frequency of vibration, ω_n , of the body. Then, the interaction may result in a resonance phenomenon, which is considered as primarily responsible for the possible failure of the structure [4]. Such a problem falls into the general area of vibration-induced oscillations that has all along constituted a main focus of applied science and is at the heart of a vast engineering literature; see the monographs [4, 11, 25] and the bibliography therein.

In the classical model employed to investigate this type of question, one regards the structure, \mathcal{B} , as a rigid body subject to a linear restoring force, while the liquid, \mathcal{L} , is described by the Navier–Stokes equations [3, 10, 23, 31]. In such a framework, the occurrence of resonance is usually explained as follows [4, Section 3.5], [2]: In the time-periodic regime of the liquid, the motion of the structure is that of a simple harmonic oscillator subject to a forced time-periodic motion caused by the action of the force, F , exerted by \mathcal{L}

on \mathcal{B} . Then, one can show that for ω close to (a multiple of) ω_n , the amplitude of the forced oscillations of \mathcal{B} becomes very large and, in the absence of a damping mechanism, tends to infinity. A significant consequence of this argument is that, *without damping*, the coupled system $\mathcal{S} := \mathcal{B} - \mathcal{L}$ cannot perform a time-periodic motion of *arbitrary* frequency ω . While seemingly reasonable, this way of thinking appears rather simplistic. First, it decouples the highly coupled system \mathcal{S} , by somehow prescribing the action of \mathcal{L} on \mathcal{B} . Moreover, it disregards the fact that F splits into two competing components: the one that, indeed, forces the oscillations of \mathcal{B} , and the other, due to viscosity effects, which tends to absorb them. As a result, if the latter prevails over the former, occurrence of resonance can be excluded, even in the *absence* of a structural damping mechanism.

Also motivated by the above considerations, very recently the authors and their collaborators started a *rigorous* and systematic analysis of this model for different flow geometries, by investigating general mathematical properties, such as existence, uniqueness, stability, and bifurcation of solutions [1, 5–9, 19, 20, 26]. For other contributions related to a similar problem, we also refer to [16, 24].

In this paper we aim at furnishing a further contribution, which may help a better understanding of the resonance phenomenon. More specifically, we consider the general case when \mathcal{B} (of arbitrary shape) is subject to a – possibly anisotropic – linear restoring force, and is immersed in a Navier–Stokes liquid filling the whole space, Ω , outside \mathcal{B} . The motion of the coupled system \mathcal{S} is driven by a time-periodic uniform velocity, $V = V(t)$, of period T impressed on the liquid at large distance from \mathcal{B} . We assume the worst-case scenario, namely, *no external damping* mechanism acts on the structure, so that the only dissipative effect is due to the viscosity of \mathcal{L} . We then ask the following question: Will the coupled system \mathcal{S} perform a time-periodic motion of period T , for *arbitrary* T and *arbitrary* magnitude of V ? A positive answer would suggest that the model used so far is probably not appropriate for resonance studies, and that other aspects should be accounted for.

The main achievement of this paper is to show that, indeed, the above question admits an affirmative answer, provided only that V has a mild degree of regularity and \mathcal{B} is of class C^2 . As expected, since we want to keep the “size” of V arbitrary, we are lead to perform this study in a suitable class of weak solutions; see Definition 3.1. The approach we use is, in principle, rather usual for time-periodic flow in exterior domains, and employs the “invading domains” technique [15, 18]. However, in the case at hand, its implementation is by no means straightforward, and presents a number of difficulties that are described next.

We recall that the above technique develops along the following steps. One picks an increasing sequence of bounded domains, $\{\Omega_n\}$, whose union coincides with Ω , and suitably reformulates the original problem in each Ω_n . Then, following an idea due to Prouse [28], a time-periodic solution in Ω_n is searched via the finite-dimensional Galerkin approximation, by showing with the help of Brouwer’s theorem that the Poincaré map, M , bringing initial conditions into corresponding solutions at time T , has a fixed point. This fact, in conjunction with appropriate uniform estimates in n , allows one to construct a

solution on each Ω_n , and eventually pass to the limit $n \rightarrow \infty$ to obtain a (weak) solution to the original problem. It should be emphasized that, for this approach to work, it is crucial that along the Galerkin approximant the total energy of the system, E , in the absence of a forcing term, be bounded by an exponentially decreasing function of time, to guarantee that M is a self-map. This means that the energy equation should contain a dissipative term proportional to E . It is exactly here that, in our case, the first difficulty arises. Actually, E is the sum of the kinetic energy ($K_{\mathcal{L}}$) of the liquid, and the kinetic ($K_{\mathcal{B}}$) and potential ($U_{\mathcal{B}}$) energies of the body. By using viscosity dissipation combined with trace theorems we obtain some damping terms for both $K_{\mathcal{L}}$ and $K_{\mathcal{B}}$ but, unsurprisingly, no damping for $U_{\mathcal{B}}$, namely, for the oscillations amplitude of \mathcal{B} . In order to get the latter, one may think of adapting a procedure seemingly introduced by Haraux in the context of non-linear wave equations [21, p. 162ff]; see also [22]. However, though this procedure formally works on the original system of equations in Ω_n , it is rather doubtful that it can be applied at the finite-dimensional level along the Galerkin approximant. Therefore, we are forced to find a fixed point of the Poincaré map M directly on that system. Since we want existence for data of arbitrary size, the fixed point should be found in a class of weak solutions. Moreover, now being in an infinite-dimensional framework, M should also possess suitable compactness properties. This type of question has been addressed by Prodi for classical Navier–Stokes equations [27]. Nevertheless, his method requires uniqueness and continuous dependence of solutions upon the data in the energy norm, properties that, to date, are only known to hold in two dimensions. As a consequence, Prodi’s method is inapplicable to our case, which produces yet another difficulty. To overcome this issue, we introduce an entirely different strategy that combines a suitable mollification of the non-linear term in the original problem in Ω_n , along with the use of time-weighted norms that vanish at $t = 0$. Incidentally, we remark that, by using this new strategy, one could extend Prodi’s result to arbitrary dimension $d \geq 2$.

Our approach employs the following steps: Thanks to the regularization procedure, we are able to show the existence of global weak solutions to the mollified problem for initial data possessing only finite energy, that however are strong (à la Prodi–Ladyzhenskaya) at any positive time; see Lemma 3.1. In this class of solutions it is easily shown that the Poincaré map is compact in the energy space. Thus, in order to prove existence of time-periodic solutions to the mollified problem, it remains to ascertain that M maps some ball in the energy space into itself. This property is shown by adapting Haraux’s argument mentioned earlier on, which gives the desired dissipation also for $U_{\mathcal{B}}$; see Lemma 3.5. In this way, by letting the mollification parameter go to 0, we finally deduce the existence of a time-periodic weak solution in every Ω_n ; see Proposition 3.1. The last step is to let $n \rightarrow \infty$, and this brings the last difficulty. Actually, for the dissipative term for $U_{\mathcal{B}}$, it does not appear possible to provide an estimate that is *uniform* in n , which means that, as $n \rightarrow \infty$, we have no control on the amplitude of the oscillations of \mathcal{B} . However, we prove a uniform control on the velocity of the center of mass of \mathcal{B} along with estimates for the time derivative of the flow velocity in suitable distributional spaces. Combining these estimates

allows us to deduce, by a (local) compactness argument, existence of time-periodic weak solutions for the original problem in the whole Ω ; see Section 3.6.

We conclude this introductory section with the following remark. In the model considered in the present paper (and precisely presented in the next section), the body \mathcal{B} can move only by translational motion and is not free to rotate. Removing this constraint and adding a corresponding restoring torque would lead to a more complete model that might present resonance phenomena, owing to possible interaction of the different degrees of freedom. Plainly, the mathematical analysis of such a model becomes far more complicated, also because of the notorious difficulty due to the presence of an unbounded coefficient in the flow equations [13]. The investigation of this topic will be the object of future work.

The plan of the paper is as follows. After introducing the mathematical formulation of the problem in Section 1, in the following Section 2 we introduce the basic function spaces and collect some preliminary results that, among other things, help us to furnish a suitable reformulation of the original problem. In Section 3.1 we state our main result in Theorem 3.1 and present the strategy we use for its proof. Sections 3.2–3.5 are dedicated to the proof of existence of time-periodic solutions to a suitable modification of the problem in an arbitrary bounded domain, strictly containing \mathcal{B} , along with uniform estimates independent of the “size” of the domain. With this result in hand, in the final Section 3.6 we are able to produce a full proof of Theorem 3.1.

1. Formulation of the problem

Consider a rigid body \mathcal{B} , occupying the closure of the bounded domain Ω_0 , completely surrounded by a Navier–Stokes liquid, \mathcal{L} , filling the entire space, Ω , outside \mathcal{B} . Then \mathcal{B} is subject to an elastic restoring force, \mathbf{R} , applied at its center of mass G . We take \mathbf{R} to be linear, but not necessarily isotropic, that is,

$$\mathbf{R} = -\hat{\mathbf{A}} \cdot \boldsymbol{\delta},$$

where $\boldsymbol{\delta} := \overrightarrow{GO}$, O is a fixed point, and $\hat{\mathbf{A}}$ is a 3×3 symmetric and positive definite real matrix (stiffness matrix). We assume that \mathcal{B} can move only by translational motion, which can be accomplished by having a suitable torque acting on it. The motion of the coupled system body-liquid is driven by a time-periodic flow of \mathcal{L} imposed at “large” spatial distances from \mathcal{B} and characterized by a uniform T -periodic velocity field $-\hat{\mathbf{V}}$, where $\hat{\mathbf{V}}$ is a bounded function of time t only, satisfying therefore

$$\hat{\mathbf{V}}(t) = \hat{\mathbf{V}}(t + T), \quad \text{for all } t \in \mathbb{R}.$$

Denote by L and M the diameter and mass of \mathcal{B} , and by ρ and μ the density and shear viscosity coefficient of the liquid. Then the T -periodic motion of the coupled system body-liquid when referred to a body-fixed frame $\mathcal{F} \equiv \{G, e_i\}$ is governed by the following set

of dimensionless equations (see for instance [13, Section 1]):

$$\left\{ \begin{array}{l} \left. \begin{array}{l} \partial_t \mathbf{v} + \lambda(\mathbf{v} - \boldsymbol{\gamma}) \cdot \nabla \mathbf{v} = \Delta \mathbf{v} - \nabla p, \\ \operatorname{div} \mathbf{v} = 0, \end{array} \right\} \quad \text{in } \Omega \times \mathbb{R}, \\ \mathbf{v}(x, t) = \boldsymbol{\gamma}(t), \quad (x, t) \in \partial\Omega \times \mathbb{R}, \\ \lim_{|x| \rightarrow \infty} \mathbf{v}(x, t) = -\mathbf{V}(t), \quad t \in \mathbb{R}, \\ \left. \begin{array}{l} \boldsymbol{\gamma} + \mathbb{A} \cdot \boldsymbol{\delta} + \varpi \int_{\partial\Omega} \mathbb{T}(\mathbf{v}, p) \cdot \mathbf{n} = \mathbf{0}, \\ \dot{\boldsymbol{\delta}} = \boldsymbol{\gamma}, \end{array} \right\} \quad \text{in } \mathbb{R}. \end{array} \right. \quad (1.1)$$

Here, \mathbf{v} and p are (non-dimensional) velocity and pressure fields of the liquid, while

$$\mathbb{A} := \frac{\rho^2 L^4}{M \mu^2} \hat{\mathbb{A}}, \quad \varpi := \frac{\rho L^3}{M}, \quad \lambda := \frac{\rho V_\infty L}{\mu}, \quad \mathbf{V} := \frac{\hat{\mathbf{V}}}{V_\infty}, \quad V_\infty := \sup_{t \in \mathbb{R}} |\hat{\mathbf{V}}(t)|.$$

Moreover,

$$\mathbb{T}(\mathbf{z}, \psi) := 2\mathbb{D}(\mathbf{z}) - \psi \mathbb{I}, \quad \mathbb{D}(\mathbf{z}) := \frac{1}{2}(\nabla \mathbf{z} + (\nabla \mathbf{z})^\top),$$

with \mathbb{I} the identity matrix, is the (dimensionless) Cauchy stress tensor, and \mathbf{n} the unit outer normal at $\partial\Omega$.

Notice that, with the above non-dimensionalization, we have

$$\sup_{t \in \mathbb{R}} |\mathbf{V}(t)| = 1. \quad (1.2)$$

Our ultimate goal is to show that, for any given $\lambda, \mathbb{A}, \varpi$, and $T > 0$, and any (sufficiently smooth) \mathbf{V} , problem (1.1) has at least one, suitably defined, T -periodic weak solution $(\mathbf{v}, \boldsymbol{\gamma}, \boldsymbol{\delta})$.

2. Preliminaries

2.1. Functional spaces and some related properties

Before describing our functional framework, we begin with some notation. We indicate by $\Omega \subset \mathbb{R}^3$ the exterior domain of class C^2 , defined as the complement of the closure of the bounded domain Ω_0 occupied by \mathcal{B} . We take the origin of coordinates in the interior of Ω_0 , and set $B_r := \{x \in \mathbb{R}^3 : |x| < r, r > 0\}$, and

$$\Omega_R := \Omega \cap B_R, \quad R > R_* := \operatorname{diam} \Omega_0.$$

As is customary, for A a domain of \mathbb{R}^3 , $L^q = L^q(A)$, $W^{m,2} = W^{m,2}(A)$, $q \in [1, \infty]$, $m \in \mathbb{N}$, are Lebesgue and Sobolev spaces with norms $\|\cdot\|_{q,A}$, and $\|\cdot\|_{m,2,A}$. By $(\cdot, \cdot)_A$

we indicate the $L^2(A)$ -scalar product. Furthermore, $D^{m,q} = D^{m,q}(A)$ is the homogeneous Sobolev space with semi-norm $\sum_{|l|=m} \|D^l u\|_{q,A}$. In all the above notation we shall typically omit the subscript “ A ”, unless confusion arises. For a Banach space X , we may, occasionally, indicate its norm by $\|\cdot\|_X$. Finally, by $L^q(I; X)$, $W^{1,q}(I; X)$, and $C^m(I; X)$, where I is a real interval, we shall denote the classical Bochner spaces.

If $A \subseteq \mathbb{R}^3$ is a domain with $A \supset \overline{\Omega_0}$, let

$$\begin{aligned} \mathcal{K} &= \mathcal{K}(A) := \{\varphi \in C_0^\infty(A) : \exists \hat{\varphi} \in \mathbb{R}^3 \text{ s.t. } \varphi(x) = \hat{\varphi} \text{ in a neighborhood of } \Omega_0\}, \\ \mathcal{C} &= \mathcal{C}(A) := \{\varphi \in \mathcal{K}(A) : \operatorname{div} \varphi = 0 \text{ in } A\}, \\ \mathcal{C}_0 &= \mathcal{C}_0(A) := \{\varphi \in \mathcal{C}(A) : \hat{\varphi} = \mathbf{0}\}. \end{aligned}$$

In $\mathcal{K}(A)$ we introduce the scalar product

$$\langle \varphi, \psi \rangle_A := \varpi^{-1} \hat{\varphi} \cdot \hat{\psi} + (\varphi, \psi)_{A \cap \Omega}, \quad \varphi, \psi \in \mathcal{K}, \tag{2.1}$$

and define

$$\begin{aligned} \mathcal{L}^2(\mathbb{R}^3) &:= \{\text{completion of } \mathcal{K}(\mathbb{R}^3) \text{ in the norm induced by (2.1)}\}, \\ \mathcal{H}(\mathbb{R}^3) &:= \{\text{completion of } \mathcal{C}(\mathbb{R}^3) \text{ in the norm induced by (2.1)}\}, \\ \mathcal{G}(\mathbb{R}^3) &:= \{\mathbf{h} \in \mathcal{L}^2(\mathbb{R}^3) : \exists p \in D^{1,2}(\Omega) \text{ s.t. } \mathbf{h} = \nabla p \text{ in } \Omega, \\ &\quad \text{and } \mathbf{h} = -\varpi \int_{\partial\Omega} p \mathbf{n} \text{ in } \Omega_0\}. \end{aligned} \tag{2.2}$$

It is shown in [29, Theorem 3.1 and Lemma 3.2] that

$$\begin{aligned} \mathcal{L}^2(\mathbb{R}^3) &= \{\mathbf{u} \in L^2(\mathbb{R}^3) : \mathbf{u} = \hat{\mathbf{u}} \text{ in } \Omega_0, \text{ for some } \hat{\mathbf{u}} \in \mathbb{R}^3\}, \\ \mathcal{H}(\mathbb{R}^3) &= \{\mathbf{u} \in \mathcal{L}^2(\mathbb{R}^3) : \operatorname{div} \mathbf{u} = 0\}, \end{aligned}$$

along with the following orthogonal decomposition [29, Theorem 3.2]:

$$\mathcal{L}^2(\mathbb{R}^3) = \mathcal{H}(\mathbb{R}^3) \oplus \mathcal{G}(\mathbb{R}^3). \tag{2.3}$$

We next define the space

$$\mathcal{D}^{1,2} = \mathcal{D}^{1,2}(\mathbb{R}^3) := \{\text{completion of } \mathcal{C}(\mathbb{R}^3) \text{ in the norm } \|\mathbb{D}(\cdot)\|_2\},$$

whose basic properties are collected in the next lemma; see [13, Lemmas 9–11].

Lemma 2.1. *The space $\mathcal{D}^{1,2}$ is a separable Hilbert space when equipped with the scalar product*

$$(\mathbb{D}(\mathbf{u}_1), \mathbb{D}(\mathbf{u}_2)), \quad \mathbf{u}_i \in \mathcal{D}^{1,2}, \quad i = 1, 2.$$

Moreover, we have the characterization

$$\mathcal{D}^{1,2} = \{\mathbf{u} \in L^6(\mathbb{R}^3) \cap D^{1,2}(\mathbb{R}^3); \operatorname{div} \mathbf{u} = 0; \mathbf{u} = \hat{\mathbf{u}} \text{ in } \Omega_0, \text{ for some } \hat{\mathbf{u}} \in \mathbb{R}^3\}.$$

Also, for each $\mathbf{u} \in \mathcal{D}^{1,2}$, it holds

$$\|\nabla \mathbf{u}\|_2 = \sqrt{2} \|\mathbb{D}(\mathbf{u})\|_2, \tag{2.4}$$

and

$$\|\mathbf{u}\|_6 \leq \kappa_0 \|\mathbb{D}(\mathbf{u})\|_2, \tag{2.5}$$

for some numerical constant $\kappa_0 > 0$. Finally, there is another positive constant κ_1 such that

$$|\hat{\mathbf{u}}| \leq \kappa_1 \|\mathbb{D}(\mathbf{u})\|_2. \tag{2.6}$$

Along with the spaces \mathcal{L}^2 , \mathcal{H} , and $\mathcal{D}^{1,2}$ defined above, we introduce suitable “local” versions of these spaces. Precisely, we set

$$\begin{aligned} \mathcal{L}^2(B_R) &:= \{\varphi \in L^2(B_R) : \varphi|_{\Omega_0} = \hat{\varphi} \text{ for some } \hat{\varphi} \in \mathbb{R}^3\}, \\ \mathcal{H}(B_R) &:= \{\varphi \in \mathcal{L}^2(B_R) : \operatorname{div} \varphi = 0, \varphi \cdot \mathbf{n}|_{\partial B_R} = 0\}, \\ \mathcal{D}^{1,2}(B_R) &:= \{\varphi \in W^{1,2}(B_R) : \operatorname{div} \varphi = 0, \varphi|_{\Omega_0} = \hat{\varphi} \text{ for some } \hat{\varphi} \in \mathbb{R}^3, \varphi|_{\partial B_R} = \mathbf{0}\}, \\ \mathcal{D}_0^{1,2}(\Omega_R) &:= \{\varphi \in \mathcal{D}^{1,2}(B_R) : \hat{\varphi} = \mathbf{0}\}. \end{aligned}$$

Then $\mathcal{H}(B_R)$ and $\mathcal{D}^{1,2}(B_R)$ are Hilbert spaces with scalar products

$$\langle \varphi_1, \varphi_2 \rangle_{B_R}, \quad \varphi_i \in \mathcal{H}(B_R); \quad (\mathbb{D}(\psi_1), \mathbb{D}(\psi_2))_{B_R}, \quad \psi_i \in \mathcal{D}^{1,2}(B_R), \quad i = 1, 2.$$

Moreover, the following decomposition holds, analogous to (2.3) [29, Theorem 3.1 and Lemma 3.2]:

$$\mathcal{L}^2(B_R) = \mathcal{H}(B_R) \oplus \mathcal{G}(B_R), \tag{2.7}$$

where $\mathcal{G}(B_R)$ is defined as in (2.2)₃, by replacing Ω with Ω_R .

Finally, the dual spaces of $\mathcal{D}^{1,2}(B_R)$ and $\mathcal{D}_0^{1,2}(\Omega_R)$ will be denoted by $\mathcal{D}^{-1,2}(B_R)$ and $\mathcal{D}_0^{-1,2}(\Omega_R)$, respectively.

Remark 2.1. The space $\mathcal{D}^{1,2}(B_R)$ can be viewed as a subspace of $W^{1,2}(\Omega) \cap \mathcal{D}^{1,2}(\mathbb{R}^3)$, by extending its generic element to 0 in $\mathbb{R}^3 \setminus B_R$. Therefore, all the properties mentioned in Lemma 2.1 continue to hold for $\mathcal{D}^{1,2}(B_R)$.

We next recall basic facts about the mollification of fields in $\mathcal{D}^{1,2}(B_R)$. Let $\mathbf{u} \in \mathcal{D}^{1,2}(B_R)$ and continue to denote by \mathbf{u} its extension to $W^{1,2}(\mathbb{R}^3) \cap \mathcal{D}^{1,2}(\mathbb{R}^3)$, in the sense of Remark 2.1. Moreover, let $\eta_0 > 0$ be small enough so that the domain

$$\Omega_{00} := \{x \in \Omega_0 : \operatorname{dist}(x, \partial\Omega) > \eta_0\} \tag{2.8}$$

is not empty. For a given $\eta \in (0, \eta_0)$, we then indicate by \mathbf{u}_η the (Friederichs) mollifier of \mathbf{u} , namely,

$$\mathbf{u}_\eta(x) := \int_{\mathbb{R}^3} k_\eta(x - y) \mathbf{u}(y) \, dy, \tag{2.9}$$

where

$$k_\eta(\xi) := \eta^{-3} k(\xi/\eta); \quad k \in C_0^\infty(B_1), \quad \int_{\mathbb{R}^3} k(x) \, dx = 1.$$

Lemma 2.2. *Let $\mathbf{u} \in \mathcal{D}^{1,2}(B_R)$ and \mathbf{u}_η be defined by (2.9) with $\eta \in (0, \eta_0)$. The following properties hold:*

- (a) $\operatorname{div} \mathbf{u}_\eta(x) = 0$, for all $x \in \mathbb{R}^3$;
- (b) $\mathbf{u}_\eta(x) = \hat{\mathbf{u}}$, for all $x \in \overline{\Omega_{00}}$, and with $\hat{\mathbf{u}} = \mathbf{u}|_{\Omega_0}$;
- (c) for any $\mathbf{w} \in \mathcal{D}^{1,2}(B_R)$,

$$\int_{\Omega_R} (\mathbf{u}_\eta - \hat{\mathbf{u}}) \cdot \nabla \mathbf{w} \cdot \mathbf{w} = 0.$$

Proof. Since $\operatorname{div} \mathbf{u}(x) = 0$, $x \in \mathbb{R}^3$ and $\mathbf{u}(x) = \hat{\mathbf{u}}$, $x \in \Omega_0$, both (a) and (b) follow from the properties of mollifiers and the definition of Ω_{00} in (2.8). Integrating by parts and using $\mathbf{w}|_{\partial B_R} = \mathbf{0}$, $\mathbf{w}|_{\partial \Omega} = \hat{\mathbf{w}}$ along with (a), we get

$$\begin{aligned} \int_{\Omega_R} (\mathbf{u}_\eta - \hat{\mathbf{u}}) \cdot \nabla \mathbf{w} \cdot \mathbf{w} &= \frac{1}{2} |\hat{\mathbf{w}}|^2 \int_{\partial \Omega} (\mathbf{u}_\eta - \hat{\mathbf{u}}) \cdot \mathbf{n} = \frac{1}{2} |\hat{\mathbf{w}}|^2 \int_{\partial \Omega} \mathbf{u}_\eta \cdot \mathbf{n} \\ &= \frac{1}{2} |\hat{\mathbf{w}}|^2 \int_{\Omega_0} \operatorname{div} \mathbf{u}_\eta = 0. \end{aligned} \quad \blacksquare$$

2.2. Reformulation of the problem with velocity fields vanishing at infinity

In order to solve Problem (1.1) formulated in Section 1, it is convenient and customary to deal with velocity fields vanishing as $|x| \rightarrow \infty$. This requires an appropriate lifting of V , which will be accomplished with the help of the following result.

Lemma 2.3. *Let $\eta_0 > 0$, Ω_{00} be as in (2.8), and set*

$$\Omega^0 := \mathbb{R}^3 \setminus \overline{\Omega_{00}}.$$

Then, for any $\varepsilon > 2/\ln(1/\eta_0)$, there exists $\mathbf{U} = \mathbf{U}(\varepsilon; x, t)$, $(x, t) \in \Omega^0 \times \mathbb{R}$ such that

- (i) $\mathbf{U}(t) \in C^\infty(\Omega^0)$, $t \in \mathbb{R}$;
- (ii) $\operatorname{div} \mathbf{U}(x, t) = 0$, $(x, t) \in \Omega^0 \times \mathbb{R}$;
- (iii) $\operatorname{supp}(\mathbf{U}(t)) \subset \{x \in \Omega^0 : 0 \leq \operatorname{dist}(x, \partial \Omega_{00}) \leq e^{-1/\varepsilon}\} =: \Omega_\varepsilon$, for all $t \in \mathbb{R}$;
- (iv) $\mathbf{U}(x, t) = \mathbf{V}(t)$, $(x, t) \in \Omega_{\frac{\varepsilon}{2}} \times \mathbb{R} \supset \partial \Omega \times \mathbb{R}$;
- (v) $\|\mathbf{U}(t)\|_{2,2} \leq c_\varepsilon |\mathbf{V}(t)|$;
- (vi) $\|\partial_t \mathbf{U}(t)\|_2 \leq c_\varepsilon |\dot{\mathbf{V}}(t)|$;
- (vii) if $\hat{\Omega} \supset \Omega_\varepsilon$, and $\mathbf{w}, \mathbf{z} \in W^{1,2}(\hat{\Omega})$ with $\mathbf{w}|_{\partial \Omega_{00}} = \mathbf{0}$, then

$$\int_{\hat{\Omega}} |\mathbf{w} \cdot \nabla \mathbf{z} \cdot \mathbf{U}| \leq c_0 \varepsilon \|\nabla \mathbf{w}\|_2 \|\nabla \mathbf{z}\|_2,$$

where c_0 is a positive constant independent of ε .

Proof. Let $\psi = \psi(r)$, $r \in [0, \infty)$ be a smooth, non-decreasing real function such that $\psi(r) = 0$ if $r \leq 1$, and $\psi(r) = 1$, if $r \geq 2$, and set

$$\phi(\varepsilon; x) = \psi(-\varepsilon \ln d(x)),$$

where $d(x) := \text{dist}(x, \partial\Omega_{00})$. Clearly,

$$\phi(\varepsilon; x) = \begin{cases} 1 & \text{if } d(x) \leq e^{-2/\varepsilon}, \\ 0 & \text{if } d(x) \geq e^{-1/\varepsilon}, \end{cases}$$

and, moreover,

$$|\nabla\phi(\varepsilon; x)| \leq \frac{c\varepsilon}{d(x)}, \tag{2.10}$$

with c independent of ε . Let

$$\mathbf{U}(x, t) := x_3 V_2(t) \mathbf{e}_1 + x_1 V_3(t) \mathbf{e}_2 + x_2 V_1(t) \mathbf{e}_3$$

and define

$$\mathbf{U}(\varepsilon; x, t) := \text{curl}(\phi(\varepsilon; x)\mathbf{U}(x, t)).$$

Since $\text{curl } \mathbf{U} = \mathbf{V}$, we get

$$\mathbf{U}(\varepsilon; x, t) = \phi(\varepsilon; x)\mathbf{V}(t) - \mathbf{U}(x, t) \times \nabla\phi(\varepsilon; x), \tag{2.11}$$

which, by the properties of ϕ and the choice of ε , shows the validity of (i)–(vi) above. Moreover, by observing that $\sup_{\Omega_\varepsilon} |\mathbf{U}| \leq c|\mathbf{V}(t)|$, and using the Schwarz inequality, (2.11), and (2.10), we deduce that¹

$$\begin{aligned} \left(\int_{\hat{\Omega}} |\mathbf{w} \cdot \nabla \mathbf{z} \cdot \mathbf{U}| \right)^2 &= \left(\int_{\Omega_\varepsilon} |\mathbf{w} \cdot \nabla \mathbf{z} \cdot \mathbf{U}| \right)^2 \leq \|\nabla \mathbf{z}\|_2^2 \int_{\Omega_\varepsilon} |\mathbf{U}|^2 |\mathbf{w}|^2 \\ &\leq c \|\nabla \mathbf{z}\|_2^2 \int_{\Omega_\varepsilon} \left(\phi^2 |\mathbf{w}|^2 + \varepsilon^2 \frac{|\mathbf{w}|^2}{d^2} \right) \\ &\leq c \|\nabla \mathbf{z}\|_2^2 \int_{\Omega_\varepsilon} (d^2 + \varepsilon^2) \frac{|\mathbf{w}|^2}{d^2} \\ &\leq c \varepsilon^2 \|\nabla \mathbf{z}\|_2^2 \int_{\Omega_\varepsilon} \frac{|\mathbf{w}|^2}{d^2}, \end{aligned}$$

where c is independent of ε . Using the Hardy inequality, see e.g. [14, Lemma III.6.3], we conclude that

$$\left(\int_{\hat{\Omega}} |\mathbf{w} \cdot \nabla \mathbf{z} \cdot \mathbf{U}| \right)^2 \leq c \varepsilon^2 \|\nabla \mathbf{z}\|_2^2 \int_{\Omega_\varepsilon} \frac{|\mathbf{w}|^2}{d^2} \leq c \varepsilon^2 \|\nabla \mathbf{z}\|_2^2 \int_{\Omega_\varepsilon} |\nabla \mathbf{w}|^2,$$

and property (vii) follows. ■

¹Recall (1.2).

With Lemma 2.3 at hand, we now rewrite (1.1) with the new unknown velocity field

$$\mathbf{u} = \mathbf{v} + \mathbf{V} - \mathbf{U}$$

so that $(\mathbf{u}, \boldsymbol{\gamma}, \boldsymbol{\delta})$ now satisfies the following problem with source terms:

$$\left\{ \begin{array}{l} \left. \begin{array}{l} \partial_t \mathbf{u} + \lambda[(\mathbf{u} - \boldsymbol{\gamma} + \mathbf{U} - \mathbf{V}) \cdot \nabla \mathbf{u} + (\mathbf{u} - \boldsymbol{\gamma}) \cdot \nabla \mathbf{U}] \\ \qquad \qquad \qquad = \operatorname{div} \mathbb{T}(\mathbf{u}, p) + \mathbf{f}, \\ \qquad \qquad \qquad \operatorname{div} \mathbf{u} = 0, \end{array} \right\} \text{ in } \Omega \times \mathbb{R}, \\ \mathbf{u}(x, t) = \boldsymbol{\gamma}(t), \qquad \qquad \qquad (x, t) \in \partial\Omega \times \mathbb{R}, \\ \lim_{|x| \rightarrow \infty} \mathbf{u}(x, t) = \mathbf{0}, \qquad \qquad \qquad t \in \mathbb{R}, \\ \left. \begin{array}{l} \dot{\boldsymbol{\gamma}} + \mathbb{A} \cdot \boldsymbol{\delta} + \varpi \int_{\partial\Omega} \mathbb{T}(\mathbf{u}, p) \cdot \mathbf{n} = \mathbf{F}, \\ \qquad \qquad \qquad \dot{\boldsymbol{\delta}} = \boldsymbol{\gamma}, \end{array} \right\} \text{ in } \mathbb{R}. \end{array} \right. \quad (2.12)$$

where

$$\begin{aligned} \mathbf{f} &:= (\mathbf{V} - \mathbf{U}) \cdot \nabla \mathbf{U} + \Delta \mathbf{U} - \partial_t \mathbf{U}, \\ \mathbf{F} &:= -\varpi \operatorname{vol}(\mathcal{B}) \dot{\mathbf{V}}, \\ p &:= p - \dot{\mathbf{V}} \cdot \mathbf{x}. \end{aligned} \quad (2.13)$$

Observe that, by the properties of \mathbf{U} listed in Lemma 2.3, and (1.2), we have, in particular, that

$$\operatorname{supp}(\mathbf{f}(t)) \subset \Omega_\varepsilon, \text{ for all } t \in \mathbb{R}; \quad \|\mathbf{f}(t)\|_2 \leq c(|\mathbf{V}(t)| + |\dot{\mathbf{V}}(t)|). \quad (2.14)$$

We thus obtain that the original problem (1.1) has been formally and equivalently reformulated as (2.12)–(2.13) where the velocity of \mathcal{L} vanishes at infinity, whereas non-zero prescribed external T -periodic forces are now acting on both \mathcal{L} and \mathcal{B} .

3. Existence of T -periodic weak solutions

3.1. Definition of T -periodic weak solutions

We first need to introduce a suitable general class of T -periodic test functions. Precisely, let A be either B_R or \mathbb{R}^3 . By $\mathcal{C}_\#(A)$, we denote the space of restriction to $[0, T]$ of functions $\boldsymbol{\varphi} \in C^1(A \times \mathbb{R})$, satisfying

- (a) $\operatorname{div} \boldsymbol{\varphi}(x, t) = 0$ for $(x, t) \in A \times \mathbb{R}$;
- (b) $\boldsymbol{\varphi}(x, t) = \hat{\boldsymbol{\varphi}}(t)$, for some $\hat{\boldsymbol{\varphi}} \in C^1(\mathbb{R})$, for x in a neighborhood of Ω_0 and $t \in \mathbb{R}$;
- (c) $\operatorname{supp}_x \boldsymbol{\varphi}(x, t) \subset A$ for all $t \in \mathbb{R}$;
- (d) $\boldsymbol{\varphi}(x, t + T) = \boldsymbol{\varphi}(x, t)$ for all $(x, t) \in A \times \mathbb{R}$.

We are then able to give the definition of a T -periodic weak solution to (2.12). Testing (2.12)₁ by arbitrary $\varphi \in \mathcal{C}_\#(\mathbb{R}^3)$, integrating by parts over $\Omega \times [0, T]$, and employing (2.12)₂₋₅, we infer

$$\begin{aligned} & \langle \mathbf{u}(T), \varphi(T) \rangle - \langle \mathbf{u}(0), \varphi(0) \rangle \\ &= - \int_0^T [-\langle \mathbf{u}, \partial_t \varphi \rangle + \lambda((\mathbf{u} - \boldsymbol{\gamma} + \mathbf{U} - \mathbf{V}) \cdot \nabla \mathbf{u} + (\mathbf{u} - \boldsymbol{\gamma}) \cdot \nabla \mathbf{U}, \varphi) \\ & \quad + 2(\mathbb{D}(\mathbf{u}), \mathbb{D}(\varphi)) + \varpi^{-1} \hat{\boldsymbol{\varphi}} \cdot \mathbb{A} \cdot \boldsymbol{\delta} - (\mathbf{f}, \varphi) - \mathbf{F} \cdot \hat{\boldsymbol{\varphi}}] dt, \end{aligned}$$

where we recall that $(\cdot, \cdot) \equiv (\cdot, \cdot)_\Omega$ is the $L^2(\Omega)$ -scalar product, whereas

$$\langle \varphi, \boldsymbol{\psi} \rangle = \varpi^{-1} \hat{\boldsymbol{\varphi}} \cdot \hat{\boldsymbol{\psi}} + (\varphi, \boldsymbol{\psi})_\Omega.$$

Furthermore, from (2.12)₆, we also have

$$\boldsymbol{\delta}(T) - \boldsymbol{\delta}(0) = \int_0^T \boldsymbol{\gamma}(t) dt.$$

Thus, if $(\mathbf{u}, \boldsymbol{\gamma}, \boldsymbol{\delta})$ is a T -periodic (sufficiently smooth) solution to (3.3), then $\boldsymbol{\gamma}$ has zero average. Since the periodicity also implies

$$\langle \mathbf{u}(T), \varphi(T) \rangle = \langle \mathbf{u}(0), \varphi(0) \rangle,$$

we deduce that any such T -periodic solution satisfies

$$\begin{aligned} & \int_0^T [-\langle \mathbf{u}, \partial_t \varphi \rangle + \lambda((\mathbf{u} - \boldsymbol{\gamma} + \mathbf{U} - \mathbf{V}) \cdot \nabla \mathbf{u} + (\mathbf{u} - \boldsymbol{\gamma}) \cdot \nabla \mathbf{U}, \varphi) \\ & \quad + 2(\mathbb{D}(\mathbf{u}), \mathbb{D}(\varphi)) + \varpi^{-1} \hat{\boldsymbol{\varphi}} \cdot \mathbb{A} \cdot \boldsymbol{\delta} - (\mathbf{f}, \varphi) - \mathbf{F} \cdot \hat{\boldsymbol{\varphi}}] dt = 0, \quad (3.1) \\ & \dot{\boldsymbol{\delta}} = \boldsymbol{\gamma}, \quad \int_0^T \boldsymbol{\gamma}(t) dt = \mathbf{0}, \end{aligned}$$

for arbitrary test functions φ in $\mathcal{C}_\#(\mathbb{R}^3)$. Conversely, with the help of the decomposition (2.3), it is easy to see that every sufficiently smooth tuple $(\mathbf{u}, \boldsymbol{\gamma}, \boldsymbol{\delta})$ obeying (3.1) is a T -periodic solution of (2.12). With this in mind, we give the following definition of a weak solution.

Definition 3.1. The triple $(\mathbf{u}, \boldsymbol{\gamma}, \boldsymbol{\delta})$ is a T -periodic weak solution to (2.12) if

- (i) $\mathbf{u} \in L^2(0, T; \mathcal{D}^{1,2}(\mathbb{R}^3))$, with $\mathbf{u}(x, t)|_{\partial\Omega} = \boldsymbol{\gamma}(t)$, a.a. $t \in [0, T]$, $\boldsymbol{\gamma} \in L^2(0, T; \mathbb{R}^3)$;
- (ii) $\boldsymbol{\delta} \in W^{1,2}(0, T; \mathbb{R}^3)$;
- (iii) $(\mathbf{u}, \boldsymbol{\gamma}, \boldsymbol{\delta})$ satisfies (3.1).

Remark 3.1. In view of Lemma 2.1, it is easy to check that the integral in (3.1) is well defined for a weak solution. Likewise, the boundary condition in (i) is meaningful in the trace sense.

3.2. Statement of the main theorem and strategy of the proof

The main contribution of this paper is expressed by the following result:

Theorem 3.1. *Suppose $V \in W^{1,2}(0, T; \mathbb{R}^3)$ is T -periodic for some $T > 0$. Then there exists at least one corresponding T -periodic weak solution to (2.12). This solution satisfies the estimate*

$$\int_0^T (\|\nabla \mathbf{u}(t)\|_2^2 + |\boldsymbol{\gamma}(t)|^2) dt \leq C \int_0^T (|V|^2 + |\dot{V}|^2) dt, \tag{3.2}$$

where the constant C depends only on Ω and the physical parameters of the body and the liquid.

Our strategy to prove Theorem 3.1 goes as follows: In a first step, we shall consider a suitable modification of problem (2.12)–(2.13) in a generic bounded domain of type Ω_R . Precisely, for any $R > 3R_*$ and such that $B_R \supset \Omega_\varepsilon$,² we shall prove that the problem

$$\left\{ \begin{array}{l} \left. \begin{array}{l} \partial_t \mathbf{u} + \lambda[(\mathbf{u} - \boldsymbol{\gamma} + \mathbf{U} - V) \cdot \nabla \mathbf{u} + (\mathbf{u} - \boldsymbol{\gamma}) \cdot \nabla U] \\ = \operatorname{div} \mathbb{T}(\mathbf{u}, p) + \mathbf{f}, \\ \operatorname{div} \mathbf{u} = 0, \end{array} \right\} \text{ in } \Omega_R \times \mathbb{R}, \\ \left. \begin{array}{l} \mathbf{u}(x, t) = \boldsymbol{\gamma}(t), \\ \mathbf{u}(x, t) = \mathbf{0}, \end{array} \right\} \begin{array}{l} (x, t) \in \partial\Omega \times \mathbb{R}, \\ (x, t) \in \partial B_R \times \mathbb{R}, \end{array} \\ \left. \begin{array}{l} \dot{\boldsymbol{\gamma}} + \mathbb{A} \cdot \boldsymbol{\delta} + \varpi \int_{\partial\Omega} \mathbb{T}(\mathbf{u}, p) \cdot \mathbf{n} = \mathbf{F}, \\ \dot{\boldsymbol{\delta}} = \boldsymbol{\gamma}, \end{array} \right\} \text{ in } \mathbb{R}, \end{array} \tag{3.3}$$

with \mathbf{f} and \mathbf{F} given in (2.13) has at least one T -periodic (weak) solution $(\mathbf{u}, \boldsymbol{\gamma}, \boldsymbol{\delta})$ that, in addition, obeys certain bounds in terms of the data, uniformly with respect to R . We then let $R \rightarrow \infty$ along a sequence and prove that the corresponding solutions will converge to a weak solution to the original problem.

Even though this approach is classical, its implementation in the present setting is by no means straightforward, due to the fact that the spring has no damping. In particular, the classical method of showing the existence of a fixed point for the Poincaré map at the finite-dimensional level (along the Galerkin approximations) fails, due to the lack of “sufficient dissipation”. We are thus lead to prove this existence at the infinite-dimensional level.

We develop the steps towards the proof of Theorem 3.1 in the next subsections. We first show existence and uniqueness of strong solutions to the initial-value problem associated to a regularized version of (3.3), where the non-linear term has been suitably mollified. Then we prove that the Poincaré map associated to this problem has a fixed point, which leads to the existence of T -periodic strong solutions. Finally, we let the mollifying

²Later on, the parameter ε will be fixed in terms of the data, so that this request is meaningful; see (3.6).

parameter tend to 0, thus obtaining the same result for the original problem (3.3) in the class of weak solutions.

The last step is to let $R \rightarrow \infty$. We are not able to provide uniform estimates on the dissipation in δ , which means that, as $R \rightarrow \infty$, we have no control on the amplitude of the oscillations. To overcome this issue, we prove the velocity $\dot{\delta}$ is uniformly bounded so that the oscillation rate remains bounded as $R \rightarrow \infty$. Combining this with an estimate for $\partial_t \mathbf{u}$ and using the equation along the approximating sequence, we eventually control the average of δ and conclude existence of time-periodic weak solutions for the original problem in the whole Ω .

All the above will be accomplished through several intermediate steps. From now on it will be tacitly understood that $V(t)$ is T -periodic.

3.3. The initial-boundary value problem

In this subsection we shall study the following initial-boundary value problem associated to a mollified version of (3.3):

$$\left\{ \begin{array}{l} \left. \begin{array}{l} \partial_t \mathbf{u} + \lambda(\mathbf{u}_\eta - \boldsymbol{\gamma} + \mathbf{U} - \mathbf{V}) \cdot \nabla \mathbf{u} + (\mathbf{u} - \boldsymbol{\gamma}) \cdot \nabla \mathbf{U} \\ \qquad \qquad \qquad = \operatorname{div} \mathbb{T}(\mathbf{u}, p) + \mathbf{f}, \\ \qquad \qquad \qquad \operatorname{div} \mathbf{u} = 0, \end{array} \right\} \text{ in } \Omega_R \times (0, \infty), \\ \mathbf{u}(x, t) = \boldsymbol{\gamma}(t), \qquad \qquad \qquad (x, t) \in \partial\Omega \times (0, \infty), \\ \mathbf{u}(x, t) = \mathbf{0}, \qquad \qquad \qquad (x, t) \in \partial B_R \times (0, \infty), \\ \mathbf{u}(x, 0) = \mathbf{u}_0, \qquad \qquad \qquad x \in \Omega_R, \\ \left. \begin{array}{l} \dot{\boldsymbol{\gamma}} + \mathbb{A} \cdot \delta + \varpi \int_{\partial\Omega} \mathbb{T}(\mathbf{u}, p) \cdot \mathbf{n} = \mathbf{F}, \\ \qquad \qquad \qquad \dot{\delta} = \boldsymbol{\gamma}, \end{array} \right\} \text{ in } (0, \infty), \\ \boldsymbol{\gamma}(0) = \boldsymbol{\gamma}_0, \delta(0) = \delta_0, \end{array} \right. \quad \text{(I-BVP)}$$

where \mathbf{u}_η is the (Friederichs) mollifier of \mathbf{u} as defined in (2.9). We shall prove existence and uniqueness of solutions to (I-BVP) in a suitable functional class. To this end, we begin to transform it in an appropriate “weak” form. Testing (I-BVP)₁ by $\boldsymbol{\psi} \in \mathcal{D}^{1,2}(B_R)$, integrating by parts, and using (I-BVP)₂₋₈, we deduce

$$\begin{aligned} \langle \mathbf{u}(t) - \mathbf{u}_0, \boldsymbol{\psi} \rangle &= - \int_0^t [\lambda((\mathbf{u}_\eta - \boldsymbol{\gamma} + \mathbf{U} - \mathbf{V}) \cdot \nabla \mathbf{u} + (\mathbf{u} - \boldsymbol{\gamma}) \cdot \nabla \mathbf{U}, \boldsymbol{\psi}) \\ &\quad + 2(\mathbb{D}(\mathbf{u}), \mathbb{D}(\boldsymbol{\psi})) + \varpi^{-1} \hat{\boldsymbol{\psi}} \cdot \mathbb{A} \cdot \delta \\ &\quad - (\mathbf{f}, \boldsymbol{\psi}) - \mathbf{F} \cdot \hat{\boldsymbol{\psi}}] \, ds, \end{aligned} \tag{3.4}$$

$$\delta(t) - \delta_0 = \int_0^t \boldsymbol{\gamma}(s) \, ds,$$

where we recall that $(\cdot, \cdot) \equiv (\cdot, \cdot)_{\Omega_R}$ is the $L^2(\Omega_R)$ -scalar product, whereas

$$\langle \boldsymbol{\varphi}, \boldsymbol{\psi} \rangle = \varpi^{-1} \hat{\boldsymbol{\varphi}} \cdot \hat{\boldsymbol{\psi}} + (\boldsymbol{\varphi}, \boldsymbol{\psi})_{\Omega_R}.$$

Definition 3.2. The triple $(\mathbf{u}, \boldsymbol{\gamma}, \boldsymbol{\delta})$ is a *weak solution* to **(I-BVP)** if, for all $t > 0$,

- (i) $\mathbf{u} \in C_w([0, t]; L^2(\Omega_R)) \cap L^2(0, t; \mathcal{D}^{1,2}(\Omega_R))$, with $\mathbf{u}(x, t)|_{\partial\Omega} = \boldsymbol{\gamma}(t)$, $\boldsymbol{\gamma} \in C([0, t]; \mathbb{R}^3)$;
- (ii) $\boldsymbol{\delta} \in C^1([0, t]; \mathbb{R}^3)$;
- (iii) $(\mathbf{u}, \boldsymbol{\gamma}, \boldsymbol{\delta})$ satisfies (3.4) for all $\boldsymbol{\psi} \in \mathcal{D}^{1,2}(B_R)$.

Remark 3.2. Taking into account classical properties of mollifiers and Remark 2.1, it is easy to check that the integral in (3.4) is well defined for a weak solution.

Remark 3.3. With the help of the decomposition (2.7), one can show by classical arguments that if $(\mathbf{u}, \boldsymbol{\gamma}, \boldsymbol{\delta})$ satisfies (3.4) and is sufficiently regular, then there exists a suitable pressure field $p = p(x, t)$ such that $(\mathbf{u}, p, \boldsymbol{\gamma}, \boldsymbol{\delta})$ is a solution to **(I-BVP)**.

The next lemma deals with the well-posedness of **(I-BVP)**. For a weak solution $(\mathbf{u}, \boldsymbol{\gamma}, \boldsymbol{\delta})$ of **(I-BVP)**, we define the energy by

$$E(t) = \frac{1}{2} [\|\mathbf{u}(t)\|_{2,\Omega_R}^2 + \varpi^{-1} (|\boldsymbol{\gamma}(t)|^2 + \boldsymbol{\delta}(t) \cdot \mathbb{A} \cdot \boldsymbol{\delta}(t))]. \tag{3.5}$$

Observe that, since $\boldsymbol{\gamma} \equiv \hat{\mathbf{u}}$, the functional $E^{\frac{1}{2}}$ defines a norm on $\mathcal{H}(B_R) \times \mathbb{R}^3$.

Lemma 3.1. Let $V \in W^{1,2}(0, T; \mathbb{R}^3)$, and let

$$\eta_0 < \exp(-8c_0\lambda),^3 \tag{3.6}$$

with η_0 as in Lemma 2.2, $\eta \in (0, \eta_0)$, and c_0 as in Lemma 2.3 (vii). Then, for any given $\mathbf{u}_0 \in \mathcal{H}(B_R)$, $\boldsymbol{\gamma}_0 \in \mathbb{R}^3$, $\boldsymbol{\delta}_0 \in \mathbb{R}^3$, with the compatibility condition $\boldsymbol{\gamma}_0 := \hat{\mathbf{u}}_0$, there exists one and only one corresponding weak solution $(\mathbf{u}, \boldsymbol{\gamma}, \boldsymbol{\delta})$ to **(I-BVP)** such that

- (i) for all $\sigma > 0$ and all $t > \sigma > 0$,
 - $\mathbf{u} \in W^{1,2}(\sigma, t; L^2(\Omega_R)) \cap L^2(\sigma, \tau; W^{2,2}(\Omega_R))$,
 - $\boldsymbol{\gamma} \in W^{1,2}(\sigma, t; \mathbb{R}^3)$,
 - $\boldsymbol{\delta} \in W^{2,2}(\sigma, t; \mathbb{R}^3)$;
- (ii) for all $\sigma > 0$ and all $t > \sigma > 0$, $\mathbf{u} \in C([\sigma, t]; \mathcal{D}^{1,2}(\Omega_R))$, and there exists $p \in L^2(\sigma, \tau; W^{1,2}(\Omega_R))$ such that $(\mathbf{u}, p, \boldsymbol{\gamma}, \boldsymbol{\delta})$ satisfies **(I-BVP)**₁ a.a. in $\Omega_R \times (0, t)$ and **(I-BVP)**_{6,7} a.a. in $(0, t)$;
- (iii) the initial conditions $(\mathbf{u}_0, \boldsymbol{\gamma}_0, \boldsymbol{\delta}_0)$ are attained by $\boldsymbol{\gamma}$ and $\boldsymbol{\delta}$ in the sense of pointwise continuity, and by \mathbf{u} in the L^2 -sense, i.e.

$$\lim_{t \rightarrow 0} \|\mathbf{u}(t) - \mathbf{u}_0\|_{2,\Omega_R} = 0; \tag{3.7}$$

³See footnote 2.

(iv) for some C_1, C_2 independent of R and η , some C depending only on the data, η , and t , the energy estimates

$$\begin{aligned}
 E(t) + C_1 \int_0^t (\|\nabla \mathbf{u}(s)\|_{2, \Omega_R}^2 + |\boldsymbol{\gamma}(s)|^2) ds \\
 \leq E(0) + C_2 \int_0^t (|V|^2 + |\dot{V}|^2) ds, \quad \text{for all } t > 0, \quad (3.8) \\
 \max_{s \in [\sigma, t]} \|\nabla \mathbf{u}(s)\|_2 \leq C, \quad \text{for all } t > \sigma,
 \end{aligned}$$

hold.

Moreover, the solution depends continuously on the initial data in the norm $E^{\frac{1}{2}}$.

The proof of Lemma 3.1 relies on Galerkin approximation. We shall look for “approximated” solutions to (3.4) of the form

$$\mathbf{u}_N(x, t) = \sum_{k=1}^N c_{kN}(t) \boldsymbol{\psi}_k(x), \quad \boldsymbol{\gamma}_N(t) = \sum_{k=1}^N c_{kN}(t) \hat{\boldsymbol{\psi}}_k, \quad \delta_N(t),$$

where $\{\boldsymbol{\psi}_k\}$ is the special basis provided in the following lemma, whose proof is given in [17].

Lemma 3.2. For any fixed $R > R_*$, the problem

$$\begin{aligned}
 \left. \begin{aligned}
 -\nabla \cdot \mathbb{T}(\boldsymbol{\psi}, \phi) &= \mu \boldsymbol{\psi}, \\
 \operatorname{div} \boldsymbol{\psi} &= 0,
 \end{aligned} \right\} \quad \text{in } \Omega_R, \\
 \boldsymbol{\psi} &= \hat{\boldsymbol{\psi}} \text{ in } \Omega_0, \quad \boldsymbol{\psi} = 0 \text{ at } \partial B_R, \\
 \mu \hat{\boldsymbol{\psi}} &= \varpi \int_{\partial \Omega} \mathbb{T}(\boldsymbol{\psi}, \phi) \cdot \mathbf{n},
 \end{aligned} \quad (3.9)$$

admits a denumerable number of positive eigenvalues $\{\mu_i\}$ clustering at infinity, and corresponding eigenfunctions $\{\boldsymbol{\psi}_i\} \subset \mathcal{D}^{1,2}(B_R) \cap W^{2,2}(\Omega_R)$ forming an orthonormal basis of $\mathcal{H}(B_R)$ that is also orthogonal in $\mathcal{D}^{1,2}(\Omega_R)$. Furthermore, the correspondent “pressure” fields satisfy $\phi_i \in W^{1,2}(\Omega_R)$, $i \in \mathbb{N}$.

We also need the following approximation result that will be used in the proof of Proposition 3.1.

Lemma 3.3. Let $R > R_*$ be given, and let $\{\boldsymbol{\psi}_k\}_{k \in \mathbb{N}}$ be the basis given in Lemma 3.2. Then, for any $\boldsymbol{\varphi} \in C_{\#}^1(B_R)$ and any $\varepsilon > 0$, there is $N = N(\boldsymbol{\varphi}, \varepsilon) \in \mathbb{N}$, and corresponding T -periodic functions $r_k \in C^1(\mathbb{R})$, $k = 1, \dots, N$, such that

$$\begin{aligned}
 \max_{t \in [0, T]} \{ &\|(\boldsymbol{\varphi}_N - \boldsymbol{\varphi})(\cdot, t)\|_{1,2} + |(\boldsymbol{\varphi}_{1N} - \boldsymbol{\varphi}_1)(t)| \\
 &+ \|\partial_t(\boldsymbol{\varphi}_N - \boldsymbol{\varphi})(\cdot, t)\|_2 + |\partial_t(\boldsymbol{\varphi}_{1N} - \boldsymbol{\varphi}_1)(t)| \} < \varepsilon
 \end{aligned}$$

with $\boldsymbol{\varphi}_N(x, t) = \sum_{k=1}^N r_k(t) \boldsymbol{\psi}_k(x)$, and $\boldsymbol{\varphi}_{1N}(t) = \sum_{k=1}^N r_k(t) \hat{\boldsymbol{\psi}}_k$.

The proof is obtained arguing as in [18, Lemma 3.1] and therefore omitted.

As already said, Lemma 3.1 will be proved with the classical Galerkin procedure. To start the process, we search for an approximated solution $(\mathbf{u}_N(x, t), \boldsymbol{\gamma}_N(t), \boldsymbol{\delta}_N(t))$ to (3.4) of the form

$$\begin{aligned} \mathbf{u}_N(x, t) &= \sum_{k=1}^N c_{kN}(t) \boldsymbol{\psi}_k(x), \\ \boldsymbol{\gamma}_N(t) &= \sum_{k=1}^N c_{kN}(t) \hat{\boldsymbol{\psi}}_k, \end{aligned} \tag{3.10}$$

where $\{\boldsymbol{\psi}_k\}$ is the basis introduced in Lemma 3.2 and the vector functions $\mathbf{c}_N(t) := \{c_{1N}(t), \dots, c_{NN}(t)\}$ and $\boldsymbol{\delta}_N(t)$ satisfy the following system of equations:

$$\begin{aligned} \frac{d}{dt} \langle \mathbf{u}_N, \boldsymbol{\psi}_i \rangle + \lambda ((\mathbf{u}_N)_\eta - \boldsymbol{\gamma}_N + \mathbf{U} - \mathbf{V}) \cdot \nabla \mathbf{u}_N + (\mathbf{u}_N - \boldsymbol{\gamma}_N) \cdot \nabla \mathbf{U}, \boldsymbol{\psi}_i \\ = -2(\mathbb{D}(\mathbf{u}_N), \mathbb{D}(\boldsymbol{\psi}_i)) - \varpi^{-1} \hat{\boldsymbol{\psi}}_i \cdot \mathbb{A} \cdot \boldsymbol{\delta}_N + (\mathbf{f}, \boldsymbol{\psi}_i) + \mathbf{F} \cdot \hat{\boldsymbol{\psi}}_i, \\ \dot{\boldsymbol{\delta}}_N = \boldsymbol{\gamma}_N, \end{aligned} \tag{3.11}$$

$i = 1, \dots, N$. Here $(\cdot, \cdot) \equiv (\cdot, \cdot)_{\Omega_R}$ and $\langle \cdot, \cdot \rangle \equiv \langle \cdot, \cdot \rangle_{B_R}$. This yields a system of first-order differential equations in normal form in the unknowns $\mathbf{c}_N, \boldsymbol{\chi}_N$. Indeed, since we have the orthogonality conditions

$$\langle \boldsymbol{\psi}_i, \boldsymbol{\psi}_j \rangle = \delta_{ij}, \tag{3.12}$$

plugging the ansatz (3.10) into (3.11) entails

$$\begin{aligned} \dot{c}_{iN} &= F_i(\mathbf{c}_N, \boldsymbol{\delta}_N), \quad i = 1, \dots, N, \\ \dot{\boldsymbol{\delta}}_N &= \boldsymbol{\gamma}_N, \end{aligned} \tag{3.13}$$

where

$$\begin{aligned} F_i := & - \sum_{k=1}^N c_{kN} [\lambda ((\mathbf{U} - \mathbf{V}) \cdot \nabla \boldsymbol{\psi}_k + (\boldsymbol{\psi}_k - \hat{\boldsymbol{\psi}}_k) \cdot \nabla \mathbf{U}, \boldsymbol{\psi}_i) + 2(\mathbb{D}(\boldsymbol{\psi}_k), \mathbb{D}(\boldsymbol{\psi}_i))] \\ & - \sum_{k=1}^N \sum_{l=1}^N c_{kN} c_{lN} (((\boldsymbol{\psi}_k)_\eta - \hat{\boldsymbol{\psi}}_k) \cdot \nabla \boldsymbol{\psi}_l, \boldsymbol{\psi}_i) - \frac{1}{\varpi} \hat{\boldsymbol{\psi}}_i \cdot \mathbb{A} \cdot \boldsymbol{\delta}_N \\ & + (\mathbf{f}, \boldsymbol{\psi}_i) + \mathbf{F} \cdot \hat{\boldsymbol{\psi}}_i. \end{aligned}$$

The initial conditions $(\mathbf{u}_0, \boldsymbol{\gamma}_0, \boldsymbol{\delta}_0)$ at the level of the coefficients read

$$c_{iN}(0) = \langle \mathbf{u}_0, \boldsymbol{\psi}_i \rangle = (\mathbf{u}_0, \boldsymbol{\psi}_i) + \varpi^{-1} \boldsymbol{\gamma}_0 \cdot \hat{\boldsymbol{\psi}}_i, \quad \boldsymbol{\delta}_N(0) = \boldsymbol{\delta}_0. \tag{3.14}$$

Since $\{\boldsymbol{\psi}_i\}$ is an orthonormal basis of $\mathcal{H}(B_R)$, multiplying the first identity in (3.14) by $\boldsymbol{\psi}_i$ and summing over the index i from 1 to N delivers a bound on the initial conditions $(\mathbf{u}_N(0), \boldsymbol{\gamma}_N(0), \boldsymbol{\delta}_N(0))$,

$$\|\mathbf{u}_N(0)\|_2^2 + \varpi^{-1} |\boldsymbol{\gamma}_N(0)|^2 \leq \|\mathbf{u}_0\|_2^2 + \varpi^{-1} |\boldsymbol{\gamma}_0|^2. \tag{3.15}$$

Now that we have settled the starting point of the argument, we turn to the proof of Lemma 3.1.

Proof of Lemma 3.1. We begin to derive three basic energy estimates for the approximated solution \mathbf{u}_N which are the approximated forms of estimates that would be obtained formally for \mathbf{u} by choosing abusively \mathbf{u} , $-t \operatorname{div} \mathbb{T}(\mathbf{u}, p)$, and $t \partial_t \mathbf{u}$ as test functions in (3.4). Once we have these estimates at hand, we can actually let N go to infinity and prove the assertions (i)–(iv) for the limit \mathbf{u} of the sequence $\{\mathbf{u}_N\}$. Then it will remain to prove the continuous dependence on the initial data implying at once the uniqueness of the solution \mathbf{u} .

Step 1: Energy estimates.

First estimate: To mimic the choice of $\boldsymbol{\psi} = \mathbf{u}$ in (3.4), we multiply both sides of (3.11)₁ by c_{iN} , sum over i , and integrate by parts over Ω_R . Using (3.10) along with Lemma 2.2 (c) we show

$$\begin{aligned} & \frac{1}{2} \frac{d}{dt} [\|\mathbf{u}_N\|_2^2 + \varpi^{-1} (|\boldsymbol{\gamma}_N|^2 + \boldsymbol{\delta}_N \cdot \mathbb{A} \cdot \boldsymbol{\delta}_N)] + 2 \|\mathbb{D}(\mathbf{u}_N)\|_2^2 \\ & = \lambda((\mathbf{u}_N - \boldsymbol{\gamma}_N) \cdot \nabla \mathbf{u}_N, \mathbf{U}) + (\mathbf{f}, \mathbf{u}_N) + \mathbf{F} \cdot \boldsymbol{\gamma}_N. \end{aligned} \tag{3.16}$$

Since η_0 satisfies (3.6) and c_0 does not depend on ε , we can choose ε in the construction of \mathbf{U} such that

$$\lambda|((\mathbf{u}_N - \boldsymbol{\gamma}_N) \cdot \nabla \mathbf{u}_N, \mathbf{U})| \leq \lambda c_0 \varepsilon \|\nabla \mathbf{u}_N\|_2^2 \leq \frac{1}{4} \|\nabla \mathbf{u}_N\|_2^2, \tag{3.17}$$

whereas, from (2.4) and (2.6) it follows that

$$2 \|\mathbb{D}(\mathbf{u}_N)\|_2^2 - \frac{1}{4} \|\nabla \mathbf{u}_N\|_2^2 \geq \frac{1}{2} \|\nabla \mathbf{u}_N\|_2^2 + \kappa |\boldsymbol{\gamma}_N|^2, \tag{3.18}$$

for some universal $\kappa > 0$.

We next estimate the last terms in (3.16). Recalling that $\mathbf{f} = (\mathbf{V} - \mathbf{U}) \cdot \nabla \mathbf{U} + \Delta \mathbf{U} - \partial_t \mathbf{U}$ and $\mathbf{F} = -\varpi \operatorname{vol}(\mathcal{B}) \dot{\mathbf{V}}$, the properties of \mathbf{U} proved in Lemma 2.3, together with the Hölder inequality, (2.13), and the Cauchy–Schwarz inequality imply

$$|(\mathbf{f}, \mathbf{u}_N)| + |\mathbf{F} \cdot \boldsymbol{\gamma}_N| \leq c(|\mathbf{V}(t)|^2 + |\dot{\mathbf{V}}(t)|^2) + \frac{1}{4} \|\nabla \mathbf{u}_N\|_2^2 + \frac{1}{2} \kappa |\boldsymbol{\gamma}_N|^2.$$

Thus, employing in (3.16) the latter, together with (3.17), (3.18), we establish that

$$\begin{aligned} & \frac{1}{2} \frac{d}{dt} [\|\mathbf{u}_N\|_2^2 + \varpi^{-1} (|\boldsymbol{\gamma}_N|^2 + \boldsymbol{\delta}_N \cdot \mathbb{A} \cdot \boldsymbol{\delta}_N)] + \frac{1}{4} \|\nabla \mathbf{u}_N\|_2^2 + \frac{1}{2} \kappa |\boldsymbol{\gamma}_N|^2 \\ & \leq c_2 (|\mathbf{V}(t)|^2 + |\dot{\mathbf{V}}(t)|^2). \end{aligned} \tag{3.19}$$

Denoting by $E_N(t)$ the energy of the approximated solution, i.e.

$$E_N(t) := \frac{1}{2} [\|\mathbf{u}_N(t)\|_2^2 + \varpi^{-1} (|\boldsymbol{\gamma}_N(t)|^2 + \boldsymbol{\delta}_N(t) \cdot \mathbb{A} \cdot \boldsymbol{\delta}_N(t))],$$

and using the bound (3.15) on the initial conditions, we infer

$$\begin{aligned}
 E_N(t) &+ c_3 \int_0^t (\|\nabla \mathbf{u}_N(s)\|^2 + |\boldsymbol{\gamma}_N(s)|^2) \, ds \\
 &\leq E_N(0) + c_2 \int_0^t (|\mathbf{V}(s)|^2 + |\dot{\mathbf{V}}(s)|^2) \, ds \\
 &\leq \frac{1}{2} (\|\mathbf{u}_0\|^2 + \varpi^{-1} (|\boldsymbol{\gamma}_0|^2 + \boldsymbol{\delta}_0 \cdot \mathbb{A} \cdot \boldsymbol{\delta}_0)) + c_2 \int_0^t (|\mathbf{V}(s)|^2 + |\dot{\mathbf{V}}(s)|^2) \, ds, \tag{3.20}
 \end{aligned}$$

where c_2 and c_3 are independent of R and η .

As a consequence of this energy bound, we deduce that whatever $\tau > 0$, there exists $C(\tau) > 0$, independent of R and η , such that

$$|\mathbf{c}_N(t)| + |\boldsymbol{\delta}_N(t)| \leq C(\tau), \quad t \in [0, \tau].$$

In particular, this bound in turn implies that the initial-value problem (3.13)–(3.14) has a unique global solution (i.e. defined for all $t > 0$).

Second estimate: To mimic the formal choice $\boldsymbol{\psi} = -t \operatorname{div} \mathbb{T}(\mathbf{u}, p)$ in (3.4), we next multiply both sides of (3.11)₁ by $t\mu_i c_{iN}$, $t > 0$, and sum over i . Integrating by parts over Ω_R and employing (3.9), we show

$$\begin{aligned}
 &\frac{1}{2} \frac{d}{dt} (t \|\nabla \mathbf{u}_N\|_2^2) + t \|\operatorname{div} \mathbb{T}(\mathbf{u}_N, p_N)\|_2^2 + \varpi t |\mathbf{S}_N|^2 - \frac{1}{2} \|\nabla \mathbf{u}_N\|_2^2 \\
 &= \lambda t ((\mathbf{u}_N)_\eta - \boldsymbol{\gamma}_N + \mathbf{U} - \mathbf{V}) \cdot \nabla \mathbf{u}_N + (\mathbf{u}_N - \boldsymbol{\gamma}_N) \cdot \nabla \mathbf{U}, \operatorname{div} \mathbb{T}(\mathbf{u}_N, p_N) \\
 &\quad + t \mathbf{S}_N \cdot \mathbb{A} \cdot \boldsymbol{\delta}_N + t(\mathbf{f}, \operatorname{div} \mathbb{T}(\mathbf{u}_N, p_N)) + t\varpi(\mathbf{F}, \mathbf{S}_N), \tag{3.21}
 \end{aligned}$$

where

$$\mathbf{S}_N := \int_{\partial\Omega} \mathbb{T}(\mathbf{u}_N, p_N) \cdot \mathbf{n}, \quad p_N := \sum_{k=1}^N c_{kN} \phi_k,$$

and ϕ_k is the “pressure” field associated to $\boldsymbol{\psi}_k$. We estimate the right-hand side of (3.21) piece by piece. Since

$$\|(\mathbf{u}_N)_\eta\|_\infty \leq c_\eta \|\mathbf{u}_N\|_2,$$

for some $c_\eta > 0$, we have

$$\lambda |((\mathbf{u}_N)_\eta \cdot \nabla \mathbf{u}_N, \operatorname{div} \mathbb{T}(\mathbf{u}_N, p_N))| \leq c_\eta \|\mathbf{u}_N\|_2^2 \|\nabla \mathbf{u}_N\|_2^2 + \frac{1}{8} \|\operatorname{div} \mathbb{T}(\mathbf{u}_N, p_N)\|_2^2. \tag{3.22}$$

Using the Hölder inequality, the properties of \mathbf{U} , (2.5), and (2.6), we also get

$$\begin{aligned}
 &\lambda |((\mathbf{u}_N - \boldsymbol{\gamma}_N) \cdot \nabla \mathbf{U}, \operatorname{div} \mathbb{T}(\mathbf{u}_N, p_N))| \\
 &\quad \leq c |\mathbf{V}(t)|^2 \|\nabla \mathbf{u}_N\|_2^2 + \frac{1}{8} \|\operatorname{div} \mathbb{T}(\mathbf{u}_N, p_N)\|_2^2, \\
 &\lambda |((-\boldsymbol{\gamma}_N + \mathbf{U} - \mathbf{V}) \cdot \nabla \mathbf{u}_N, \operatorname{div} \mathbb{T}(\mathbf{u}_N, p_N))| \\
 &\quad \leq c (\|\nabla \mathbf{u}_N\|_2^2 + |\mathbf{V}(t)|^2) \|\nabla \mathbf{u}_N\|_2^2 + \frac{1}{8} \|\operatorname{div} \mathbb{T}(\mathbf{u}_N, p_N)\|_2^2. \tag{3.23}
 \end{aligned}$$

We now set

$$g(t) := 1 + \|\mathbf{u}_N(t)\|_{1,2}^2 + |\mathbf{V}(t)|^2,$$

and

$$G(t) := \|\nabla \mathbf{u}_N(t)\|_2^2 + t|\boldsymbol{\delta}_N(t)|^2 + t\|\mathbf{f}(t)\|_2^2 + t|\mathbf{F}(t)|^2,$$

and we observe that from the energy estimate (3.20), recalling also (2.14), we obtain for an arbitrary $\tau > 0$,

$$\int_0^\tau (g(t) + G(t)) dt \leq F(\tau),$$

where F depends only on the norm of the initial data, the $W^{1,2}$ -norm of $\mathbf{V}(t)$, τ , and η . Finally, we make use of the classical estimate for the Stokes problem,

$$\|\mathbf{u}_N\|_{2,2} \leq c_5(\|\operatorname{div} \mathbb{T}(\mathbf{u}_N, p_N)\|_2 + |\hat{\mathbf{u}}_N|) \leq c_6(\|\operatorname{div} \mathbb{T}(\mathbf{u}_N, p_N)\|_2 + \|\nabla \mathbf{u}_N\|_2), \tag{3.24}$$

where, in the last inequality, we have used (2.6), and the constant c_6 depends on R . Then, using the Cauchy–Schwarz inequality one more time, we conclude from (3.21)–(3.24) that

$$\frac{d}{dt}(t\|\nabla \mathbf{u}_N\|_2^2) + c_7 t\|\mathbf{u}_N\|_2^2 \leq c_8 g(t)(t\|\nabla \mathbf{u}_N\|_2^2) + c_9 G(t), \tag{3.25}$$

the constants depending on R .

As a result, using Grönwall’s lemma in (3.25) entails

$$\sup_{t \in [0, \tau]} (t\|\nabla \mathbf{u}_N(t)\|_2) + \int_0^\tau (t\|\mathbf{u}_N(t)\|_{2,2}^2) dt \leq H_1(\tau), \quad \text{for all } \tau > 0, \tag{3.26}$$

where H_1 has the same property as F and, in addition, also depends on R .

Third estimate: We finally mimic the formal choice $\boldsymbol{\psi} = t\partial_t \mathbf{u}$ in (3.4) multiplying both sides of (3.11)₁ by $t\dot{c}_{iN}$, summing over i , and integrating by parts over Ω_R as necessary. Taking into account (3.10), (2.4), and Remark 2.1 again, we show

$$\begin{aligned} & \frac{1}{2} \frac{d}{dt}(t\|\nabla \mathbf{u}_N\|_2^2) + t\|\partial_t \mathbf{u}_N\|_2^2 + \varpi^{-1} t|\dot{\boldsymbol{\gamma}}_N|^2 - \frac{1}{2}\|\nabla \mathbf{u}_N\|_2^2 \\ &= \lambda t((\mathbf{u}_N)_\eta - \boldsymbol{\gamma}_N + \mathbf{U} - \mathbf{V}) \cdot \nabla \mathbf{u}_N + (\mathbf{u}_N - \boldsymbol{\gamma}_N) \cdot \nabla \mathbf{U}, \partial_t \mathbf{u}_N \\ & \quad - \frac{t}{\varpi} \dot{\boldsymbol{\gamma}}_N \cdot \mathbb{A} \cdot \boldsymbol{\delta}_N + t(\mathbf{f}, \partial_t \mathbf{u}_N) + t\mathbf{F} \cdot \dot{\boldsymbol{\gamma}}_N. \end{aligned} \tag{3.27}$$

We may now proceed to estimate the right-hand side of (3.27) exactly as we did in (3.23), (3.22), with $\partial_t \mathbf{u}_N$ in place of $\operatorname{div} \mathbb{T}$, to show that (3.27) implies the following further bound:

$$\int_0^\tau t(\|\partial_t \mathbf{u}_N(t)\|_2^2 + |\dot{\boldsymbol{\gamma}}_N(t)|^2) dt \leq H_2(\tau), \quad \text{for all } \tau > 0, \tag{3.28}$$

where H_2 has the same property as H_1 .

Step 2: Convergence of the sequence $\{\mathbf{u}_N, \boldsymbol{\gamma}_N, \boldsymbol{\delta}_N\}$. We warn the reader that, throughout this step of the proof, convergence is understood up to the choice of an ad hoc subsequence if needed, and all extracted subsequences will still be denoted by $\{\mathbf{u}_N, \boldsymbol{\gamma}_N, \boldsymbol{\delta}_N\}$ for simplicity. Integrating both sides of (3.11)₁ over the interval $[t_1, t_2]$, $0 \leq t_1 < t_2$, and taking into account (3.12), we get

$$\begin{aligned} & \langle \mathbf{u}_N(t_2) - \mathbf{u}_N(t_1), \boldsymbol{\psi}_i \rangle \\ &= - \int_{t_1}^{t_2} \{ \lambda((\mathbf{u}_N)_\eta - \boldsymbol{\gamma}_N + \mathbf{U} - \mathbf{V}) \cdot \nabla \mathbf{u}_N + (\mathbf{u}_N - \boldsymbol{\gamma}_N) \cdot \nabla \mathbf{U}, \boldsymbol{\psi}_i \} \\ & \quad + 2(\mathbb{D}(\mathbf{u}_N), \mathbb{D}(\boldsymbol{\psi}_i)) + \varpi^{-1} \mathbb{A} \cdot \boldsymbol{\delta}_N \cdot \hat{\boldsymbol{\psi}}_i - (\mathbf{f}, \boldsymbol{\psi}_i) - \mathbf{F} \cdot \hat{\boldsymbol{\psi}}_i \} dt. \end{aligned} \tag{3.29}$$

Employing (3.29) and the uniform bound (3.20), by means of a classical argument [12, Section 3], we can show the existence of

$$\mathbf{u} \in L^\infty(0, t; \mathcal{H}(B_R)), \quad \text{for all } t > 0, \tag{3.30}$$

such that for all $\boldsymbol{\psi} \in \mathcal{D}^{1,2}(B_R)$, and setting $\boldsymbol{\gamma} = \hat{\mathbf{u}}$,

$$\lim_{N \rightarrow \infty} [(\mathbf{u}_N(t), \boldsymbol{\psi}) + \varpi^{-1} \boldsymbol{\gamma}_N(t) \cdot \hat{\boldsymbol{\psi}}] = (\mathbf{u}(t), \boldsymbol{\psi}) + \varpi^{-1} \boldsymbol{\gamma}(t) \cdot \hat{\boldsymbol{\psi}}, \tag{3.31}$$

uniformly in $t \in [0, \tau]$, for all $\tau > 0$. This implies, in particular,

$$\lim_{N \rightarrow \infty} \boldsymbol{\gamma}_N(t) = \boldsymbol{\gamma}(t), \quad \text{uniformly in } t \in [0, \tau], \tag{3.32}$$

which furnishes

$$\boldsymbol{\gamma} \in C([0, \tau]; \mathbb{R}^3), \tag{3.33}$$

and also, by (3.13)₂,

$$\boldsymbol{\delta}_N \rightarrow \boldsymbol{\delta} \text{ in } C^1([0, \tau]; \mathbb{R}^3), \quad \dot{\boldsymbol{\delta}}(t) = \boldsymbol{\gamma}(t), \quad t \in [0, \tau]. \tag{3.34}$$

From the uniformity in N of the estimates (3.20), (3.26), and (3.28), it is routine to deduce that

$$\mathbf{u}_N \rightarrow \mathbf{u} \begin{cases} \text{weakly in } L^2(0, t; W^{1,2}(\Omega_R)), \\ \text{weak}^* \text{ in } L^\infty(0, t; L^2(\Omega_R)), \end{cases} \tag{3.35}$$

for all $t > 0$ and so, combining (3.31) with (3.35)₁, we also deduce (we refer to [12, Section 5] for details)

$$\mathbf{u}_N \rightarrow \mathbf{u}, \text{ strongly in } L^2(0, t; L^2(\Omega_R)), \text{ for all } t > 0. \tag{3.36}$$

In particular, from (3.35) (3.20), (3.32), and (3.34) we infer the validity of (3.8)₁. Likewise, from (3.26) and (3.28) it follows that

$$\mathbf{u}_N \rightarrow \mathbf{u} \text{ weakly in } W^{1,2}(\sigma, \tau; \mathcal{L}^2(B_R)) \text{ and } L^2(\sigma, \tau; W^{2,2}(\Omega_R)) \text{ for all } \tau > \sigma > 0.$$

The latter implies, on the one hand, by a well-known interpolation theorem, $\mathbf{u} \in C([\sigma, \tau]; W^{1,2}(\Omega_R))$, and, on the other hand, in combination with (3.26), the estimate (3.8)₂.

We next pass to the limit $N \rightarrow \infty$ in (3.29). Using (3.31), (3.35), and (3.36) together with the properties of the base $\{\boldsymbol{\psi}_i\}$ (see Lemma 3.2) and classical arguments from [12, Section 5], we can show that (3.29) continues to hold with $(\mathbf{u}_N, \boldsymbol{\gamma}_N, \boldsymbol{\delta}_N)$ replaced by $(\mathbf{u}, \boldsymbol{\gamma}, \boldsymbol{\delta})$ and $\boldsymbol{\psi}_i$ replaced by arbitrary $\boldsymbol{\psi} \in \mathcal{D}^{1,2}(B_R)$.

All in all, $(\mathbf{u}, \boldsymbol{\gamma}, \boldsymbol{\delta})$ satisfies (3.4), for all $t > 0$ and in view of (3.33) and (3.30), this also implies

$$\mathbf{u} \in C_w(0, t; L^2(\Omega_R)). \tag{3.37}$$

In order to complete the existence part of the proof, it remains to consider the sense in which the initial conditions $(\mathbf{u}_0, \boldsymbol{\gamma}_0, \boldsymbol{\delta}_0)$ are satisfied and in fact it remains only to show (3.7). The energy estimate (3.8) and the continuity of $(\boldsymbol{\gamma}, \boldsymbol{\delta})$ at $t = 0$ imply

$$\limsup_{t \rightarrow 0^+} \|\mathbf{u}(t)\|_2 \leq \|\mathbf{u}_0\|_2,$$

which entails (3.7) as an immediate consequence of the weak continuity (3.37).

Step 3: Continuous dependence on the initial data. Let $(\mathbf{u}_i, \boldsymbol{\gamma}_i, \boldsymbol{\delta}_i)$, $i = 1, 2$, be two weak solutions corresponding to the same V . Setting $\mathbf{u} = \mathbf{u}_1 - \mathbf{u}_2$, $\boldsymbol{\gamma} = \boldsymbol{\gamma}_1 - \boldsymbol{\gamma}_2$, $\boldsymbol{\delta} = \boldsymbol{\delta}_1 - \boldsymbol{\delta}_2$ we have for arbitrary $\tau > \sigma > 0$,

$$\left\{ \begin{array}{l} \partial_t \mathbf{u} + \lambda((\mathbf{u} + \mathbf{u}_2)_\eta - \boldsymbol{\gamma} + \mathbf{U} - V) \cdot \nabla \mathbf{u} \\ + (\mathbf{u} - \boldsymbol{\gamma}) \cdot \nabla \mathbf{U} + \mathbf{u}_\eta \cdot \nabla \mathbf{u}_2 = \operatorname{div} \mathbb{T}(\mathbf{u}, p), \\ \operatorname{div} \mathbf{u} = 0, \end{array} \right\} \quad \text{in } \Omega_R \times [\sigma, \tau],$$

$$\left\{ \begin{array}{l} \mathbf{u}(x, t) = \boldsymbol{\gamma}(t), \\ \mathbf{u}(x, t) = \mathbf{0}, \\ \mathbf{u}(x, 0) = \mathbf{u}_1(x, 0) - \mathbf{u}_2(x, 0), \end{array} \right\} \quad \begin{array}{l} (x, t) \in \partial\Omega \times [\sigma, \tau], \\ \text{a.a. } (x, t) \in \partial B_R \times [\sigma, \tau], \\ x \in \Omega_R, \end{array} \tag{3.38}$$

$$\left\{ \begin{array}{l} \dot{\boldsymbol{\gamma}} + \mathbb{A} \cdot \boldsymbol{\delta} + \varpi \int_{\partial\Omega} \mathbb{T}(\mathbf{v}, p) \cdot \mathbf{n} = \mathbf{0}, \\ \dot{\boldsymbol{\delta}} = \boldsymbol{\gamma}, \end{array} \right\} \quad \text{in } (0, \infty),$$

$$\boldsymbol{\gamma}(0) = \boldsymbol{\gamma}_1(0) - \boldsymbol{\gamma}_2(0), \quad \boldsymbol{\delta}(0) = \boldsymbol{\delta}_1(0) - \boldsymbol{\delta}_2(0),$$

for some $p \in L^2(\sigma, \tau; W^{1,2}(\Omega_R))$. Testing (3.38)₁ by \mathbf{u} , integrating by parts over Ω_R , and taking into account (3.38)₂₋₅ we get

$$\frac{1}{2} \frac{dE}{dt} + 2\|\mathbb{D}(\mathbf{u})\|_2^2 = \lambda[(\mathbf{u} - \boldsymbol{\gamma}) \cdot \nabla \mathbf{u}, \mathbf{U}] - (\mathbf{u}_\eta \cdot \nabla \mathbf{u}_2, \mathbf{u}), \quad t \in [\sigma, \tau], \tag{3.39}$$

where E is given in (3.5). Arguing as in (3.17), (3.18) from (3.39) we infer, in particular,

$$\frac{dE}{dt} \leq -\lambda(\mathbf{u}_\eta \cdot \nabla \mathbf{u}_2, \mathbf{u}), \quad t \in [\sigma, \tau].$$

By the property of the mollifier, we have $\|\mathbf{u}_\eta\|_\infty \leq c_\eta \|\mathbf{u}\|_2$, so that

$$|(\mathbf{u}_\eta \cdot \nabla \mathbf{u}_2, \mathbf{u})| \leq c_\eta \|\nabla \mathbf{u}_2\|_2 \|\mathbf{u}\|_2^2 \leq 2c_\eta \|\nabla \mathbf{u}_2\|_2 E.$$

Combining the last two displayed equations and using Grönwall’s lemma, we get

$$E(\tau) \leq E(\sigma) \exp\left(2c_\eta \int_0^\tau \|\nabla \mathbf{u}_2(s)\|_2 \, ds\right).$$

If we let $\sigma \rightarrow 0$ and use properties (i) and (ii) of a weak solution along with (3.7), we get

$$\lim_{\sigma \rightarrow 0} E(\sigma) = E(0),$$

yielding the claimed continuous dependence property and, in particular, uniqueness for the Cauchy problem. ■

We end this subsection on the initial-boundary value problem with the following important estimate that will be used later on in Proposition 3.1:

Lemma 3.4. *Let $V \in W^{1,2}(0, T; \mathbb{R}^3)$. Then, every weak solution to (I-BVP) satisfies for all $t > 0$ and all $R > R_*$,*

$$\frac{d\mathbf{u}}{dt} \in L^1(0, t; \mathcal{D}^{-1,2}(B_R)).$$

Moreover, there exists a constant $C = C(R)$ independent of η such that

$$\begin{aligned} & \left\| \frac{d\mathbf{u}}{dt} \right\|_{L^1(0, T; \mathcal{D}^{-1,2}(B_R))} \\ & \leq C \int_0^t (|V(s)| + |\dot{V}(s)| + |\delta(s)|^2 + \|\nabla \mathbf{u}(s)\|_2 + \|\nabla \mathbf{u}(s)\|_2^2) \, ds. \end{aligned} \tag{3.40}$$

Proof. From (3.4)₁ we deduce, for a.a. $t > 0$ and all $\psi \in \mathcal{D}^{1,2}(B_R)$,

$$\frac{d}{dt} \langle \mathbf{u}(t), \psi \rangle = G_\psi(t), \tag{3.41}$$

where

$$\begin{aligned} G_\psi(t) = & -\lambda((\mathbf{u}_\eta - \boldsymbol{\gamma} + \mathbf{U} - V) \cdot \nabla \mathbf{u} + (\mathbf{u} - \boldsymbol{\gamma}) \cdot \nabla \mathbf{U}, \psi) \\ & - 2(\mathbb{D}(\mathbf{u}), \mathbb{D}(\psi)) - \varpi^{-1} \hat{\boldsymbol{\psi}} \cdot \mathbb{A} \cdot \boldsymbol{\delta} + (\mathbf{f}, \psi) + \mathbf{F} \cdot \hat{\boldsymbol{\psi}}. \end{aligned}$$

Employing the Hölder inequality several times, along with (2.4)–(2.6) and classical properties of the mollifier, we show

$$|G_\psi(t)| \leq C(R)(|V(t)| + |\dot{V}(t)| + |\delta(t)|^2 + \|\nabla \mathbf{u}(t)\|_2 + \|\nabla \mathbf{u}(t)\|_2^2) \|\mathbb{D}(\psi)\|_2. \tag{3.42}$$

Thus, from (3.41)–(3.42) and the assertions (i)–(ii) of Definition 3.2, we deduce that there exists $g \in L^1(0, T; \mathcal{D}^{-1,2}(B_R))$ such that

$$\frac{d}{dt} \mathbf{u}(t) = g(t)$$

in the sense of distributions, as well as the validity of (3.40). ■

3.4. A key lemma on the total dissipation

As explained in the introduction, we shall obtain the existence of a periodic solution as a fixed point of the Poincaré map M in the energy space. To prove that M maps some ball in the energy space into itself, we next adapt an argument originally introduced in the context of non-linear wave equations [21, p. 162], [22]. The basic idea, reformulated in our case, consists in perturbing the energy by adding the term $\zeta \varpi^{-1} \boldsymbol{\gamma} \cdot \boldsymbol{\delta}$ for the solid, to recover some dissipation in $\boldsymbol{\delta}$. To be compatible with the fluid, we need to add another ad hoc contribution in the energy functional. To this aim, we first construct (classically) a solenoidal extension of $\boldsymbol{\delta} = \boldsymbol{\delta}(t)$. Set

$$\begin{aligned} \mathbf{H}(x, t) &:= x_3 \delta_2(t) \mathbf{e}_1 + x_1 \delta_3(t) \mathbf{e}_2 + x_2 \delta_1(t) \mathbf{e}_3, \\ \mathbf{H}(x, t) &:= \operatorname{curl}(\psi(|x|)\mathbf{H}(x, t)) \equiv \psi(|x|)\boldsymbol{\delta}(t) - \mathbf{H}(x, t) \times \nabla \psi(|x|), \end{aligned}$$

where ψ is a smooth function that takes value 1 in a neighborhood of Ω_0 and vanishes for $|x| \geq 2R_*$. Clearly,

$$\begin{aligned} \mathbf{H}(x, t) &= \boldsymbol{\delta}(t), \quad (x, t) \in \partial\Omega \times (0, \infty), \\ \operatorname{div} \mathbf{H}(x, t) &= 0, \quad (x, t) \in \Omega_R \times (0, \infty), \\ \sup_{x \in \Omega_R} |\mathbf{H}(x, t)| &\leq c|\boldsymbol{\delta}(t)|, \\ \sup_{x \in \Omega_R} |\partial_t \mathbf{H}(x, t)| &\leq c|\dot{\boldsymbol{\delta}}(t)| = c|\boldsymbol{\gamma}(t)|. \end{aligned} \tag{3.43}$$

We then introduce the map from $C^1([0, t]; \mathbb{R}^3)$ to $L^2(0, t; \mathcal{D}^{1,2}(\Omega_R))$ that associates the solenoidal extension $\mathbf{H}(t)$ we just built to $\boldsymbol{\delta}(t)$. We could write this field as $\mathbf{H}_\boldsymbol{\delta}(t)$, but we just keep the notation $\mathbf{H}(t)$ since no confusion arises. Let $(\mathbf{u}, \boldsymbol{\gamma}, \boldsymbol{\delta})$ be the weak solution of (I-BVP) determined in Lemma 3.1. For $\zeta \in (0, \infty)$, we define the following modified energy functional:

$$E_\zeta(\mathbf{u}, \boldsymbol{\gamma}, \boldsymbol{\delta}) := E(\mathbf{u}, \boldsymbol{\gamma}, \boldsymbol{\delta}) + 2\zeta[(\mathbf{u}, \mathbf{H}) + \varpi^{-1} \boldsymbol{\gamma} \cdot \boldsymbol{\delta}],$$

where E is given in (3.5). With the help of (3.43) and the Cauchy–Schwarz inequality, it is easy to check that there is $\zeta_1 = \zeta_1(\varpi, R_*, \mathbb{A})$ such that, if $\zeta \leq \zeta_1$, then

$$\frac{1}{2}E \leq E_\zeta \leq 2E. \tag{3.44}$$

As a consequence, $E^{\frac{1}{2}}$ and $E_\zeta^{\frac{1}{2}}$ are equivalent norms in $\mathcal{L}^2(B_R)$ and the set $\{(\mathbf{u}, \boldsymbol{\delta}) \in \mathcal{H}(B_R) \times \mathbb{R}^3 : E_\zeta(\mathbf{u}, \hat{\mathbf{u}}, \boldsymbol{\delta}) \leq \rho^2\}$ is convex for ζ small. This fact is needed in the next subsection, in order to obtain a fixed point of the Poincaré map via the Schauder theorem.

We next show that there exist $\zeta_0, \rho_0 > 0$, depending on the parameters of the problem such that, for the norm $E_{\zeta_0}^{\frac{1}{2}}$, the Poincaré map maps the ball of radius ρ_0 in $\mathcal{L}^2(B_R) \times \mathbb{R}^3 \times \mathbb{R}^3$ into itself. To this end, we set

$$\mathcal{V} := \int_0^T (|V|^2 + |\dot{V}|^2) dt.$$

Lemma 3.5. *Let V satisfy the assumption of Lemma 3.1 and let $(\mathbf{u}, \boldsymbol{\gamma}, \boldsymbol{\delta})$ be the corresponding unique solution to (I-BVP) given in Lemma 3.1. There exist $\zeta_0, \rho_0 > 0$, depending on $R, \mathcal{V}, \lambda, \mathbb{A},$ and ϖ , such that if $E_{\zeta_0}(0) \leq \rho_0$, then $E_{\zeta_0}(T) \leq \rho_0$ as well. In addition, the following estimate holds:*

$$E_{\zeta_0}(T) + \frac{1}{2}\zeta_0\varpi^{-1} \int_0^T \boldsymbol{\delta}(t) \cdot \mathbb{A} \cdot \boldsymbol{\delta}(t) dt \leq E_{\zeta_0}(0) + C_0\mathcal{V}, \tag{3.45}$$

where C_0 is independent of η .

Proof. For sufficiently small $\sigma > 0$, consider (I-BVP) with $t \in [\sigma, T]$. Testing (I-BVP)₁ by \mathbf{u} and using exactly the same argument leading to (3.19) we obtain

$$\frac{dE}{dt} + \frac{1}{2}\|\nabla\mathbf{u}\|_2^2 + \kappa|\boldsymbol{\gamma}|^2 \leq C_1(|V|^2 + |\dot{V}|^2), \tag{3.46}$$

where, here and in the rest of the proof, by $C_i, i = 1, \dots,$ and, later on, ζ_2 , we denote generic positive constants depending, at most, on $\lambda, \mathbb{A}, \varpi,$ and the (spatial) support of \mathbf{H} , but not on η . Furthermore, again from (I-BVP)₁, tested this time on \mathbf{H} , we get

$$\begin{aligned} &\frac{d}{dt}[(\mathbf{u}, \mathbf{H}) + \varpi^{-1}\boldsymbol{\gamma} \cdot \boldsymbol{\delta}] \\ &= (\mathbf{u}, \partial_t \mathbf{H}) + \varpi^{-1}|\boldsymbol{\gamma}|^2 - \lambda((\mathbf{u}_\eta - \boldsymbol{\gamma} + \mathbf{U} - \mathbf{V}) \cdot \nabla\mathbf{u} + (\mathbf{u} - \boldsymbol{\gamma}) \cdot \nabla\mathbf{U}, \mathbf{H}) \\ &\quad - 2(\mathbb{D}(\mathbf{u}), \mathbb{D}(\mathbf{H})) - \varpi^{-1}\boldsymbol{\delta} \cdot \mathbb{A} \cdot \boldsymbol{\delta} + (\mathbf{f}, \mathbf{H}) + \varpi^{-1}\mathbf{F} \cdot \boldsymbol{\delta}. \end{aligned}$$

If we employ (2.4)–(2.6), the property of the mollifier, assertion (v) of Lemma 2.3, (3.43), and recalling that $\sup_t |V(t)| = 1$, we can estimate the right-hand side in the following way:

$$\begin{aligned} |(\mathbf{u}, \partial_t \mathbf{H})| &\leq C_2\|\nabla\mathbf{u}\|_2|\boldsymbol{\gamma}|, \\ 2|(\mathbb{D}(\mathbf{u}), \mathbb{D}(\mathbf{H}))| + \lambda|((\mathbf{u}_\eta - \boldsymbol{\gamma} + \mathbf{U} - \mathbf{V}) \cdot \nabla\mathbf{u} + (\mathbf{u} - \boldsymbol{\gamma}) \cdot \nabla\mathbf{U}, \mathbf{H})| \\ &\leq C_2|\boldsymbol{\delta}|\|\nabla\mathbf{u}\|_2(1 + \|\nabla\mathbf{u}\|_2), \\ |(\mathbf{f}, \mathbf{H}) + \varpi^{-1}\mathbf{F} \cdot \boldsymbol{\delta}| &\leq C_3|\boldsymbol{\delta}|(\|\mathbf{f}\|_2 + |\mathbf{F}|) \leq C_4|\boldsymbol{\delta}|(|V| + |\dot{V}|). \end{aligned} \tag{3.47}$$

We also make use of the estimate

$$|(\mathbf{H}, \partial_t \mathbf{H})| \leq C_2|\boldsymbol{\delta}||\boldsymbol{\gamma}|.$$

From the latter, (3.46)–(3.47), and the Cauchy–Schwarz inequality we deduce that there is $\zeta_2 \in (0, \zeta_1)$ such that if

$$\zeta \leq \zeta_2,$$

then

$$\frac{dE_\zeta}{dt} + \frac{1}{8}\|\nabla\mathbf{u}\|_2^2 + C_6|\boldsymbol{\gamma}|^2 + \frac{1}{2}\zeta\varpi^{-1}\boldsymbol{\delta} \cdot \mathbb{A} \cdot \boldsymbol{\delta} \leq C_7(\zeta|\boldsymbol{\delta}|\|\nabla\mathbf{u}\|_2^2 + |V|^2 + |\dot{V}|^2), \tag{3.48}$$

and, in addition, (3.44) holds.

We claim there exist $\zeta_0 \in (0, \zeta_2)$ and $\rho_0 > 0$ such that $E_\zeta(0) \leq \rho_0$ implies $E_\zeta(T) \leq \rho_0$ for $\zeta \leq \zeta_0$. Imposing $E_\zeta(0) \leq \rho_0$ with

$$\rho_0 \geq \mathcal{V},$$

we deduce from (3.8) and (3.44) that

$$E(t) \leq E(0) + C_2 \mathcal{V} \leq 2E_\zeta(0) + C_2 \mathcal{V} \leq (2 + C_2)\rho_0,$$

yielding the estimate

$$|\delta(t)| \leq C_8 \rho_0^{\frac{1}{2}}.$$

Plugging this bound into (3.48), we infer

$$\frac{dE_\zeta}{dt} + \frac{1}{8} \|\nabla \mathbf{u}\|_2^2 + C_6 |\boldsymbol{\gamma}|^2 + \frac{1}{2} \zeta \varpi^{-1} \boldsymbol{\delta} \cdot \mathbb{A} \cdot \boldsymbol{\delta} \leq \frac{1}{16C_9} \zeta \rho_0^{\frac{1}{2}} \|\nabla \mathbf{u}\|_2^2 + C_7 (|\mathbf{V}|^2 + |\dot{\mathbf{V}}|^2),$$

and therefore

$$\frac{dE_{\zeta_0}}{dt} + \frac{1}{16} \|\nabla \mathbf{u}\|_2^2 + C_6 |\boldsymbol{\gamma}|^2 + \frac{1}{2} \zeta_0 \varpi^{-1} \boldsymbol{\delta} \cdot \mathbb{A} \cdot \boldsymbol{\delta} \leq C_7 (|\mathbf{V}|^2 + |\dot{\mathbf{V}}|^2), \tag{3.49}$$

with

$$\zeta_0 := C_9 \rho_0^{-\frac{1}{2}}. \tag{3.50}$$

It is classical to show, from (2.5) and the Hölder inequality, that $\|\nabla \mathbf{u}\|_2 \geq C_{10} R^{-1} \|\mathbf{u}\|_2$. Therefore, we infer, with the help of (3.44),

$$\begin{aligned} & \frac{1}{16} \|\nabla \mathbf{u}\|_2^2 + C_6 |\boldsymbol{\gamma}|^2 + \frac{1}{2} \zeta_0 \varpi^{-1} \boldsymbol{\delta} \cdot \mathbb{A} \cdot \boldsymbol{\delta} \\ & \geq \frac{\zeta_0}{2} \left(C_{11}^{-1} R^{-2} \rho_0^{\frac{1}{2}} \|\mathbf{u}\|_2^2 + \frac{1}{\varpi} (C_{11}^{-1} \rho_0^{\frac{1}{2}} |\boldsymbol{\gamma}|^2 + \boldsymbol{\delta} \cdot \mathbb{A} \cdot \boldsymbol{\delta}) \right) \\ & \geq \zeta_0 E \geq \frac{1}{2} \zeta_0 E_{\zeta_0}, \end{aligned}$$

if $\rho_0 \geq C_{11}^2 (R^4 + 1)$. Replacing the latter in (3.49) entails

$$\frac{dE_{\zeta_0}}{dt} + \frac{1}{2} \zeta_0 E_{\zeta_0} \leq C_7 (|\mathbf{V}|^2 + |\dot{\mathbf{V}}|^2), \quad \text{for all } t \in [\sigma, T],$$

which, once integrated between σ and T , furnishes

$$E_{\zeta_0}(T) \leq E_{\zeta_0}(\sigma) \exp\left(-\frac{1}{2} \zeta_0 T\right) + C_7 \mathcal{V}.$$

However, from the assertion (iii) of Lemma 3.1, we easily show that

$$\lim_{\sigma \rightarrow 0} E_{\zeta_0}(\sigma) = E_{\zeta_0}(0) \tag{3.51}$$

and the claimed result that $E_{\zeta_0}(T) \leq \rho_0$ follows by imposing that ρ_0 satisfies the condition

$$\rho_0 \geq \frac{C_7 \mathcal{V}}{1 - \exp(-\frac{1}{2}\zeta_0 T)}. \tag{3.52}$$

It is readily checked that (3.52) certainly holds, provided we choose ρ_0 greater than some quantity depending only on C_9, C_7, \mathcal{V} , and T . In fact, taking into account the choice of ζ_0 made in (3.50) and setting $x = C_9 T / 2\rho_0^{\frac{1}{2}}$, (3.52) is equivalent to

$$\frac{1}{C_7 \mathcal{V}} \left(\frac{2}{C_9 T} \right)^2 \geq x \frac{x}{1 - e^{-x}},$$

which is true, provided x is less than a suitable quantity with the properties stated above.

Finally, the estimate (3.45) follows by integrating (3.49) over $[\sigma, T]$, letting $\sigma \rightarrow 0$, and using (3.51). ■

3.5. Approximated solutions in bounded domains

With the help of what we have shown so far, we are now in a position to prove the existence of a T -periodic weak solution to (3.3). To this end, we begin to give the definition of a T -periodic weak solution to problem (3.3). This is done exactly as we did in the case of Definition 3.1, by replacing \mathbb{R}^3 and Ω by B_R and Ω_R , respectively.

Definition 3.3. The triple $(\mathbf{u}, \boldsymbol{\gamma}, \boldsymbol{\delta})$ is a T -periodic weak solution to (3.3) if

- (i) $\mathbf{u} \in L^2(0, T; \mathcal{D}^{1,2}(B_R))$, with $\mathbf{u}(x, t)|_{\partial\Omega} = \boldsymbol{\gamma}(t)$, a.a. $t \in [0, T]$, $\boldsymbol{\gamma} \in L^2(0, T; \mathbb{R}^3)$;
- (ii) $\boldsymbol{\delta} \in W^{1,2}(0, T; \mathbb{R}^3)$;
- (iii) $(\mathbf{u}, \boldsymbol{\gamma}, \boldsymbol{\delta})$ satisfies the following equations (with $(\cdot, \cdot) \equiv (\cdot, \cdot)_{\Omega_R}$ and $\langle \cdot, \cdot \rangle \equiv \langle \cdot, \cdot \rangle_{B_R}$):

$$\int_0^T \left[- \langle \mathbf{u}, \partial_t \boldsymbol{\varphi} \rangle + \lambda((\mathbf{u} - \boldsymbol{\gamma} + \mathbf{U} - \mathbf{V}) \cdot \nabla \mathbf{u} + (\mathbf{u} - \boldsymbol{\gamma}) \cdot \nabla \mathbf{U}, \boldsymbol{\varphi}) + 2(\mathbb{D}(\mathbf{u}), \mathbb{D}(\boldsymbol{\varphi})) + \varpi^{-1} \hat{\boldsymbol{\varphi}} \cdot \mathbb{A} \cdot \boldsymbol{\delta} - (\mathbf{f}, \boldsymbol{\varphi}) - \mathbf{F} \cdot \hat{\boldsymbol{\varphi}} \right] dt = 0, \tag{3.53}$$

$$\dot{\boldsymbol{\delta}} = \boldsymbol{\gamma}, \quad \int_0^T \boldsymbol{\gamma}(t) dt = \mathbf{0},$$

whatever $\boldsymbol{\varphi} \in \mathcal{C}_\#(B_R)$.

Remark 3.4. It is easy to show that if $(\mathbf{u}, \boldsymbol{\gamma}, \boldsymbol{\delta})$ is a sufficiently regular T -periodic solution to (3.3), then it satisfies (3.53) and that, with the help of the decomposition (2.7), the converse is also true.

Proposition 3.1. Let $\mathbf{V} \in W^{1,2}(0, T; \mathbb{R}^3)$. Then, for any $R > 3R_*$ there is at least one T -periodic weak solution to (3.3) in B_R . This solution satisfies

$$\int_0^T (\|\nabla \mathbf{u}(t)\|_{2, \Omega_R}^2 + |\boldsymbol{\gamma}(t)|^2) dt \leq C \mathcal{V}, \tag{3.54}$$

where the constant C is independent of R . Moreover, given $R_0 > 3R_*$, there exists a constant C_1 depending on R_0 but independent of R such that for all $R \geq R_0$,

$$\int_0^T \|\mathbf{u}(t)\|_{2,\Omega_{R_0}}^2 \leq C_1 \mathcal{V}, \quad \left\| \frac{d\mathbf{u}}{dt} \right\|_{L^1(0,T;\mathcal{D}_0^{-1,2}(\Omega_{R_0}))} \leq C_1(\mathcal{V}^{\frac{1}{2}} + \mathcal{V}). \tag{3.55}$$

Proof. Set

$$\mathcal{S}_{\rho_0} := \{(\mathbf{u}, \boldsymbol{\delta}) \in \mathcal{H}(B_R) \times \mathbb{R}^3 : E_{\zeta_0}(\mathbf{u}, \hat{\mathbf{u}}, \boldsymbol{\delta}) \leq \rho_0^2\},$$

with ζ_0 and ρ_0 as in Lemma 3.5. Clearly, \mathcal{S}_{ρ_0} is a closed convex subset of $\mathcal{H}(B_R) \times \mathbb{R}^3$. Next, let $\mathbf{s}(t) := (\mathbf{u}(t), \boldsymbol{\gamma}(t), \boldsymbol{\delta}(t))$ be the solution to (I-BVP) determined in Lemma 3.1, and consider the map

$$M: \mathbf{s}(0) \mapsto \mathbf{s}(T).$$

By Lemma 3.5, M maps \mathcal{S}_{ρ_0} into itself. By Lemma 3.1, M is also continuous, with $\mathbf{u}(T) \in W^{1,2}(\Omega_R)$, thus furnishing that $M(\mathcal{S}_{\rho_0})$ is compact. As a result, by the Schauder fixed-point theorem we conclude that the ‘‘mollified’’ problem (I-BVP) has at least one T -periodic solution.

Our next goal is to prove that, if we let $\eta \rightarrow 0$ along a sequence $\{\eta_n\}$, the sequence of corresponding T -periodic (strong) solutions $(\mathbf{u}_n, \boldsymbol{\gamma}_n, \boldsymbol{\delta}_n)$ converge to a T -periodic weak solution to (3.3). To this end, we remark that from (3.8) and (3.45) it follows that

$$\begin{aligned} \int_0^T (\|\nabla \mathbf{u}_n(t)\|_2^2 + |\boldsymbol{\gamma}_n(t)|^2) dt &\leq C_1 \mathcal{V}, \\ \int_0^T \boldsymbol{\delta}_n(t) \cdot \mathbb{A} \cdot \boldsymbol{\delta}_n(t) dt &\leq C_2 \mathcal{V}, \end{aligned} \tag{3.56}$$

where C_1 is independent of R and η , and C_2 depends on R but is independent of η . Multiplying both sides of (I-BVP)₁ – written for these solutions – by the test function $r(t)\boldsymbol{\psi}$, we show that for all $\boldsymbol{\psi} \in \mathcal{D}^{1,2}(B_R)$ and all smooth r ,

$$\begin{aligned} &\langle \mathbf{u}_n(t), r(t)\boldsymbol{\psi} \rangle - \langle \mathbf{u}_n(0), r(0)\boldsymbol{\psi} \rangle \\ &= - \int_0^t \left[- \langle \mathbf{u}_n, r'\boldsymbol{\psi} \rangle \right. \\ &\quad + \lambda(((\mathbf{u}_n)_{\eta_n} - \boldsymbol{\gamma}_n + U - V) \cdot \nabla \mathbf{u}_n + (\mathbf{u}_n - \boldsymbol{\gamma}_n) \cdot \nabla U, r\boldsymbol{\psi}) \\ &\quad \left. + 2(\mathbb{D}(\mathbf{u}_n), \mathbb{D}(r\boldsymbol{\psi})) + \varpi^{-1}r\hat{\boldsymbol{\psi}} \cdot \mathbb{A} \cdot \boldsymbol{\delta}_n - (\mathbf{f}, r\boldsymbol{\psi}) - \mathbf{F} \cdot r\hat{\boldsymbol{\psi}} \right] ds. \end{aligned}$$

Now, by a standard procedure (see e.g. [12, Section 2]), we use this relation for $t = T$, Lemma 3.3, and the functional properties of $(\mathbf{u}_n, \boldsymbol{\gamma}_n, \boldsymbol{\delta}_n)$ to deduce that for all $\boldsymbol{\varphi} \in \mathcal{C}_\#(B_R)$, it holds that

$$\begin{aligned} &\int_0^T \left[- \langle \mathbf{u}_n, \partial_t \boldsymbol{\varphi} \rangle + \lambda(((\mathbf{u}_n)_{\eta_n} - \boldsymbol{\gamma}_n + U - V) \cdot \nabla \mathbf{u}_n + (\mathbf{u}_n - \boldsymbol{\gamma}_n) \cdot \nabla U, \boldsymbol{\varphi}) \right. \\ &\quad \left. + 2(\mathbb{D}(\mathbf{u}_n), \mathbb{D}(\boldsymbol{\varphi})) + \varpi^{-1}\hat{\boldsymbol{\varphi}} \cdot \mathbb{A} \cdot \boldsymbol{\delta}_n - (\mathbf{f}, \boldsymbol{\varphi}) - \mathbf{F} \cdot \hat{\boldsymbol{\varphi}} \right] ds = 0. \end{aligned} \tag{3.57}$$

We also recall from (3.4)₂,

$$\dot{\delta}_n(t) = \gamma_n(t), \quad t \in [0, T], \quad \int_0^T \gamma_n(t) dt = \mathbf{0}. \tag{3.58}$$

By (3.56), (3.58), and (2.6), we infer that there exist

$$(\mathbf{u}, \boldsymbol{\gamma}, \boldsymbol{\delta}) \in L^2(0, T; \mathcal{D}^{1,2}(B_R)) \times L^2(0, T; \mathbb{R}^3) \times W^{1,2}(0, T; \mathbb{R}^3),$$

with $\hat{\mathbf{u}}(t) = \boldsymbol{\gamma}(t)$ such that, as $n \rightarrow \infty$,⁴

$$\begin{aligned} \mathbf{u}_n &\rightharpoonup \mathbf{u}, && \text{weakly in } L^2(0, T; \mathcal{D}^{1,2}(B_R)), \\ \boldsymbol{\gamma}_n &\rightharpoonup \boldsymbol{\gamma}, && \text{weakly in } L^2(0, T; \mathbb{R}^3), \\ \boldsymbol{\delta}_n &\rightharpoonup \boldsymbol{\delta}, && \text{weakly in } W^{1,2}(0, T; \mathbb{R}^3), \text{ and in } C([0, T]; \mathbb{R}^3). \end{aligned} \tag{3.59}$$

From Lemma 3.4, (3.56), (3.59)₁, and the Simon compactness theorem [30], we also get

$$\mathbf{u}_n \rightarrow \mathbf{u}, \quad \text{strongly in } L^2(0, T; \mathcal{H}(B_R)), \tag{3.60}$$

which implies

$$\boldsymbol{\gamma}_n \rightarrow \boldsymbol{\gamma}, \quad \text{strongly in } L^2(0, T; \mathbb{R}^3). \tag{3.61}$$

We now pass to the limit $n \rightarrow \infty$ in (3.57)–(3.58). Employing (3.59)–(3.61), it is not difficult to show that, in doing so, we can replace everywhere in (3.57)–(3.58), \mathbf{u}_n , $\boldsymbol{\gamma}_n$, and $\boldsymbol{\delta}_n$ with \mathbf{u} , $\boldsymbol{\gamma}$, and $\boldsymbol{\delta}$, respectively, with (3.58)₁ holding for a.a. $t \in [0, T]$. The only point that deserves a little care is the convergence of the non-linear term:

$$I_n := \int_0^T ((\mathbf{u}_n)_{\eta_n} - \boldsymbol{\gamma}_n) \cdot \nabla \mathbf{u}_n, \boldsymbol{\varphi} dt \rightarrow \int_0^T ((\mathbf{u} - \boldsymbol{\gamma}) \cdot \nabla \mathbf{u}, \boldsymbol{\varphi}) dt =: I. \tag{3.62}$$

To show (3.62), we first observe that, using the Schwarz inequality,

$$\begin{aligned} |I_n - I| &\leq \max_{t,x} |\boldsymbol{\varphi}(x, t)| \int_0^T (\|(\mathbf{u}_n - \mathbf{u})_{\eta_n}\|_2^2 + \|(\mathbf{u})_{\eta_n} - \mathbf{u}\|_2^2 + |\boldsymbol{\gamma}_n - \boldsymbol{\gamma}|^2) \|\nabla \mathbf{u}_n\|_2^2 dt \\ &\quad + \left| \int_0^T ((\mathbf{u} - \boldsymbol{\gamma}) \cdot (\nabla \mathbf{u}_n - \nabla \mathbf{u}), \boldsymbol{\varphi}) dt \right|. \end{aligned}$$

Then, using (3.59)₁, (3.60), (3.61), and classical properties of the mollifier, we infer the convergence in (3.62) holds. Finally, (3.54) is established by letting $n \rightarrow \infty$ in (3.56)₁ and using (3.59)_{1,2}.

It remains to prove (3.55). The first inequality in (3.55) is an obvious consequence of (2.5) and (3.54). To show the second one, we choose in (3.53) $\boldsymbol{\varphi} = \psi \boldsymbol{\phi}$, for arbitrary $\boldsymbol{\phi} \in \mathcal{C}_0(B_{R_0})$ and $\psi \in C_0^\infty((0, T); \mathbb{R})$. We thus obtain

$$\int_0^T \psi'(t) (\mathbf{u}(t), \boldsymbol{\phi}) dt = \int_0^T G_{\boldsymbol{\phi}}(t) \psi(t) dt, \tag{3.63}$$

⁴In what follows we shall not make notational distinction between sequences and subsequences.

where

$$G_\phi(t) = -\lambda((\mathbf{u} - \boldsymbol{\gamma} + \mathbf{U} - \mathbf{V}) \cdot \nabla \mathbf{u} + (\mathbf{u} - \boldsymbol{\gamma}) \cdot \nabla \mathbf{U}, \phi) - 2(\mathbb{D}(\mathbf{u}), \mathbb{D}(\phi)) + (\mathbf{f}, \phi).$$

Arguing exactly as in the proof of Lemma 3.4, we then prove that there is a constant C depending on R_0 such that

$$|G_\phi(t)| \leq C(R_0)(|\mathbf{V}(t)| + \|\nabla \mathbf{u}(t)\|_2 + \|\nabla \mathbf{u}(t)\|_2^2)\|\mathbb{D}(\phi)\|_2, \tag{3.64}$$

for all $\phi \in \mathcal{C}_0(B_{R_0})$. Henceforth, since $\mathcal{C}_0(B_{R_0})$ is dense in $\mathcal{D}_0^{1,2}(\Omega_{R_0})$, we infer

$$G_\phi(t) = [\mathbf{g}(t), \phi]$$

for some $\mathbf{g}(t) \in \mathcal{D}_0^{-1,2}(\Omega_{R_0})$, where $[\cdot, \cdot]$ denotes the duality pairing

$$\mathcal{D}_0^{-1,2}(\Omega_{R_0}) \leftrightarrow \mathcal{D}_0^{1,2}(\Omega_{R_0}).$$

This, in combination with (3.63) and the arbitrariness of ψ , furnishes

$$\frac{d}{dt}(\mathbf{u}, \phi) = [\mathbf{g}(t), \phi],$$

in the sense of distribution. Furthermore, in view of (3.54) and (3.64), we have

$$\int_0^T |G_\phi(t)| dt \leq C(\mathcal{V}^{\frac{1}{2}} + \mathcal{V})\|\mathbb{D}(\phi)\|_2,$$

and the desired property is then proved. ■

3.6. T -periodic weak solutions for the original problem

This last subsection is dedicated to the proof of Theorem 3.1. Let $\{\Omega_n \equiv \Omega_{R_n}\}$, $R_1 > 3R_*$, be a sequence of “invading domains”, namely,

$$\Omega_{n-1} \subset \Omega_n, \quad n \in \mathbb{N}; \quad \bigcup_{n=1}^\infty \Omega_n = \Omega,$$

and let $\{s_n \equiv (\mathbf{u}_n, \boldsymbol{\gamma}_n, \delta_n)\}$ be the sequence of corresponding T -periodic weak solutions determined in Proposition 3.1. For each n , we extend \mathbf{u}_n to 0 outside Ω_n and continue to denote by \mathbf{u}_n its extension. Consequently, by Remark 2.1, $\{\mathbf{u}_n\} \subset W^{1,2}(\Omega) \cap \mathcal{D}^{1,2}(\mathbb{R}^3)$. From the bound (3.54), we deduce that there is a subsequence of $\{(\mathbf{u}_n, \boldsymbol{\gamma}_n)\}$, again denoted by the same symbol, with $\mathbf{u}_n|_{\Omega_0} = \boldsymbol{\gamma}_n$, and functions $(\mathbf{u}, \boldsymbol{\gamma}) \in L^2(0, T; \mathcal{D}^{1,2}(\mathbb{R}^3)) \times L^2(0, T; \mathbb{R}^3)$ such that

$$\begin{aligned} \mathbf{u}_n &\rightharpoonup \mathbf{u}, && \text{weakly in } L^2(0, T; \mathcal{D}^{1,2}(\mathbb{R}^3)), \\ \boldsymbol{\gamma}_n &\rightharpoonup \boldsymbol{\gamma}, && \text{weakly in } L^2(0, T; \mathbb{R}^3), \end{aligned} \tag{3.65}$$

and for which (3.2) holds. Fix $R_0 > R_*$ arbitrarily. From (3.55), (3.65)₁ and the Aubin–Lions–Simon theorem [30], we can extract another subsequence, again denoted by $\{u_n\}$, such that

$$u_n \rightarrow u, \quad \text{strongly in } L^2(0, T; L^2(\Omega_{R_0})). \tag{3.66}$$

Also, from a classical trace inequality, we get

$$|\gamma_n - \gamma| \leq c_\epsilon \|u_n - u\|_2 + \epsilon \|\nabla(u_n - u)\|_2,$$

where $\epsilon > 0$ is arbitrary [14, Exercise II.4.1], so that, from the latter, (3.65), and (3.66) we deduce

$$\gamma_n \rightarrow \gamma, \quad \text{strongly in } L^2(0, T; \mathbb{R}^3). \tag{3.67}$$

Since for all $n \in \mathbb{N}$ we have

$$\int_0^T \gamma_n(t) dt = \mathbf{0},$$

we deduce that γ has zero average. Now set

$$\tilde{\delta}_n := \delta_n - \bar{\delta}_n \quad \left(\bar{w} := T^{-1} \int_0^T w(t) dt \right).$$

Recalling that $\dot{\delta}_n = \gamma_n$, the Poincaré–Wirtinger inequality yields

$$\int_0^T |\tilde{\delta}_n - \tilde{\delta}_m|^2 \leq T^2 \int_0^T |\gamma_n - \gamma_m|^2,$$

so that $\{\tilde{\delta}_n\}$ is a Cauchy sequence in $L^2(0, T; \mathbb{R}^3)$. It follows that there exists $\tilde{\delta} \in W^{1,2}(0, T; \mathbb{R}^3)$ with zero average such that $\tilde{\delta} = \gamma$ and

$$\tilde{\delta}_n \rightarrow \tilde{\delta}, \quad \text{strongly in } W^{1,2}(0, T; \mathbb{R}^3). \tag{3.68}$$

The last subsequence we have selected may depend on R_0 . However, covering Ω , with an increasing sequence of bounded domains and using the Cantor diagonal method, we may extract a further subsequence for which all the above properties, and in particular (3.66), hold for all R_0 .

In order to complete the proof of the theorem, it remains to show that the limiting functions determined above satisfy the weak formulation of (3.3). In particular, we still need to prove the convergence of the sequence of averages $\{\bar{\delta}_n\}$. From the weak formulation (3.53)₁ satisfied by $(u_n, \gamma_n, \delta_n)$ and the arbitrariness of R it follows that for any fixed at will $\varphi \in \mathcal{C}_\#(\mathbb{R}^3)$ the sequence $\{s_n\}$ obeys the following equation for all sufficiently large $n \in \mathbb{N}$:

$$\begin{aligned} & \int_0^T [\langle u_n, \partial_t \varphi \rangle - \lambda((u_n - \gamma_n + U - V) \cdot \nabla u_n + (u_n - \gamma_n) \cdot \nabla U, \varphi) \\ & \quad - 2(\mathbb{D}(u_n), \mathbb{D}(\varphi)) - \varpi^{-1} \hat{\varphi} \cdot \mathbb{A} \cdot \tilde{\delta}_n + (f, \varphi) + F \cdot \hat{\varphi}] dt \\ & \quad = \varpi^{-1} T \tilde{\varphi} \cdot \mathbb{A} \cdot \bar{\delta}_n. \end{aligned} \tag{3.69}$$

The convergences proved for the sequences $\{\mathbf{u}_n, \boldsymbol{\gamma}_n, \tilde{\boldsymbol{\delta}}_n\}$ in (3.65)₁, (3.66) – valid for all R_0 – and (3.67)–(3.68), and an argument similar to that used in the proof of (3.62), we show that, as $n \rightarrow \infty$, the left-hand side of (3.69) converges to the same quantity with $(\mathbf{u}_n, \boldsymbol{\gamma}_n, \tilde{\boldsymbol{\delta}}_n)$ replaced by $(\mathbf{u}, \boldsymbol{\gamma}, \tilde{\boldsymbol{\delta}})$, for any arbitrarily fixed $\boldsymbol{\varphi}$. This implies that $\overline{\boldsymbol{\delta}}_n \rightarrow \tilde{\boldsymbol{\delta}}$, for some $\tilde{\boldsymbol{\delta}} \in \mathbb{R}^3$. We then deduce that the triple $(\mathbf{u}, \boldsymbol{\gamma}, \boldsymbol{\delta} = \tilde{\boldsymbol{\delta}} + \boldsymbol{\delta})$ satisfies (3.1)₁ whatever the test function $\boldsymbol{\varphi}$ taken in $\mathcal{C}_\#(\mathbb{R}^3)$. This concludes the proof. ■

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