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A uniqueness result for the continuity equation in two dimensions

Dedicated to Constantine Dafermos on the occasion of his 70th birthday

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Abstract. We characterize the autonomous, divergence-free vector fields b on the plane such that the Cauchy problem for the continuity equation $\partial_t u + \operatorname{div}(bu) = 0$ admits a unique bounded solution (in the weak sense) for every bounded initial datum; the characterization is given in terms of a property of Sard type for the potential f associated to b. As a corollary we obtain uniqueness under the assumption that the curl of b is a measure. This result can be extended to certain nonautonomous vector fields b with bounded divergence.

Keywords. Continuity equation, transport equation, uniqueness of weak solutions, weak Sard property, disintegration of measures, coarea formula

1. Introduction

In this paper we consider the continuity equation

$$\partial_t u + \operatorname{div}(bu) = 0, \tag{1.1}$$

where $b:[0,T)\times\mathbb{R}^d\to\mathbb{R}^d$ is a bounded, time-dependent vector field on \mathbb{R}^d (with T possibly equal to $+\infty$). We recall that (1.1) is the equation satisfied by the density u of a continuous distribution of particles moving according to the velocity field b, that is, u(t,x) is the number of particles per unit volume at time t and position x, and the trajectory of each particle satisfies the ordinary differential equation

$$\dot{x} = b(t, x). \tag{1.2}$$

Throughout this paper, the vector field b will not be any more regular than just bounded and with bounded (distributional) divergence. Accordingly, solutions of the

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¹ In particular b is far from being Lipschitz in the space variable, which is the minimal regularity required to apply the method of characteristics to (1.1).

Cauchy problem for (1.1) are intended in the weak (or distributional) sense: a function $u:[0,T)\times\mathbb{R}^d\to\mathbb{R}$ solves (1.1) with the initial condition $u(0,\cdot)=u_0$ if

$$\int_{0}^{T} \int_{\mathbb{R}^{d}} (\partial_{t} \varphi + b \cdot \nabla \varphi) u \, dx \, dt + \int_{\mathbb{R}^{d}} \varphi(0, \cdot) u_{0} \, dx = 0$$

$$\forall \varphi \in C_{c}^{\infty}([0, T) \times \mathbb{R}^{d}).$$

$$(1.3)$$

Concerning the existence of solutions, it can be shown that if the divergence of b is bounded from below then, for every bounded initial datum u_0 , a solution of the Cauchy problem for (1.1) exists for all times in the future and is bounded for finite times.²

We focus therefore on the problem of uniqueness, and precisely on the following question: *Under which assumption on b does the Cauchy problem for* (1.1) *admit a unique bounded solution for every bounded initial datum* u_0 ?

In the fundamental paper [13], R. J. DiPerna and P.-L. Lions have proved uniqueness under the assumption that b is in the Sobolev class $W^{1,1}$ (locally in space), and later L. Ambrosio [4] extended this uniqueness result to vector fields b of class BV (locally in space).³

When b is a divergence-free autonomous vector field on the plane (d=2), uniqueness has been proved with the BV or Sobolev regularity of b replaced by various assumptions on the direction of b ([7, 17], see also [10, 8] for the bounded divergence case). In the opposite direction, N. Depauw [11] gave an example of a time-dependent, divergence-free, bounded vector field b on the plane for which there is no uniqueness (see also [1, 9]).

We can now turn to the main result of this paper. Let $b: \mathbb{R}^2 \to \mathbb{R}^2$ be a bounded, divergence-free, autonomous vector field on the plane; then b admits a Lipschitz potential $f: \mathbb{R}^2 \to \mathbb{R}$, that is, $b = \nabla^\perp f$, where $\nabla^\perp := (-\partial_2, \partial_1)$. In Theorem 4.7 we prove that the Cauchy problem for (1.1) admits a unique bounded solution for every bounded initial datum *if and only if* the potential f has what we call the *weak Sard property*.

We recall that a differentiable function $f: \mathbb{R}^2 \to \mathbb{R}$ has the (*strong*) Sard property if the image under f of the critical set S (the set of all x where $\nabla f(x) = 0$) is negligible, that is, $\mathcal{L}^1(f(S)) = 0$.⁴ The weak Sard property is a measure-theoretic version of the

² The bound div $b \ge -m$ implies that the Jacobian determinant J of the flux associated to (1.2) satisfies $J \ge e^{-mt}$, and therefore, according to the mechanical interpretation given above, the particle density u should satisfy $||u(t,\cdot)||_{\infty} \le e^{mt} ||u_0||_{\infty}$. This estimate can indeed be proved for smooth b, and then used to obtain solutions for nonsmooth b by approximation.

³ Both uniqueness results, besides being deeply interesting *per se*, had relevant applications to other problems, among which we mention the Boltzmann [14] and Vlasov–Poisson equations [15], and the Keyfitz–Kranzer system [5]. The mechanical interpretation given above suggests that uniqueness for (1.1) should be connected with uniqueness of solutions of (1.2) with initial condition $x(0) = x_0$ for a *generic* initial point x_0 , but not necessarily for every x_0 . This relation can be made rigorous in terms of uniqueness of the regular Lagrangian flow associated to (1.2) (see [4]), and explains why the regularity assumptions on b in the theorems of DiPerna–Lions and Ambrosio are distinctly weaker than those required for the uniqueness of (1.2) at every initial point.

⁴ Sard's theorem states that functions of class C^2 have the Sard property (see [20], or [18, Chapter 3, Theorem 1.3]), while this is not always the case for functions of class C^1 .

Sard property; the precise formulation requires a few additional definitions, and will be given in §2.13. Conditions on f (and therefore on b) that imply the weak Sard property are discussed in detail in §2.15; here we just recall that this property holds whenever b is non-zero a.e. (§2.15(i)), or is of class BV;⁵ indeed it suffices that the curl of b is a measure (§2.15(vi)).

We stress once again that while the results in the literature give only sufficient conditions for uniqueness, Theorem 4.7 gives a *necessary and sufficient* condition. Accordingly, the examples of functions without the weak Sard property constructed in the companion paper [2] immediately yield examples of divergence-free autonomous vector fields in the plane for which there is no uniqueness (Corollary 4.8). Note that the weak Sard property, like the Sard property, is completely nonlinear in character, and is satisfied by a generic divergence-free vector field (in the category sense, see §2.15(i)).

We conclude this introduction with an outline of the proof of Theorem 4.7. An essential ingredient is the description of (generic) level sets of Lipschitz functions on the plane given in [2]; we refer to that paper for a detailed discussion of all related measure-theoretic and real-analytic issues, and in particular for counterexamples.

The first step is a dimension-reduction argument (Theorem 3.10) that can be summarized as follows: a bounded function u solves the continuity equation (1.1) if and only if it solves a suitable one-dimensional variant of the same equation on *every* nontrivial connected component of the level set $E_y := f^{-1}(y)$ for *almost every* y. In other words, the problem of uniqueness for equation (1.1) is reduced to the uniqueness for the corresponding equations on the nontrivial connected components of a (generic) level set E_y .

The proof of Theorem 3.10 relies on the notion of disintegration of a measure with respect to the level sets of a function and on the coarea formula for Lipschitz functions, and it works in every space dimension. The rest of the proof of Theorem 4.7 is strictly two-dimensional.

The next key fact is that, for a.e. y, every connected component C of E_y is a *rectifiable simple curve* (as shown in [2]), and more precisely it is possible to choose an interval I and a Lipschitz parametrization $\gamma: I \to C$ so that, under the change of variable $x = \gamma(s)$, the equation on C becomes

$$\partial_t(u(1+\lambda)) + \partial_s u = 0, \tag{1.4}$$

where λ is a suitable singular measure on I. Note that, due to the particular choice of γ , the vector field b no longer appears in the equation.

Next we observe that the Cauchy problem for (1.4) admits a unique bounded solution for every bounded initial datum if and only if the measure λ is trivial. Indeed, if $\lambda = 0$

⁵ Thus we recover the uniqueness results in [7, 17] in full generality, and those in [4, 13] limited to our particular class of b.

⁶ The notion of "solving the equation on (a subset of) a level set" is properly defined in §3.6; a connected component is *nontrivial* if it contains more than one point. Note that this statement is a close relative of the method of characteristics: indeed the vector field b, being orthogonal to ∇f , is tangent to the level set E_y at a.e. point (with respect to the length measure) and for a.e. y. Hence these level sets are the proper replacement for characteristic curves.

⁷ The precise definition of λ and the weak formulation of (1.4) are given in Lemma 4.4.

equation (1.4) reduces to $\partial_t u + \partial_s u = 0$, for which uniqueness is well-known. To understand why the converse holds, assume that I is the real line and λ is a Dirac mass at 0; if u represents the density of a distribution of particles, then equation (1.4) means that each particle moves at constant speed 1 from left to right, except when it reaches the point 0, where it may stop for any given period of time. Therefore, if u_0 is an initial datum with support contained in $(-\infty, 0)$, a solution of (1.4) with initial condition $u(0, s) = u_0(s)$ is given by the function $u: [0, +\infty) \times \mathbb{R} \to \mathbb{R}$ defined by

$$u(t,s) := \begin{cases} u_0(s-t) & \text{if } s \neq 0, \\ 0 & \text{if } s = 0 \end{cases}$$

(no particle stops at 0), while another solution is

$$\tilde{u}(t,s) := \begin{cases} u_0(s-t) & \text{if } s < 0, \\ \int_{-t}^0 u_0(\tau) \, d\tau & \text{if } s = 0, \\ 0 & \text{if } s > 0 \end{cases}$$

(all particles stop at 0).

We conclude that uniqueness for (1.1) holds if and only if $\lambda=0$ for every nontrivial connected component C of a.e. level set E_y , which in turn is equivalent to the weak Sard property (Lemma 2.14). Without entering into details, we just mention that the connection between the weak Sard property, which is related to the critical set S, and the measures λ lies in the fact that these measures are given by the disintegration (with respect to the level sets of f) of the restriction of Lebesgue measure to $S \cap E^*$, where E^* is the union of all nontrivial connected components of all levels sets of f.

As one might expect, Theorem 4.7 can be extended in many different ways. In Theorem 5.2 we consider the case of a vector field on the plane of the form $b = a \nabla^{\perp} f$, where f is a Lipschitz function as above and a a scalar function depending also on time. Further extensions are mentioned in §6.1. It is indeed conceivable to apply the strategy outlined above even in higher space dimension (see §6.2).

This paper is organized as follows: Section 2 contains the measure-theoretic and real-analytic results used in the rest of the paper: the disintegration of a measure with respect to the level sets of a function, the coarea formula for Lipschitz functions, the description of connected components of the level sets of a Lipschitz function on the plane, the definition of weak Sard property and some related results. Section 3 is devoted to the dimension-reduction argument (Theorem 3.10); Sections 4 and 5 contain the proof of the main result (Theorem 4.7) and of the generalization mentioned above (Theorem 5.2). Section 6 contains some additional remarks, and finally the Appendix contains a measurable selection lemma used in the proof of Theorem 3.10.

2. Measure-theoretic preliminaries

We begin this section by recalling the notion of disintegration of a measure, focusing in particular on the disintegration of the Lebesgue measure with respect to Lipschitz maps (§2.7 and Lemma 2.8). We then restrict our attention to Lipschitz functions on the plane (Lemma 2.11); in this setting we introduce the notion of weak Sard property (§2.13), prove a characterization of this property in terms of disintegration of the Lebesgue measure (Lemma 2.14), and list some conditions which imply it (§2.15). We conclude with a change-of-variable formula used in the proof of Lemma 4.5.

2.1. Basic notation. In this paper we follow the standard notation of measure theory. Sets and functions are always assumed to be Borel measurable, and measures are always defined on the Borel σ -algebra of some locally compact, separable metric space X. Unless otherwise stated, a measure on X is assumed to be positive and locally finite. We write $\lambda \ll \mu$ when the measure λ is absolutely continuous with respect to the measure μ , and $\lambda \perp \mu$ when λ is singular with respect to μ . We say that the measure μ is supported on the set E if $\mu(X \setminus E) = 0$.

Given a function $\rho: X \to [0, +\infty]$, we denote by $\rho \cdot \mu$ the measure on X defined by

$$(\rho \cdot \mu)(E) := \int_{E} \rho \, d\mu$$

for every Borel set E contained in X. Given a set A contained in X, we write $\mathbf{1}_A: X \to \{0, 1\}$ for the characteristic function of A, and therefore $\mathbf{1}_A \cdot \mu$ is the restriction of μ to A.

Given a metric space Y and a map $f: X \to Y$, the push-forward of μ by f is the measure $f_\#\mu$ on Y defined by

$$[f_{\#}\mu](E) := \mu(f^{-1}(E))$$

for every Borel set E contained in Y. Thus

$$\int_{Y} \varphi \, d(f_{\#}\mu) = \int_{X} \varphi \circ f \, d\mu$$

for every function $\varphi: Y \to [0, +\infty]$, and therefore also for every real- or vector-valued φ in $L^1(f_\#\mu)$.

As usual, \mathscr{L}^d is the Lebesgue measure on (domains in) \mathbb{R}^d while \mathscr{H}^d is the d-dimensional Hausdorff measure on every metric space—the usual d-dimensional volume for d-dimensional surfaces of class C^1 in some Euclidean space. When the measure is not specified, it is assumed to be the Lebesgue measure, and we often write $\int g(x) \, dx$ instead of $\int g \, d\mathscr{L}^d$.

2.2. Borel families of measures. Let Y be a metric space and $\{\mu_y : y \in Y\}$ a family of measures on a locally compact, separable metric space X. We say that such a family, or more precisely the map $y \mapsto \mu_y$, is Borel if the function

$$y \mapsto \int_{X} \varphi \, d\mu_{y} \tag{2.1}$$

⁸ Recall that $\lambda \ll \mu$ if $\lambda(E) = 0$ whenever $\mu(E) = 0$, or equivalently (in the context of this paper) when λ can be written as $\lambda = \rho \cdot \mu$ for a suitable density ρ ; λ is singular with respect to μ if it is supported on a μ -negligible set, or equivalently if λ and μ are supported on disjoint sets.

⁹ Note that E need not be closed, and does not necessarily contain the support of μ (defined in the usual distributional sense).

is Borel for every test function $\varphi: X \to \mathbb{R}$ which is continuous and compactly supported. It then follows that the function in (2.1) is Borel measurable also for every Borel test function φ on $X \times Y$ such that the integral at the right-hand side of (2.1) is well-defined.

- **2.3. Disintegration of measures.** Let X and Y be locally compact, separable metric spaces, μ a measure on X, $f: X \to Y$ a Borel map, and ν a measure on Y such that $f_\#\mu \ll \nu$. Then there exists a Borel family $\{\mu_y: y \in Y\}$ of measures on X such that
- (i) μ_y is supported on the level set $E_y := f^{-1}(y)$ for every $y \in Y$;
- (ii) μ can be decomposed as $\mu = \int_{V} \mu_{V} d\nu(y)$, which means that

$$\mu(A) = \int_{Y} \mu_{y}(A) d\nu(y) \tag{2.2}$$

for every Borel set A contained in X.

Any family $\{\mu_y\}$ satisfying (i) and (ii) is called a *disintegration* of μ with respect to f and ν . The disintegration is unique in the sense that given another disintegration $\{\tilde{\mu}_y\}$ there holds $\mu_{\nu} = \tilde{\mu}_{\nu}$ for ν -a.e. $\nu \in Y$.

The existence and uniqueness of the disintegration is a standard result in case X is compact, μ is finite, and $\nu := f_{\#}\mu$ (see for instance [12, Chapter III, §70 and §71]). The statement given above can be easily derived from this particular case.

- **2.4. Properties of disintegration.** We list here some general properties of disintegration that will be used (often tacitly) throughout the paper.
 - (i) Formula (2.2) implies that

$$\int_{X} \varphi \, d\mu = \int_{Y} \left[\int_{E_{y}} \varphi \, d\mu_{y} \right] d\nu(y) \tag{2.3}$$

for every Borel function $\varphi: X \to [0, +\infty]$.

If φ is a real-valued or vector-valued function in $L^1(\mu)$, by applying identity (2.3) with $|\varphi|$ in place of φ we find that φ belongs to $L^1(\mu_y)$ for ν -a.e. y, and the function $y \mapsto \int |\varphi| \, d\mu_y$ is in $L^1(\nu)$. Thus both sides of (2.3) make sense, and the equality holds even for such φ .

- (ii) If A is a set in X with $\mu(A) < \infty$, formula (2.2) implies that $\mu_y(A)$ is finite for ν -a.e. y. In particular, if μ is finite then μ_y is finite for ν -a.e. y.
- (iii) Formula (2.2) shows that a set A in X is μ -negligible if and only if it is μ_y -negligible for ν -a.e. y. We infer the following: (a) $\mu(f^{-1}(N)) = 0$ for every set N such that $\nu(N) = 0$; (b) μ is supported on a set F if and only if μ_y is supported on $F \cap E_y$ for ν -a.e. y; (c) if P(x) is a proposition that depends on the variable $x \in X$, then P(x) holds for μ -a.e. x if and only if it holds for μ_y -a.e. x and ν -a.e. y.

¹⁰ This is equivalent to the notion of Borel measurability for maps with values in the space of locally finite measures on X, when the latter has been endowed with the weak* topology (as dual of the space of continuous functions with compact support).

(iv) If μ' is a measure on the metric space X', we write $\mu \otimes \mu'$ for the product measure on $X \times X'$ and we consider the function $\tilde{f}: X \times X' \to Y$ defined by $\tilde{f}(x, x') = f(x)$. Then the disintegration of $\mu \otimes \mu'$ with respect to \tilde{f} and ν is $\{\mu_{\nu} \otimes \mu'\}$.

2.5. Lipschitz maps. For the rest of this section d, k are integers such that $0 \le k < d$, and $f : \mathbb{R}^d \to \mathbb{R}^{d-k}$ is a Lipschitz map. For every $y \in \mathbb{R}^{d-k}$ we write E_y for the level set $f^{-1}(y)$.

By the Rademacher theorem f is differentiable at almost every point in \mathbb{R}^d , and at all such points we define the Jacobian

$$J := [\det(\nabla f \cdot \nabla^t f)]^{1/2}.$$

We define the *critical set* of f to be the set S of all points in \mathbb{R}^d where either f is not differentiable or J=0, that is, the rank of ∇f is strictly less than d-k.

A connected component C of E_y is a connected subset of E_y which is maximal with respect to inclusion. Notice that every such connected component is a closed set. For every $y \in \mathbb{R}^{d-k}$ we denote by \mathscr{C}_y the family of all connected components C of E_y such that $\mathscr{H}^k(C) > 0$; we then define E_y^* as the union of all C in \mathscr{C}_y , and E^* as the union of all E_y^* with $y \in \mathbb{R}^{d-k}$. Note that the sets E^* and E_y^* are Borel (see [2, Proposition 6.1]).

2.6. Lemma. In the context of the previous subsection, the following statements hold:

- (i) if $\mu := J \cdot \mathcal{L}^d$ then $f_{\#}\mu$ is absolutely continuous with respect to \mathcal{L}^{d-k} and its disintegration with respect to f and \mathcal{L}^{d-k} is $\mu_y := \mathbf{1}_{E_y} \cdot \mathcal{H}^k$;
- (ii) for a.e. $y \in \mathbb{R}^{d-k}$ we have $\mathcal{H}^k(E_y \cap S) = 0$, and for every bounded set A in \mathbb{R}^d we have $\mathcal{H}^k(A \cap E_y) < +\infty$ and

$$\int_{A\cap E_{\nu}}\frac{1}{J}\,d\mathcal{H}^{k}<+\infty;$$

- (iii) if $\tilde{\mu} := \mathbf{1}_{\mathbb{R}^d \setminus S} \cdot \mathcal{L}^d$ then $f_{\#}\tilde{\mu}$ is absolutely continuous with respect to \mathcal{L}^{d-k} , and its disintegration with respect to f and \mathcal{L}^{d-k} is given by $\tilde{\mu}_{\gamma} := (1/J)\mathbf{1}_{E_{\gamma}} \cdot \mathcal{H}^k$;
- (iv) the family \mathscr{C}_y is countable and $\mathscr{H}^k(E_y \setminus E_y^*) = 0$ for \mathscr{L}^{d-k} -a.e. y;
- (v) $\mathcal{L}^d(\mathbb{R}^d \setminus (E^* \cup S)) = 0.$

Proof. Statement (i) is just a reformulation of the coarea formula for Lipschitz maps (see for instance [16, §3.2.11], [21, §10], or [19, Corollary 5.2.6]).

Using statement (i) and formula (2.2) we obtain

$$0 = \mu(S) = \int_{\mathbb{R}^{d-k}} \mu_{y}(S) \, dy = \int_{\mathbb{R}^{d-k}} \mathscr{H}^{k}(E_{y} \cap S) \, dy,$$

and therefore $\mathcal{H}^k(E_v \cap S) = 0$ for a.e. y. Similarly,

$$+\infty > \mu(A) = \int_{\mathbb{R}^{d-k}} \mu_y(A) \, dy = \int_{\mathbb{R}^{d-k}} \mathscr{H}^k(A \cap E_y) \, dy$$

implies that $\mathscr{H}^k(E_y \cap A)$ is finite for a.e. y. The last part of the statement follows by applying (2.3) with $\varphi := (1/J)\mathbf{1}_{A\setminus S}$:

$$+\infty > \mathscr{L}^d(A \setminus S) = \int \varphi \, d\mu = \int_{\mathbb{R}^{d-k}} \left[\int_{E_y} \varphi \, d\mu_y \right] dy = \int_{\mathbb{R}^{d-k}} \left[\int_{A \cap E_y} \frac{1}{J} \, d\mathscr{H}^k \right] dy.$$

Similarly, statement (iii) follows by applying (2.3) with $\varphi := (1/J)\mathbf{1}_{A\setminus S}$ and A an arbitrary Borel set in \mathbb{R}^d .

Statement (iv) is proved in [2, Theorem 2.5(iii)].

Statement (v) follows from statements (iii) and (iv):

$$\mathcal{L}^{d}(\mathbb{R}^{d} \setminus (E^{*} \cup S)) = \tilde{\mu}(\mathbb{R}^{d} \setminus E^{*}) = \int_{\mathbb{R}^{d-k}} \tilde{\mu}_{y}(\mathbb{R}^{d} \setminus E^{*}) \, dy$$
$$= \int_{\mathbb{R}^{d-k}} \left[\int_{E_{y} \setminus E_{y}^{*}} \frac{1}{J} \, d\mathcal{H}^{k} \right] dy = 0. \qquad \Box$$

- **2.7. Disintegration of Lebesgue measure.** We take f as above and choose a measure v^s on \mathbb{R}^{d-k} so that v^s is singular with respect to \mathcal{L}^{d-k} and $f_\#\mathcal{L}^d \ll v$ with $v := \mathcal{L}^{d-k} + v^s$. For the rest of this section we denote by $\{\mu_y\}$ the disintegration of \mathcal{L}^d with respect to f and v.
- **2.8. Lemma.** In the context of the previous subsection, the following statements hold:
- (i) for v^s -a.e. y the measure μ_v is supported on $S \cap E_v$;
- (ii) for \mathcal{L}^{d-k} -a.e. y the measure μ_{ν} can be decomposed as

$$\mu_{y} = (1/J)\mathbf{1}_{E_{y}\backslash S}\cdot \mathscr{H}^{k} + \mu_{y}^{s}$$

with μ_y^s a measure supported on $E_y \cap S$; moreover μ_y^s is singular with respect to $(1/J)\mathbf{1}_{E_y \setminus S} \cdot \mathcal{H}^k$, and the latter measure agrees with $(1/J)\mathbf{1}_{E_y} \cdot \mathcal{H}^k$.

Proof. Note that \mathcal{L}^d can be written as the sum of $\mu^1 := \mathbf{1}_S \cdot \mathcal{L}^d$ and $\mu^2 := \mathbf{1}_{\mathbb{R}^d \setminus S} \cdot \mathcal{L}^d$, and therefore its disintegration (with respect to f and ν) is obtained by summing the disintegrations of μ^1 and μ^2 .

Now μ^1 is supported on S and then μ^1_y is supported on $S \cap E_y$ for ν -a.e. y; on the other hand, Lemma 2.6(iii) implies that $\mu^2_y = 0$ for ν^s -a.e. y and $\mu^2_y = (\mathbf{1}_{E_y}/J) \cdot \mathscr{H}^k$ for \mathscr{L}^{d-k} -a.e. y. Putting together these facts we immediately obtain statement (i) and the first part of (ii).

Finally, μ_y^s is singular with respect to $(\mathbf{1}_{E_y\setminus S}/J)\cdot \mathcal{H}^k$ because these measures are supported on the disjoint sets $E_y\setminus S$ and $E_y\cap S$. The latter measure agrees with $(\mathbf{1}_{E_y}/J)\cdot \mathcal{H}^k$ because $\mathcal{H}^k(E_y\cap S)=0$ (Lemma 2.6(ii)).

2.9. Lipschitz functions on the plane. For the rest of this section we assume d=2 and k=1, that is, f is a Lipschitz function on \mathbb{R}^2 . Therefore $J=|\nabla f|$, S is the set of all points where either f is not differentiable or $\nabla f=0$, and \mathscr{C}_y is the family of all

nontrivial connected components of the level set E_y , namely those which contain more than one point.

2.10. Simple curves. We denote by $[0, L]^*$ the space obtained by identifying the endpoints of the interval [0, L], endowed with the distance

$$d(x, y) := |x - y| \wedge (L - |x - y|).$$

The canonical measure on $[0, L]^*$ is the Lebesgue measure, which we denote simply by \mathcal{L} .

We say that a set C in \mathbb{R}^d is a closed, simple curve with finite length if there exist L > 0 and a Lipschitz bijection $\gamma : [0, L]^* \to C$. We call any such γ a parametrization of C.

- **2.11. Lemma.** Assume that f has compact support, and take μ_y , μ_y^s as in Lemma 2.8. The following statements hold for \mathcal{L}^1 -a.e. $y \in \mathbb{R}$ and every connected component $C \in \mathcal{C}_y$:
- (i) *C* is a closed simple curve with finite length;
- (ii) there exists a parametrization $\gamma:[0,L]^* \to C$ such that $\dot{\gamma} = \nabla^{\perp} f(\gamma) \neq 0$ a.e. (recall that $\nabla^{\perp}:=(-\partial_2,\partial_1)$);
- (iii) the push-forward of \mathcal{L} according to γ is $(1/J)\mathbf{1}_C \cdot \mathcal{H}^1$;
- (iv) there exists a measure λ on $[0, L]^*$ such that λ is singular with respect to \mathcal{L} , and the push-forward of λ according to γ is the restriction of μ_{ν}^s to C;
- (v) the push-forward of $\mathcal{L} + \lambda$ according to γ is the restriction of μ_{ν} to C.
- **2.12. Remark.** (i) The assumption that f has compact support was made for the sake of simplicity, and it can be easily removed. In that case statement (i) should be modified, allowing for connected components that are simple curves with end points at infinity.
- (ii) It is always possible to choose the parametrization γ in such a way that the measure λ has no atom at 0, that is, $\lambda(\{0\}) = 0$. This assumption makes certain proofs much simpler, and will be tacitly made in the following.

Proof of Lemma 2.11. By Lemma 2.8(ii) we can assume that $\mathcal{H}^1(S \cap C) = 0$ and

$$\int_{C} \frac{1}{J} d\mathcal{H}^{1} < +\infty. \tag{2.4}$$

Then statement (i) is proved in [2, Theorem 2.5(iv)].

The same theorem shows that C admits a parametrization $\varphi:[0,L']^*\to C$ such that

$$J(\varphi) \neq 0$$
 and $\dot{\varphi} = \frac{\nabla^{\perp} f(\varphi)}{|\nabla f(\varphi)|} = \frac{\nabla^{\perp} f(\varphi)}{J(\varphi)}$ a.e. (2.5)

Thus the parametrization in (ii) is given by $\gamma := \varphi \circ \tau^{-1}$ where $\tau : [0, L'] \to [0, L]$ is the function defined by

$$\tau(t) := \int_0^t \frac{1}{J(\varphi(s))} \, ds \quad \text{and} \quad L := \int_0^{L'} \frac{1}{J(\varphi(s))} \, ds.$$

The definition of τ is well-posed because the integral that defines L is finite, and more precisely it agrees with the integral on the left-hand side of (2.4) (because $|\dot{\varphi}| = 1$ a.e.). Since $J \leq m$, where m denotes the Lipschitz constant of f, we have

$$\tau(t_2) - \tau(t_1) \ge \frac{1}{m}(t_2 - t_1)$$

whenever $t_1 < t_2$, which implies that τ^{-1} is a Lipschitz function, and therefore so is γ . Moreover τ is of class $W^{1,1}$ and $\dot{\tau} = 1/J(\varphi) \neq 0$ a.e.; using this fact one obtains

$$(\tau^{-1})' = J(\gamma) = |\nabla f(\gamma)|$$
 a.e.

This equation and (2.5) yield $\dot{\gamma} = \nabla^{\perp} f(\gamma)$ a.e., concluding the proof of statement (ii). Statement (iii) follows by the fact that $|\dot{\gamma}| = |\nabla f(\gamma)| = J(\gamma)$ a.e.

To prove statement (iv), let λ be the push-forward according to γ^{-1} of $\mathbf{1}_C \cdot \mu_y^s$. In order to prove that $\lambda \perp \mathscr{L}$, note that μ_y^s is supported on $S \cap C$, and therefore λ is supported on $\gamma^{-1}(S \cap C)$. Moreover $S \cap C$ is \mathscr{H}^1 -negligible, and therefore $\gamma^{-1}(S \cap C) = \tau(\varphi^{-1}(S \cap C))$ is \mathscr{L} -negligible because both φ^{-1} and τ map negligible sets to negligible sets.

Statement (v) follows from statements (iii) and (iv).

2.13. The weak Sard property. Let $f: \mathbb{R}^2 \to \mathbb{R}$ be a Lipschitz function. ¹¹ We say that f has the *weak Sard property* if the push-forward according to f of the restriction of the Lebesgue measure to $S \cap E^*$ (where S and E^* have been defined in §2.5) is singular, that is.

$$f_{\#}(\mathbf{1}_{S\cap E^*}\cdot\mathscr{L}^2)\perp\mathscr{L}^1.$$
 (2.6)

- **2.14. Lemma.** The following statements are equivalent:
- (a) f has the weak Sard property;
- (b) $\mu_{\mathbf{y}}^{s}(E_{\mathbf{y}}^{*}) = 0$ for \mathcal{L}^{1} -a.e. $\mathbf{y} \in \mathbb{R}$;
- (c) $\mu_{\nu}^{s}(C) = 0$ for \mathcal{L}^{1} -a.e. $y \in \mathbb{R}$ and every $C \in \mathcal{C}_{\nu}$.

Proof. Take ν as in §2.7 and μ_y^s as Lemma 2.8. Set $\tilde{\mu} := \mathbf{1}_{S \cap E^*} \cdot \mathcal{L}^2$, and let $\{\tilde{\mu}_y\}$ be the disintegration of $\tilde{\mu}$ with respect to f and ν . Thus (2.6) can be restated as

$$\tilde{\mu}_{y} = 0$$
 for \mathcal{L}^{1} -a.e. y.

On the other hand Lemma 2.8(ii) implies that $\tilde{\mu}_y = \mathbf{1}_{S \cap E^*} \cdot \mu_y^s$ for \mathcal{L}^1 -a.e. y, and therefore $\tilde{\mu}_y = 0$ if and only if $0 = \mu_y^s(S \cap E^*) = \mu_y^s(E_y^*)$ (for the second equality use that μ_y^s is supported on $E_y \cap S$ and $E_y \cap E^* = E_y^*$). We have thus proved that (2.6) holds if and only if $\mu_y^s(E_y^*) = 0$ for \mathcal{L}^1 -a.e. y, that is, statements (a) and (b) are equivalent.

The equivalence of (b) and (c) is immediate because E_y^* is the union of all $C \in \mathcal{C}_y$, and \mathcal{C}_y is countable for a.e. y (Lemma 2.6(iv)).

The weak Sard property can be defined in the more general context of Lipschitz maps from \mathbb{R}^d into \mathbb{R}^{d-k} , but only the case d=2, k=1 is relevant to this paper.

2.15. On the weak Sard property. We list here some conditions on f which imply the weak Sard property. Indeed all of them imply the slightly stronger property

$$f_{\#}(\mathbf{1}_{S} \cdot \mathcal{L}^{2}) \perp \mathcal{L}^{1}.$$
 (2.7)

Since in Theorem 4.7 the function f is the potential associated to a given vector field b by the relation $\nabla^{\perp} f = b$, we will express the following conditions also in terms of b.

- (i) The weak Sard property is implied by the condition $\mathcal{L}^2(S \cap E^*) = 0$, which in turn is implied by $\nabla f \neq 0$ a.e., or, equivalently, $b \neq 0$ a.e. Note that the set of all vector fields b that satisfy the last condition is residual 12 in the Banach space X of divergence-free, bounded vector fields endowed with the L^{∞} norm (this is a reformulation of [2, Proposition 4.10]). In particular the set of all b whose potential f has the weak Sard property is residual in X.
- (ii) The (strong) Sard property implies the weak Sard property. Indeed the measure $f_{\#}(\mathbf{1}_{S\cap E^*}\cdot \mathscr{L}^2)$ is supported on the set f(S), and if this set is negligible (Sard property), then the measure must be singular (weak Sard property).
- (iii) A function f on \mathbb{R}^2 is said to have the *Lusin property of order* k if there exists a sequence of functions f_n of class C^k on \mathbb{R}^2 and Borel sets A_n such that $f = f_n$ on A_n , and the sets A_n cover almost all of \mathbb{R}^2 . A Lipschitz function with the Lusin property of order 2 satisfies (2.7) and therefore has the weak Sard property.

Let indeed S_n be the critical set of f_n . Then $S \cap A_n$ is contained in $S_n \cap A_n$ up to a negligible set, and S is contained in the union of $S_n \cap A_n$ up to a negligible set. Hence the measure $\mathbf{1}_{S \cap E^*} \cdot \mathscr{L}^2$ is supported on the union of $S_n \cap A_n$, and its push-forward according to f is supported on the union of all $f(S_n \cap A_n) = f_n(S_n \cap A_n)$, which is negligible because it is contained in the union of the negligible sets $f_n(S_n)$.

(iv) Following [23, §3.5.5] we say that f admits an L^1 -Taylor expansion of order 2 at x if $f(x+h) = P_x(h) + R_x(h)$ for every h, where P_x is a polynomial of degree 2, and the remainder R_x satisfies

$$\int_{B(r)} |R_x(h)| \, dh = o(r^2)$$

(the barred integral stands for the average over the ball B(r) with centre 0 and radius r). If f admits an L^1 -Taylor expansion of order 2 at a.e. point, then it has the Lusin property of order 2 (use the L^p -version of Whitney extension theorem given in [23, Theorem 3.6.3]), and therefore also the weak Sard property.

- (v) If f is locally in the Sobolev class $W^{2,1}$ (or, equivalently, b is locally in the class $W^{1,1}$) then f admits an L^1 -Taylor expansion of order 2 at a.e. point [23, Theorem 3.4.2], and therefore it has the weak Sard property.
- (vi) The proof of [23, Theorem 3.4.2] can be easily modified to show that f admits an L^1 -Taylor expansion of order 2 at a.e. point—which implies the weak Sard property—even when the second order distributional derivative of f is locally a vector-valued measure (that is, when b is locally of class BV). In [3, Proposition 4.4] we prove that the same

¹² A set in a topological space is *residual* if it contains a countable intersection of open dense sets.

¹³ Thus Lusin's theorem states that a Borel function has the Lusin property of order 0.

conclusion holds under the weaker assumption that the distributional Laplacian of f (that is, the curl of b) is locally a signed measure.

(vii) In [2, Section 4] we construct a function on the plane without the weak Sard property and of class $C^{1,\alpha}$ for every $\alpha < 1$.

We conclude this section with a particular change of variable that will be used in the proof of Lemma 4.5.

2.16. A particular change of variable. In what follows, I is the interval [0, L], \mathcal{L} is the Lebesgue measure on I, λ is a measure on I which is singular with respect to \mathcal{L} , and \mathcal{A} is the set of atoms of λ (points with positive measure). We set

$$\widehat{L} := (\mathcal{L} + \lambda)(I)$$
 and $\widehat{I} := [0, \widehat{L}],$

and denote by $\widehat{\mathscr{L}}$ the Lebesgue measure on \widehat{I} .

We denote by $\hat{\sigma}$ the multi-valued function from I to \widehat{I} that to every $s \in I$ associates the interval

$$\hat{\sigma}(s) := [\hat{\sigma}_{-}(s), \hat{\sigma}_{+}(s)]$$

where

$$\hat{\sigma}_{-}(s) := (\mathcal{L} + \lambda)([0, s))$$
 and $\hat{\sigma}_{+}(s) := (\mathcal{L} + \lambda)([0, s]),$

and define $\sigma: \widehat{I} \to I$ as the inverse of $\widehat{\sigma}$.

Then $\hat{\sigma}$ is surjective on I, strictly increasing, and single-valued for every $s \notin \mathcal{A}$ because $\hat{\sigma}_-$ and $\hat{\sigma}_+$ are strictly increasing, and $\hat{\sigma}_-(s) = \hat{\sigma}_+(s)$ whenever $s \notin \mathcal{A}$. Moreover

$$s_2 - s_1 \le \hat{s}_2 - \hat{s}_1 \tag{2.8}$$

for every $s_1, s_2 \in I$ with $s_1 < s_2$, and every $\hat{s}_1 \in \hat{\sigma}(s_1), \hat{s}_2 \in \hat{\sigma}(s_2)$.

Accordingly, σ is surjective from \widehat{I} onto I, single-valued and 1-Lipschitz because of (2.8), constant on the interval $\sigma(s)$ for every $s \in \mathscr{A}$ and strictly increasing at every point outside $\sigma(\mathscr{A})$.

- **2.17. Lemma.** Let F be an \mathcal{L} -negligible set in I which supports the measure λ , and set $\widehat{F} := \widehat{\sigma}(F)$. Then:
- (i) the push-forward of the measure $\widehat{\mathcal{L}}$ according to σ is $\mathcal{L} + \lambda$, and this means that for every function $g: I \to [-\infty, +\infty]$ which is either positive or belongs to $L^1(\mathcal{L} + \lambda)$ we have

$$\int_{I} g \, d(\mathcal{L} + \lambda) = \int_{\widehat{I}} g \circ \sigma \, d\widehat{\mathcal{L}}; \qquad (2.9)$$

(ii) the derivative of σ agrees with $\mathbf{1}_{\widehat{I}\setminus\widehat{F}}$ a.e. in \widehat{I} .

Proof. Let $E := (s_1, s_2)$ be an open interval in I. By the definition of σ and $\hat{\sigma}$ we have

$$\sigma^{-1}(E) = \hat{\sigma}(E) = (\hat{\sigma}_{+}(s_1), \hat{\sigma}_{-}(s_2))$$

and therefore

$$\widehat{\mathscr{L}}(\sigma^{-1}(E)) = \widehat{\sigma}_{-}(s_2) - \widehat{\sigma}_{+}(s_1) = (\mathscr{L} + \lambda)(E).$$

Then a standard argument shows that the identity $\widehat{\mathscr{L}}(\sigma^{-1}(E)) = (\mathscr{L} + \lambda)(E)$ holds for every open set E, and hence also for every Borel set E. Thus (i) is proved.

In order to prove (ii), we first notice that by the choice of F we have $\mathbf{1}_F \cdot (\mathcal{L} + \lambda) = \lambda$ and $\mathbf{1}_{I \setminus F} \cdot (\mathcal{L} + \lambda) = \mathcal{L}$, and therefore statement (i) yields

$$\sigma_{\#}(\mathbf{1}_{\widehat{F}}\cdot\widehat{\mathscr{L}}) = \mathscr{L} \quad \text{and} \quad \sigma_{\#}(\mathbf{1}_{\widehat{I}\setminus\widehat{F}}\cdot\widehat{\mathscr{L}}) = \lambda.$$
 (2.10)

Since λ is singular with respect to \mathscr{L} , it is well known that the density of λ with respect to \mathscr{L} is equal to 0 at \mathscr{L} -a.e. point, and equal to $+\infty$ at λ -a.e. point. This means that the derivatives of the monotone functions $\hat{\sigma}_-$ and $\hat{\sigma}_+$ are equal to 1 at \mathscr{L} -a.e. point, and equal to $+\infty$ at λ -a.e. point. Therefore, taking into account (2.10), we conclude that the derivative of σ is equal to 1 at $\widehat{\mathscr{L}}$ -a.e. point in $\widehat{I} \setminus \widehat{F}$, and equal to 0 at $\widehat{\mathscr{L}}$ -a.e. point in \widehat{F} .

- **2.18. Remark.** The definitions of σ and $\hat{\sigma}$ and Lemma 2.17 can be extended with few obvious modification to the case where the intervals [0, L] and $[0, \widehat{L}]$ are replaced by the quotients $[0, L]^*$ and $[0, \widehat{L}]^*$ (see §2.10).
- **2.19. List of notation.** We recall here the notation that will be frequently used in the rest of the paper. Given a Lipschitz map $f : \mathbb{R}^d \to \mathbb{R}^{d-k}$ we have defined:

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J Jacobian of f (§2.5); S critical set of f (§2.5); E_y := f^{-1}(y) with y \in \mathbb{R}^{d-k}, level set of f (§2.5); \mathscr{C}_y collection of connected components C of E_y with \mathscr{H}^k(C) > 0 (§2.5); E^* union of all C with C \in \mathscr{C}_y and y \in \mathbb{R}^{d-k} (§2.5); v := \mathscr{L}^{d-k} + v^s, a measure such that f_\#\mathscr{L}^d \ll v (§2.7); \mu_y disintegration of \mathscr{L}^d with respect to f and v corresponding to the level set E_y (§2.7); \mu_y^s disintegration of the restriction of \mathscr{L}^d to S with respect to f and v (Lemma 2.8).
```

In the specific case of a Lipschitz function $f: \mathbb{R}^2 \to \mathbb{R}$ we have:

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J = |\nabla f| (§2.9);

S set of points where f is not differentiable or \nabla f = 0 (§2.9);

\mathscr{C}_y collection of nontrivial connected components of E_y (§2.9);

\gamma : [0, L]^* \to C, parametrization of C \in \mathscr{C}_y given in Lemma 2.11;

\mathscr{L} Lebesgue measure on [0, L]^* (§2.10);

\lambda measure on [0, L]^* defined in Lemma 2.11.
```

3. Dimension reduction

We now begin to consider the question of uniqueness for the Cauchy problem for the continuity equation (1.1). In this section, we still work in general space dimension. The existence of a Lipschitz function associated to the vector field as in §3.1 gives the possi-

bility to relate the uniqueness for the Cauchy problem in \mathbb{R}^d to the uniqueness for suitably defined Cauchy problems on the level sets of that Lipschitz function (Theorem 3.10).

3.1. Assumptions on the vector field. In this section $b:[0,T)\times\mathbb{R}^d\to\mathbb{R}^d$ is a bounded vector field and we assume that there exists a Lipschitz function $f:\mathbb{R}^d\to\mathbb{R}^{d-k}$ with compact support such that

$$\nabla f(x) \cdot b(t, x) = 0 \quad \text{for a.e. } (t, x) \in [0, T) \times \mathbb{R}^d. \tag{3.1}$$

We take S, E_y , μ_y and so on as in Section 2 (see §2.19). Further assumptions on b and f will be introduced when needed.

3.2. The uniqueness problem. By linearity, the uniqueness for the Cauchy problem for (1.1) is equivalent to the fact that every solution in the sense of distributions with initial value $u_0 = 0$ is trivial, that is, a.e. null in $[0, T) \times \mathbb{R}^d$. Therefore, in this and in the following sections, we only consider solutions of the Cauchy problem

$$\begin{cases} \partial_t u + \operatorname{div}(bu) = 0, \\ u(0, \cdot) = 0. \end{cases}$$
(3.2)

We recall that a bounded function $u:[0,T)\times\mathbb{R}^d\to\mathbb{R}$ solves (3.2) in the sense of distributions if

$$\int_0^T \!\! \int_{\mathbb{R}^d} (\partial_t \varphi + b \cdot \nabla \varphi) u \, dx \, dt = 0 \quad \forall \varphi \in C_c^{\infty}([0, T) \times \mathbb{R}^d). \tag{3.3}$$

- **3.3. Remark.** In the following we will use (often tacitly) that the test functions in (3.3) can be equivalently taken in any of the following classes:
- (i) $\varphi: [0, T) \times \mathbb{R}^d \to \mathbb{R}$ Lipschitz with compact support;
- (ii) φ of the form $\varphi_1(t) \varphi_2(x)$, with $\varphi_1 \in C_c^{\infty}([0, T))$ and $\varphi_2 \in C_c^{\infty}(\mathbb{R}^d)$;
- (iii) φ of the form $\varphi_1(t) \varphi_2(x)$, with $\varphi_1 : [0, T) \to \mathbb{R}$ Lipschitz with compact support and $\varphi_2 : \mathbb{R}^d \to \mathbb{R}$ Lipschitz with compact support.
- **3.4. Remark.** Here and in the rest of the paper, u and b are functions defined at every point of their domain, and not equivalence classes. ¹⁴ The reason is that it is more convenient to use the disintegration formula (2.3) with functions defined everywhere rather than almost everywhere.
- **3.5. Measure solutions.** For technical reasons it is convenient to introduce the following notion of generalized solution of (3.2). We say that a *signed* locally finite measure η on $[0, T) \times \mathbb{R}^d$ is a *measure solution* of (3.2) if 15

$$\int_{[0,T)\times\mathbb{R}^d} (\partial_t \varphi + b \cdot \nabla \varphi) \, d\eta = 0 \quad \forall \varphi \in C_c^{\infty}([0,T)\times\mathbb{R}^d). \tag{3.4}$$

Accordingly, we avoid the notations $u \in L^{\infty}$ and $b \in L^{\infty}$.

¹⁵ This notion is similar but not equivalent to that of measure solution considered for instance in [6].

3.6. Solutions on level sets. Take $y \in \mathbb{R}^{d-k}$ and let C be a subset of the level set E_y . We say that a bounded function $u:[0,T)\times C\to \mathbb{R}$ is a *solution of* (3.2) *on* C if the measure $\eta=u\mathbf{1}_C\cdot \mathcal{L}^1\otimes \mu_y$ is a measure solution in the sense of §3.5, that is,

$$\int_0^T \int_C (\partial_t \varphi + b \cdot \nabla \varphi) u \, d\mu_y \, dt = 0 \quad \forall \varphi \in C_c^{\infty}([0, T) \times \mathbb{R}^d). \tag{3.5}$$

It is important to keep in mind that this notion of solution does not involve only the function u and the set C, but also the disintegration measure μ_v .

The following lemma shows that any bounded solution of (3.2) is also a solution on a generic level set.

3.7. Lemma. Let $u:[0,T)\times\mathbb{R}^d\to\mathbb{R}$ be a bounded solution of (3.2). Then u solves (3.2) on the level set E_y for v-a.e. $y\in\mathbb{R}^{d-k}$.

Proof. Fix a function $\rho \in C_c^{\infty}(\mathbb{R}^{d-k})$. We consider in (3.3) Lipschitz test functions of the form $\varphi(t,x)\rho(f(x))$, where $f:\mathbb{R}^d\to\mathbb{R}^{d-k}$ is the Lipschitz function introduced in §3.1 (recall Remark 3.3(i)). After immediate computations we obtain

$$0 = \int_0^T \int_{\mathbb{R}^d} \left[\rho(f) \partial_t \varphi + \rho(f) b \cdot \nabla \varphi + \varphi \langle \nabla f \cdot b; \nabla \rho(f) \rangle \right] u \, dx \, dt$$
$$= \int_0^T \int_{\mathbb{R}^d} \rho(f) (\partial_t \varphi + b \cdot \nabla \varphi) u \, dx \, dt = 0,$$

where the last equality follows from (3.1).

Using (2.3), the above integral can be rewritten as

$$0 = \int_{\mathbb{R}^{d-k}} \int_0^T \int_{E_y} \rho(f) (\partial_t \varphi + b \cdot \nabla \varphi) u \, d\mu_y(x) \, dt \, d\nu(y)$$
$$= \int_{\mathbb{R}^{d-k}} \rho(y) \left[\int_0^T \int_{E_y} (\partial_t \varphi + b \cdot \nabla \varphi) u \, d\mu_y(x) \, dt \right] d\nu(y) = 0.$$

The arbitrariness of ρ gives the existence of a ν -negligible set $N_{\varphi} \subset \mathbb{R}^{d-k}$ with the property that

$$\int_0^T \!\! \int_{E_y} (\partial_t \varphi + b \cdot \nabla \varphi) u \, d\mu_y(x) \, dt = 0 \quad \text{for all } y \notin N_{\varphi}.$$

Finally, we consider a countable dense set $\mathscr{D}\subset C_c^\infty([0,T)\times\mathbb{R}^d)$ and set $N=\bigcup_{\varphi\in\mathscr{D}}N_\varphi$. Since \mathscr{D} is countable the set N is ν -negligible. Moreover, for all $\varphi\in D$ we have

$$\int_0^T\!\!\int_{E_y} (\partial_t \varphi + b \cdot \nabla \varphi) u \, d\mu_y(x) \, dt = 0 \quad \text{for all } y \not\in N;$$

this fact and the density of \mathscr{D} in $C_c^{\infty}([0,T)\times\mathbb{R}^d)$ imply the result.

3.8. Lemma. Take $y \in \mathbb{R}^{d-k}$ with $y \neq 0$ and let $u : [0, T) \times E_y \to \mathbb{R}$ be a bounded solution of (3.2) on E_y . Then u solves (3.2) on every connected component C of E_y .

Proof. The fact that u is a solution of (3.2) on E_v means that (recall §3.6)

$$\int_0^T \!\! \int_{E_y} (\partial_t \varphi + b \cdot \nabla \varphi) u \, d\mu_y \, dt = 0 \quad \forall \varphi \in C_c^{\infty}([0, T) \times \mathbb{R}^d). \tag{3.6}$$

Given a connected component C of E_y , it is possible to find a decreasing sequence of bounded open sets U_n contained in \mathbb{R}^d such that the boundaries ∂U_n do not intersect E_y , and the intersection of the closures \overline{U}_n is C (a proof of this fact is briefly sketched in [2, §2.8]).

Let ε_n be the distance between the sets ∂U_n and E_y (it is positive because these sets are closed and disjoint). Given a smooth convolution kernel ρ supported in the unit ball of \mathbb{R}^d we define

$$\chi_n := \mathbf{1}_{U_n} * \rho_{\varepsilon_n/2}.$$

Considering in (3.6) test functions of the form $\varphi(t, x)\chi_n(x)$, after standard computations we obtain

$$\int_0^T \!\! \int_{E_y} \chi_n(\partial_t \varphi + b \cdot \nabla \varphi) u \, d\mu_y \, dt + \int_0^T \!\! \int_{E_y} \varphi b \cdot \nabla \chi_n \, u \, d\mu_y \, dt = 0. \tag{3.7}$$

Now we pass to the limit as $n \to +\infty$ in this equality. The definition of χ_n and the properties of the sets U_n give that $\nabla \chi_n = 0$ on E_y for every n, so that the second integral in (3.7) vanishes; concerning the first integral, we know that χ_n converges to $\mathbf{1}_C$ everywhere in E_y , and therefore, using the dominated convergence theorem,

$$\int_{0}^{T} \int_{C} (\partial_{t} \varphi + b \cdot \nabla \varphi) u \, d\mu_{y} \, dt = 0$$

for all $\varphi \in C_c^{\infty}([0, T) \times \mathbb{R}^d)$, that is, u solves (3.2) on C.

- **3.9. Lemma.** Given a bounded function $u:[0,T)\times\mathbb{R}^d\to\mathbb{R}$, the following statements are equivalent:
- (a) u solves (3.2);
- (b) there exists an \mathcal{L}^{d-k} -negligible set $N \subset \mathbb{R}^{d-k}$ such that N supports v^s , and setting $F := f^{-1}(N) \cup (\mathbb{R}^d \setminus E^*)$ with E^* given in §2.5 we have:
 - (b1) u solves (3.2) on C for every $C \in \mathcal{C}_{v}$ and every $y \notin N$;
 - (b2) $\mathbf{1}_F u \ solves \ (3.2).$

Proof. Step 1: (a) \Rightarrow (b). Since $v = \mathcal{L}^{d-k} + v^s$, Lemma 3.7 implies that u solves (3.2) on E_y for \mathcal{L}^{d-k} -a.e. y. Therefore, recalling Lemma 2.6(ii), we can choose an \mathcal{L}^{d-k} -negligible set N so that

- (i) $0 \in N$;
- (ii) N supports v^s ;

- (iii) u solves (3.2) on E_y for every $y \notin N$;
- (iv) $\mathscr{H}^k(E_v) < +\infty$ for every $y \notin N$.

Then (b1) follows from Lemma 3.8, and it remains to prove (b2).

Let $y \notin N$ be fixed for the time being. Because of (iv), \mathscr{C}_y is countable, and therefore, by summing (3.5) over all $C \in \mathscr{C}_y$ and recalling that their union is E_y^* , we obtain

$$\int_0^T \left[\int_{E_v^*} (\partial_t \varphi + b \cdot \nabla \varphi) u \, d\mu_y \right] dt = 0 \quad \forall \varphi \in C_c^{\infty}([0, T) \times \mathbb{R}^d), \tag{3.8}$$

that is, u solves (3.2) on E_y^* . By integrating (3.8) over all $y \notin N$ with respect to v we obtain

$$0 = \int_0^T \left[\int_{\mathbb{R}^{d-k} \setminus N} \int_{E_y^*} (\partial_t \varphi + b \cdot \nabla \varphi) u \, d\mu_y \, d\nu(y) \right] dt$$
$$= \int_0^T \left[\int_{\mathbb{R}^{d} \setminus F} (\partial_t \varphi + b \cdot \nabla \varphi) u \, dx \right] dt.$$

The second equality in the previous formula follows by the fact that $\{\mu_y\}$ is the disintegration of \mathcal{L}^d with respect to f and v, and by the fact that E_y^* is equal to $(\mathbb{R}^d \setminus F) \cap E_y$ for all $y \notin N$.

Thus we have proved that $\mathbf{1}_{\mathbb{R}^d\setminus F}u$ solves (3.2), and by linearity we deduce that also $\mathbf{1}_Fu = u - \mathbf{1}_{\mathbb{R}^d\setminus F}u$ solves (3.2), that is, we obtain (b2).

Step 2: (b) \Rightarrow (a). Let N' be the set of all $y \notin N$ such that $\mathscr{H}^k(E_y) = +\infty$. We know by Lemma 2.6(ii) that $\mathscr{L}^{d-k}(N') = 0$, and therefore $\nu(N') = 0$ (because $\nu = \mathscr{L}^{d-k} + \nu^s$ and ν^s is supported on N).

Fix $y \notin N \cup N'$. By (b1), u solves (3.2) on every $C \in \mathscr{C}_y$, and since $\mathscr{H}^k(E_y)$ is finite there are only countably many such C. Hence, proceeding as in Step 1, we find that u solves (3.2) on E_y^* and $\mathbf{1}_{\mathbb{R}^d \setminus F} u$ solves (3.2). Finally (b2) and the linearity of (3.2) imply that $u = \mathbf{1}_F u + \mathbf{1}_{\mathbb{R}^d \setminus F} u$ solves (3.2).

We can now state the main result of this section, namely that, under the additional assumption (3.9), uniqueness for the Cauchy problem (3.2) in \mathbb{R}^d is equivalent to uniqueness on the connected components C in \mathcal{C}_y for a.e. y.

3.10. Theorem. Take b and f as in §3.1 and assume that

$$b(t, x) = 0$$
 for a.e. $(t, x) \in [0, T) \times S$. (3.9)

Then the following statements are equivalent:

- (a) if $u : [0, T) \times \mathbb{R}^d \to \mathbb{R}$ is a bounded solution of (3.2) then u = 0 for a.e. $(t, x) \in [0, T) \times \mathbb{R}^d$;
- (b) for \mathcal{L}^{d-k} -a.e. $y \in \mathbb{R}^{d-k}$ and every $C \in \mathcal{C}_y$ the following implication holds: if $u : [0,T) \times C \to \mathbb{R}$ is a bounded solution of (3.2) on C, then u = 0 for \mathcal{L}^1 -a.e. $t \in [0,T)$ and μ_y -a.e. $x \in C$.

Proof. Step 1: (b) \Rightarrow (a). Let $u:[0,T)\times\mathbb{R}^d\to\mathbb{R}$ be a bounded solution of (3.2). We want to show that u=0 a.e.

From Lemma 3.9 we deduce the existence of an \mathscr{L}^{d-k} -negligible set N such that v^s is supported in N and u solves (3.2) on C for every $C \in \mathscr{C}_y$ and every $y \notin N$. By assumption (b) we can find an \mathscr{L}^{d-k} -negligible set N' such that u=0 for \mathscr{L}^1 -a.e. $t\in [0,T)$ and μ_y -a.e. $x\in C$, for every $C\in \mathscr{C}_y$ and every $y\notin N\cup N'$. By Lemma 2.6(ii) we can also find an \mathscr{L}^{d-k} -negligible set N'' such that for $y\notin N''$ we have $\mathscr{H}^k(E_y)<+\infty$ and therefore \mathscr{C}_y is countable.

The previous considerations imply that u = 0 for \mathcal{L}^1 -a.e. $t \in [0, T)$ and μ_y -a.e. $x \in E_y^*$ for every $y \notin N \cup N' \cup N''$. Recalling that N supports v^s we obtain

$$u = 0$$
 a.e. in $[0, T) \times (E^* \setminus f^{-1}(N \cup N' \cup N''))$. (3.10)

By Lemma 3.9 we know that $\mathbf{1}_F u$ solves (3.2), where $F := f^{-1}(N) \cup (\mathbb{R}^d \setminus E^*)$. Lemma 2.8(i) together with the fact that N is an \mathcal{L}^{d-k} -negligible set supporting v^s implies that $f^{-1}(N) \subset S$ up to negligible sets. Using Lemma 2.6(v) we deduce that $F \subset S$ up to negligible sets, thus assumption (3.9) implies that b = 0 for a.e. $(t, x) \in [0, T) \times F$. Therefore, the fact that $\mathbf{1}_F u$ solves (3.2) can be rewritten as

$$\int_0^T \!\! \int_F \partial_t \varphi \, u \, dx \, dt = 0 \quad \forall \varphi \in C_c^{\infty}([0, T) \times \mathbb{R}^d). \tag{3.11}$$

Now, since every $\phi \in C_c^{\infty}((0,T) \times \mathbb{R}^d)$ can be written as $\phi = \partial_t \varphi$ for some $\varphi \in C_c^{\infty}([0,T) \times \mathbb{R}^d)$, we see that (3.11) implies that

$$u = 0$$
 a.e. in $[0, T) \times F$. (3.12)

Taking into account (3.10) and (3.12), to complete the proof it suffices to show that $f^{-1}((N' \cup N'') \setminus N)$ is negligible. This follows from the fact that $v((N' \cup N'') \setminus N) = 0$ (recall §2.4(iii)), which in turn follows from the fact that N' and N'' are \mathcal{L}^{d-k} -negligible and v^s is supported on N.

Step 2: (a) \Rightarrow (b). Let H be the set of all $y \in \mathbb{R}^{d-k}$ such that there exist $C_y \in \mathscr{C}_y$ and $\tilde{u}_y : [0, T) \times C_y \to \mathbb{R}$ a nontrivial bounded solution of (3.2) on C_y .

We have to show that $\mathcal{L}^{d-k}(H) = 0$.

Assume for contradiction that $\mathcal{L}^{d-k}(H) > 0$. Then we can find a constant $m < +\infty$ and a subset $H' \subset H$ with $\mathcal{L}^{d-k}(H') > 0$ for which the functions \tilde{u}_y above satisfy $|\tilde{u}_y| \leq m$ on $[0, T) \times C_y$ for every $y \in H'$.

At this point we claim that the function u defined by

$$u(t,x) := \begin{cases} \tilde{u}_y(t,x) & \text{if } x \in C_y \text{ for some } y \in H', \\ 0 & \text{otherwise,} \end{cases}$$

is a nontrivial bounded solution of (3.2), in contradiction with assumption (a) of the present theorem.

Indeed, if the function u were Borel, we could use the implication (b) \Rightarrow (a) in Lemma 3.9 to prove the claim. Unfortunately this is not the case, and therefore we must proceed in a (slightly) different way.

We take m as above and consider the set \mathscr{C}^* of all couples (y, η) with $y \in \mathbb{R}^{d-k}$ and η a nontrivial signed measure on $[0, T) \times \mathbb{R}^d$ of the form

$$\eta = u \cdot \mathcal{L}^1 \otimes \mu_{\nu} \tag{3.13}$$

where u solves (3.2) on E_y and satisfies $|u| \le m$ everywhere.

The projection G of \mathscr{C}^* on \mathbb{R}^{d-k} contains H' because for every $y \in H'$ the couple $(y, \tilde{\eta}_y)$ with $\tilde{\eta}_y := \tilde{u}_y \mathbf{1}_{C_y} \cdot \mathscr{L}^1 \otimes \mu_y$ belongs to \mathscr{C}^* . Hence $\mathscr{L}^{d-k}(G) > 0$, and by Corollary 7.2 we can find a Borel set $G' \subset G$ with

$$\nu(G') \geq \mathcal{L}^{d-k}(G') = \mathcal{L}^{d-k}(G) > 0$$

and a Borel family $\{\eta_v : y \in G'\}$ of measures such that $(y, \eta_v) \in \mathscr{C}^*$ for every $y \in G'$.

We now denote by u_y the bounded solution of (3.2) on E_y associated to η_y by (3.13), and we set

$$\eta := \int_{G'} \eta_y \, d\nu(y) = \int_{G'} u_y \cdot \mathcal{L}^1 \otimes \mu_y \, d\nu(y).$$

It is clear that η is a nontrivial measure solution of (3.2), and we deduce from §2.4(iii), (iv) that η is absolutely continuous with respect to $\mathcal{L}^1 \otimes \mathcal{L}^d$,

Thus we can find a function $u:[0,T)\times\mathbb{R}^d\to\mathbb{R}$ such that $\eta=u\cdot\mathcal{L}^1\otimes\mathcal{L}^d$. Moreover, passing through the disintegration of η with respect to f, we recover that for \mathcal{L}^{d-k} -a.e. $y\in G'$ we have $u=u_y$ a.e. with respect to $\mathcal{L}^1\otimes\mu_y$, and u=0 a.e. in $[0,T)\times A$, where A is the complement of the union of E_y over all $y\in G'$.

Since $|u_y| \le m$ everywhere for every $y \in G'$ we conclude that u is a nontrivial solution of (3.2) which satisfies $|u| \le m$ a.e., and this contradicts (a).

- **3.11. Remark.** (i) The converse of Lemma 3.7 also holds: if a bounded function $u:[0,T)\times\mathbb{R}^d\to\mathbb{R}$ solves (3.2) on the level set E_y for v-a.e. $y\in\mathbb{R}^{d-k}$, then u is a bounded solution of (3.2). This is an immediate consequence of (2.3).
- (ii) The converse of Lemma 3.8 does not hold. It is not true in general that a bounded function $u:[0,T)\times E_y\to\mathbb{R}$ which solves (3.2) on every connected component of the level set E_y is also a solution of (3.2) on the level set E_y . This is related to the fact that the connected components of a level set which are \mathscr{H}^k -negligible can in principle be more than countable.
- (iii) Assumption (3.9) in Theorem 3.10 is used in the proof to show that every solution of (3.2) is a.e. null on the set F defined in Lemma 3.9. Presumably we could dispense with assumption (3.9); however, proving the corresponding variant of Theorem 3.10 would require a more refined disintegration formula, to take into account all the connected components of the level sets, including those which are \mathcal{H}^k -negligible (see the previous remark). We decided not to pursue this refinement because we did not see any specific use for it (in the two-dimensional context considered in Sections 4 and 5 assumption (3.9) is always satisified).

4. Uniqueness in the autonomous, divergence-free case

4.1. Assumptions on the vector field. Throughout this section we assume that d=2 and $b: \mathbb{R}^2 \to \mathbb{R}^2$ is a bounded, autonomous, divergence-free vector field with compact support.

It follows that there exists a Lipschitz function $f: \mathbb{R}^2 \to \mathbb{R}$ with compact support such that

$$b = \nabla^{\perp} f \quad \text{a.e. in } \mathbb{R}^2.$$
 (4.1)

Such f is unique and is called the potential associated to b; we then take S, E_y , μ_y and so on as in Section 2 (see §2.19). Note that b and f satisfy the assumptions of Theorem 3.10, namely (3.1) and (3.9).

Theorem 3.10 shows that uniqueness for the Cauchy problem (3.2) is equivalent to uniqueness for the same problem on every nontrivial connected component C of a generic level set E_y of f. Using the parametrization of C given in Lemma 2.11, we show that solving (3.2) on C is equivalent to solving the Cauchy problem with zero initial datum for the equation $\partial_t(u(1+\lambda)) + \partial_s u = 0$ on the parametrization domain $[0, L]^*$ (Lemma 4.4), and for that problem there is uniqueness if and only if $\lambda = 0$ (Lemma 4.5). In Theorem 4.7 we put together these results and prove that if b is taken as above then uniqueness for (3.2) is equivalent to the weak Sard property for the potential f.

- **4.2.** Assumptions on the set C. In the following C is a nontrivial connected component of a level set E_v of f. Taking Lemma 2.11 into account, we assume in addition that
- (i) C is a simple curve with finite length;
- (ii) $\gamma: [0, L]^* \to C$ is a one-to-one Lipschitz parametrization of C such that $\dot{\gamma} \neq 0$ a.e. with respect to \mathcal{L} (the Lebesgue measure on $[0, L]^*$);
- (iii) the restriction of μ_{γ}^{s} to C agrees with the push-forward according to γ of some measure λ on $[0, L]^*$ which is singular with respect to \mathscr{L} .

The next lemma is a particular case of [2, Corollary 7.4], and states that the test functions of the form $\varphi := \phi \circ \gamma$ with $\phi \in C_c^{\infty}(\mathbb{R}^2)$ are dense (in a suitable sense) in the class of Lipschitz functions on $[0, L]^*$.

4.3. Lemma (see [2, Corollary 7.4]). Let a be a function in $L^1([0, L])$ and μ a signed measure on [0, L] such that the functional

$$\Lambda(\varphi) := \int_0^L \dot{\varphi} a \, d\mathcal{L} + \int_0^L \varphi \, d\mu$$

satisfies $\Lambda(\varphi) = 0$ for every φ of the form $\varphi := \varphi \circ \gamma$ with $\varphi \in C_c^{\infty}(\mathbb{R}^2)$. Then $\Lambda(\varphi) = 0$ for every Lipschitz function φ on $[0, L]^*$.

4.4. Lemma. Given a bounded function $u:[0,T)\times C\to \mathbb{R}$, the following statements are equivalent:

Since div b = 0, by rotating b by 90° clockwise we obtain a curl-free vector field, which is therefore the gradient of a Lipschitz function.

- (a) u solves (3.2) on C;
- (b) the function $\tilde{u}(t, s) := u(t, \gamma(s))$ solves

$$\begin{cases} \partial_t(\tilde{u}(1+\lambda)) + \partial_s \tilde{u} = 0, \\ \tilde{u}(0,\cdot) = 0, \end{cases}$$
(4.2)

in the sense of distributions on $[0, T) \times [0, L]^*$, that is,

$$\int_0^T \left[\int_0^L (\partial_t \varphi + \partial_s \varphi) \tilde{u} \, ds + \int_0^L \partial_t \varphi \tilde{u} \, d\lambda \right] dt = 0 \tag{4.3}$$

for every $\varphi \in C_c^{\infty}([0,T) \times [0,L]^*)$. 17

Proof. Step 1: (a) \Rightarrow (b). Since $\mu_y = \gamma_\#(\mathcal{L} + \lambda)$ and b = 0 for $(\gamma_\# \lambda)$ -a.e. point of C, the fact that u solves (3.2) on C amounts to

$$\int_0^T \!\! \int_C (\dot{\varphi}_1 \varphi_2 + \varphi_1 b \cdot \nabla \varphi_2) u \, d(\gamma_\# \mathcal{L}) \, dt + \int_0^T \!\! \int_C \dot{\varphi}_1 \varphi_2 u \, d(\gamma_\# \lambda) \, dt = 0 \qquad (4.4)$$

for all $\varphi_1 \in C_c^{\infty}([0, T))$ and all $\varphi_2 \in C_c^{\infty}(\mathbb{R}^2)$.

We set $\tilde{\varphi}_2(s) = \varphi_2(\gamma(s))$ and compute

$$\dot{\tilde{\varphi}}_2(s) = \nabla \varphi_2(\gamma(s)) \cdot \dot{\gamma}(s) = \nabla \varphi_2(\gamma(s)) \cdot \nabla^{\perp} f(\gamma(s)) \quad \text{for } \mathcal{L}\text{-a.e. } s \in [0, L]^*. \tag{4.5}$$

Using (4.5) and the change of variable $x = \gamma(s)$, identity (4.4) can be rewritten as

$$\int_0^L \left[\int_0^T \varphi_1 \tilde{u} \, dt \right] \dot{\tilde{\varphi}}_2 \, ds + \int_0^L \left[\int_0^T \dot{\varphi}_1 \tilde{u} \, dt \right] \tilde{\varphi}_2 (ds + d\lambda) = 0 \tag{4.6}$$

for all $\varphi_1 \in C_c^{\infty}([0, T))$ and all $\tilde{\varphi}_2$ of the form $\tilde{\varphi}_2 = \varphi_2 \circ \gamma$ with $\varphi_2 \in C_c^{\infty}(\mathbb{R}^d)$.

Then Lemma 4.3 implies that (4.6) holds also for all Lipschitz functions $\tilde{\varphi}_2$ on $[0, L]^*$, which is enough to deduce (4.3) and obtain (b).

Step 2: (b) \Rightarrow (a). To show that u solves (3.2) on C—that is, it satisfies (3.5)—it suffices to observe that (4.3) holds also for all test functions of the form $\varphi(t, \gamma(s))$ with φ : $[0, T) \times C \to \mathbb{R}$ a Lipschitz function with compact support, and then apply the change of variable $x = \gamma(s)$ to the integral on the left-hand side as in the previous step.

- **4.5. Lemma.** Take L > 0 and let λ be any measure on $[0, L]^*$ which is singular with respect to \mathcal{L} . Then the following statements are equivalent:
- (a) if $v : [0, T) \times [0, L]^* \to \mathbb{R}$ is a bounded solution of (4.2) then v = 0 for \mathcal{L}^1 -a.e. $t \in [0, T)$ and $(\mathcal{L} + \lambda)$ -a.e. $s \in [0, L]^*$;
- (b) $\lambda = 0$.

 $^{^{17}}$ $C_c^{\infty}([0,T)\times[0,L]^*)$ corresponds to the space of smooth functions on $[0,T)\times\mathbb{R}$ which are L-periodic in the second variable and compactly supported in the first variable.

Proof. Step 1: (b) \Rightarrow (a). If $\lambda = 0$, then (4.2) reduces to

$$\begin{cases} \partial_t v + \partial_s v = 0, \\ v(0, \cdot) = 0, \end{cases}$$

which has only the trivial solution.¹⁸

Step 2: (a) \Rightarrow (b). We show that for $\lambda \neq 0$ the problem (4.2) has a nontrivial bounded solution. More precisely, we construct two distinct bounded solutions of the equation

$$\partial_t(v(1+\lambda)) + \partial_s v = 0 \tag{4.7}$$

with initial condition $v(0, \cdot) = \mathbf{1}_A$, where A is an \mathcal{L} -negligible compact set contained in (0, L) with $\lambda(A) > 0.19$

It is immediate to check that $v(t, s) := \mathbf{1}_A(s)$ is a stationary solution of (4.7) with initial condition $v(0, \cdot) = \mathbf{1}_A$.

Consider now the function $\sigma: [0, L]^* \to [0, \widehat{L}]^*$ and its inverse $\hat{\sigma}$ defined in §2.16 (see also Remark 2.18); since the (distributional) derivative of $\hat{\sigma}$ is $1 + \lambda$, a *formal* computation shows that the change of variable $s = \sigma(\hat{s})$ reduces the equation

$$\partial_t w + \partial_{\hat{s}} w = 0 \tag{4.8}$$

to (4.7). This suggests a way to construct the second solution we are looking for: set

$$v(t,s) := \begin{cases} w(t,\hat{\sigma}(s)) & \text{for } s \notin \mathcal{A}, \\ \int_{\hat{\sigma}(s)} w(t,\hat{s}) \, d\hat{s} & \text{for } s \in \mathcal{A}, \end{cases}$$
(4.9)

where \mathscr{A} is the set of atoms of λ (cf. §2.16), 20 and $w:[0,T)\times[0,\widehat{L}]^*\to\mathbb{R}$ is the (unique) bounded solution of (4.8) with initial condition $w(0,\cdot)=\mathbf{1}_{\hat{\sigma}(A)}$, which means that w satisfies

$$\int_{0}^{T} \int_{0}^{\widehat{L}} (\partial_{t} \varphi + \partial_{\widehat{s}} \varphi) w \, d\widehat{s} \, dt = \int_{0}^{\widehat{L}} \varphi(0, \cdot) \mathbf{1}_{\widehat{\sigma}(A)} \, d\widehat{s}$$
 (4.10)

for every $\varphi \in C_c^{\infty}([0,T) \times [0,\widehat{L}]^*)$. We claim that v is a solution of (4.7) with initial condition $v(0,\cdot) = \mathbf{1}_A$ and is not stationary, and in particular it differs from the stationary solution

Step 3: v solves (4.7) with initial condition $v(0,\cdot) = \mathbf{1}_A$. Let F be an \mathcal{L} -negligible set which supports λ . Then the claim can be rewritten as follows: for every test function $\varphi \in C_c^{\infty}([0,T)\times[0,L]^*)$ we have

$$\int_{0}^{T} \int_{0}^{L} (\partial_{t} \varphi + \mathbf{1}_{I \setminus F} \partial_{s} \varphi) v \, d(\mathcal{L} + \lambda) \, dt = \int_{0}^{L} \varphi(0, \cdot) \mathbf{1}_{A} \, d(\mathcal{L} + \lambda). \tag{4.11}$$

This is well-known with $[0, L]^*$ replaced by \mathbb{R} ; the same proof works also for $[0, L]^*$.

Such an A can be found because the measure λ is supported in (0, L) (cf. §2.12(ii)).

Recall that $\hat{\sigma}$ is single-valued for $s \notin \mathcal{A}$, and $\hat{\sigma}(s)$ is the interval $[\hat{\sigma}_{-}(s), \hat{\sigma}_{+}(s)]$ for $s \in \mathcal{A}$ (cf. §2.16).

We apply the change of variable $s = \sigma(\hat{s})$; letting $\hat{v}(t, \hat{s}) := v(t, \sigma(\hat{s}))$ and $\hat{\varphi}(t, \hat{s}) := \varphi(t, \sigma(\hat{s}))$ and using (2.9), identity (4.11) becomes

$$\int_0^T \int_0^{\widehat{L}} (\partial_t \hat{\varphi} + \partial_{\widehat{s}} \hat{\varphi}) \hat{v} \, d\hat{s} \, dt = \int_0^{\widehat{L}} \hat{\varphi}(0, \cdot) \mathbf{1}_{\hat{\sigma}(A)} \, d\hat{s} \tag{4.12}$$

(note that $\partial_t \hat{\varphi} = \partial_t \varphi$ and $\partial_{\hat{s}} \hat{\varphi} = \partial_s \varphi \mathbf{1}_{\widehat{I} \setminus \hat{\sigma}(F)}$).

If \hat{v} were equal to w, then (4.12) would follow by writing (4.10) with the test function $\hat{\varphi}$ (recall Remark 3.3). Unfortunately, \hat{v} agrees with w only on the complement of $\hat{\sigma}(\mathscr{A})$. However, we can recover (4.12) from (4.10) by showing that for every $t \in [0, T)$,

$$\int_{\hat{\sigma}(\mathscr{A})} (\partial_t \hat{\varphi} + \partial_{\hat{s}} \hat{\varphi}) \hat{v} \, d\hat{s} = \int_{\hat{\sigma}(\mathscr{A})} (\partial_t \hat{\varphi} + \partial_{\hat{s}} \hat{\varphi}) w \, d\hat{s}.$$

Indeed,

$$\int_{\hat{\sigma}(\mathscr{A})} (\partial_t \hat{\varphi} + \partial_{\hat{s}} \hat{\varphi}) \hat{v} \, d\hat{s} = \sum_{s \in \mathscr{A}} \int_{\hat{\sigma}(s)} (\partial_t \hat{\varphi} + \partial_{\hat{s}} \hat{\varphi}) \hat{v} \, d\hat{s} = \sum_{s \in \mathscr{A}} \partial_t \varphi(t, s) \int_{\hat{\sigma}(s)} \hat{v} \, d\hat{s}
= \sum_{s \in \mathscr{A}} \partial_t \varphi(t, s) \int_{\hat{\sigma}(s)} w \, d\hat{s} = \int_{\hat{\sigma}(\mathscr{A})} (\partial_t \hat{\varphi} + \partial_{\hat{s}} \hat{\varphi}) w \, d\hat{s}, \quad (4.13)$$

where the second and the last equalities follow from the fact that $\partial_{\hat{s}}\hat{\varphi}(t,\hat{s}) = 0$ and $\partial_t\hat{\varphi}(t,\hat{s}) = \partial_t\varphi(t,s)$ for all \hat{s} in the interval $\hat{\sigma}(s)$, while the third equality follows from that fact that in each of these intervals the function $\hat{v}(t,\hat{s})$ coincides with the average of $w(t,\hat{s})$ over the same interval (see (4.9)).

Step 4: v is not stationary. This claim can be proved in many ways. We choose to look at the maximum a(t) of the support of the function $v(t, \cdot)$ (computed with respect to the measure $\mathcal{L} + \lambda$), and show that this function is strictly increasing at t = 0.

Let $\hat{a}(t)$ be the maximum of the support of the function $w(t,\cdot)$ (computed with respect to the Lebesgue measure). Since w solves (4.8), a continuity equation with a regular (in fact, constant) vector field, it is propagated along characteristics, and therefore \hat{a} is strictly increasing in t, at least as long as it does not reach the value L.

Moreover $a = \sigma(\hat{a})$. Therefore, recalling the properties of σ (see §2.16), we deduce that a is strictly increasing at t = 0 provided that $\hat{a}(0)$ does not belong to any of the intervals $\sigma(s)$ with $s \in \mathcal{A}$, or, if it does, then it agrees with the supremum of such an interval. Finally, one can check that this condition is always satisfied (recall that A is closed).

4.6. Remark. Concerning the previous proof, it should be noted that the change of variable $s = \sigma(\hat{s})$ reduces to (4.7) *every* equation of the form

$$\partial_t v + \partial_{\hat{s}}(g_1 v + g_2) = 0 \tag{4.14}$$

where the functions g_1 and g_2 satisfy the following conditions: g_1 is constant a.e. in $\hat{\sigma}(F)$, $g_1 = 1$ and $g_2 = 0$ a.e. outside $\hat{\sigma}(F)$. More precisely, if w is a solution of (4.14), then the function v defined by (4.9) is a solution of (4.7). In particular, taking $g_1 := 1$ and $g_2 := -\mathbf{1}_A$ we recover the stationary solution $v(t, s) := \mathbf{1}_A(s)$.

4.7. Theorem. Take b and f as in $\S4.1$. Then the following statements are equivalent:

- (a) if $u : [0, T) \times \mathbb{R}^2 \to \mathbb{R}$ is a bounded solution of (3.2) then u = 0 for a.e. $(t, x) \in [0, T) \times \mathbb{R}^2$;
- (b) the potential f has the weak Sard property (defined in §2.13).

Proof. Since b and f satisfy the assumptions of Theorem 3.10, the uniqueness for the Cauchy problem (3.2) (statement (a)) is equivalent to uniqueness for the same problem on every $C \in \mathcal{C}_y$ for a.e. y.

By Lemma 2.11, for a.e. y, every connected component $C \in \mathcal{C}_y$ satisfies the assumptions in §4.2, and therefore Lemmas 4.4 and 4.5 imply that the uniqueness of (3.2) on C is equivalent to $\lambda = 0$, which means that the restriction of μ_y^s to C vanishes.

Finally Lemma 2.14 states that $\mathbf{1}_C \cdot \mu_y^s = 0$ for every $C \in \mathcal{C}_y$ and a.e. y if and only if f has the weak Sard property.

4.8. Corollary. There exists a divergence-free autonomous vector field b on the plane which belongs to $C^{0,\alpha}$ for every $\alpha < 1$, and for which the Cauchy problem (3.2) has nontrivial bounded solutions.

Proof. In [2, §4.8] we construct a function $f': \mathbb{R}^2 \to \mathbb{R}$ of class $C^{1,\alpha}$ for every $\alpha < 1$ which does not have the weak Sard property. We modify f' so as to make it compactly supported, and set $b := \nabla^{\perp} f'$. Then Theorem 4.7 shows that for such b the Cauchy problem (3.2) admits nontrivial bounded solutions.

5. Uniqueness in the bounded divergence case

5.1. Assumptions on the vector field. In this section we consider a bounded, time-dependent vector field $b: [0, T) \times \mathbb{R}^2 \to \mathbb{R}^2$ of the form

$$b(t, x) = a(t, x) \nabla^{\perp} f(x) \quad \text{a.e. in } [0, T) \times \mathbb{R}^2, \tag{5.1}$$

where $f: \mathbb{R}^2 \to \mathbb{R}$ is a Lipschitz function with compact support and $a: [0, T) \times \mathbb{R}^2 \to \mathbb{R}$ is a bounded function. We then take S, E_y, μ_y and so on as in Section 2 (see §2.19).

We assume moreover that b has bounded divergence in the sense of distributions, that is, there exists a bounded function $g:[0,T)\times\mathbb{R}^2\to\mathbb{R}$ such that

$$\int_0^T \!\! \int_{\mathbb{R}^2} (b \cdot \nabla \varphi + g\varphi) \, dx \, dt = 0 \quad \forall \varphi \in C_c^{\infty}((0, T) \times \mathbb{R}^2). \tag{5.2}$$

The next result extends Theorem 4.7 and provides a characterization of uniqueness within the special class of bounded vector fields considered in the previous subsection.

- **5.2. Theorem.** Take b, f and a as in §5.1. The following statements hold:
- (i) If f has the weak Sard property then there is uniqueness for the Cauchy problem (3.2), that is, every bounded solution $u : [0, T) \times \mathbb{R}^2 \to \mathbb{R}$ of (3.2) satisfies u = 0 for a.e. $(t, x) \in [0, T) \times \mathbb{R}^2$.

- (ii) Conversely, if f does not have the weak Sard property and there exists $\delta > 0$ such that $a \geq \delta$ a.e. on $[0, T) \times \mathbb{R}^2$, then (3.2) admits a nontrivial bounded solution.
- **5.3. Remark.** (i) Observe that, for given b and f, the value of a is determined (a.e.) in $[0, T) \times (\mathbb{R}^2 \setminus S)$ by equation (5.1), but can be freely chosen on $[0, T) \times S$.
- (ii) The class of vector fields considered in this section clearly includes the one considered in Section 4. The structural assumption (5.1) is however very rigid, in that it allows only for vector fields whose direction does not depend on time, and does not even include all autonomous vector fields with bounded divergence.²¹
- (iii) The role of the nondegeneracy assumption $a \ge \delta > 0$ in statement (ii) of Theorem 5.2 will be discussed in detail in Remark 5.5 at the end of this section. Here we just summarize a few conclusions:
- (a) Some nonegeneracy assumption on a is needed: if b (or equivalently a) vanishes on a neighbourhood of the critical set S, then uniqueness for (3.2) holds regardless of whether f satisfies the weak Sard property or not.
- (b) It is not enough to assume that a is strictly positive (a > 0).
- (c) The assumption $a \ge \delta > 0$ can be clearly replaced by $a \le -\delta < 0$, 22 but none of these is optimal. Yet we could not find weaker conditions that are easily expressed only in terms of a and f.
- (iv) Consider two different couples f_1 , a_1 and f_2 , a_2 that decompose b as in (5.1), that is, $b = a_1 \nabla^{\perp} f_1 = a_2 \nabla^{\perp} f_2$. Under the assumption that both a_1 and a_2 are positive and bounded away from zero, f_1 has the weak Sard property if and only if f_2 does (because by Theorem 5.2 both conditions are equivalent to the uniqueness for (3.2)). More generally, if a_2 is positive and bounded away from zero then the weak Sard property of f_1 implies that of f_2 .

We first prove a lemma that characterizes the derivative of a along a nontrivial connected component C of a generic level set E_y . The rest of the proof of Theorem 5.2 is a more or less straightforward modification of the proof of Theorem 4.7, and we will only summarize the main steps.

5.4. Lemma. For a.e. $y \in \mathbb{R}$, every connected component $C \in \mathcal{C}_y$ satisfies the assumptions in §4.2, and

$$\partial_s \tilde{a} = \tilde{g} (1 + \lambda) \tag{5.3}$$

in the sense of the distributions on $(0, T) \times [0, L]^*$, where we have set $\tilde{a}(t, s) := a(t, \gamma(s))$ and $\tilde{g}(t, s) := g(t, \gamma(s))$ and $\gamma, \lambda, \mathcal{L}$ are taken as in §4.2.

Moreover there exists a function $\alpha:[0,T)\times[0,\widehat{L}]^*\to\mathbb{R}$ which is bounded, uniformly Lipschitz in the second variable, and satisfies $\tilde{a}(t,s)=\alpha(t,\hat{\sigma}(s))$ for a.e. (t,s), where \hat{L} and $\hat{\sigma}:[0,L]^*\to[0,\widehat{L}]^*$ are given in §2.16 (see also Remark 2.18).

In particular when $\lambda = 0$ we have $\widehat{L} = L$, $\widehat{\sigma}(s) = s$, and therefore $\widetilde{a} = \alpha$ a.e.

The vector field b(x) := x cannot be written in the form (5.1) in any neighbourhood of 0: if it were, all level sets of f would contain the point 0, which is clearly impossible for a nonconstant Lipschitz function.

Just apply the result with a and f replaced by -a and -f.

Sketch of proof. Step 1: proof of (5.3). Following the proofs of Lemmas 3.7 and 3.8 we find that for every $C \in \mathcal{C}_v$ and a.e. y,

$$\int_0^T \!\! \int_C (b \cdot \nabla \varphi + g\varphi) \, d\mu_y \, dt = 0 \quad \forall \varphi \in C_c^{\infty}((0, T) \times \mathbb{R}^2). \tag{5.4}$$

Moreover by Lemma 2.11 the connected component C satisfies the requirements in §4.2 and therefore, arguing as in the proof of Lemma 4.4, equation (5.4) translates into

$$\int_0^T \left[\int_0^L (\tilde{a}\partial_s \varphi + g\varphi) \, d\mathcal{L} + \int_0^L \tilde{g}\varphi \, d\lambda \right] dt = 0 \quad \forall \varphi \in C_c^{\infty}((0,T) \times [0,L]^*),$$

which is precisely the weak formulation of (5.3).

Step 2: construction of α . For every $t \in [0, T)$ and $\hat{s} \in [0, \widehat{L}]$ we set

$$\alpha(t, \hat{s}) := c(t) + \int_0^{\hat{s}} \tilde{g}(t, \sigma(\cdot)) d\mathcal{L},$$

and we choose c(t) so that the integrals of $\tilde{a}(t,\cdot)$ and $\alpha(t,\hat{\sigma}(\cdot))$ over [0,L] are the same.

The function α is clearly well-defined on $[0, T) \times [0, \widehat{L}]$, bounded, and Lipschitz in the second variable. Moreover (5.3) and formula (2.9) imply that for a.e. t,

$$0 = \int_0^L \tilde{g}(t, \cdot) d(\mathcal{L} + \lambda) = \int_0^{\hat{L}} \tilde{g}(t, \sigma(\cdot)) d\mathcal{L}$$

and therefore $\alpha(t,0) = \alpha(t,\widehat{L})$. Hence $\alpha(t,\cdot)$ is well-defined and Lipschitz also as a function on $[0,\widehat{L}]^*$.

Step 3. Formula (2.9) yields

$$\alpha(t, \hat{\sigma}(s)) = c(t) + \int_0^s \tilde{g}(t, \cdot) d(\mathcal{L} + \lambda)$$

for every t and every s such that $\hat{\sigma}$ is single-valued at s. This identity and (5.3) imply that

$$\partial_s \alpha(t, \hat{\sigma}(s)) = \tilde{g}(1+\lambda) = \partial_s \tilde{a}$$
 (5.5)

in the sense of distributions on $(0, T) \times [0, L]^*$. In turn, (5.5) and the choice of c(t) imply that $\alpha(t, \hat{\sigma}(s)) = \tilde{a}(t, s)$ for a.e. (t, s).

Sketch of proof of Theorem 5.2. Step 1. Since b and f satisfy the assumptions of Theorem 3.10, the uniqueness for the Cauchy problem (3.2) is equivalent to uniqueness for the same problem on every $C \in \mathcal{C}_y$ for a.e. y. Moreover, in view of Lemma 5.4, we can restrict ourselves to those C which satisfy the assumptions in §4.2 and in Lemma 5.4.

Step 2. A straightforward modification of Lemma 4.4 shows that the uniqueness for the Cauchy problem (3.2) on C is equivalent to the uniqueness for

$$\begin{cases} \partial_t(\tilde{u}(1+\lambda)) + \partial_s(\tilde{a}\tilde{u}) = 0, \\ \tilde{u}(0,\cdot) = 0, \end{cases}$$
 (5.6)

on $[0, T) \times [0, L]^*$, where \tilde{a} is defined as in Lemma 5.4.

Step 3: proof of statement (i). Since f has the weak Sard property, by Lemma 2.14 we can assume that the restriction of μ_{ν}^{s} to C vanishes, which means that λ vanishes.

By the previous steps, it suffices to show that for $\lambda = 0$ the only bounded solution of (5.6) is the trivial one. Indeed (5.6) reduces to the standard continuity equation

$$\begin{cases} \partial_t \tilde{u} + \partial_s(\tilde{a}\tilde{u}) = 0, \\ \tilde{u}(0, \cdot) = 0. \end{cases}$$
 (5.7)

Since \tilde{a} agrees a.e. with the function α given in Lemma 5.4, and α is Lipschitz in the second variable, the fact that \tilde{u} is trivial follows by a standard result.

Step 4: proof of statement (ii). In view of the first two steps, it suffices to show that if $\lambda \neq 0$ then the problem (5.6) admits a nontrivial bounded solution.

To do this, we strictly follow the second part of the proof of Lemma 4.5 and construct two distinct bounded solutions of the equation

$$\partial_t(v(1+\lambda)) + \partial_s(\tilde{a}v) = 0, \tag{5.8}$$

with the same initial condition $v(0, \cdot) = \mathbf{1}_A$, where A is chosen as in that proof.

One solution is the stationary one, and the other one is the function v given by formula (4.9) by taking as w the (unique) bounded solution of

$$\partial_t w + \partial_{\hat{s}}(\alpha w) = 0 \tag{5.9}$$

on $[0, T) \times [0, \widehat{L}]^*$ with initial condition $w(0, \cdot) = \mathbf{1}_{\widehat{\sigma}(A)}$, where α is the function constructed in Lemma 5.4.

The fact that v solves (5.8) and is different from the stationary solution can be proved as in Lemma 4.5.

5.5. Remark. The assumption $a \ge \delta > 0$ in statement (ii) of Theorem 5.2 is used only once in the proof above, and precisely to show that the solution v of (5.8) built out of a certain solution w of (5.9) is not stationary. Assuming that a, and therefore α , is positive and bounded away from zero is clearly sufficient for this purpose (even though not necessary).

In the rest of this subsection we argue that just assuming a strictly positive is not enough. We take a function f of class C^1 on \mathbb{R}^2 without the weak Sard property (see [2, Section 4]) and construct a function a = a(x) so that \tilde{a} is Lipschitz, positive, and vanishes only on the closed set $\gamma^{-1}(S)$. We claim that the Cauchy problem (3.2) with $b := a \nabla^{\perp} f$ admits only the trivial solution, despite the fact that f does not have the weak Sard property and a is strictly positive.²³

As pointed out in the proof above, this is equivalent to showing that (5.6) admits only the trivial solution. Now \tilde{a} vanishes on $\gamma^{-1}(S)$ by construction, and since λ is supported on $\gamma^{-1}(S)$, it follows that \tilde{a} vanishes λ -a.e. Therefore the equation in (5.6) can be rewritten as $\partial_t v + \partial_s(\tilde{a}v) = 0$ with $v := \tilde{u}(1 + \lambda)$, and it is known that for this equation there is uniqueness also among measure solutions because \tilde{a} is Lipschitz.

²³ By construction we only imposed that a is strictly positive (a.e.) outside S, but remember that a can be freely modified in S without affecting b. The fact that b has bounded divergence follows by a suitable converse of Lemma 5.4.

6. Additional remarks

6.1. Extensions of Theorems 4.7 and 5.2. (i) In both statements, the assumption that b has compact support is made for the sake of simplicity, and can be easily removed. In that case, one should also consider nontrivial connected components C of the level sets of f which are simple curves with end points at infinity.

(ii) The continuity equation (1.1) can be modified so as to include a (possibly non-linear) source term on the right-hand side, that is,

$$\partial_t u + \operatorname{div}(bu) = h(t, x, u),$$

where the function $h:[0,T)\times\mathbb{R}^2\times\mathbb{R}\to\mathbb{R}$ is bounded. The key point is clearly that uniqueness holds for the corresponding one-dimensional equation

$$\partial_t v + \partial_s(\tilde{a}v) = \tilde{h}(t, s, v)$$

on $[0, T) \times [0, L]^*$.

- (iii) Uniqueness can be shown in the class of weak solutions that are integrable in space and time (instead of being bounded). The key point is that uniqueness holds for the corresponding one-dimensional equation $\partial_t v + \partial_s(\tilde{a}v) = 0$ among solutions which are integrable in space and time.
- **6.2. Extension to higher dimension.** It is possible to extend Theorem 5.2 to higher dimensions, and more precisely to bounded, time-dependent vector fields b on \mathbb{R}^d with bounded divergence which satisfy the following structural assumption: there exists a Lipschitz map $f: \mathbb{R}^d \to \mathbb{R}^{d-1}$ such that $\nabla f \cdot b = 0$ a.e., and b = 0 at a.e. point where the rank of ∇f is not maximal.

In this case, the uniqueness for the Cauchy problem (3.2) should be proved under the following assumptions: (i) f satisfies a suitable version of the weak Sard property, and (ii) for a.e. y the level set E_y of f does not contain triods, and therefore its nontrivial connected components are simple curves.

It seems that both assumptions are not only sufficient for uniqueness, but also necessary. Note that assumption (ii) is automatically satisfied when d=2 (Lemma 2.11(i)), while for d>2 it is satisfied when f is of class $C^{1,1/2}$ (see [2, Lemma 2.16]), but may fail when f is of class $C^{1,\alpha}$ with $\alpha<1/(d-1)$ (see [2, Section 3]).

Since also in this case the uniqueness for the Cauchy problem (3.2) turns out to be equivalent to the uniqueness for the same problem on the nontrivial connected components of a generic level set of f, the lack of triods is (presumably) also necessary for uniqueness.

6.3. Renormalization property. The key step in the uniqueness proof by DiPerna and Lions [13] is proving the renormalization property for weak solutions.

This property is not used in this paper, but it can be easily obtained as a consequence of our dimension-reduction technique. Let us assume for simplicity that b is divergence-free and autonomous; in this case the renormalization property for a weak solution u of (1.1) simply means that $\beta(u)$ is a weak solution of the same equation for every smooth function $\beta: \mathbb{R} \to \mathbb{R}^{24}$ when the potential f associated to b has the weak Sard property

²⁴ If u is of class C^1 , this follows by a straightforward computation.

this fact follows from the renormalization property for the one-dimensional equation $\partial_t v + \partial_s v = 0$.

- **6.4. Regular Lagrangian flow.** Under the weak regularity assumptions on b considered in this paper, the meaningful notion of flow for the ordinary differential equation $\dot{x} = b(t, x)$ is that of regular Lagrangian flow (see [4]): we say that a Borel map Φ : $[0, T) \times \mathbb{R}^2 \to \mathbb{R}^2$ is a *regular Lagrangian flow* associated to b if
- (i) for a.e. $x \in \mathbb{R}^2$ the map $t \mapsto \Phi(t, x)$ solves $\dot{x} = b(t, x)$ in the integral sense;
- (ii) there exists a constant c such that for every t the push-forward of \mathcal{L}^2 according to $\Phi_t := \Phi(t,\cdot)$ satisfies $(\Phi_t)_\# \mathcal{L}^2 \le c \mathcal{L}^2$.

Theorems 4.7 and 5.2, together with the abstract theory of regular Lagrangian flows developed in [6], give the following result: if b is taken as in §4.1 or §5.1, and the corresponding function f has the weak Sard property, then there exists a unique regular Lagrangian flow for b.

6.5. Strong locality of the divergence operator. It is well-known that the (distributional) gradient is *strongly local* for Sobolev functions, in the sense that the following implication holds for every Sobolev function *u* and every Borel set *A* in the domain of *u*:

$$u = \text{const. a.e. on } A \implies \nabla u = 0 \text{ a.e. on } A.$$

It follows immediately that also every first-order differential operator is strongly local on Sobolev spaces.

However, this is no longer true on larger spaces; in particular the divergence operator is not strongly local on the space of vector fields with bounded divergence. Indeed, for vector fields b on \mathbb{R}^2 of the form

$$b = a\nabla^{\perp}f,\tag{6.1}$$

where f is a Lipschitz function on \mathbb{R}^2 and a a bounded function on \mathbb{R}^2 , the strong locality of the divergence is strictly related to the weak Sard property of f. More precisely:

- **6.6. Proposition.** Let b, f and a be given as above, and assume that b has bounded divergence. Then the following statements hold for every $v \in \mathbb{R}^2$:
- (i) If the function $f(x) v \cdot x$ has the weak Sard property, then for every Borel set A in \mathbb{R}^2 we have the implication

$$b = v \text{ a.e. on } A \Rightarrow \operatorname{div} b = 0 \text{ a.e. on } A.$$
 (6.2)

(ii) Conversely, if $f(x) - v \cdot x$ does not have the weak Sard property then there exist b as above and a Borel set A such that the implication (6.2) fails.

In [2, Section 5] we construct an explicit example of a bounded vector field b on the plane whose divergence is bounded, nontrivial, and supported in the set where b = 0, and then use it to construct a Lipschitz function without the weak Sard property.

Sketch of proof of Proposition 6.6. With no loss of generality we can assume v = 0. We first prove statement (i). Let g be the divergence of b, that is,

$$\int_{\mathbb{R}^2} (b \cdot \nabla \varphi + g\varphi) \, dx = 0 \quad \forall \varphi \in C_c^{\infty}(\mathbb{R}^2).$$

Starting from this identity and arguing as in the proof of Lemma 3.9 we find an \mathcal{L}^1 -negligible set N in \mathbb{R} such that the following hold: the singular measure v^s defined in §2.7 is supported on N, for every $y \notin N$ and every $C \in \mathcal{C}_v$ we have

$$\int_{C} (b \cdot \nabla \varphi + g\varphi) \, d\mu_{y} = 0 \quad \forall \varphi \in C_{c}^{\infty}(\mathbb{R}^{2}), \tag{6.3}$$

and setting $F := f^{-1}(N) \cup (\mathbb{R}^2 \setminus E^*)$ we have

$$\int_{\mathbb{R}} (b \cdot \nabla \varphi + g\varphi) \, dx = 0 \quad \forall \varphi \in C_c^{\infty}(\mathbb{R}^2). \tag{6.4}$$

Since F is contained (up to a negligible subset) in the critical set S, we find that b=0 a.e. on F, and therefore (6.4) implies that

$$g = 0$$
 a.e. in F . (6.5)

Moreover, choosing a parametrization γ of C as in §4.2 and recalling that the measure λ which appears there is null because of the weak Sard property (Lemma 2.14), we can rewrite (6.3) as

$$\int_0^L (\tilde{a}\dot{\varphi} + \tilde{g}\varphi) \, ds = 0 \quad \forall \varphi \in C^{\infty}([0, L]^*), \tag{6.6}$$

where we have set $\tilde{a}(s) := a(\gamma(s))$ and $\tilde{g}(s) := g(\gamma(s))$. This means that \tilde{g} is the distributional derivative of \tilde{a} .

Now, the assumption b=0 a.e. on A implies that $\tilde{a}=0$ a.e. on $\gamma^{-1}(A)$, and therefore the strong locality of derivatives of Sobolev functions yields $\tilde{g}=0$ a.e. in $\gamma^{-1}(A)$, that is,

$$g = 0 \quad \mu_{\nu} \text{-a.e. in } A \cap C. \tag{6.7}$$

Since (6.7) holds for every $C \in \mathcal{C}_y$ and for ν -a.e. y, the disintegration formula for the Lebesgue measure (cf. §2.7) yields

$$g = 0$$
 a.e. in $A \setminus F$,

and together with (6.5) this implies that g = 0 a.e. in A, which concludes the proof of statement (i).

We now prove statement (ii). Since f does not have the weak Sard property, there exists a set G in \mathbb{R} with positive measure such that for every $y \in G$ there exists $C \in \mathcal{C}_y$ for which the measure λ is nontrivial (recall Lemma 2.14).

For every such C, consider the usual parametrization γ and choose two bounded functions \tilde{a} and \tilde{g} on $[0, L]^*$ so that

$$\partial_s \tilde{a} = \tilde{g}(1+\lambda)$$

in the sense of distributions on $[0, L]^*$, and \tilde{g} is not λ -a.e. null.

Now define functions g and a on \mathbb{R}^2 so that g = a = 0 out of the union of all C given above, and $a(\gamma(s)) = \tilde{a}(s)$ and $g(\gamma(s)) = \tilde{g}(s)$ for $(\mathcal{L} + \lambda)$ -a.e. s and every C given above.

It can be verified that g is the distributional divergence of $b := a\nabla^{\perp} f$, and the restriction of g to the singular set S is not a.e. null. Hence the implication (6.2) fails for A := S.

7. Appendix: a measurable selection lemma

In this appendix we prove a measurable selection lemma used in the proof of Theorem 3.10. For the notation we refer to Section 2. Given m, T > 0, we denote by \mathscr{C} the set of all couples (y, η) such that $y \in \mathbb{R}^{d-k}$ and η is a real-valued measure on $[0, T) \times \mathbb{R}^d$ of the form

$$\eta = u \cdot \mathcal{L}^1 \otimes \mu_{\nu},$$

where $u:[0,T)\times E_y\to\mathbb{R}$ satisfies $|u|\leq m$ everywhere and solves (3.2) on E_y , or equivalently η is a measure solution of (3.2) in the sense of §3.5.

We denote by \mathscr{M} the Banach space of real-valued measures on $[0,T)\times\mathbb{R}^d$. In the measure-theoretic issues that will be considered below, the word "Borel" in relation to \mathscr{M} refers to the σ -algebra generated by the weak* topology of \mathscr{M} (that is, the one induced by the standard duality with the space of continuous functions on $[0,T)\times\mathbb{R}^d$ vanishing at infinity). The key point is that the restriction of the weak* topology to any closed ball of \mathscr{M} is compact and metrizable, and therefore Polish (separable and completely metrizable).

7.1. Proposition. The set \mathscr{C} defined above is a Borel subset of $\mathbb{R}^{d-k} \times \mathscr{M}$.

Proof. Step 1. We rewrite \mathscr{C} as the set of all (y, η) such that

$$|\eta|(E) \le m\mathcal{L}^1 \otimes \mu_{\mathcal{V}}(E) \quad \forall E \subset [0, T) \times \mathbb{R}^d,$$
 (7.1)

where $|\eta|$ denotes the variation of the real-valued measure η , and

$$\int (\partial_t \varphi + b \cdot \nabla \varphi) \, d\eta = 0 \quad \forall \varphi \in C_c^{\infty}([0, T) \times \mathbb{R}^d). \tag{7.2}$$

Indeed (7.1) means that the measure η is absolutely continuous with respect to $\mathcal{L}^1 \otimes \mu_y$ and that the density u of the former measure with respect to the latter satisfies $|u| \leq m$ almost everywhere, while (7.2) means that η is a measure solution of (3.2).

Step 2: the set of all (y, η) which satisfy (7.1) is a Borel subset of $\mathbb{R}^{d-k} \times \mathcal{M}$. Note that (7.1) can be rewritten as

$$\int \varphi \, d\eta \le m \int |\varphi| \, dt \, d\mu_y \quad \forall \varphi \in C_0([0, T) \times \mathbb{R}^d). \tag{7.3}$$

Moreover it suffices that the inequality in (7.3) is satisfied for all $\varphi \in \mathcal{D}$, where \mathcal{D} is a given *countable* dense subset of $C_c([0,T)\times\mathbb{R}^d)$. In other words (7.3), and therefore also (7.1), can be rewritten as $\Lambda_{\varphi}(y,\eta) \leq 0$ for all $\varphi \in \mathcal{D}$, where

$$\Lambda_{\varphi}(y,\eta) := \int \varphi \, d\eta - m \int |\varphi| \, dt \, d\mu_y. \tag{7.4}$$

Now, the first integral in (7.4) is a continuous function of η (by the definition of the topology on \mathcal{M}), while the second is a Borel function of y (because the family $\{\mu_y\}$ is Borel, cf. §2.2 and §2.3), and therefore Λ_{φ} is a Borel function on $\mathbb{R}^{d-k} \times \mathcal{M}$.

Hence the set of all (y, η) satisfying $\Lambda_{\varphi}(y, \eta) \leq 0$ is Borel for every given φ , and so is the set of all (y, η) satisfying $\Lambda_{\varphi}(y, \eta) \leq 0$ for all $\varphi \in \mathcal{D}$.

Step 3: the set of all (y, η) which satisfy (7.2) is a Borel subset of $\mathbb{R}^{d-k} \times \mathcal{M}$. Note that in (7.2) we can equivalently require that the equality holds just for all test functions $\varphi \in \mathcal{D}'$, where \mathcal{D}' is a given *countable* dense subset of $C_c^{\infty}([0, T) \times \mathbb{R}^d)$. Therefore, arguing as in Step 2, we only need to show that for every $\varphi \in C_c^{\infty}([0, T) \times \mathbb{R}^d)$ the integral on the left-hand side of (7.2) is a Borel function of η .

We actually prove that for every bounded Borel function $a:[0,T)\times\mathbb{R}^d\to\mathbb{R}$ the integral $\int a\,d\eta$ is a Borel function of η . Let indeed \mathscr{F} be the class of all bounded Borel real functions a on $[0,T)\times\mathbb{R}^d$ such that this is true. One easily checks the following:

- (i) \mathcal{F} is a vector space;
- (ii) if f is the pointwise limit of an (increasing) sequence of uniformly bounded functions in \mathscr{F} , then f belongs to \mathscr{F} ;
- (iii) \mathscr{F} contains all functions in $C_0([0,T)\times\mathbb{R}^d)$ and therefore also all bounded continuous functions on $[0,T)\times\mathbb{R}^d$ by statement (ii).

Then the functional version of the monotone class theorem (see for instance [12, Chapter I, Theorem 21]) implies that \mathscr{F} contains all bounded Borel functions.

7.2. Corollary. Let \mathscr{C}^* be the set of all $(y, \eta) \in \mathscr{C}$ with $\eta \neq 0$, and let G be the projection of \mathscr{C}^* on \mathbb{R}^{d-k} . Then G is Lebesgue measurable, and there exist a Borel set $G' \subset G$ such that $\mathscr{L}^{d-k}(G \setminus G') = 0$ and a Borel family $\{\eta_y : y \in G'\}$ such that $(y, \eta_y) \in \mathscr{C}^*$ for every $y \in G'$.

Proof. Since \mathscr{C} is a Borel set in $\mathbb{R}^{d-k} \times \mathscr{M}$, and so is the set \mathscr{I} of all couples (y,0) with $y \in \mathbb{R}^{d-k}$, then $\mathscr{C}^* = \mathscr{C} \setminus \mathscr{I}$ is a Borel set in $\mathbb{R}^{d-k} \times \mathscr{M}$, too.

In particular \mathscr{C}^* belongs to the product σ -algebra generated by the σ -algebra of Lebesgue measurable sets on \mathbb{R}^{d-k} times the Borel σ -algebra on \mathscr{M} . Therefore, by von Neumann's selection theorem (cf. [22, Corollary 5.5.8]), the projection G is Lebesgue measurable, and there exists a Lebesgue measurable map that to every $y \in G$ associates η_y such that $(y, \eta_y) \in \mathscr{C}^*$.

To conclude, we choose a Borel subset G' of G such that the restriction of this selection map to G' is Borel, and $\mathcal{L}^{d-k}(G \setminus G') = 0$.

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References

- Aizenman, M.: On vector fields as generators of flows: a counterexample to Nelson's conjecture. Ann. of Math. 107, 287–296 (1978) Zbl 0394.28012 MR 0482853
- [2] Alberti, G., Bianchini, S., Crippa, G.: Structure of level sets and Sard-type properties of Lipschitz maps. Ann. Scuola Norm. Sup. Pisa 12, 863–902 (2013)
- [3] Alberti, G., Bianchini, S., Crippa, G.: On the L^p -differentiability of certain classes of functions. Rev. Mat. Iberoamer., to appear
- [4] Ambrosio, L.: Transport equation and Cauchy problem for BV vector fields. Invent. Math. 158, 227–260 (2004) Zbl 1075.35087 MR 2096794
- [5] Ambrosio, L., Bouchut, F., De Lellis, C.: Well-posedness for a class of hyperbolic systems of conservation laws in several space dimensions. Comm. Partial Differential Equations 29, 1635–1651 (2004) Zbl 1072.35116 MR 2103848
- [6] Ambrosio, L., Crippa, G.: Existence, uniqueness, stability and differentiability properties of the flow associated to weakly differentiable vector fields. In: Transport Equations and Multi-D Hyperbolic Conservation Laws, Lect. Notes Unione Mat. Ital. 5, Springer, Berlin, 3–57 (2008) Zbl 1155.35313 MR 2409676
- Bouchut, F., Desvillettes, L.: On two-dimensional Hamiltonian transport equations with continuous coefficients. Differential Integral Equations 14, 1015–1024 (2001) Zbl 1028.35042 MR 1827101
- [8] Colombini, F., Crippa, G., Rauch, J.: A note on two-dimensional transport with bounded divergence. Comm. Partial Differential Equations 31, 1109–1115 (2006) Zbl 1122.35143 MR 2254607
- [9] Colombini, F., Luo, T., Rauch, J.: Uniqueness and nonuniqueness for nonsmooth divergencefree transport. In: Séminaire Équations aux Dérivées Partielles 2002–2003, exp. XXII, École Polytech., Palaiseau (2003) Zbl 1065.35089 MR 2030717
- [10] Colombini, F., Rauch, J.: Uniqueness in the Cauchy problem for transport in \mathbb{R}^2 and \mathbb{R}^{1+2} . J. Differential Equations **211**, 162–167 (2005) Zbl 1134.35311 MR 2121112
- [11] Depauw, N.: Nonunicité du transport par un champ de vecteurs presque BV. In: Séminaire Équations aux Dérivées Partielles 2002–2003, exp. XIX, École Polytech., Palaiseau (2003) Zbl 1063.35054 MR 2030714
- [12] Dellacherie, C., Meyer, P.-A.: Probabilities and Potential. North-Holland Math. Stud. 29, North-Holland, Amsterdam (1978) Zbl 0494.60001 MR 0521810
- [13] DiPerna, R. J., Lions, P.-L.: Ordinary differential equations, transport theory and Sobolev spaces. Invent. Math. 98, 511–547 (1989) Zbl 0696.34049 MR 1022305
- [14] DiPerna, R. J., Lions, P.-L.: On the Cauchy problem for Boltzmann equations: global existence and weak stability. Ann. of Math. (2) 130, 321–366 (1989) Zbl 0698.45010 MR 1014927
- [15] DiPerna, R. J., Lions, P.-L.: Global weak solutions of Vlasov–Maxwell systems. Comm. Pure Appl. Math. 42, 729–757 (1989) Zbl 0698.35128 MR 1003433
- [16] Federer, H.: Geometric Measure Theory. Grundlehren Math. Wiss. 153, Springer, New York (1969); Reprinted in the series Classics in Mathematics, Springer, Berlin (1996) Zbl 0176.00801 MR 0257325

[17] Hauray, M.: On two-dimensional Hamiltonian transport equations with $L^p_{\rm loc}$ coefficients. Ann. Inst. H. Poincaré Anal. Non Linéaire **20**, 625–644 (2003) Zbl 1028.35148 MR 1981402

- [18] Hirsch, M. W.: Differential Topology. Grad. Texts in Math. 33, Springer, New York (1976) Zbl 0356.57001 MR 0448362
- [19] Krantz, S. G., Parks, H. R.: Geometric Integration Theory. Birkhäuser Boston, Boston (2008) Zbl 1149.28001 MR 2427002
- [20] Sard, A.: The measure of the critical values of differentiable maps. Bull. Amer. Math. Soc. 48, 883–890 (1942) Zbl 0063.06720 MR 0007523
- [21] Simon, L.: Lectures on Geometric Measure Theory. Proc. Centre Math. Anal. 3, Austral. National Univ., Canberra (1983) Zbl 0546.49019 MR 0756417
- [22] Srivastava, S. M.: A Course on Borel Sets. Grad. Texts in Math. 180, Springer, New York (1998) Zbl 0903.28001 MR 1619545
- [23] Ziemer, W. P.: Weakly Differentiable Functions, Sobolev Spaces and Functions of Bounded Variation. Grad. Texts in Math. 120, Springer, New York (1989) Zbl 0692.46022 MR 1014685