

Weighted Weyl estimates near an elliptic trajectory

Thierry Paul and Alejandro Uribe

Abstract. Let ψ_j^{\hbar} and E_j^{\hbar} denote the eigenfunctions and eigenvalues of a Schrödinger-type operator H_{\hbar} with discrete spectrum. Let $\psi_{(x,\xi)}$ be a coherent state centered at a point (x, ξ) belonging to an elliptic periodic orbit, γ of action S_{γ} and Maslov index σ_{γ} . We consider “weighted Weyl estimates” of the following form: we study the asymptotics, as $\hbar \rightarrow 0$ along any sequence

$$\hbar = \frac{S_{\gamma}}{2\pi l - \alpha + \sigma_{\gamma}},$$

$l \in \mathbb{N}$, $\alpha \in \mathbb{R}$ fixed, of

$$\sum_{|E_j - E| \leq c\hbar} |(\psi_{(x,\xi)}, \psi_j^{\hbar})|^2.$$

We prove that the asymptotics depend strongly on α -dependent arithmetical properties of c and on the angles θ of the Poincaré mapping of γ . In particular, under irrationality assumptions on the angles, the limit exists for a non-open set of full measure of c 's. We also study the regularity of the limit as a function of c .

1. Introduction and results.

Consider a Schrödinger operator $H = -\hbar^2 \Delta + V(x)$ with V smooth, either on $M = \mathbb{R}^m$ (in which case we assume V tends to infinity at infinity and therefore H has discrete spectrum) or on a compact Riemannian

manifold, M . In [7] we considered “trace formulae” associated to projectors on coherent states in the following sense. For $(x, \xi) \in \mathbb{R}^{2m}$ and $a \in \mathcal{S}(\mathbb{R}^m)$ define the coherent state $\psi_{x\xi}^a$ as:

$$(1) \quad \psi_{(x,\xi)}^a(y) = \rho(y-x) (2\pi\hbar)^{-3m/4} 2^{-m/4} e^{-ix\xi/2\hbar} e^{i\xi y/\hbar} \hat{a}\left(\frac{y-x}{\sqrt{\hbar}}\right).$$

Here ρ is a cut-off function near zero and \hat{a} is the Fourier transform of a , (in the manifold case $(x, \xi) \in T^*M$ and the above definition is in local coordinates near x). Let ψ_j and E_j the eigenfunctions and eigenvalues of H . Then if φ is a Schwartz function whose Fourier transform is compactly supported and $E = |\xi|^2 + V(x)$, we have

$$(2) \quad \sum_j \varphi\left(\frac{E_j - E}{\hbar}\right) |(\psi_{(x,\xi)}, \psi_j)|^2 \sim \sum_{j=0} c_j^\varphi(x, \xi) \hbar^{-m+1/2+j},$$

for $\hbar \rightarrow 0$. (If $E \neq |\xi|^2 + V(x)$, the left-hand side tends to 0 rapidly in \hbar .) Although the form of the asymptotic expansion does not depend on (x, ξ) , the coefficient $c_0(x, \xi)$ is highly sensitive to the point (x, ξ) being periodic or not with respect to the classical flow. In case (x, ξ) is either not periodic or is on a *hyperbolic* trajectory, we proved in [7] (using a Tauberian theorem) that, for every $c \in \mathbb{R}$,

$$(3) \quad \sum_{|E_j - E| \leq c\hbar} |(\psi_{(x,\xi)}, \psi_j)|^2 = \chi_{[-c,c]}^{x, \xi}(x, \xi) \hbar^{-m+1/2} + o(\hbar^{-m+1/2}),$$

as $\hbar \rightarrow 0$ possibly along certain sequence. (Here $\chi_{[-c,c]}$ is the characteristic function of the interval $[-c, c]$.) The main goal of this paper is to study the case where (x, ξ) belongs to an elliptic closed trajectory.

Our results are related to the existence of quasi-modes near an elliptic trajectory. Recall that if H is as before and γ is a closed elliptic trajectory of the Hamiltonian $|\xi|^2 + V(x)$ with energy E , period T_γ , action S_γ , Maslov index σ_γ and Poincaré mapping of angles θ_j , $j = 1, \dots, m-1$, then one can construct (see [9], [3], [8], [7]) quasi-modes of H (namely solutions of the Schrödinger equation modulo a remainder), microlocalized near γ , of quasi-energies

$$(4) \quad E_{QM}^{k,l} = E + \frac{\hbar}{T_\gamma} \left(\left(2\pi l - \frac{S_\gamma}{\hbar} \right) + \sum_{j=1}^{m-1} \left(k_j + \frac{1}{2} \right) \theta_j + \sigma_\gamma \right),$$

for $(k, l) \in \mathbb{Z}^m$, l large. The remainder is $O(\hbar^2)$ uniformly as

$$\left| 2\pi l - \frac{S\gamma}{\hbar} \right| \quad \text{and} \quad |k| := \sum k_j$$

remain bounded. The existence of these quasi-modes implies that part of the spectral density of H concentrates near the quasi-energies defined by (4), but this doesn't say anything about $E_{QM}^{k,l}$ as $|k| \rightarrow \infty$ and does not involve the rest of the spectrum. The results of this paper will indicate that the rescaled localized spectral density

$$(5) \quad \sum_j \delta\left(\frac{E_j - \lambda}{\hbar}\right) |(\psi_{(x,\xi)}, \psi_j)|^2$$

(which is the rescaled spectral density microlocalized at the point in phase space (x, ξ)) has a certain semiclassical limit whose singularities are indeed precisely the quasi-energies (4), and this time with no restriction on $|k|$.

We will now state our results, valid for more general quantum Hamiltonians: Let $H_\hbar = \sum_{l=0}^L \hbar^l P_l(x, D_x)$ where P_l is a differential operator of order l on \mathbb{R}^m (or M) of principal symbol P_l^0 , sub-principal symbol P_l^{-1} (formally P_l is regarded as acting on half-densities) and smooth coefficients. Let $\mathcal{H}(x, \xi) = \sum_{l=0}^L P_l^0(x, \xi)$ and $\mathcal{H}_{\text{sub}}(x, \xi) = \sum_{l=0}^L P_l^{-1}(x, \xi)$ be the principal and sub-principal symbols of H_\hbar . We assume that P_L is elliptic, \mathcal{H} is positive, and in case $M = \mathbb{R}^m$, that \mathcal{H} tends polynomially to infinity at infinity. We will also suppose for simplicity that $\mathcal{H}_{\text{sub}}(x, \xi) = 0$.

Let E_j^\hbar and ψ_j^\hbar denote the eigenvalues and eigenvectors of H_\hbar . Let us suppose that (x, ξ) belongs to an elliptic trajectory of period T_γ , action S_γ , Maslov index σ_γ and Poincaré mapping of angles $\theta = (\theta_1, \dots, \theta_{m-1})$. We will use throughout the notations

$$k = (k_1, \dots, k_{m-1}) \in \mathbb{N}^{m-1},$$

$$(6) \quad k \theta := \sum_{j=1}^{m-1} k_j \theta_j \quad \text{and} \quad \left(k + \frac{1}{2}\right) \theta := \sum_{j=1}^{m-1} \left(k_j + \frac{1}{2}\right) \theta_j.$$

Theorem 1.1. *Assume that $\theta_1/(2\pi), \dots, \theta_{m-1}/(2\pi)$ are rational. Then, for every $\alpha \in [0, 2\pi)$, as $\hbar \rightarrow 0$ along the sequence*

$$(7) \quad \hbar = \frac{S_\gamma}{2\pi l - \alpha + \sigma_\gamma}, \quad l \in \mathbb{N},$$

one has

$$(8) \quad \sum_{|E_j - E| \leq c\hbar} |(\psi_{(x,\xi)}, \psi_j)|^2 = \hbar^{-m+1/2} \mathcal{L}_\alpha(c) + o(\hbar^{-m+1/2}),$$

for all c such that

$$(9) \quad c \neq \pm \frac{1}{T_\gamma} \left(2\pi j + \left(k + \frac{1}{2} \right) \theta + \alpha \right), \quad \text{for all } j \in \mathbb{Z}, k \in \mathbb{N}^{m-1}.$$

Moreover, as a function of c the limit $\mathcal{L}_\alpha(c)$ is a step function constant on the intervals defined by (9).

Next we consider the irrational case:

Theorem 1.2. *Assume that $1, \theta_1/(2\pi), \dots, \theta_{m-1}/(2\pi)$ are linearly independent over the rationals. Then there exists a set \mathcal{M}^α of values of c , of full Lebesgue measure, such that for all $c \in \mathcal{M}^\alpha$*

$$(10) \quad \sum_{|E_j - E| \leq c\hbar} |(\psi_{(x,\xi)}, \psi_j)|^2 = \hbar^{m-1/2} \mathcal{L}_\alpha(c) + o(\hbar^{m-1/2}),$$

for \hbar as in (7). Moreover, as a function of c , $\mathcal{L}_\alpha(c)$ is locally Lipschitz on \mathcal{M}^α in the sense that for all $c \in \mathcal{M}^\alpha$ there exists $\beta_c > 0$ such that,

$$(11) \quad |\mathcal{L}_\alpha(c') - \mathcal{L}_\alpha(c)| \leq \beta_c |c' - c|, \quad \text{for all } c' \in \mathcal{M}^\alpha.$$

Finally there exists a rapidly decreasing family $\{g_k\}_{k \in \mathbb{N}^{m-1}}$ (related to the microlocalization of the symbol a of $\psi_{(x,\xi)}$) such that

$$(12) \quad \{c : \text{for all } k \in \mathbb{N}^{m-1} \quad |1 - e^{i(cT_\gamma + (k+1/2)\theta + \alpha)}| > \varepsilon g_k\} \subset \mathcal{M}^\alpha,$$

for all $\varepsilon > 0$. (For a precise definition of the set \mathcal{M}^α see Lemma 3.3.).

REMARK. In the rational case the discontinuities of the function \mathcal{L}_α are located exactly at the values of the $E_{QM}^{k,l}$ defined before by (4), for the values of \hbar given by (7). In the irrational case in order to prove that $\mathcal{L}_\alpha(c)$ exists we need that c be at some distance from the quasi-energies $E_{QM}^{k,l}$ (unless the symbol a of the quasi-mode is chosen very judiciously, in which case we can work with c in the complement of the set of all quasi-energies). In all cases this suggests that the weighted spectral measure, (5), in the semi-classical limit, is particularly singular exactly at the values of the $E_{QM}^{k,l}$ defined before. We hope to provide a rigorous proof of a precise statement of this elsewhere.

The paper is organized as follows: In Section 3 we prove the existence of the functions \mathcal{L}_α which are studied in Section 4. In Section 5 we finish the proof of the main Theorems, using a Tauberian argument that we recall in Section 2. Finally, in the appendix we review and extend slightly a result on Hölder continuity of function such as \mathcal{L}_α using wavelets.

2. A Tauberian lemma.

In this section we refine the Tauberian lemma of [2] and [7]. Consider an expression of the following form

$$(13) \quad \Upsilon_{E,\hbar}^w(\varphi) = \sum_j w_j(\hbar) \varphi\left(\frac{E_j(\hbar) - E}{\hbar}\right),$$

defined for all $\varphi \in \mathcal{R}$ where \mathcal{R} will henceforth denote the set of all Schwartz functions on the line with compactly supported Fourier transform.

Let \mathcal{M}^α a subset of \mathbb{R}^+ of full Lebesgue measure in a bounded interval.

We introduce the following notations. Fix a positive function $f \in \mathcal{R}$ satisfying $f(0) = 1$ and $\hat{f}(0) = 1$. For every $a > 0$, define

$$(14) \quad f_a(r) := a^{-1} f\left(\frac{r}{a}\right)$$

and for every $a > 0$ and $c > 0$

$$(15) \quad \varphi_{a,c} := f_a * \chi_{[-c,c]},$$

where $\chi_{[-c,c]}$ is the characteristic function of the interval $[-c, c]$.

The Tauberian lemma in question is:

Theorem 2.1 (See [2] and [7]). *Let \mathcal{M}^α a subset of \mathbb{R}^+ of full Lebesgue measure in a bounded interval. Suppose $w_j(\hbar)$, $E_j(\hbar)$, E and Υ_\hbar^w itself satisfy all of the following:*

1) *There exists a positive function $\omega(\hbar)$, defined on an interval $(0, \hbar_0)$, and a functional \mathcal{F}_0 on \mathcal{R} , such that for all $\varphi \in \mathcal{R}$*

$$(16) \quad \Upsilon_{E, \hbar}^w(\varphi) = \mathcal{F}_0(\varphi) \omega(\hbar) + o(\omega(\hbar)), \quad \hbar \rightarrow 0.$$

2) *for all $c \in \mathcal{M}^\alpha$ the limit*

$$\mathcal{L}_\alpha(c) = \lim_{a \rightarrow 0} \mathcal{F}_0(\varphi_{a,c})$$

exists.

3) *\mathcal{L}_α is a continuous function on \mathcal{M}^α .*

4) *There exists a $k \in \mathbb{Z}$ such that $\hbar^k = \mathcal{O}(\omega(\hbar))$, $\hbar \rightarrow 0$.*

5) *There exists an $\varepsilon > 0$ such that for every φ there is a constant C_φ such that for all $E' \in [E - \varepsilon, E + \varepsilon]$*

$$(17) \quad |\Upsilon_{E', \hbar}^w(\varphi)| \leq C_\varphi \omega(\hbar)$$

(rough uniformity in E).

6) *The $w_j(\hbar)$ are non-negative and bounded: there exists a constant $C \geq 0$ such that for all j and all $\hbar, 0 < \hbar < \hbar_0$*

$$(18) \quad 0 \leq w_j(\hbar) \leq C.$$

7) *The eigenvalues $E_j(\hbar)$ satisfy the following rough estimate: for each C_1 there exist constants C_2, N_0 such that for all k*

$$(19) \quad \#\{j : E_j(\hbar) \leq C_1 + k\hbar\} \leq C_2(\hbar^{-1}k)^{N_0}.$$

Define the weighted counting function by

$$(20) \quad N_{E,c}^w(\hbar) = \sum_{j : |x_j(\hbar)| \leq c} w_j(\hbar),$$

where

$$(21) \quad x_j(\hbar) := \frac{E_j(\hbar) - E}{\hbar}.$$

Then the conclusion is: for all $c \in \mathcal{M}^\alpha$,

$$(22) \quad N_{E,c}^w(\hbar) = \mathcal{L}_\alpha(c) \omega(\hbar) + o(\omega(\hbar)), \quad \hbar \longrightarrow 0.$$

PROOF. Except for the fact that the set \mathcal{M}^α of allowed c 's is not \mathbb{R}^+ , this theorem is precisely [2, Theorem 6.3]. Proceeding exactly as in the proof of the [2, inequalities (188)], one shows that for all $R > 0$, for all $N \in \mathbb{N}$ exists $C > 0$, $C_N > 0$ such that for all $a \in (0, R)$ and for all η , $0 < \eta < c$,

$$(23) \quad \begin{aligned} \frac{1}{\omega(\hbar)} \left(1 - C \frac{a}{\eta}\right) N_{E,c-\eta}(\hbar) &\leq \frac{1}{\omega(\hbar)} \Upsilon_{E,\hbar}(\varphi_{a,c}) \\ &\leq \frac{1}{\omega(\hbar)} N_{E,c+\eta}(\hbar) + C_N \left(\frac{a}{\eta}\right)^N. \end{aligned}$$

Let $c \in \mathcal{M}^\alpha$ be given. We begin by observing that by the first of the inequalities (23)

$$(24) \quad \frac{1}{\omega(\hbar)} N_{E,c}(\hbar) \leq \frac{1}{\omega(\hbar)} \Upsilon_{E,\hbar}(\varphi_{a,c+\eta}) + C_1 \frac{a}{\eta},$$

where we have also used the fact that $N_{E,c}(\hbar)/\omega(\hbar)$ is bounded (a trivial consequence of (16)). For every η such that $0 < \eta < c$ one can take the limit in (24) as $\hbar \longrightarrow 0$ to obtain that

$$(25) \quad \limsup_{\hbar \rightarrow 0} \frac{1}{\omega(\hbar)} N_{E,c}(\hbar) \leq \mathcal{F}_0(\varphi_{a,c+\eta}) + C_1 \frac{a}{\eta}.$$

If we now assume that $\eta + c \in \mathcal{M}^\alpha$ we can take the limit as $a \longrightarrow 0$ to obtain

$$(26) \quad \limsup_{\hbar \rightarrow 0} \frac{1}{\omega(\hbar)} N_{E,c}(\hbar) \leq \mathcal{L}_\alpha(c + \eta).$$

By the assumption that \mathcal{M}^α has full measure, we can find a sequence $\{\eta_j\}$ such that for all j , $c + \eta_j \in \mathcal{M}^\alpha$ and $\eta_j \longrightarrow 0$. Taking the limit in

(26) of $\mathcal{L}_\alpha(c + \eta_j)$ as $j \rightarrow \infty$ and using the fact that \mathcal{L}_α is continuous at c we obtain

$$(27) \quad \limsup_{\hbar \rightarrow 0} \frac{1}{\omega(\hbar)} N_{E,c}(\hbar) \leq \mathcal{L}_\alpha(c).$$

A similar argument starting with the second inequality (23) shows that

$$(28) \quad \liminf_{\hbar \rightarrow 0} \frac{1}{\omega(\hbar)} N_{E,c}(\hbar) \geq \mathcal{L}_\alpha(c),$$

which finishes the proof.

3. The existence of $\mathcal{L}_\alpha(c)$.

In this section we prove the existence of the coefficients $\mathcal{L}_\alpha(c)$ in the limits (8) and (10) (see (36) below).

Lemma 3.1. *There exists a rapidly decreasing family of non-negative numbers, $\{c_k\}_{k \in \mathbb{N}^{m-1}}$, such that for all $\varphi \in \mathcal{R}$ the first coefficient $c_k^\varphi(x, \xi)$ in (2) can be written as*

$$(29) \quad c_0^\varphi(x, \xi) = \sum_{n=-\infty}^{+\infty} \sum_{k \in \mathbb{N}^{m-1}} \hat{\varphi}(n T_\gamma) c_k e^{in((k+1/2)\theta + \alpha)}.$$

PROOF. In [7] we proved that the first coefficient $c_0^\varphi(x, \xi)$ in (2) can be written as

$$(30) \quad \begin{aligned} & 2^{2n} \pi^{(3n+1)/2} c_0^\varphi(x, \xi) \\ &= \sum_{n=-\infty}^{+\infty} \hat{\varphi}(n T_\gamma) e^{in S_\gamma / \hbar + \sigma_\gamma} \int_{-\infty}^{+\infty} (a, Z((s \dot{x}, s \dot{\xi})) U^n a) ds, \end{aligned}$$

where $(\dot{x}, \dot{\xi})$ is the tangent vector to the classical flow at (x, ξ) , Z is the Weyl/Heisenberg operator defined by

$$(31) \quad Z(e, f)(a)(\eta) = e^{-ief/2} e^{ie\eta} a(\eta - f)$$

and U is the metaplectic representation of the linearized flow at time T_γ . (We should point out that in the manifold case a defines intrinsically a

smooth vector in the metaplectic representation of $T_{(x,\xi)}(T^*M)$, and U and Z are operators in that representation space.) Denoting by S the linearized flow at time T_γ , we also showed that one can find a symplectic mapping R such that $R^{-1}SR$ is block-diagonal of the form

$$(32) \quad R^{-1}SR = \begin{pmatrix} 1 & \mu & 0 \\ 0 & 1 & 0 \\ 0 & 0 & A_\theta \end{pmatrix},$$

where $\mu \in \mathbb{R}$ and A_θ is the direct sum of rotations of angles $\theta_1, \dots, \theta_{m-1}$. Furthermore, the transformation R maps the vector $(s\dot{x}, s\dot{\xi})$ to the vector $(s, 0)$.

Let us denote $a' := \text{Mp}(R)^{-1}a$ and $V := \text{Mp}(R^{-1}SR)$, where $\text{Mp}(R)$ denotes the metaplectic representation of the mapping R . Then, letting $Z(s) := Z((s\dot{x}, s\dot{\xi}))$ and

$$W(s) := \text{Mp}(R)^{-1}Z(s)\text{Mp}(R) = Z(s, 0, 0, 0),$$

one has

$$(a, Z(s)U^n a) = (a', W(s)V^n a').$$

Denote the variables of a' by (η_1, η_2) where $\eta_1 \in \mathbb{R}$ and $\eta_2 \in \mathbb{R}^{m-1}$, and let $e^{i\theta(D_{\eta_2}^2 + \eta_2^2)/2}$ denote the direct sum of the propagators of one-dimensional Harmonic oscillators at times $\theta_1, \dots, \theta_{m-1}$, acting on a' by acting on the η_2 variables. If $e^{in\mu\partial_{\eta_1}^2/2}$ denotes the metaplectic quantization of

$$(33) \quad \begin{pmatrix} 1 & n\mu \\ 0 & 1 \end{pmatrix},$$

we get that (30) becomes

$$\begin{aligned} & 2^{2n} \pi^{(3n+1)/2} c_0^\varphi(x, \xi) \\ &= \sum_{n=-\infty}^{+\infty} \hat{\varphi}(n T_\gamma) e^{in\alpha} \\ & \quad \cdot \int \overline{a'(\eta)} (e^{in\mu\partial_{\eta_1}^2/2} e^{in\theta(D_{\eta_2}^2 + \eta_2^2)/2} (a')) (\eta_1 - s, \eta_2) d\eta ds. \end{aligned}$$

The integral over ds is a convolution and the integral over $d\eta_1$ is the integral of that convolution. Therefore, using the Fourier inversion

formula plus the fact that on the Fourier transform side the operator $e^{in\mu\partial_{\eta_1}^2/2}$ is multiplication by $e^{-in\mu\zeta^2/2}$ (ζ being the dual variable), one gets

$$(34) \quad \begin{aligned} & 2^{2n} \pi^{(3n+1)/2} c_0^\varphi(x, \xi) \\ &= \sum_{n=-\infty}^{+\infty} \hat{\varphi}(n T_\gamma) e^{in\alpha} \int \overline{a'^{\wedge}(0, \eta_2)} e^{in\theta(D_{\eta_2}^2 + \eta_2^2)/2} a'^{\wedge}(0, \eta_2) d\eta_2, \end{aligned}$$

where a'^{\wedge} is the Fourier transform of a' with respect to η_1 . Let $b(x) := a'^{\wedge}(0, x)$ and let us decompose b on the Hermite basis, h_k , of eigenfunctions of the harmonic oscillator

$$(35) \quad b = \sum_{k \in \mathbb{N}^{m-1}} b_k h_k.$$

Then, letting $c_k := |b_k|^2$ we get (29) and the family $\{c_k\}$ is non-negative. It is also rapidly decreasing since the function b is Schwartz.

REMARK. For a given quantum Hamiltonian H , the coefficients $\{c_k\}$ depend only on the symbol a of the coherent state. Observe that the proof shows that given any rapidly decreasing family $\{c_k\}$ one can find an a giving rise to it.

We next prove the existence of the limit

$$(36) \quad \mathcal{L}_\alpha(c) := \lim_{a \rightarrow 0} c_0^{(f_a * \chi_{[-c, c]})}(x, \xi),$$

for f as in the Tauberian lemma and $c_0^\varphi(x, \xi)$ as in (29). Let $\phi_a(c) := c_0^{(f_a * \chi_{[-c, c]})}(x, \xi)$, that is

$$(37) \quad \phi_a(c) := c + \sum_{n \neq 0, k} \hat{f}(a n) \frac{\sin(n c T_\gamma)}{n T_\gamma} c_k e^{in((k+1/2)\theta + \alpha)}.$$

We must then prove that the limit $\mathcal{L}_\alpha(c) = \lim_{a \rightarrow 0} \phi_a(c)$ exists.

To lighten up the notation a bit, let us define

$$(38) \quad d_k := \left(k + \frac{1}{2}\right)\theta + \alpha, \quad k \in \mathbb{N}^{m-1},$$

keeping in mind the notation (6). Let $0 < a < 1$, then

$$\begin{aligned}
 \phi_1(c) - \phi_a(c) &= \frac{1}{T_\gamma} \sum_{(n,k) \in \mathbb{Z} \times \mathbb{N}^{m-1}} \sin(cnT_\gamma) c_k e^{ind_k} \int_a^1 \hat{f}'(tn) dt \\
 (39) \qquad &= \frac{1}{T_\gamma} \sum_{(n,k)} \left(\frac{e^{incT_\gamma} - e^{-incT_\gamma}}{2i} \right) c_k e^{ind_k} \int_a^1 \hat{f}'(tn) dt.
 \end{aligned}$$

Applying the Poisson summation formula to the series over n , we get (after a calculation)

$$\begin{aligned}
 \phi_1(c) - \phi_a(c) &= \frac{-\pi}{T_\gamma} \sum_{(j,k) \in \mathbb{Z} \times \mathbb{N}^{m-1}} c_k \int_a^1 \left(g\left(\frac{1}{t}(2\pi j + cT_\gamma + d_k)\right) \right. \\
 (40) \qquad &\qquad \qquad \left. - g\left(\frac{1}{t}(2\pi j - cT_\gamma + d_k)\right) \right) \frac{dt}{t},
 \end{aligned}$$

where $g(x) := x f(x)$.

Lemma 3.2. *Define*

$$\mathcal{M}_0^\alpha = \left\{ c \in \mathbb{R} : \text{for all } (j,k) \in \mathbb{Z} \times \mathbb{N}^{m-1}, c \neq \pm \frac{1}{T_\gamma} (2\pi j + d_k) \right\}.$$

If $\theta_1/(2\pi), \dots, \theta_{m-1}/(2\pi)$ are rational and $c \in \mathcal{M}_0^\alpha$, then each of the limits

$$(41) \qquad \lim_{a \rightarrow 0} \sum_{(j,k) \in \mathbb{Z} \times \mathbb{N}^{m-1}} c_k \int_a^1 g\left(\frac{1}{t}(2\pi j \pm cT_\gamma + d_k)\right) \frac{dt}{t},$$

exists (and is finite). Moreover, the convergence is locally uniform in c .

PROOF. By the rationality assumption the complement of \mathcal{M}_0^α is discrete. Therefore, if $c \in \mathcal{M}_0^\alpha$ there exists ε such that

$$0 < \varepsilon \leq |2\pi j \pm cT_\gamma + d_k|, \qquad \text{for all } (j,k) \in \mathbb{Z} \times \mathbb{N}^{m-1}.$$

The function g is rapidly decreasing: for all $N \in \mathbb{N} \exists C_N > 0$ such that for all $x \in \mathbb{R}$, $|g(x)| \leq C_N (1 + |x|)^{-N}$. Therefore

$$\begin{aligned}
 (42) \qquad \left| g\left(\frac{2\pi j \pm cT + d_k}{t}\right) \right| &\leq C_N \frac{t^N}{t^N + (2\pi j \pm cT + d_k)^N} \\
 &\leq C_N \frac{t^N}{(2\pi j \pm cT + d_k)^N},
 \end{aligned}$$

and so for all $(j, k) \in \mathbb{Z} \times \mathbb{N}^{m-1}$ and for all $a \in (0, 1)$

$$(43) \quad \int_a^1 \left| g\left(\frac{1}{t}(2\pi j \pm cT_\gamma + d_k)\right) \right| \frac{dt}{t} \leq \frac{C_N}{N} \frac{1 - a^{N+1}}{|2\pi j \pm cT_\gamma + d_k|^N}.$$

This shows that each of the integrals in the series (41) extends to a continuous function of $a \in [0, 1)$. Moreover, since the family

$$M_{k,j} := \frac{c_k}{|2\pi j \pm cT_\gamma + d_k|^N}, \quad (j, k) \in \mathbb{Z} \times \mathbb{N}^{m-1}$$

is absolutely convergent (for N sufficiently large) and it dominates the absolute values of the terms of (41), we are done.

We now turn to the irrational case.

Lemma 3.3. *Assume that $1, \theta_1/(2\pi), \dots, \theta_{m-1}/(2\pi)$ are linearly independent over the rationals. Let*

$$(44) \quad \mathcal{M}_\pm^\alpha := \left\{ c \in \mathcal{M}_0^\alpha : \sum_{k \in \mathbb{N}^{m-1}} c_k \left(\pm \left(d_k + \frac{cT}{2\pi} \right) \right)^{-2} < \infty \right\},$$

where $\{x\}$ denotes the fractional part of x , and let

$$(45) \quad \mathcal{M}^\alpha := \mathcal{M}_+^\alpha \cap \mathcal{M}_-^\alpha.$$

Then, if $c \in \mathcal{M}^\alpha$, each of the limits

$$\lim_{a \rightarrow 0} \sum_{j,k} c_k \int_a^b g\left(\frac{1}{t}(2\pi j \pm cT_\gamma + d_k)\right) \frac{dt}{t}$$

exists and is finite. Moreover, the convergence is locally uniform in c .

PROOF. It is enough to consider one of the series above, say the one with the plus sign. Let $c \in \mathcal{M}^\alpha$ and define

$$O^+ := \{(j, k) \in \mathbb{Z} \times \mathbb{N}^{m-1} : 2\pi j + cT_\gamma + d_k > 0\},$$

and

$$O^- := \{(j, k) \in \mathbb{Z} \times \mathbb{N}^{m-1} : 2\pi j + cT_\gamma + d_k < 0\}.$$

Since $c \in \mathcal{M}_0^\alpha$, $\mathbb{Z} \times \mathbb{N}^{m-1} = O^+ \cup O^-$. Recalling that $g(x) = x f(x)$ and that f as well as the c_k are non-negative, we see that the terms with $(j, k) \in O^\pm$ have the sign \pm and therefore each of

$$\sum_{(j,k) \in O^\pm} c_k \int_a^1 g\left(\frac{1}{t}(2\pi j + cT_\gamma + d_k)\right) \frac{dt}{t}$$

is a decreasing function of a . It therefore suffices to show that

$$\lim_{a \rightarrow 0} \sum_{(j,k) \in O^+} c_k \int_a^1 g\left(\frac{1}{t}(2\pi j + cT_\gamma + d_k)\right) \frac{dt}{t} < \infty$$

and similarly for the series over O^- .

Specializing (43) to $N = 2$, we see that exists $C > 0$ such that for all $a \in (0, 1)$ and for all $(j, k) \in O^+$

$$(46) \quad \int_a^1 g\left(\frac{2\pi j + cT + d_k}{t}\right) \frac{dt}{t} \leq \frac{C}{(2\pi j + cT + d_k)^2}.$$

(The last denominator is not zero if $(j, k) \in O^+$.) Therefore, the Lemma will be proved provided we show the convergence of the double series of scalars

$$(47) \quad \sum_{(j,k) \in O^+} M_{k,j},$$

where

$$M_{k,j} = c_k \left(j + \frac{cT + d_k}{2\pi}\right)^{-2},$$

that is

$$(48) \quad M_{k,j} = \frac{c_k}{(j + k\xi + \beta)^2},$$

where

$$(49) \quad \xi = \left(\frac{\theta_1}{2\pi}, \dots, \frac{\theta_{m-1}}{2\pi}\right) \quad \text{and} \quad \beta = \frac{1}{2\pi} \left(cT + \alpha + \sum_{j=1}^{m-1} \frac{\theta_j}{2}\right).$$

Since the terms in (47) are positive, we can prove its convergence by first summing over j with k fixed, and then summing over $k \in \mathbb{N}^{m-1}$. Observe that

$$(50) \quad (j, k) \in O^+ \quad \text{if and only if} \quad j \geq [-k\xi - \beta] + 1,$$

where $[x]$ denotes the greatest integer less than or equal to x . For every k consider the series

$$(51) \quad \sum_{j=-[k\xi+\beta]}^{+\infty} M_{k,j} .$$

(If $x \notin \mathbb{Z}$, then $[-x] = -[x] - 1$, and since $c \in \mathcal{M}_0^\alpha$, for all $k \in \mathbb{N}^{m-1}$, $k\xi + \beta \notin \mathbb{Z}$.) Comparing this series with the integral

$$(52) \quad \int_{-[k\xi+\beta]}^{+\infty} \frac{dx}{(x + k\xi + \beta)^2} ,$$

we find that

$$(53) \quad \sum_{j=-[k\xi+\beta]}^{+\infty} M_{k,j} \leq M_{k,[-k\xi+\beta]} + \frac{c_k}{-[k\xi + \beta] + k\xi + \beta} ,$$

or with the notation $\{x\} =$ fractional part of $x = x - [x]$,

$$(54) \quad \sum_{j=-[k\xi+\beta]+1}^{+\infty} M_{k,j} \leq \frac{c_k}{\{k\xi + \beta\}^2} + \frac{c_k}{\{k\xi + \beta\}} .$$

Therefore convergence of (47) follows from the convergence of

$$\sum_{k \in \mathbb{N}^{m-1}} \frac{c_k}{\{k\xi + \beta\}^2} .$$

But since by assumption $c \in \mathcal{M}_+^\alpha$, this series converges.

In conclusion we have shown that $\mathcal{L}_\alpha(c)$ exists for c as defined by the Lemmas.

REMARKS. In the irrational case:

1) To find examples of numbers c in \mathcal{M}_\pm^α , it suffices to find a family $\{g_k\}_{k \in \mathbb{N}^{m-1}}$ of positive numbers such that $\sum g_k^{-2} c_k < \infty$. Then if

$$(55) \quad \left| 2\pi j \pm c T_\gamma + \left(k + \frac{1}{2}\right)\theta + \alpha \right| > g_k , \quad \text{for all } (j, k) \in \mathbb{Z} \times \mathbb{N}^{m-1} ,$$

then $c \in \mathcal{M}^\alpha$. Defining $\hat{g}_k = \varepsilon g_k$, $\varepsilon > 0$, one see that still $\sum \hat{g}_k^{-2} c_k < \infty$ and therefore associated to \hat{g}_k (by (55)) is a subset of \mathcal{M}_α whose intersection with any interval I has a co-measure in I arbitrary small as $\varepsilon \rightarrow 0$; therefore \mathcal{M}^α has full measure.

2) The set \mathcal{M}^α is related to the rate of decay of the c_k (that is to the properties of the symbol, a , of the coherent states), as well as to irrationality properties of $\theta/(2\pi)$. At one extreme, we can choose a such that only finitely-many of the coefficients c_k are non-zero (see the remark following Lemma 3.1). In that case $\mathcal{M}^\alpha = \mathcal{M}_0^\alpha$ is just the complement of the set of quasi-energies of the quasi-modes associated with the trajectory.

4. Properties of the function \mathcal{L}_α .

Having established the existence of the function $\mathcal{L}_\alpha(c)$, we now derive some of its properties.

Rational case. Let us go back to the identity $\mathcal{L}_\alpha(c) = \lim_{a \rightarrow 0} \phi_a(c)$ where ϕ_a is defined in (37). Applying in (37) the Poisson summation formula to the series over n with k fixed one obtains

$$(56) \quad \mathcal{L}_\alpha(c) = \lim_{a \rightarrow 0} \frac{1}{a} \left(F_c * f \left(\frac{\cdot}{a} \right) \right) (0),$$

where

$$(57) \quad F_c(y) = \int_{-c}^c \sum_{j,k} c_k \delta(T_\gamma(x-y) - 2\pi j - d_k) dx.$$

For each $c > 0$ the function F_c is a step function; indeed

$$(58) \quad F_c = \sum_{j,k} c_k \chi_{[-c-(2\pi j+d_k)/T, c-(2\pi j+d_k)/T]}.$$

Since $f(\cdot/a)/a \rightarrow \delta$, we obtain

$$(59) \quad \mathcal{L}_\alpha(c) = \sum_{\{j,k : -cT < 2\pi j + d_k < cT\}} c_k, \quad \text{for all } c \in \mathcal{M}_0^\alpha,$$

which is clearly a step function (*i.e.* a locally constant function) of $c \in \mathcal{M}_0^\alpha$.

Irrational case. To study the function $\mathcal{L}_\alpha(c)$ on \mathcal{M}^α as defined by (36), we will use a wavelet decomposition.

Let $g \in L^2$ be a function satisfying $\int g(x) dx = 0$ and $\int x g(x) dx = 0$. If it exists, the wavelet coefficient of $\mathcal{L}_\alpha(c) - c$ is

$$(60) \quad T(a, b) = \frac{1}{a} \int g\left(\frac{x-b}{a}\right) (\mathcal{L}_\alpha(x) - x) dx .$$

Plugging in (60) the expression

$$(61) \quad \mathcal{L}_\alpha(x) - x = \sum_{\substack{n \neq 0 \\ k}} \frac{\sin(n x T_\gamma)}{n T_\gamma} e^{i n d_k} c_k ,$$

one finds, supposing \hat{g} even

$$(62) \quad T(a, b) = \frac{1}{2i} \sum_{\substack{n \neq 0 \\ k}} \frac{1}{n T_\gamma} \hat{g}(a n) \sin(n b T_\gamma) c_k e^{i n d_k} .$$

The following result shows that such a decomposition is indeed valid.

Proposition 4.1. *Let g as before, \hat{g} being compactly supported and even, and let us suppose that φ is a compactly supported function satisfying*

$$(63) \quad \int \bar{\varphi}(a) \hat{g}(a) \frac{da}{a} = \int \bar{\hat{g}}(-a) \hat{\varphi}(-a) \frac{da}{a} = 1 .$$

Then, for all $c \in \mathcal{M}^\alpha$,

$$(64) \quad \mathcal{L}_\alpha(c) - c = \lim_{\varepsilon \rightarrow 0} \int_\varepsilon^{+\infty} \frac{da}{a} \int_{-\infty}^{+\infty} \varphi\left(\frac{c-b}{a}\right) T(a, b) db ,$$

where

$$(65) \quad T(a, b) = \frac{1}{2i} \sum_{\substack{n \neq 0 \\ k}} \frac{1}{n T_\gamma} \hat{g}(a n) \sin(n b T_\gamma) c_k e^{i n d_k} .$$

PROOF.

$$\begin{aligned}
 & \int_{\varepsilon}^{+\infty} \frac{da}{a} \int_{-\infty}^{+\infty} db \varphi\left(\frac{c-b}{a}\right) T(a, b) \\
 &= \int_{\varepsilon}^{+\infty} \frac{da}{a} \sum_{\substack{n \neq 0 \\ k}} \frac{1}{n T_{\gamma}} (\hat{\varphi}(a n) e^{inc T_{\gamma}} - \hat{\varphi}(-a n) e^{-inc T_{\gamma}}) \\
 & \quad \cdot e^{ind_k} \hat{g}(a n) c_k \\
 (66) \quad &= \int_{\varepsilon}^{+\infty} \frac{da}{a} \sum_{\substack{n \neq 0 \\ k}} \hat{\varphi}(a n) \hat{g}(a n) \sin(n c T_{\gamma}) e^{in((k+1/2)\theta+\alpha)} c_k \\
 &= \sum_{\substack{n \neq 0 \\ k}} \psi(\varepsilon n) \sin(n c T_{\gamma}) e^{ind_k} c_k,
 \end{aligned}$$

where

$$\psi(\varepsilon) := \int_{\varepsilon}^{+\infty} \frac{da}{a} \hat{\varphi}(a) \hat{g}(a).$$

Noting that $\psi'(a) = \hat{\varphi}(a) \hat{g}(a)/a$ is compactly supported and $\psi(0) = 1$ by hypothesis one get the result, thanks to Lemma 3.3.

The next result, thanks to the result of the Appendix will enable us to prove the Lipschitz continuity on \mathcal{M}^{α} .

Proposition 4.2.

$$(67) \quad T(a, b) = O(a), \quad \text{near } 0 \text{ almost everywhere and uniformly in } b.$$

PROOF. Since $\int x g(x) dx = \int g(x) dx = 0$, $g'(0) = 0$. So one can find a C^{∞} function f such that $\hat{g}(\xi) = \xi f(\xi)$ and $f(0) = 0$. Then

$$(68) \quad T(a, b) = a \sum_{\substack{n \neq 0 \\ k}} f(a n) \sin(b n T_{\gamma}) e^{ind_k} c_k,$$

and it is easy to check, by the same argument as in Lemma 3.3, that if $b \in \mathcal{M}^{\alpha}$,

$$\sum_{\substack{n \neq 0 \\ k}} f(a n) \sin(b n T_{\gamma}) e^{in((k+1/2)\theta+\alpha)} c_k$$

is bounded.

5. End of proofs.

The convergence statements in both theorems are immediate consequences of the Tauberian lemma of Section 2, applied to the following objects

$$(69) \quad \Upsilon_{\hbar}(a, c) = \sum_j w_j(\hbar) \varphi\left(\frac{E_j(\hbar) - E}{\hbar}\right),$$

where

$$(70) \quad w_j(\hbar) = |(\psi_{(x, \xi)}, \psi_j)|^2.$$

The weighted counting function is therefore

$$(71) \quad \sum_{\substack{j \\ |E_j(\hbar) - E| \leq c\hbar}} |(\psi_{(x, \xi)}, \psi_j)|^2.$$

The functional of the Tauberian lemma is

$$(72) \quad \mathcal{F}_0(\varphi) := c_0^\varphi(x, \xi)$$

as defined by (29). We must check that the above objects satisfy the assumptions of the Tauberian lemma.

a) Theorem 1.1. It is easy to see that the functional \mathcal{F}_0 defined where $c_0^\varphi(x, \xi)$ is defined by (29) satisfies the hypothesis 2 of the Tauberian Lemma of Section 2 if we take for \mathcal{M}^α the set defined by (9). Moreover the other hypotheses are satisfied as in [7]. Then just apply the Tauberian Lemma.

b) Theorem 1.2. The Lipschitz continuity of \mathcal{F}_0 is an immediate consequence of Proposition 4.2 together with Theorem A.1 below. The fact that \mathcal{M}^α is of full Lebesgue measure, is a classical result of Diophantine analysis (recall that the sequence $\{g_k\}$ in the remark 1, Section 3 is rapidly decreasing).

Appendix. Wavelets and Hölder continuity.

In this appendix we will prove an easy extension of results of [6], [5] and [4].

Let \mathcal{M}^α a bounded subset of \mathbb{R} of full Lebesgue measure.

Theorem A.1. *Let g be a continuously differentiable compactly supported function. Let f defined and bounded on \mathcal{M}^α . Let us suppose that f admits a “scale-space coefficient $T(a, b)$ ” decomposition with respect to g , namely*

$$(73) \quad f(x) = \int_0^\infty \int_{-\infty}^{+\infty} g\left(\frac{x-b}{a}\right) T(a, b) \frac{da}{a} db, \quad \text{for all } x \in \mathcal{M}^\alpha.$$

Let us suppose moreover that

$$(74) \quad T(a, b) = o(a^\alpha),$$

near 0 almost everywhere and uniformly in b . Then F is α -Hölder continuous on \mathcal{M}^α ; by this we mean

$$(75) \quad |f(x_1) - f(x_2)| = O_{x_1}(|x_2 - x_1|^\alpha), \quad \text{for all } x_1, x_2 \in \mathcal{M}^\alpha.$$

PROOF. The proof is absolutely equivalent to the one in [4], so we will only sketch it. Let us write first:

$$(76) \quad \begin{aligned} f(x) &= \left(\int_0^1 \frac{da}{a} + \int_1^\infty \frac{da}{a} \right) \int db g\left(\frac{x-b}{a}\right) T(a, b) \\ &= f_s(x) + f_l(x), \end{aligned}$$

f_l is obviously C^∞ . We concentrate on f_s .

Let $x_1, x_2 \in \mathcal{M}^\alpha$, $x_1 < x_2$, we cut f_s in three pieces.

$$(77) \quad \begin{aligned} f_s(x_1) - f_s(x_2) &= \int_0^{x_2-x_1} \frac{da}{a} \int db g\left(\frac{x_2-b}{a}\right) T(a, b) \\ &\quad - \int_0^{x_1-x_1} \frac{da}{a} \int db g\left(\frac{x_2-b}{a}\right) T(a, b) \\ &\quad + \int_{x_2-x_1}^1 da \int db \left(\frac{1}{a} g\left(\frac{x_2-b}{a}\right) - \frac{1}{a} g\left(\frac{x_1-b}{a}\right) \right) \\ &\quad \cdot T(a, b) \end{aligned}$$

$$(78) \quad =: T_1 - T_2 + T_3.$$

We now analyze each term:

- T_1 and T_2 . Since $T(a, b) = O(a^\alpha)$ almost everywhere, we have

$$(79) \quad \begin{aligned} |T_i| &= \int_0^{x_2-x_1} \frac{da}{a} \int db \left| \frac{1}{a} g\left(\frac{x_i-b}{a}\right) \right| C a^\alpha \\ &= O(|x_2-x_1|^\alpha) \|g\|_{L_1} \frac{C}{\alpha}. \end{aligned}$$

- T_3 . If g is continuously differentiable let us write

$$(80) \quad g\left(\frac{x_2-b}{a}\right) - g\left(\frac{x_1-b}{a}\right) = \frac{x_2-x_1}{a} g'\left(\frac{x'-b}{a}\right)$$

with $x_1 \leq x' \leq \bar{x}_2$. So

$$(81) \quad \begin{aligned} |T_3| &\leq \int_{x_2-x_1}^1 \frac{da}{a} \int db \left| \frac{1}{a^2} g'\left(\frac{x'-b}{a}\right) \right| |T(a, b)| |x_2-x_1| \\ &= O(|x_2-x_1|) \|g'\|_{L_1} \int_{x_2-x_1}^1 \frac{da}{a} a^{\alpha-1} \\ &= O(|x_2-x_1|^\alpha). \end{aligned}$$

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Thierry Paul
CEREMADE, URA 749 CNRS
Université Paris-Dauphine
Place de Lattre de Tassigny
75775 Paris Cedex 16, FRANCE
paulth@ceremade.dauphine.fr

and

Alejandro Uribe*
Mathematics Department
University of Michigan
Ann Arbor, Michigan 48109, U.S.A.
uribe@math.lsa.umich.edu

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